

New Directions: Exploring Charm

1. Null tests and more for e^+e^- and pp colliders
2. U-spin/CP puzzle and pion form factor: a light Z' ?

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We look for New Physics in Flavor Changing Neutral Currents (FCNCs), because they are suppressed in SM by 1. loop-factor 2. CKM-mixing 3. GIM –mass degeneracies of quarks in loop; very strong in Charm FCNCs, where d, s, b are in loop

$$\mathcal{H}_{\text{eff}} = -4 \frac{G_F}{\sqrt{2}} V_{tb} V_{ts}^* \sum_i C_i(\mu) O_i(\mu), \quad G_F/\sqrt{2} = g^2/(8m_W^2)$$

$$O_9 = \bar{s}_L \gamma_\mu b_L \bar{\ell} \gamma^\mu \ell \quad (\text{vector}) \quad O_{10} = \bar{s}_L \gamma_\mu b_L \bar{\ell} \gamma^\mu \gamma_5 \ell \quad (\text{axial-vector})$$

Ongoing precision programs in B -physics. ($b \rightarrow s \ell \ell, \gamma$ global fit, $b \rightarrow d \ell \ell, \gamma$ [2209.04457](#)) \longrightarrow Pursue charm program

Testing the Standard Model with $c \rightarrow u$ FCNCs of mesons and baryons:

- $c \rightarrow u\gamma$ $\text{Br} \sim 10^{-6} - 10^{-4}$
- $c \rightarrow u\mu\mu, uee$ $\text{Br} \sim 10^{-7} - 10^{-6}$
- $c \rightarrow u\nu\bar{\nu}, a, Z', \dots$ $\text{Br} \lesssim 10^{-5}$

Probe various different types of physics (dipole couplings, 4-fermion operators, light NP, ..)

Complementary to kaon and B -physics – charm is unique probe of flavor in the up-sector

[0112235](#), [1510.00965](#), [1805.08516](#), [2011.09478](#),... Adolph, de Boer, Bause, Bharucha, Brod, Bigi, Burdman, d'Ambrosio, Cata, Fajfer, Feldmann, Gisbert, Golowich, Golz, Hewett, Kosnic, Magorsch, Meinel, Pakvasa, Petrov, Seidel, Singer, Silva, Solomonidi, Zwicky, ...

radiative $c \rightarrow u\gamma$: $D \rightarrow V\gamma$, $V = \rho, \dots$, $D \rightarrow P_1P_2\gamma$,

$D \rightarrow A\gamma$, $A = K_1, \dots$, $D \rightarrow P_1P_2P_3\gamma$, $\Lambda_c \rightarrow p\gamma$, $\Xi_c^0 \rightarrow \Lambda(\rightarrow p\pi)\gamma, \dots$

$B(D^0 \rightarrow \rho^0\gamma) = (1.77 \pm 0.31) \cdot 10^{-5}$ Belle'16, Cabibbo-favored modes:

$B(\Lambda_c \rightarrow \Sigma\gamma) < 2.6 \cdot 10^{-4}$, $B(\Xi_c^0 \rightarrow \Xi^0\gamma) < 1.8 \cdot 10^{-4}$ Belle 2206.12517

$B(\Lambda_c \rightarrow \Sigma\gamma) < 4.4 \cdot 10^{-4}$ BESIII 2212.07214

semileptonic $c \rightarrow ull^{(\prime)}$: $D \rightarrow \pi\mu\mu$, $D \rightarrow \mu\mu$, $D \rightarrow P_1P_2\ell\ell$,

$\Lambda_c \rightarrow p\ell\ell$, $\Xi_c^0 \rightarrow \Lambda(\rightarrow p\pi^-)\ell\ell, \dots$ $B(D \rightarrow \pi\pi\mu\mu) \simeq 9.6 \cdot 10^{-7}$ LHCb'18 ,

$B(\Lambda_c \rightarrow p\mu\mu) \lesssim 7.7 \cdot 10^{-8}$ LHCb'17, $B(\Xi_c^0 \rightarrow \Xi^0\mu^+\mu^-) < 6.5 \cdot 10^{-5}$,

$B(\Xi_c^0 \rightarrow \Xi^0e^+e^-) < 9.9 \cdot 10^{-5}$ Belle 2312.02580

dineutrinos/MET/ALPs $c \rightarrow u\nu\bar{\nu}$: $D \rightarrow \pi\nu\bar{\nu}$, $D \rightarrow \nu\bar{\nu}$, $D \rightarrow P_1P_2\nu\bar{\nu}$,

$\Lambda_c \rightarrow p\nu\bar{\nu}$, $\Xi_c^0 \rightarrow \Lambda(\rightarrow p\pi^-)\nu\bar{\nu}, \dots$ $B(D^0 \rightarrow \text{nothing}) < 9.4 \cdot 10^{-5}$

Belle'16, $B(D^0 \rightarrow \pi^0\nu\bar{\nu}) < 2.1 \cdot 10^{-4}$ BESIII 2112.14236

On April 10, 2024 BES III arXiv:2404.05973,

$$B(D_s^+ \rightarrow K^+ \pi^0 e^+ e^-) < 7.1 \cdot 10^{-5}$$

April 23, 2024, Cern seminar "from unseen to rare", search for

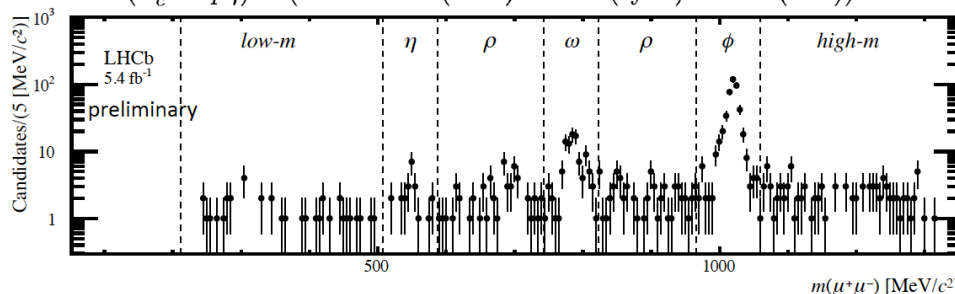
$$\Lambda_c \rightarrow p \mu \mu \text{ at LHCb, LHCb-PAPER-2024-005}$$

- Rare mode, extrapolating to full mass region from low- and high masses:
 - signal region (extrapolated): $B(\Lambda_c^+ \rightarrow p \mu^+ \mu^-) < 7.3 (8.2) \times 10^{-8}$ at 90% (95%) CL.
 - **low-mass:** $B(\Lambda_c^+ \rightarrow p \mu^+ \mu^-) < 6.5 (7.8) \times 10^{-8}$ at 90% (95%) CL
 - **high-mass:** $B(\Lambda_c^+ \rightarrow p \mu^+ \mu^-) < 11.8 (13.1) \times 10^{-8}$ at 90% (95%) CL
- Resonant mode mass regions (assuming no interference):

$$B(\Lambda_c^+ \rightarrow p \omega) = (9.82 \pm 1.23 \text{ (stat.)} \pm 0.72 \text{ (syst.)} \pm 2.79 \text{ (ext.)}) \times 10^{-4}$$

$$B(\Lambda_c^+ \rightarrow p \rho) = (1.52 \pm 0.34 \text{ (stat.)} \pm 0.14 \text{ (syst.)} \pm 0.24 \text{ (ext.)}) \times 10^{-3}$$

$$B(\Lambda_c^+ \rightarrow p \eta) = (1.67 \pm 0.69 \text{ (stat.)} \pm 0.23 \text{ (syst.)} \pm 0.34 \text{ (ext.)}) \times 10^{-3}$$



TH Progress: New BSM strategies for $|\Delta c| = |\Delta u| = 1$

SM tests in rare charm decays are **null tests** based on approximate symmetries of the SM: **GIM, CP, cLFC, LFU, LNC, $SU(3)_F$**

GIM-suppression very efficient: In SM, everything follows from tree-level W -exchange plus RGE, μ_b -matching 1707.00988:

$$O_7 = \bar{u}_L \sigma_{\mu\nu} c_R F^{\mu\nu}, \quad O_9 = \bar{u}_L \gamma_\mu c_L \bar{\ell} \gamma^\mu \ell,$$

$$O_{10} = \bar{u}_L \gamma_\mu c_L \bar{\ell} \gamma^\mu \gamma_5 \ell, \quad O_\nu = \bar{u}_L \gamma_\mu c_L \bar{\nu} \gamma^\mu (1 - \gamma_5) \nu$$

$$C_7^{\text{eff SM}} \lesssim 0.01, \quad C_9^{\text{eff SM}} \lesssim 0.05 \quad \text{1510.00311} \quad C_\nu^{\text{SM}} = 0, \quad C_{10}^{\text{SM}} = 0$$

SM-Phenomenology completely dominated by 4-quark operators;

Classify as CF $\sim V_{cs}^* V_{ud}$, **SCS** $\sim V_{cq}^* V_{uq}$, $V_{cq}^* V_{uq}$, $q = d, s$, DCS $\sim V_{cd}^* V_{us}$

$c \rightarrow u\nu\bar{\nu}$ transitions: all of them are excellent nulltests of the SM due to GIM, very clean

$D^+, D_s \rightarrow M\nu\bar{\nu}$ has BGD from $D^+, D_s \rightarrow \tau(\rightarrow M\nu)\bar{\nu}$; reducible via cuts

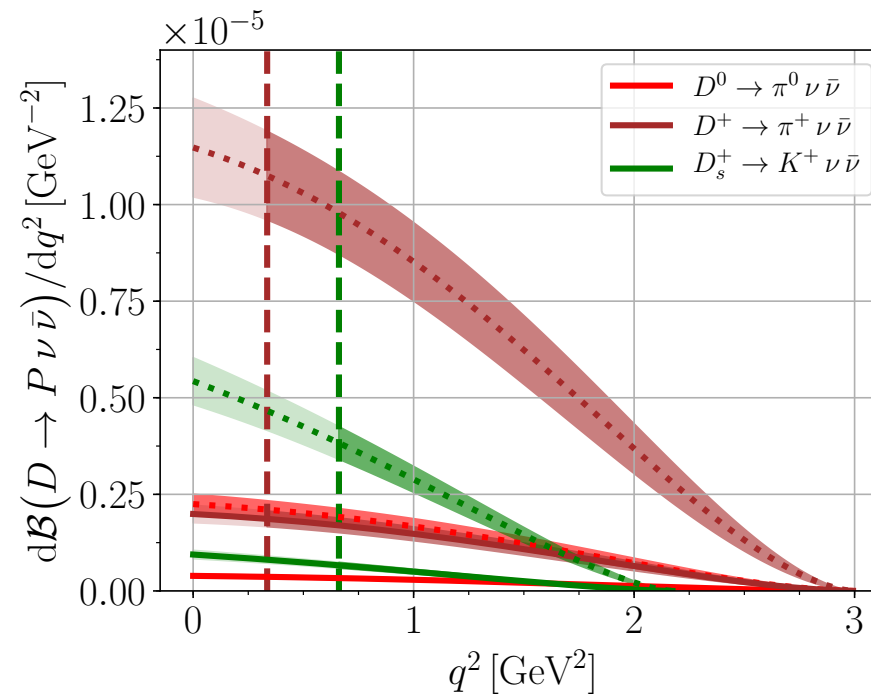


Figure 1: Differential branching ratios for $D^0 \rightarrow \pi^0 \nu \bar{\nu}$, $D^+ \rightarrow \pi^+ \nu \bar{\nu}$ and $D_s^+ \rightarrow K^+ \nu \bar{\nu}$ in red, brown and green for the LU (cLFC) limit in solid (dotted) lines. **plot shows BSM distributions** The uncertainty bands are due to the form factors, here from lattice ETM'19, **τ background removable by q^2 -cut (vertical dashed lines)** [plot from 2010.02225](#)

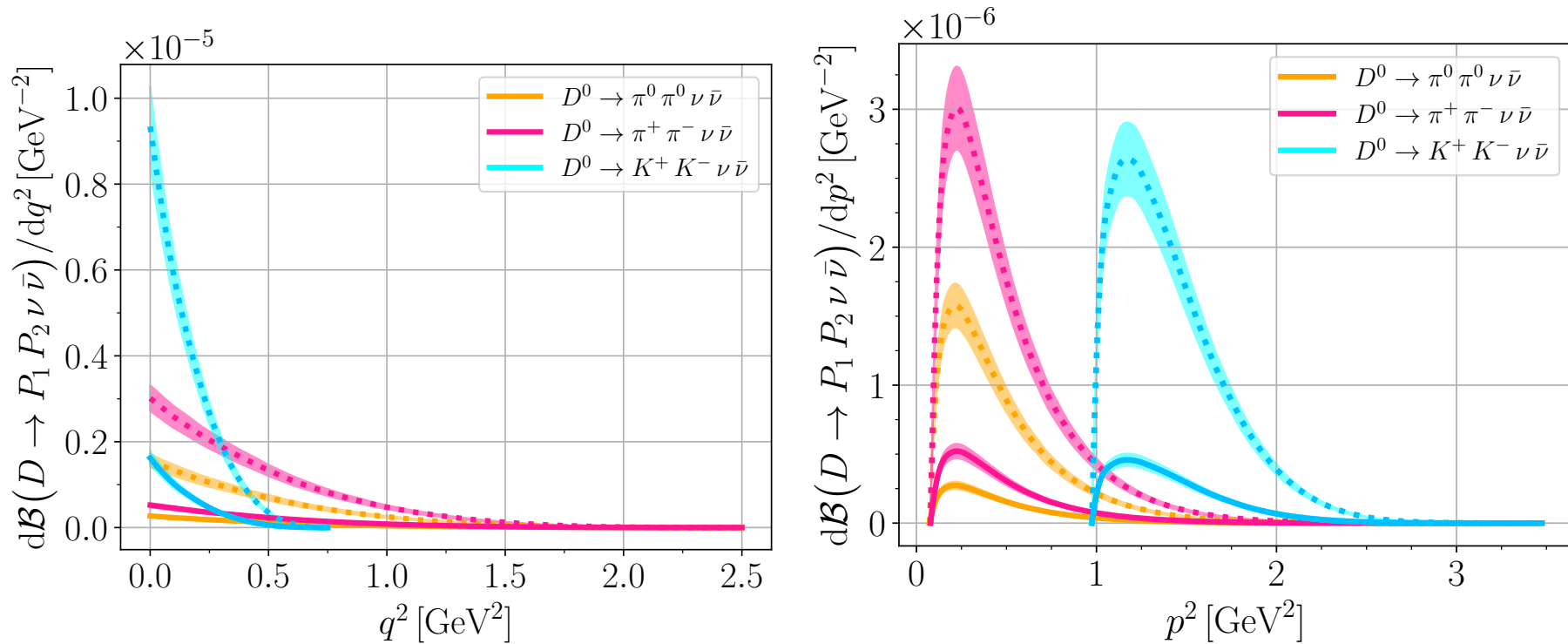


Figure 2: Differential q^2 and p^2 branching ratios for $D^0 \rightarrow \pi^0 \pi^0 \nu \bar{\nu}$, $D^0 \rightarrow \pi^+ \pi^- \nu \bar{\nu}$ and $D^0 \rightarrow K^+ K^- \nu \bar{\nu}$ decays in orange, deep pink and cyan, respectively for the LU (cLFC) limit in solid (dotted) lines. The differential branching ratio of $D^0 \rightarrow K^+ K^- \nu \bar{\nu}$ is multiplied by a factor 100 to be visible in the plots. **Can we do better with $D \rightarrow \pi\pi$ form factors?** here from HH χ PT Lee et al '92. with illustrative 10 % uncertainty. [plots from 2010.02225](#)

$D \rightarrow PP\nu\bar{\nu}$ predom. $C - C'$, complementary to $D \rightarrow P\nu\bar{\nu}$: $C + C'$

Leading semileptonic 4-fermion operators at scale above v (SMEFT) contributing to dineutrino modes $q \rightarrow q' \nu \bar{\nu}$

$$\mathcal{L}_{\text{eff}} \supset \frac{C_{\ell q}^{(1)}}{v^2} \bar{Q} \gamma_\mu Q \bar{L} \gamma^\mu L + \frac{C_{\ell q}^{(3)}}{v^2} \bar{Q} \gamma_\mu \tau^a Q \bar{L} \gamma^\mu \tau^a L + \frac{C_{\ell u}}{v^2} \bar{U} \gamma_\mu U \bar{L} \gamma^\mu L + \frac{C_{\ell d}}{v^2} \bar{D} \gamma_\mu D \bar{L} \gamma^\mu L. \quad (1)$$

$SU(2)_L \times U(1)_Y$ gauge invariance links up and down quarks, $Q = (u, d)$ and left-handed neutrinos and charged leptons $L = (\nu, \ell)$.

$$C_L^U = K_L^D = C_{\ell q}^{(1)} + C_{\ell q}^{(3)}, \quad C_R^U = K_R^U = C_{\ell u}, \quad C_L^D = K_L^U = C_{\ell q}^{(1)} - C_{\ell q}^{(3)}, \quad C_R^D = K_R^D = C_{\ell d}.$$

LH contribution to $c \rightarrow u \nu \bar{\nu}$ (C_L^U) identical to $s \rightarrow d \ell \bar{\ell}$ (K_L^D) etc

RH contribution to $c \rightarrow u \nu \bar{\nu}$ (C_R^U) identical to $c \rightarrow u \ell \bar{\ell}$ (K_R^U)

L,R denotes left or right handed quark currents; only SM-like light neutrinos.

Dineutrino vs lepton specific limits

left-handed $c \rightarrow u\nu\bar{\nu}$ (C_L^U)

Process	WC	ee	$\mu\mu$	$\tau\tau$	$e\mu$	$e\tau$	$\mu\tau$
$pp \rightarrow \ell^+\ell^-$	$\mathcal{K}_{L,R}^{sd\ell\ell'}$	3.8	2.3	5.37	2.0	6.1	6.6
$K \rightarrow \ell^+\ell^- + \pi$	$\mathcal{K}_{L,R}^{sd\ell\ell'} \cdot 10^2$	5	1.6	-	0.066	-	-
$K \rightarrow \nu\bar{\nu} + \pi$	$\mathcal{K}_R^{sd\ell\ell'} \cdot 10^2$	[-1.9,0.7]	[-1.9,0.7]	[-1.9,0.7]	1.1	1.1	1.1
prelim. $pp \rightarrow \nu\bar{\nu} + X$	$\mathcal{K}_L^{sd\ell\ell'}$	4.1 (3.3)	4.1 (3.3)	4.1 (3.3)	2.9 (2.2)	2.9 (2.2)	2.9 (2.2)
prelim. $pp \rightarrow \nu\bar{\nu} + X$	$\mathcal{K}_R^{sd\ell\ell'}$	5.7 (4.6)	5.7 (4.6)	5.7 (4.6)	4.1 (3.3)	4.1 (3.3)	4.1 (3.3)

Table 1: Limits on $sd\ell\ell'$ WCs. The first row is based on 2304.12837, while the second and third row are taken from 2007.05001 and the **NEW** ones from $pp \rightarrow \nu\bar{\nu} + X$ are from 2403.17063 (recasting ATLAS 2102.10874) with projections for $3000 fb^{-1}$ in parentheses, using $SU(2)_L$.

- synergy between high p_T flavor and rare decays
- rare kaon decays are most powerful when existing (no decays into taus)
- **prelim.** $pp \rightarrow MET + j$ similar level as Drell-Yan and better for taus

Dineutrino vs lepton specific limits

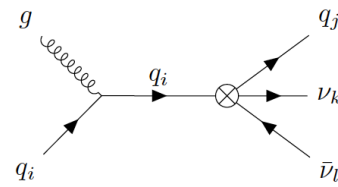
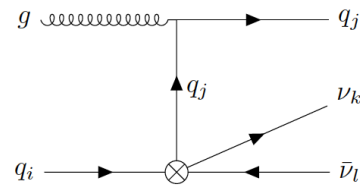
right-handed $c \rightarrow u\nu\bar{\nu}$ (C_R^U)

Process	WC	ee	$\mu\mu$	$\tau\tau$	$e\mu$	$e\tau$	
$pp \rightarrow \ell^+\ell^-$	$\mathcal{K}_{L,R}^{cull'}$	2.9	1.6	5.6	1.6	4.7	
$D \rightarrow \ell^+\ell^- + \pi$	$\mathcal{K}_{L,R}^{cull'}$	4.0	0.9	-	2.2	-	
$K \rightarrow \nu\bar{\nu} + \pi$	$\mathcal{K}_L^{cull'} \cdot 10^2$	$[-1.9, 0.7]$	$[-1.9, 0.7]$	$[-1.9, 0.7]$	1.1	1.1	
prelim. $pp \rightarrow \nu\bar{\nu} + X$	$\mathcal{K}_L^{cull'}$	5.7 (4.6)	5.7 (4.6)	5.7 (4.6)	4.1 (3.3)	4.1 (3.3)	4.1 (3.3)
prelim. $pp \rightarrow \nu\bar{\nu} + X$	$\mathcal{K}_R^{cull'}$	4.1 (3.3)	4.1 (3.3)	4.1 (3.3)	2.9 (2.2)	2.9 (2.2)	2.9 (2.2)

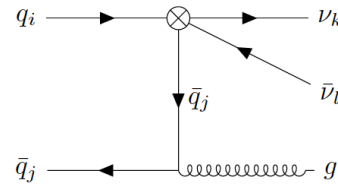
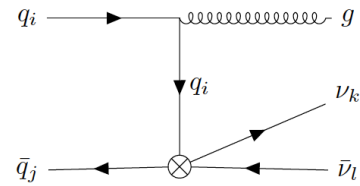
Table 2: Limits on $cull'$ WCs. The first three rows are from 2007.05001, while the ones from $pp \rightarrow \nu\bar{\nu} + X$ are from 2403.17063

- best bound on $\mathcal{K}_R^{cu\mu\mu}$ for dimuons from LHCb

Dineutrino vs lepton specific limits



(b)



(d)

energy enhanced by \hat{s}/Λ_{NP}^2 relative to SM p_T -distribution,
analytically shown in 2403.17063 , also for gluon dipole operators

Upper limits $\mathcal{B}^{\max}(h_c \rightarrow F\nu\bar{\nu})$ depend on lepton flavor structure (LFV,cLFCLFU) 2010.02225

$h_c \rightarrow F$	$\mathcal{B}_{\text{LU}}^{\max}$ [10^{-7}]	$\mathcal{B}_{\text{cLFC}}^{\max}$ [10^{-6}]	\mathcal{B}^{\max} [10^{-6}]	$N_{\text{LU}}^{\max}/\eta_{\text{eff}}$	$N_{\text{cLFC}}^{\max}/\eta_{\text{eff}}$	$N^{\max}/\eta_{\text{eff}}$
$D^0 \rightarrow \pi^0$	6.1	3.5	13	47 k (395 k)	270 k (2.3 M)	980 k (8.3 M)
$D^+ \rightarrow \pi^+$	25	14	52	77 k (650 k)	440 k (3.7 M)	1.6 M (14 M)
$D_s^+ \rightarrow K^+$	4.6	2.6	9.6	6 k (50 k)	34 k (290 k)	120 k (1.1 M)
$D^0 \rightarrow \pi^0\pi^0$	1.5	0.8	3.1	11 k (95 k)	64 k (540 k)	230 k (2.0 M)
$D^0 \rightarrow \pi^+\pi^-$	2.8	1.6	5.9	22 k (180 k)	120 k (1.0 M)	450 k (3.8 M)
$D^0 \rightarrow K^+K^-$	0.03	0.02	0.06	0.2 k (1.9 k)	1.3 k (11 k)	4.8 k (40 k)
$\Lambda_c^+ \rightarrow p^+$	18	11	39	14 k (120 k)	82 k (700 k)	300 k (2.6 M)
$\Xi_c^+ \rightarrow \Sigma^+$	36	21	76	28 k (240 k)	160 k (1.4 M)	590 k (5.0 M)

to date only a single limit exists: $B(D^0 \rightarrow \pi^0\nu\bar{\nu}) < 2.1 \cdot 10^{-4}$ BESIII 2112.14236

ULs for Belle II w 50 ab^{-1} (FCC-ee@Z w $N(c\bar{c}) = 550 \cdot 10^9$) via SMEFT and dilepton data 2010.02225

Rare decays of charm baryons $\Lambda_c, \Xi_c, \Omega_c$

theory and observables: 2107.13010, 2202.02331 **highlights for BSM searches: GIM ($C_{10}^{\text{SM}} = 0$), angular distributions, CP, cLFV, LFU**

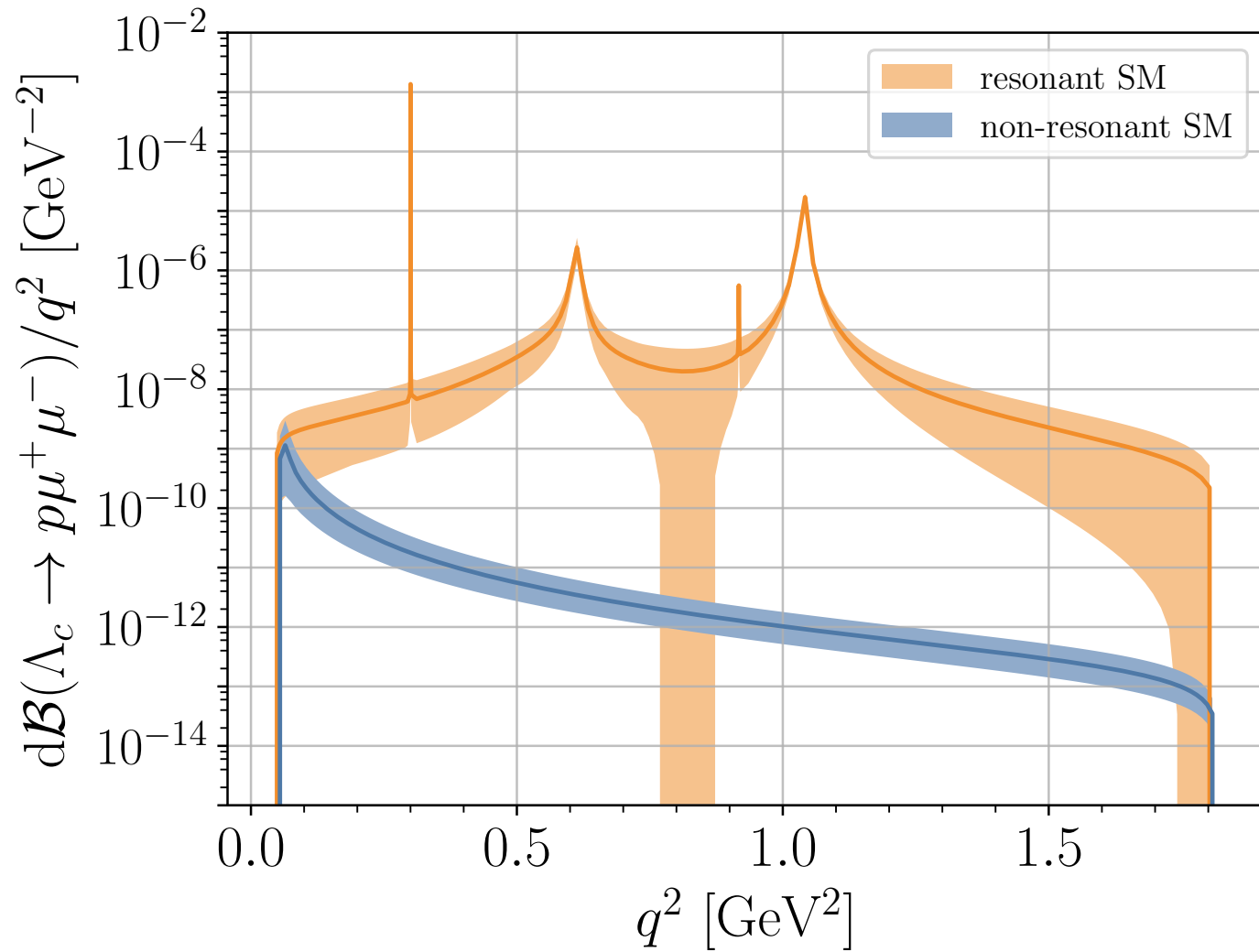
Differential angular distribution for unpolarized $\Lambda_c \rightarrow p\mu\mu$, (polarized Λ_c in 2202.02331) reads:

$$\frac{d^2\Gamma}{dq^2 d\cos\vartheta_\ell} = \frac{3}{2} (K_{1ss} \sin^2\vartheta_\ell + K_{1cc} \cos^2\vartheta_\ell + K_{1c} \cos\vartheta_\ell)$$

→ 3 observables: branching ratio (−), longitudinal pol. fraction F_L (+), Forward-Backward asymmetry $A_{\text{FB}}^\ell \propto K_{1c} \propto C_{10}$. (++)

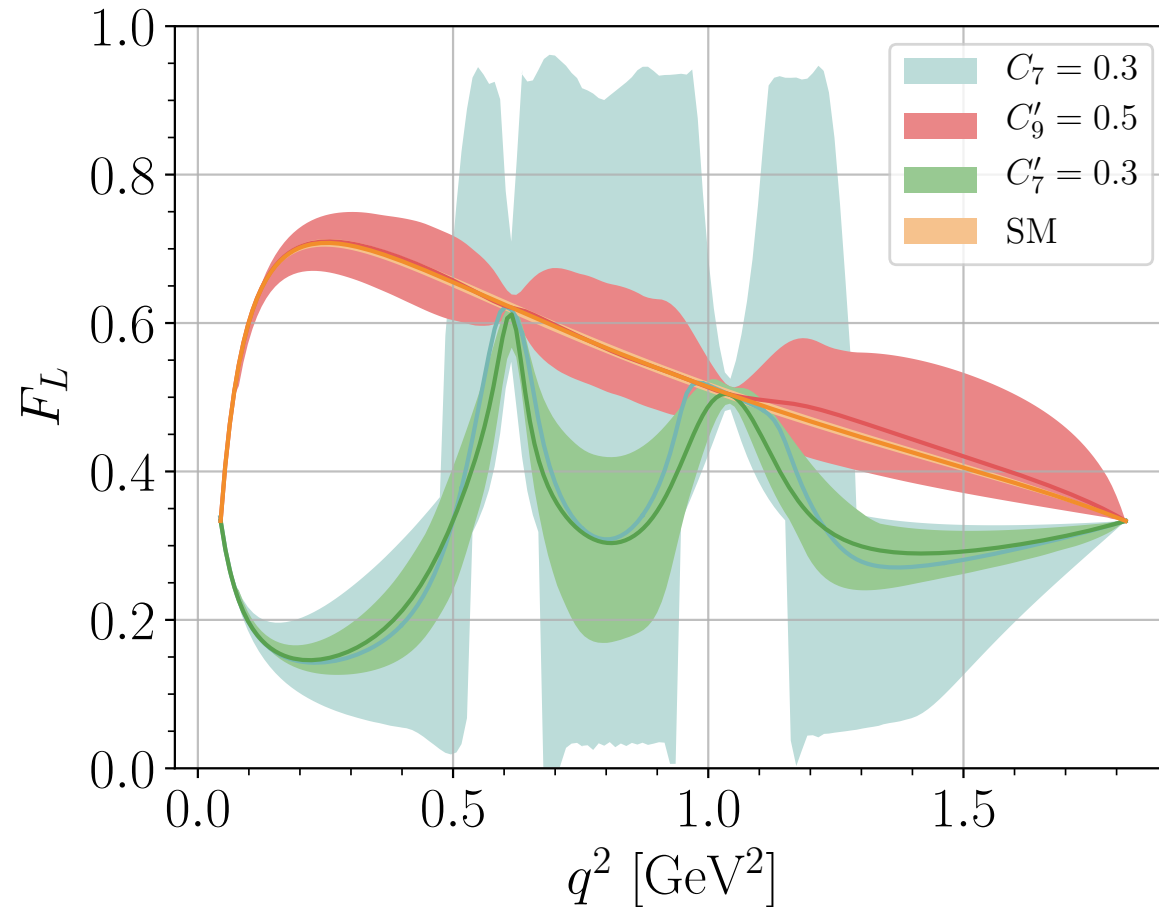
$\Lambda_c \rightarrow p$ form factors from lattice 1712.05783 – $SU(3)_F$ -relations to other baryonic decays 2203.14982

$$-\sqrt{6}h_{\perp}^{\Xi_c^0 \rightarrow \Lambda} = \sqrt{2}h_{\perp}^{\Xi_c^0 \rightarrow \Sigma^0} = h_{\perp}^{\Xi_c^+ \rightarrow \Sigma^+} = h_{\perp}^{\Lambda_c \rightarrow p}; \text{Endpoint relations (at } q^2 = \text{max): 2107.12993}$$



2107.13010

Longitudinal polarization: (+)



2107.13010

Sensitivity to dipole coefficients! $F_L(4m_\ell^2) = F_L(q_{max}^2) = 1/3$

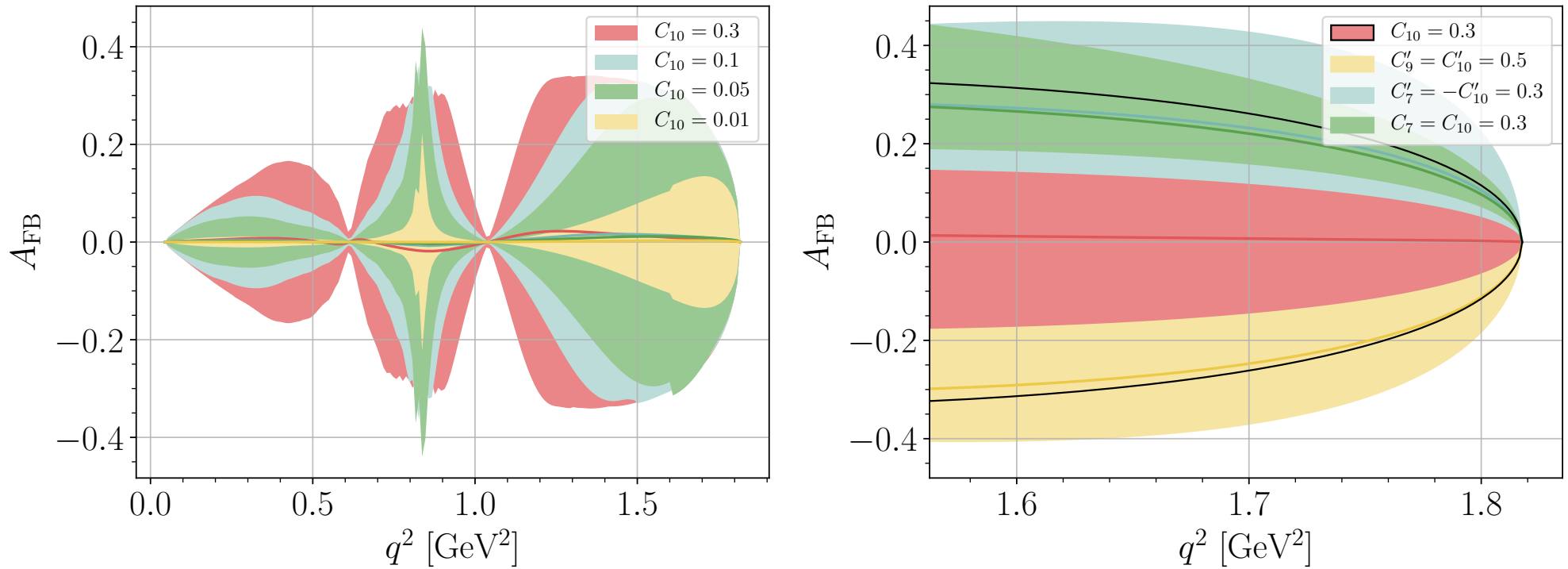


Figure 3: The forward-backward asymmetry A_{FB} of $\Lambda_c \rightarrow p \mu^+ \mu^-$ decays for different values of C_{10} in the full q^2 -region (left panel) and for various BSM contributions in the high q^2 region (right panel)

$A_{\text{FB}} \propto C_{10}$ clean null test of SM (GIM); Three more GIM-based null tests in 4-body decays

$\Xi_c^+ \rightarrow \Sigma^+ (\rightarrow p \pi^0) \ell^+ \ell^-$, $\Xi_c^0 \rightarrow \Lambda^0 (\rightarrow p \pi^-) \ell^+ \ell^-$, $\Omega_c^0 \rightarrow \Xi^0 (\rightarrow \Lambda^0 \pi^0) \ell^+ \ell^-$, 2202.02331

Full $D \rightarrow P_1 P_2 l^+ l^-$ angular distribution

Learn, e.g., from B -physics literature [1406.6681](#), earlier works in charm [1209.4235](#)

$$\frac{d^5 \Gamma(D \rightarrow P_1 P_2 l^+ l^-)}{dq^2 dp^2 d \cos \vartheta_{P_1} d \cos \vartheta_l d \varphi} = \frac{1}{2\pi} \left[\sum_i \underbrace{c_i(\vartheta_l, \varphi)}_{\text{known}} \underbrace{I_i(q^2, p^2, \cos \vartheta_{P_1})}_{\text{SM, BSM}} \right]$$

$$c_1 = 1, \quad c_2 = \cos 2\vartheta_l, \quad c_3 = \sin^2 \vartheta_l \cos 2\varphi, \quad c_4 = \sin 2\vartheta_l \cos \varphi, \quad c_5 = \sin \vartheta_l \cos \varphi, \quad c_6 = \cos \vartheta_l, \\ c_7 = \sin \vartheta_l \sin \varphi, \quad c_8 = \sin 2\vartheta_l \sin \varphi, \quad c_9 = \sin^2 \vartheta_l \sin 2\varphi.$$

I_i : angular observables; contain SM and possibly BSM contributions.

branching ratio

$$\frac{d^3 \Gamma}{dq^2 dp^2 d \cos \vartheta_{P_1}} = 2 \left(I_1 - \frac{I_2}{3} \right). \quad (2)$$

Angular distributions, such as forward-backward asymmetry in the leptons, $A_{\text{FB}} \propto I_6$

$$I_6 = \frac{1}{2} \left[\int_0^1 d \cos \vartheta_l - \int_{-1}^0 d \cos \vartheta_l \right] \frac{d^4 \Gamma}{dq^2 dp^2 d \cos \vartheta_{P_1} d \cos \vartheta_l}. \quad (3)$$

$$I_7 = \left[\int_0^\pi d\varphi - \int_\pi^{2\pi} d\varphi \right] \frac{d^4 \Gamma}{dq^2 dp^2 d \cos \vartheta_{P_1} d\varphi}, \quad (4)$$

$$I_5 = \left[\int_{-\pi/2}^{\pi/2} d\varphi - \int_{\pi/2}^{3\pi/2} d\varphi \right] \frac{d^4 \Gamma}{dq^2 dp^2 d \cos \vartheta_{P_1} d\varphi}, \quad (5)$$

$$I_8 = \frac{3\pi}{8} \left[\int_0^\pi d\varphi - \int_\pi^{2\pi} d\varphi \right] \left[\int_0^1 d \cos \vartheta_l - \int_{-1}^0 d \cos \vartheta_l \right] \frac{d^5 \Gamma}{dq^2 dp^2 d \cos \vartheta_{P_1} d \cos \vartheta_l d\varphi}, \quad (6)$$

$$I_9 = \frac{3\pi}{8} \left[\int_0^{\pi/2} d\varphi - \int_{\pi/2}^\pi d\varphi + \int_\pi^{3\pi/2} d\varphi - \int_{3\pi/2}^{2\pi} d\varphi \right] \frac{d^4 \Gamma}{dq^2 dp^2 d \cos \vartheta_{P_1} d\varphi}. \quad (7)$$

Full $D \rightarrow P_1 P_2 l^+ l^-$ angular distribution

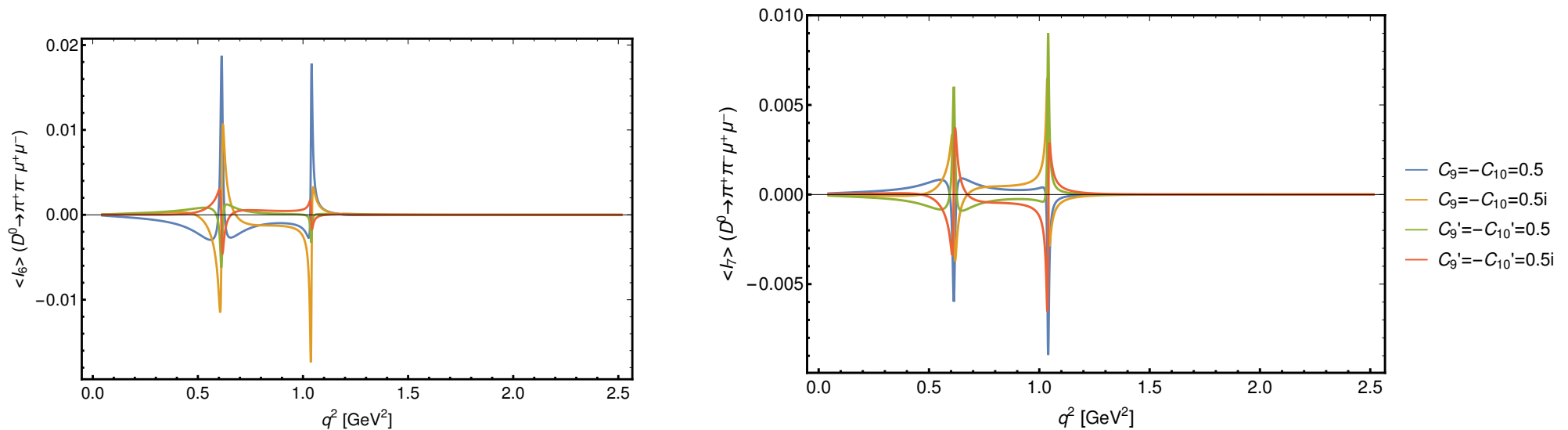
L, R : lepton current handedness, H_k : transversity amplitudes

$$\begin{aligned}
 I_1 &= \frac{1}{16} \left[|H_0^L|^2 + (L \rightarrow R) + \frac{3}{2} \sin^2 \vartheta_{P_1} \{ |H_\perp^L|^2 + |H_\parallel^L|^2 + (L \rightarrow R) \} \right], \\
 I_2 &= -\frac{1}{16} \left[|H_0^L|^2 + (L \rightarrow R) - \frac{1}{2} \sin^2 \vartheta_{P_1} \{ |H_\perp^L|^2 + |H_\parallel^L|^2 + (L \rightarrow R) \} \right], \\
 I_3 &= \frac{1}{16} \left[|H_\perp^L|^2 - |H_\parallel^L|^2 + (L \rightarrow R) \right] \sin^2 \vartheta_{P_1}, \\
 I_4 &= -\frac{1}{8} \left[\text{Re}(H_0^L H_\parallel^{L*}) + (L \rightarrow R) \right] \sin \vartheta_{P_1}, \\
 I_5 &= -\frac{1}{4} \left[\text{Re}(H_0^L H_\perp^{L*}) - (L \rightarrow R) \right] \sin \vartheta_{P_1}, \\
 I_6 &= \frac{1}{4} \left[\text{Re}(H_\parallel^L H_\perp^{L*}) - (L \rightarrow R) \right] \sin^2 \vartheta_{P_1}, \\
 I_7 &= -\frac{1}{4} \left[\text{Im}(H_0^L H_\parallel^{L*}) - (L \rightarrow R) \right] \sin \vartheta_{P_1}, \\
 I_8 &= -\frac{1}{8} \left[\text{Im}(H_0^L H_\perp^{L*}) + (L \rightarrow R) \right] \sin \vartheta_{P_1}, \\
 I_9 &= \frac{1}{8} \left[\text{Im}(H_\parallel^{L*} H_\perp^L) + (L \rightarrow R) \right] \sin^2 \vartheta_{P_1}.
 \end{aligned} \tag{8}$$

$I_{5,6,7}$ vanish due to minus signs (red) in absence of axial vector couplings.

Full $D \rightarrow P_1 P_2 l^+ l^-$ angular distribution

In charm, due to GIM, dynamics dominated by $SU(3)_C \times U(1)_{em}$: all vector-like: $I_{5,6,7}^{SM} = 0$ (proportional to $C_{10 SM}^{(l)} \lesssim 10^{-3} - 10^{-4}$) 1805.08516



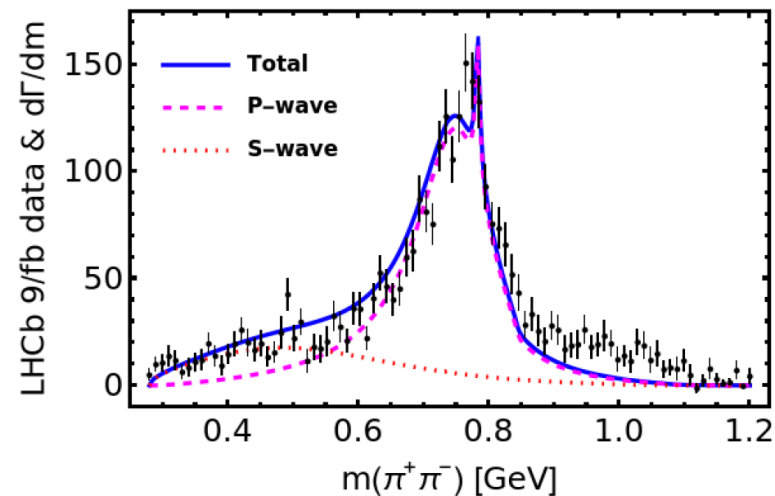
Largest BSM effects from interference with SM; peaks at ρ/ω and Φ .

Model-independent BSM effects up to few %.

Angular asymmetries already measured by LHCb 2111.03327, Null tests close to BSM sensitivity

Fit to spectra from 2312.07501, S-wave resonance contribution

$\sigma \rightarrow \pi\pi$



A Puzzle in hadronic charm CPX/UX

$\Delta A_{CP} = A_{CP}(D \rightarrow K^+ K^-) - A_{CP}(D \rightarrow \pi^+ \pi^-)$ from $A_{CP}(\pi^+ \pi^-)$?

CP & U-Spin puzzle [2207.08539](#), [2210.16330](#) two approx symmetries challenged

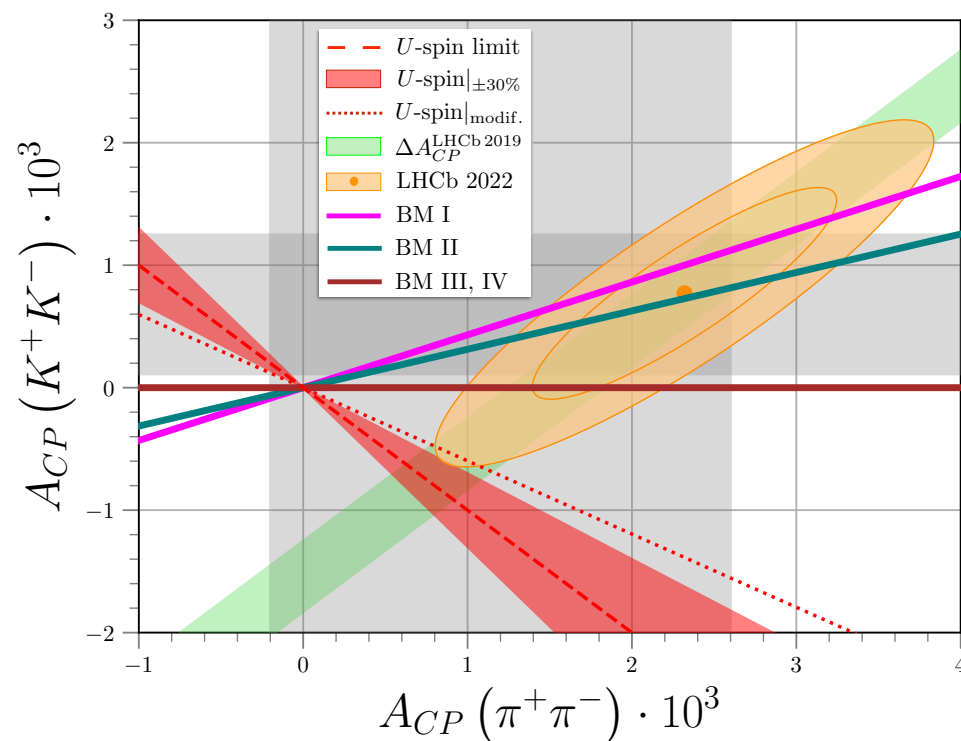


Fig from [2210.16330](#), LHCb result (orange ellipse) from [2209.03179](#) SM theory recently [2305.11951](#), [2312.13245](#)

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CP violation suppressed in charm
 t predominantly tree, h : higher order: loops, FSI, rescattering

$$A = A(D^0 \rightarrow \pi^+ \pi^-)^{\text{SM}} = \Sigma t + V_{cb}^* V_{ub} h e^{i\delta},$$
$$\bar{A} = A(\bar{D}^0 \rightarrow \pi^+ \pi^-)^{\text{SM}} = \Sigma^* t + V_{cb} V_{ub}^* h e^{i\delta},$$

where $\Sigma = (V_{cd}^* V_{ud} - V_{cs}^* V_{us})/2$ and $V_{cs}^* V_{us} \simeq -V_{cd}^* V_{ud} \gg V_{cb}^* V_{ub}$.

$$A_{\text{CP SM}}^{\pi^- \pi^+} = \frac{|A|^2 - |\bar{A}|^2}{|A|^2 + |\bar{A}|^2} \simeq 2 \cdot \text{Im}\left(\frac{V_{cb}^* V_{ub}}{V_{cd}^* V_{ud}}\right) \frac{h}{t} \sin \delta \simeq 1.2 \cdot 10^{-3} \frac{h}{t} \sin \delta$$

Data $A_{\text{CP}}^{\pi^- \pi^+} = (23.2 \pm 6.1) \cdot 10^{-4}$ implies $h/t \sin \delta \sim 2$

SM: FSI-interactions (Pich, Solomonidi, Silva) 2305.11951: $h/t \sin \delta \lesssim 0.2$, however 2203.04056 (Bediaga et al) , 1706.07780

(Khodjamirian, Petrov) $h/t \sin \delta \lesssim 0.1$. LCSR (Lenz, Piscopo et al) 2312.13245 $h/t \sin \delta \lesssim 0.1$.

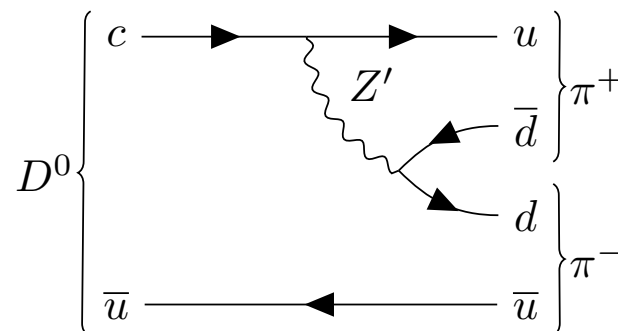
We need the U-spin breaking on top.

A light Z' in hadronic charm CPX/UX

Is large $A_{\text{CP}}^{\pi^- \pi^+} = (23.2 \pm 6.1) \cdot 10^{-4}$ and $A_{\text{CP}}^{K^- K^+}$ SM-like even explainable?

Single solution known [2210.16330](#) break U-spin explicitly, couple to charm, only $SU(2)_L$ -singlets to avoid Kaons, anomaly-free $SU(3)_C \times SU(2)_L \times U(1)_Y \times U(1)'$

Model	F_{Q_i}			F_{u_i}			F_{d_i}			F_{L_i}			F_{e_i}			F_{ν_i}		
BM I	0	0	0	9	-16	7	20	-11	-9	15	-6	-9	-16	0	16	6	12	-18
BM II	0	0	0	-19	9	10	20	-8	-12	4	1	-5	15	2	-17	8	2	-10
BM III	0	0	0	G	$-F$	0	F	$-G$	0	0	0	0	0	$-G$	F	0	G	$-F$
BM IV	0	0	0	$-F_u$	F_u	0	F_d	0	$-F_d$	0	0	0	F_e	0	$-F_e$	F_ν	$-F_\nu$	0



A light Z' in hadronic charm CPX/UX

To avoid the constraints from $D - \bar{D}$ -mixing, the $\bar{u}_R \gamma_\mu c_R$ coupling has to be small. To still get the effect in 4-quark operators $\sim \bar{u}_R \gamma_\mu c_R \bar{d}_R \gamma^\mu d_R$ the Z' has to be light.

Sub 20 GeV Z' (CMS ISR constraints), leptophobic (LHCb $A \rightarrow \mu\mu$ search), BM I,II dead, BM III $F \gg G$, BM IV $F_e \ll F_d$

Model	F_{Q_i}			F_{u_i}			F_{d_i}			F_{L_i}			F_{e_i}			F_{ν_i}		
BM III	0	0	0	G	$-F$	0	F	$-G$	0	0	0	0	0	$-G$	F	0	G	$-F$
BM IV	0	0	0	$-F_u$	F_u	0	F_d	0	$-F_d$	0	0	0	F_e	0	$-F_e$	F_ν	$-F_\nu$	0

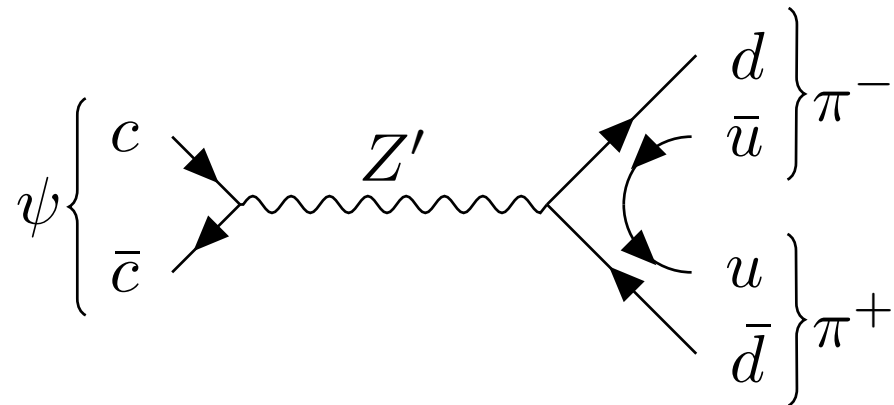
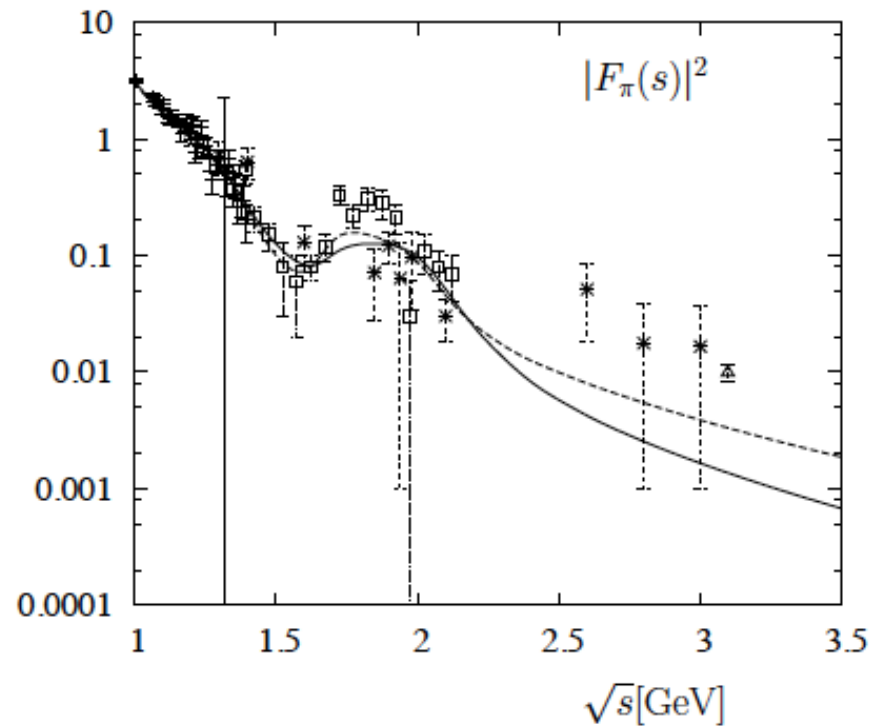
Signatures in low mass dijets, $J/\Psi/\Psi'$, Υ decays,
 $A_{CP}(D \rightarrow \pi^0 \pi^0), A_{CP}(D \rightarrow \pi^+ \pi^0) \sim A_{CP}(D \rightarrow \pi^+ \pi^-)$.

A light Z' in hadronic charm CPX/UX

Model	light quarks	b	c	e	μ	τ	ν_R
BM III $ _{M_{Z'}=2.5 \text{ GeV}}$	75	0	0	0	0	0	25
BM III $ _{M_{Z'}=15 \text{ GeV}}$	38	0	37	0	0	12	13
BM III-s $ _{M_{Z'}=2.5 \text{ GeV}}$	86	0	0	0	0	0	14
BM III-s $ _{M_{Z'}=15 \text{ GeV}}$	75	0	0	0	0	12	13
BM IV $ _{M_{Z'}=5 \text{ GeV}}$	79	0	21	0	0	0	0
BM IV $ _{M_{Z'}=15 \text{ GeV}}$	54	28	18	0	0	0	0

Table 3: Tree-level branching fractions in % for the different Z' decay modes to fermion-antifermion pairs. Results for BM III and BM III-s are given in the limit $G \ll F$. In BM IV, branching ratios depend on the different charge assignments $F_{u,d,e,\nu}$, see main text for details. The branching ratios shown in this table are obtained from $F_u = 985$, $F_d = 1393$, $F_e = 1$ and $F_\nu = 0$. Branching ratios in all BMs differ perceptibly between the low and high $M_{Z'}$ windows, as the decays $Z' \rightarrow b\bar{b}$, $c\bar{c}$, $\tau^+\tau^-$ are kinematically forbidden or suppressed in the few GeV range. Corrections to branching ratios from kinetic mixing are generically $\lesssim 10^{-7}$.

What else can the hadrophilic light Z' do for you?



plot from hep-ph/0409080, discrepancy of pion form factor (from $ee \rightarrow \pi\pi$ -scattering plus theory) with $Br(J/\psi \rightarrow \pi^+\pi^-)$ (triangle).

Simultaneous explanation with charm CP data in BM III for $M_{Z'} \sim 3$ GeV, BM IV for $M_{Z'} \sim (5 - 6)$ GeV [2210.16330](#)

- Rare charm decays can test the SM and signal new physics using null tests of the SM
- Very little experimentally explored in rare charm decays – lots of blanks in PDG, but plenty of recent progress.
- BSM effects in $|\Delta c| = |\Delta u| = 1$ can be large.
- U-spin/CP puzzle hints at combined breaking of SM symmetries. Possible connection with other low energy puzzles.
- Complementary search to K, B -decays.
- Plenty of Terra Incognita in flavor phenomenology — explore synergies

”charm is the new beauty”

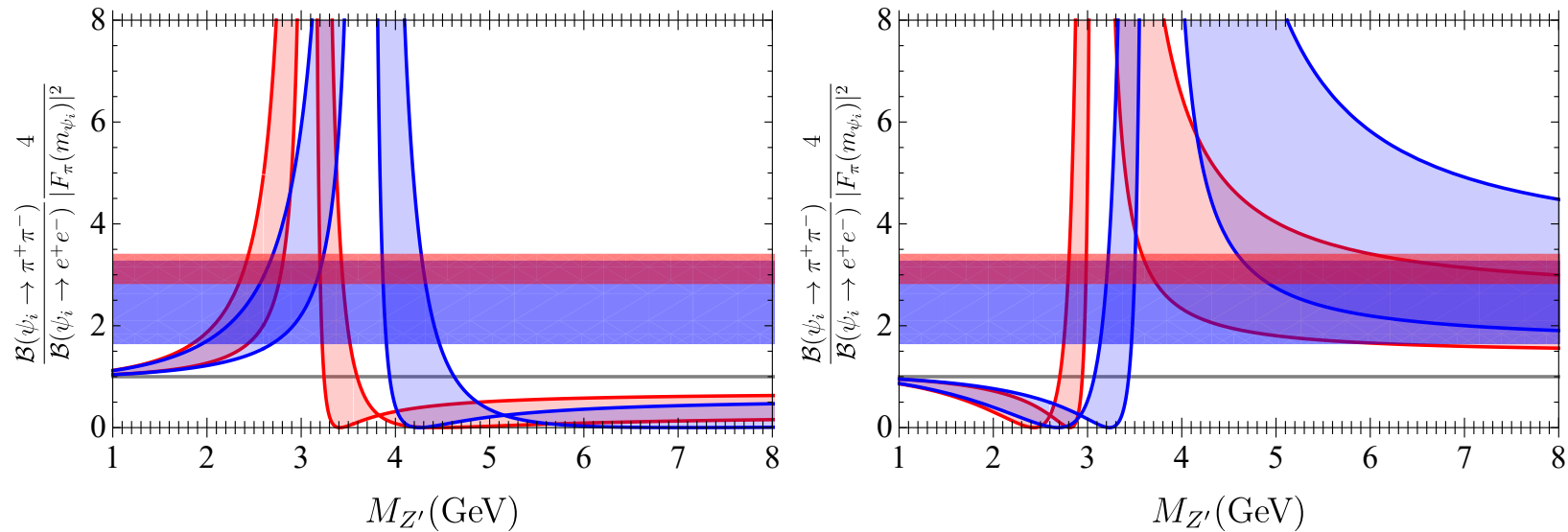


Figure 4: Constraints from charmonium decays. Horizontal red (blue) bands denote the left-hand side of (??) from 1 sigma ranges of J/ψ -data with $|F_\pi(m_{J/\psi})| = 0.056$ (ψ' decays with $|F_\pi(m_{\psi'})| = 0.04$). Curves correspond to the predictions of BM III with $F \gg G$ (left) and BM IV (right). The SM prediction via photon-exchange is shown by the grey line.