

Universality of energy-momentum response in kinetic theories

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XD, S. Ochsenfeld, S. Schlichting, PLB(2023)138161, arXiv:2306.09094

Introduction

Scales and emergence in nature

- A flock of birds shows (possibly universal) collective behavior



Starling



Sparrow

Microscopic



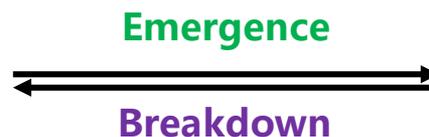
Flock of birds

Macroscopic

Scales and emergence in quantum many-body physics

- Any microscopic theory is expected to have its hydrodynamic description at macroscopic level

Quantum field theory



Hydrodynamics

- Universality in emergence? Breakdown of macroscopic theory?

Hydrodynamics

Hydrodynamic theory

- Macroscopic theory at long wave-length, low frequency limit

$$kl_{mfp} < 1, \quad \omega t_{mfp} < 1$$

- Modern construction as an order expansion of gradients close to equilibrium

$$\langle T^{\mu\nu} \rangle = \underbrace{T_{(0)}^{\mu\nu}}_{\text{Ideal hydro}} + \underbrace{T_{(1)}^{\mu\nu}(\nabla) + T_{(2)}^{\mu\nu}(\nabla^2) + \dots}_{\text{Viscous hydro}}$$

Hydrodynamic modes

small gradients

$\nabla \sim k$ small wave number

vanishing frequency

$\omega(k \rightarrow 0) = 0$

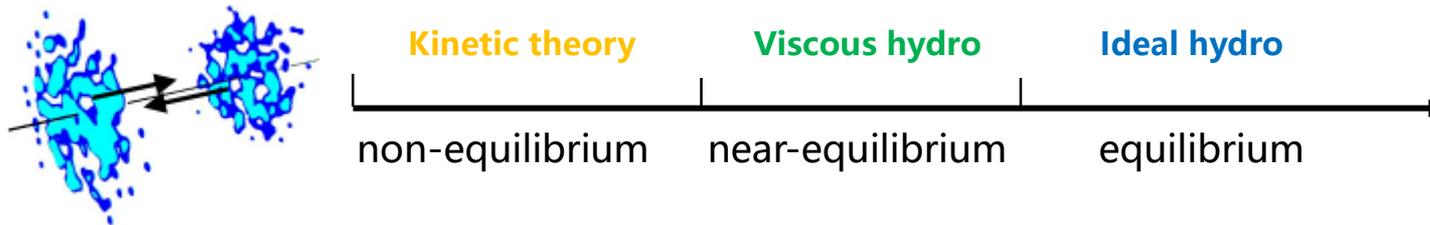
Hydrodynamics in relativistic heavy-ion collisions (HICs)

- Successful description of near-equilibrium quark-gluon plasma (QGP)
- Medium background for hard probes (jets, heavy quarks, quarkonium, etc..)
- Energy deposition to the hydrodynamic medium from hard probes as well

Non-hydrodynamics

Non-equilibrium dynamics beyond hydrodynamics in HICs

- Equilibration of jets in near-equilibrium QGP Kinetic theory
- Pre-equilibrium stage of HICs



Non-hydrodynamic modes: anything not hydrodynamic ...

large gradients

$\nabla \sim k$ not small

non-vanishing frequency

$\omega(k \rightarrow 0) \neq 0$

How to study non-hydrodynamic modes

- Any correlator in quantum field theory out of hydrodynamic region
- Specifically, effective kinetic theory with linear response

Effective kinetic theory

Linearized effective kinetic theory

- Consider an **Effective Kinetic Theory** (set of coupled Boltzmann equations)

$$p^\mu \partial_\mu f_a(t, \mathbf{x}, p) = -C_a^{LO\ 2\leftrightarrow 2, 1\leftrightarrow 2}[f](t, \mathbf{x}, p)$$

- Decompose distribution into spatially homogeneous background $f(t, p)$ and inhomogeneous perturbation $\delta f(t, \mathbf{x}, p)$

$$f_a(t, \mathbf{x}, p) = \underbrace{f_a(t, T, p)}_{\text{Background}} + \underbrace{\delta f_a(t, \mathbf{x}, p)}_{\text{Perturbation}}$$

- Fourier transform: **gradients** \rightarrow **wavenumber \mathbf{k}**

$$\delta f_a(t, \mathbf{k}, p) = \int \frac{d^3 \mathbf{x}}{(2\pi)^3} e^{-i\mathbf{x}\cdot\mathbf{k}} \delta f_a(t, \mathbf{x}, p)$$

- Results in a **Linearized Effective Kinetic Theory** with a wavenumber \mathbf{k}

$$(p^0 \partial_t + i\mathbf{p} \cdot \mathbf{k}) \delta f_a(t, \mathbf{k}, p) = -\delta C_a^{LO\ 2\leftrightarrow 2, 1\leftrightarrow 2}[f](t, \mathbf{k}, p)$$

Linear response of kinetic theory

Energy-momentum tensor

- Background

$$T^{\mu\nu}(t) = \int \frac{d^3p}{(2\pi)^3} \frac{p^\mu p^\nu}{p} \sum_a v_a f_a(t, p)$$

- Perturbation

$$\delta T^{\mu\nu}(t, k) = \int \frac{d^3p}{(2\pi)^3} \frac{p^\mu p^\nu}{p} \sum_a v_a \delta f_a(t, k, p)$$

Response function

- Response function in terms of **time t** and **wavenumber k**

$$G_{\alpha\beta}^{\mu\nu}(t, k) = \frac{\delta T^{\mu\nu}(t, k)}{\delta T^{\alpha\beta}(0, k)}$$

- Consider initial condition with scalar type perturbation

$$f_a(t_0, p) = f_a^{eq}(T, p) \text{ Background} \quad \delta f_a(t_0, k, p) = -\frac{\delta T}{T} \partial_p f_a^{eq}(T, p) \text{ Perturbation}$$

- Sound channel $G_{00}^{00}(t, k) = \frac{\delta T^{00}(t, k)}{\delta T^{00}(0, k)}$ (not the only one, but they are related)

Universal scales

Universal scales

- Response functions are in terms of \mathbf{t} and \mathbf{k}
- Different interaction strengths give different relaxation times $\tau_R \propto \eta/s$
- Rescaled & dimensionless time and wave-number

$$\bar{t} = t \frac{sT}{\eta} \quad \bar{k} = k \frac{\eta}{sT}$$

Hydrodynamization time

$$\bar{t}_H \approx 4\pi$$

Response function in universal scales

$$G(t, k) \rightarrow G(\bar{t}, \bar{k})$$

- 1st order hydrodynamics response function can be formulated in universal scales

$$G_{\text{hydro}}^{1\text{st}}(t, k) = \cos(c_s k t) e^{-\Gamma k^2 t} \quad \text{with} \quad \Gamma = \frac{2}{3} \frac{\eta}{sT}$$

$$\text{So that} \quad G_{\text{hydro}}^{1\text{st}}(t, k) = \underbrace{\cos(c_s \bar{k} \bar{t})}_{\text{Dispersion}} \underbrace{e^{-\frac{2}{3} \bar{k}^2 \bar{t}}}_{\text{Damping}}$$

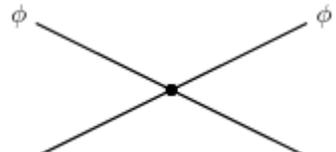
Universality among kinetic theories

- Relaxation time approximation (RTA) ■ ϕ -4 scalar theory (SCL)
- Yang-Mills kinetic theory (YM) ■ Quantum chromodynamic kinetic theory (QCD)

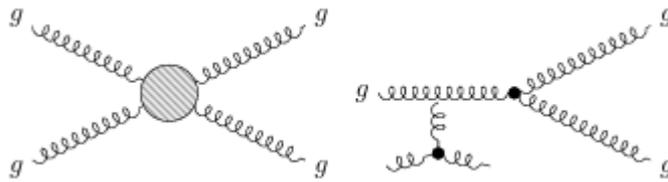
Collision integrals in kinetic theories

$$C_a^{RTA}[f](t, x, p) = \frac{p^\mu u_\mu}{\tau_R} [f(t, x, p) - f_{eq}(t, x, p)]$$

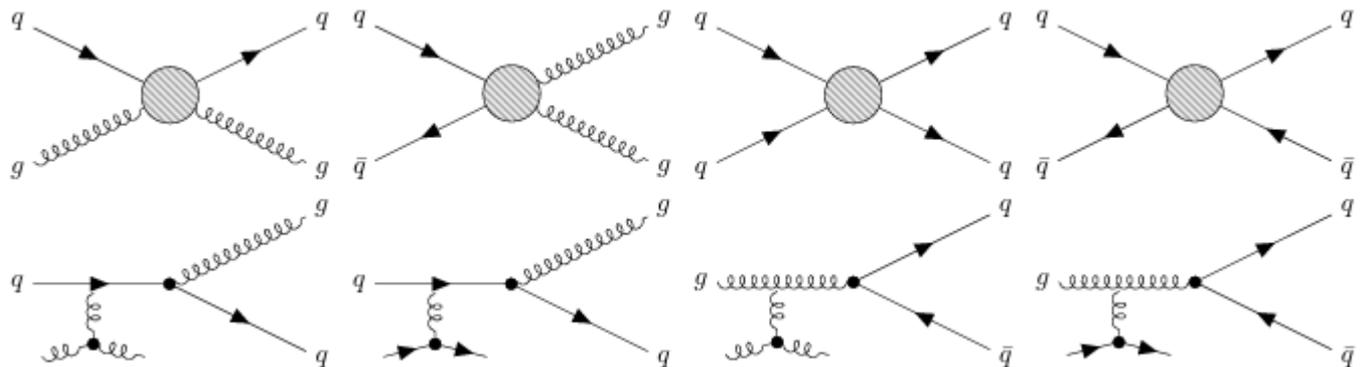
$$C_a^{SCL}[f](t, x, p) =$$



$$C_a^{YM}[f](t, x, p) =$$



$$C_a^{QCD}[f](t, x, p) =$$

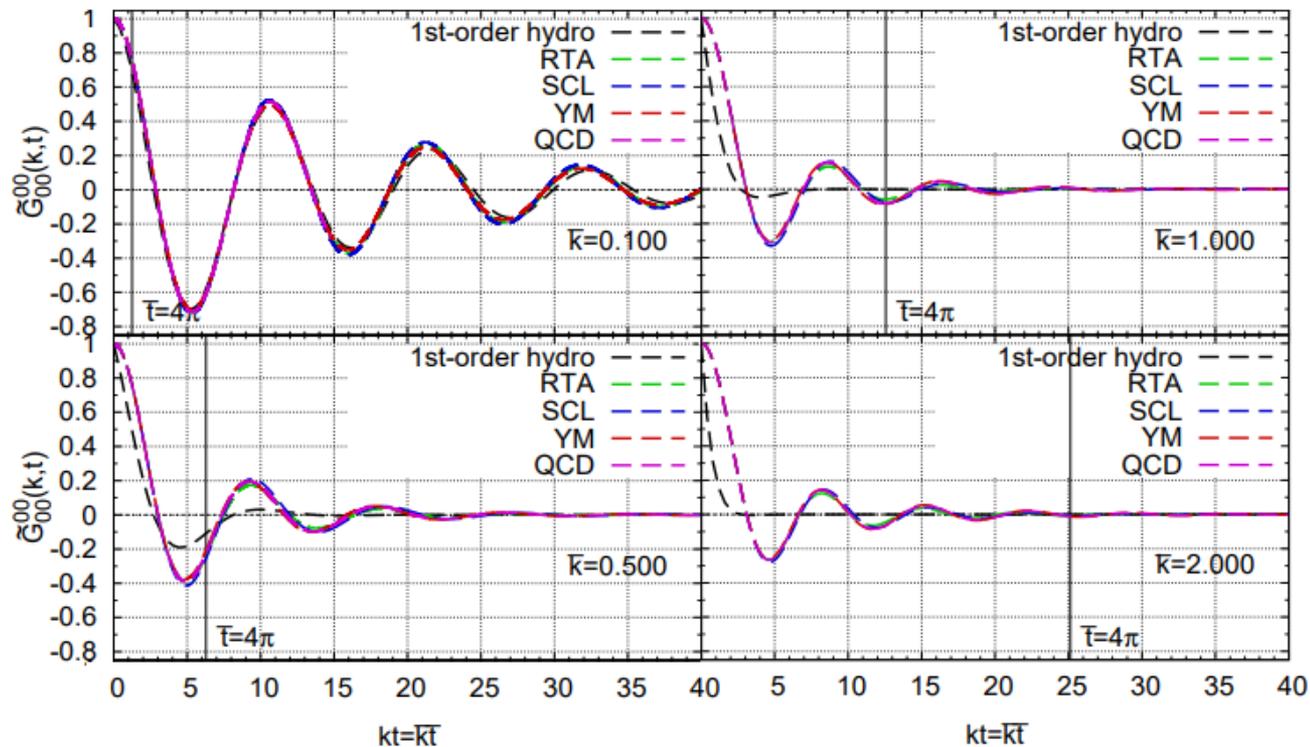


Universality among kinetic theories

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Response functions from kinetic theories

- Expected to reproduce hydrodynamics at small k (long wave-length limit)
- Universality among different kinetic theories even at large k and early time



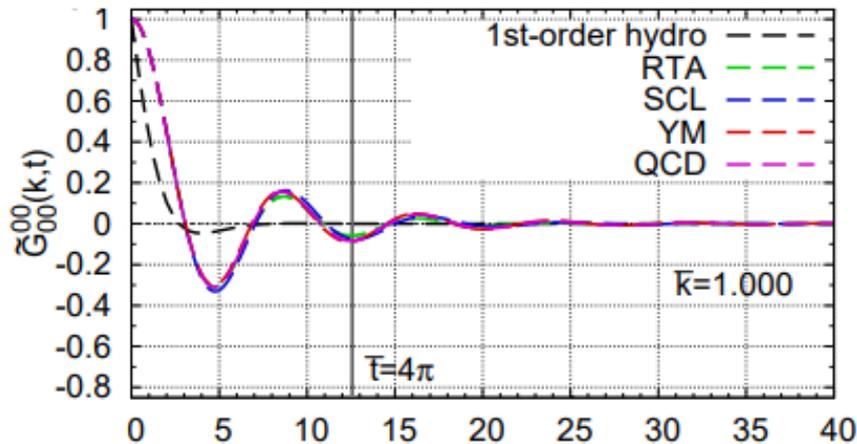
Sound and non-hydrodynamic modes

Fitting response functions in kinetic theory

- With real (dispersion/oscillation) and imaginary (damping) frequencies

$$G_{S,n}(t, k) \sim Z_k \exp[-i(\omega_k t + \phi_k)]$$

$$\omega_k = \text{Re}[\omega_k] + i\text{Im}[\omega_k]$$



non-hydro mode

$$G_n(\bar{t} < \bar{t}_H, k)$$



Hydrodynamization time

$$\bar{t}_H = 4\pi$$

$$kt = \bar{k}\bar{t}$$

sound mode

$$G_s(\bar{t} > \bar{t}_H, k)$$

Remarks

- Negative frequency gives the same mode as positive frequency

Sound modes appear in pair

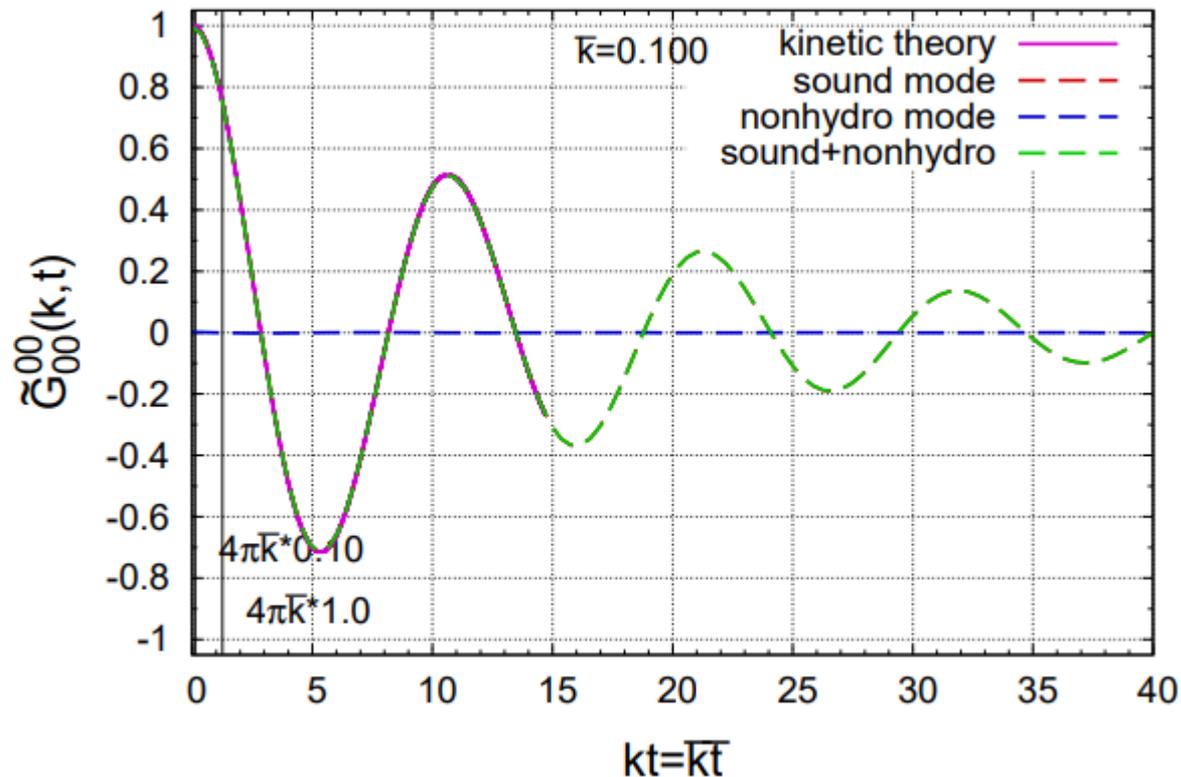
- Expected many/infinite number of non-hydrodynamic modes

We represent them with a single non-hydrodynamic mode

Sound and non-hydrodynamic modes

Fitting QCD response functions ($\bar{k}=0.1$)

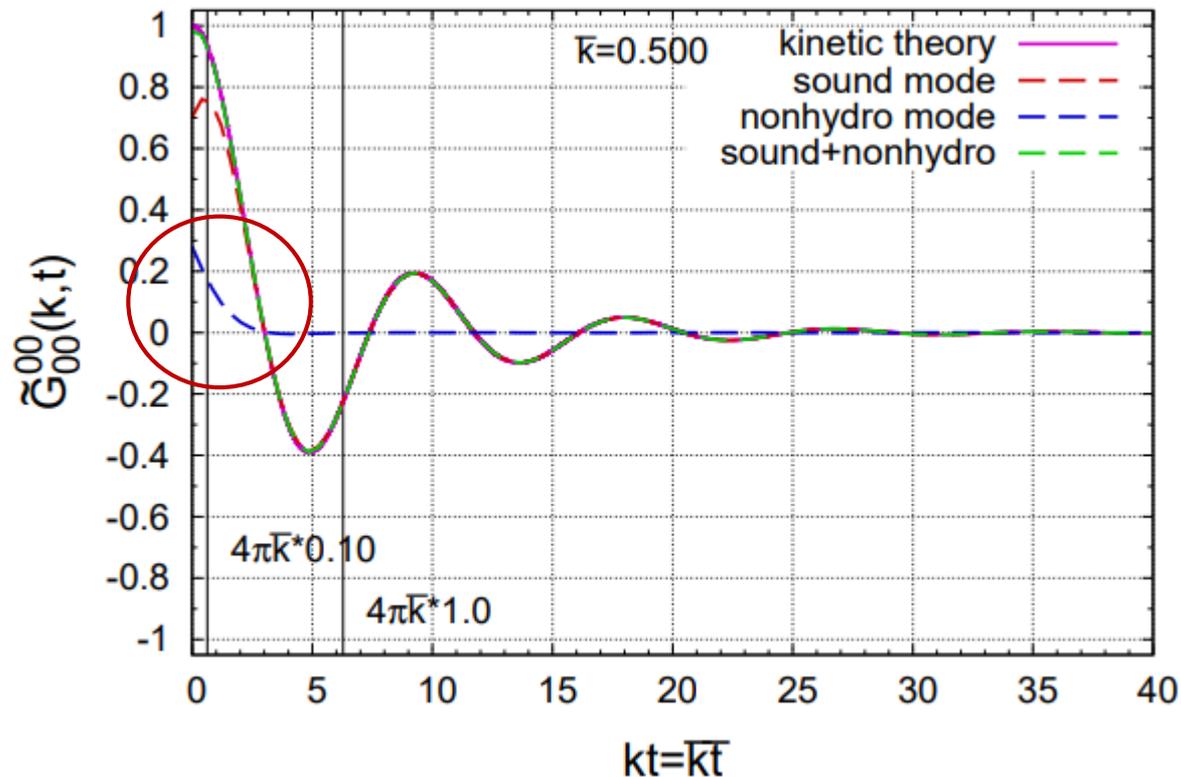
- Sound mode dominates



Sound and non-hydrodynamic modes

Fitting QCD response functions ($\bar{k}=0.5$)

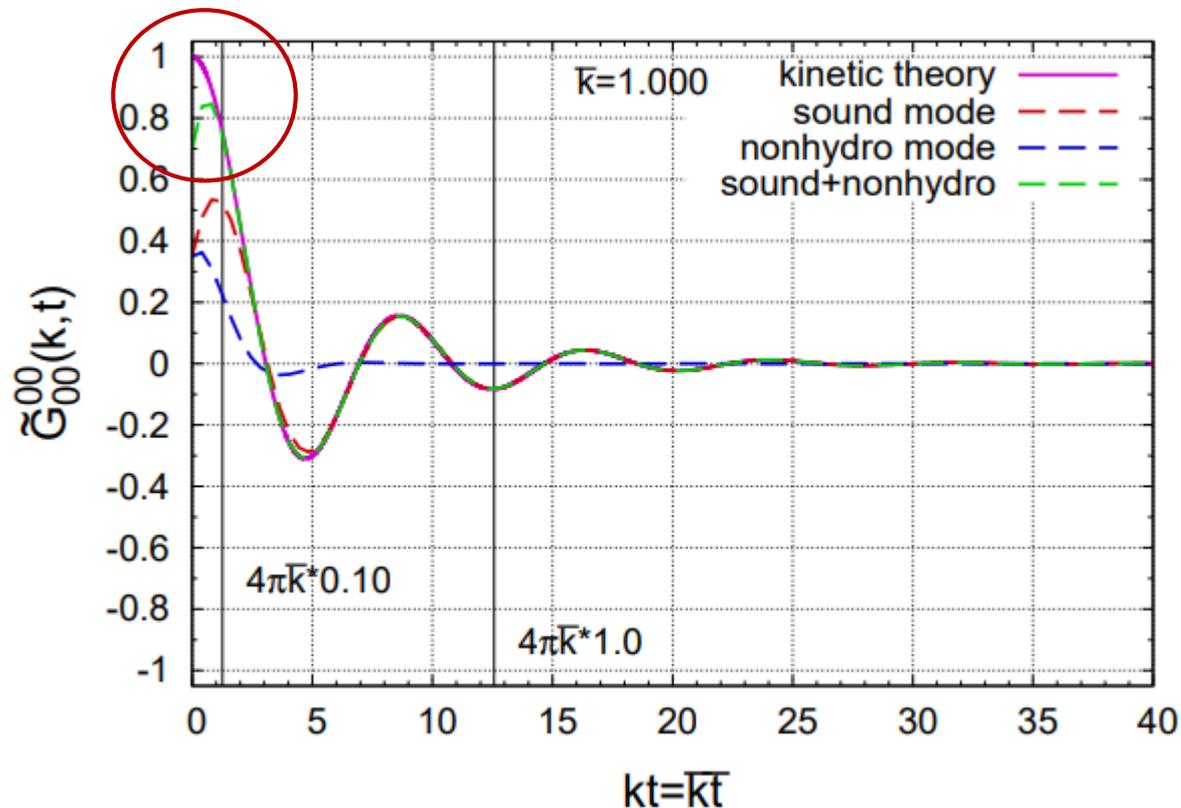
- Non-hydrodynamic mode appears



Sound and non-hydrodynamic modes

Fitting QCD response functions ($\bar{k}=1.0$)

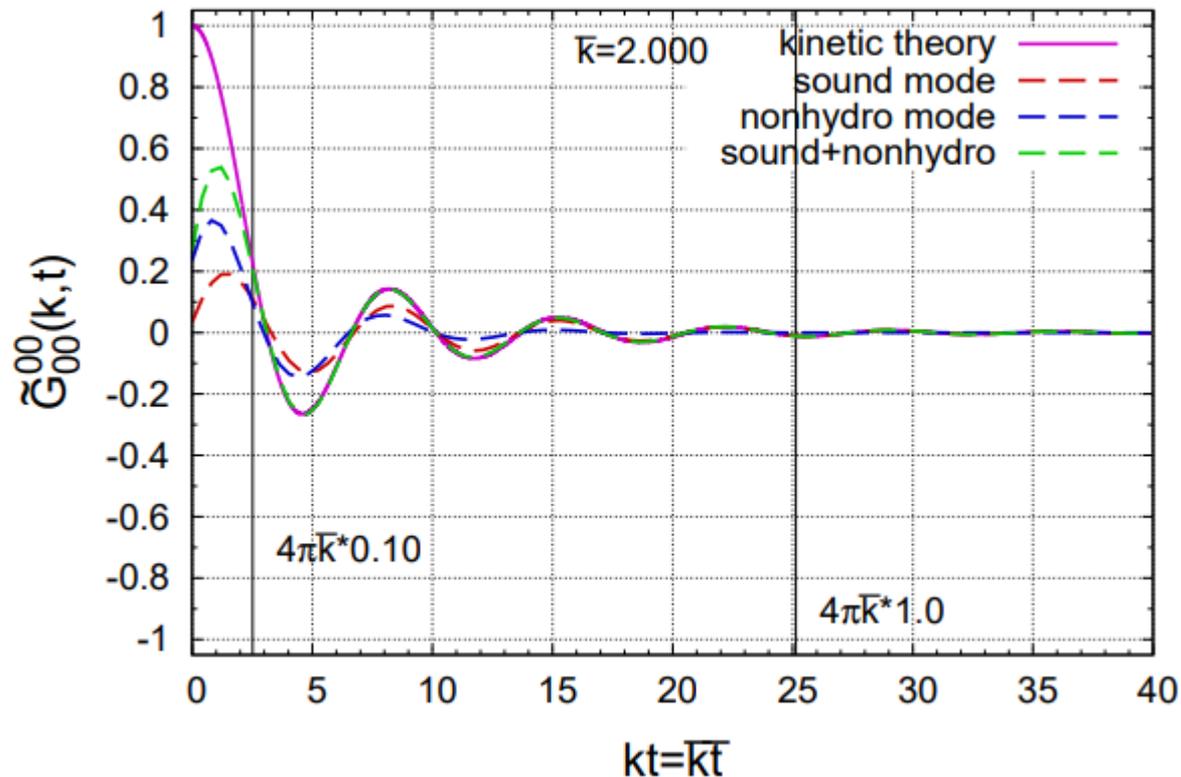
- More non-hydrodynamic modes appear at early time



Sound and non-hydrodynamic modes

Fitting QCD response functions ($\bar{k}=2.0$)

- Non-hydrodynamic mode takes over the domination by sound mode



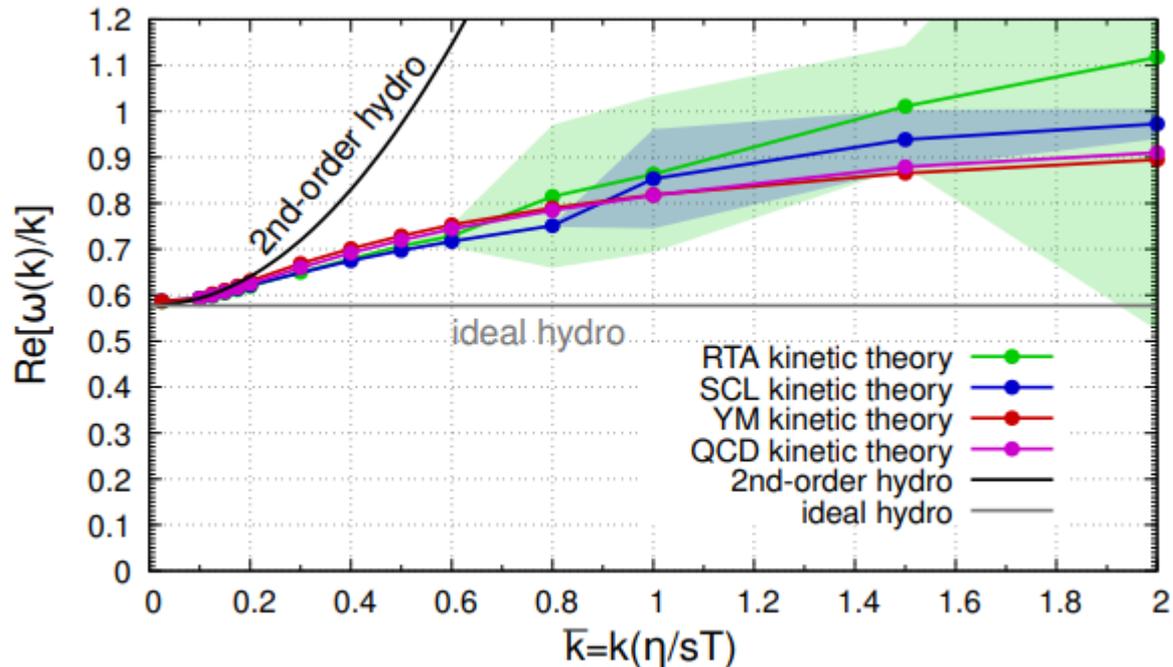
Dispersion relation

Dispersion relations among kinetic theories

- With real frequencies

$$G_{S,n}(t, k) \sim Z_k \exp[-i(\omega_k t + \phi_k)]$$

$$\omega_k = \text{Re}[\omega_k] + i\text{Im}[\omega_k]$$



- Universality of sound modes among kinetic theories at various k
- Kinetic theories converge to 2nd-order hydrodynamics at small k

$$\omega_{hydro}^{2nd}(k) = c_s k - i\Gamma k^2 + \frac{\Gamma}{c_s} \left(c_s^2 \tau_\pi - \frac{\Gamma}{2} \right) k^3 \quad \text{with} \quad \Gamma = \frac{2}{3} \frac{\eta}{sT}$$

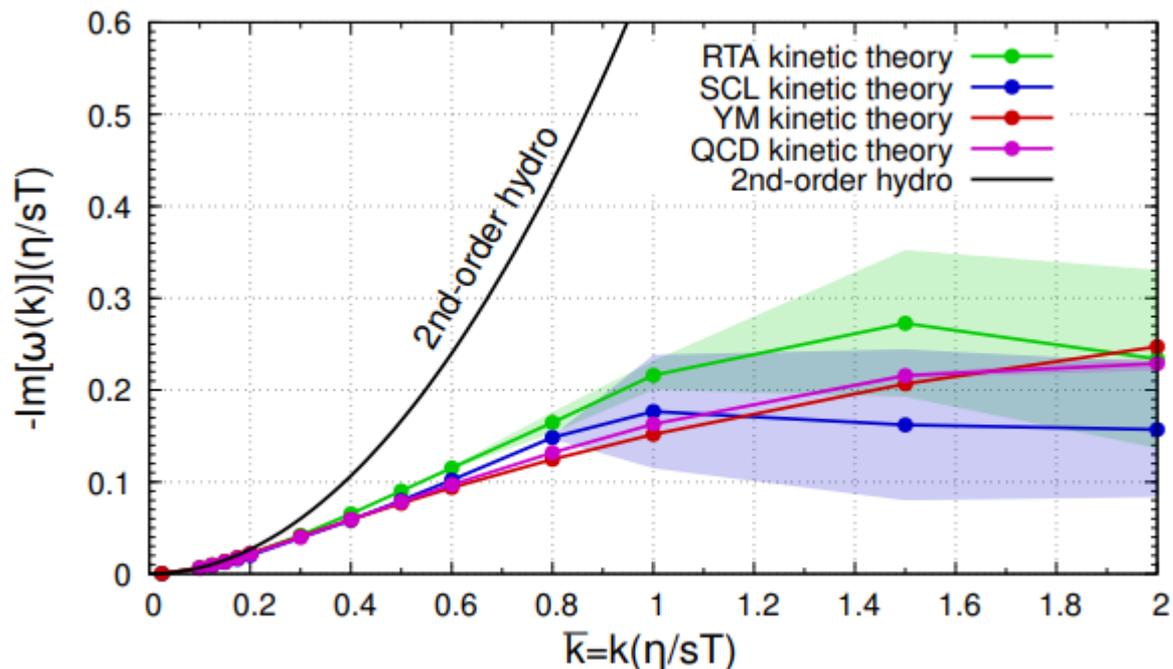
Damping relation

Damping relations among kinetic theories

- With imaginary frequencies

$$G_{S,n}(t, k) \sim Z_k \exp[-i(\omega_k t + \phi_k)]$$

$$\omega_k = \text{Re}[\omega_k] + i\text{Im}[\omega_k]$$



- Universality of sound modes among kinetic theories at various k
- Kinetic theories converge to 2nd-order hydrodynamics at small k

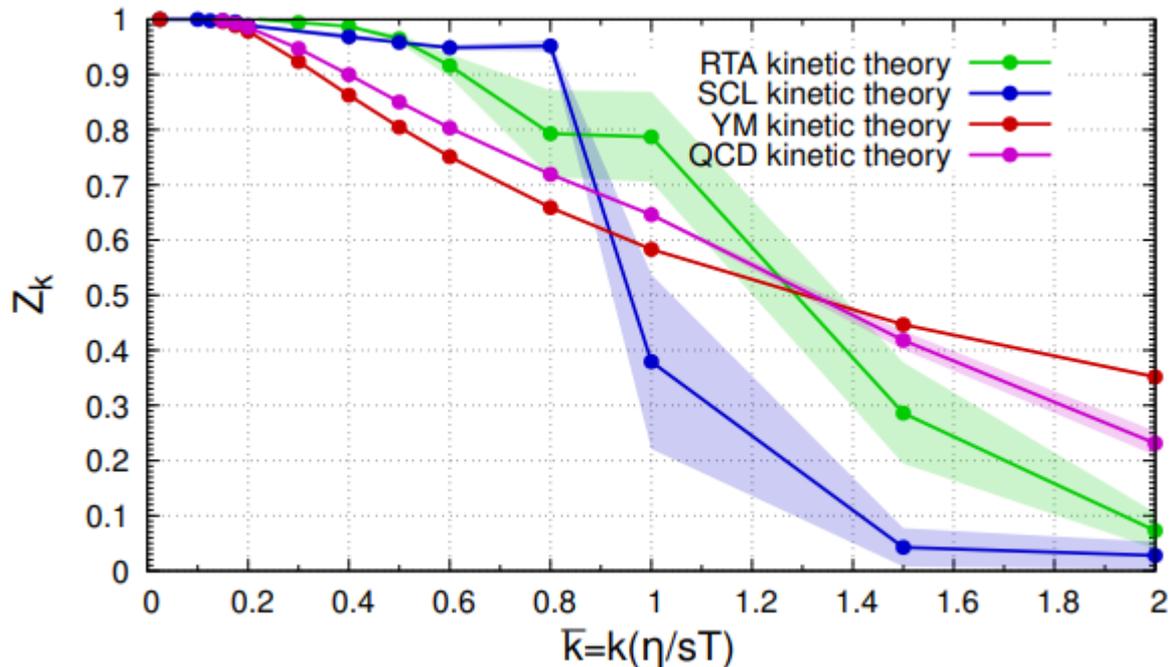
$$\omega_{hydro}^{2nd}(k) = c_s k - i\Gamma k^2 + \frac{\Gamma}{c_s} \left(c_s^2 \tau_\pi - \frac{\Gamma}{2} \right) k^3 \quad \text{with} \quad \Gamma = \frac{2}{3} \frac{\eta}{sT}$$

Residue

Residue for sound & non-hydro modes among kinetic theories

- Non-equilibrium plasma described by **sound mode** + **non-hydrodynamic mode**

$$G_{S,n}(t, k) \sim Z_k \exp[-i(\omega_k t + \phi_k)]$$

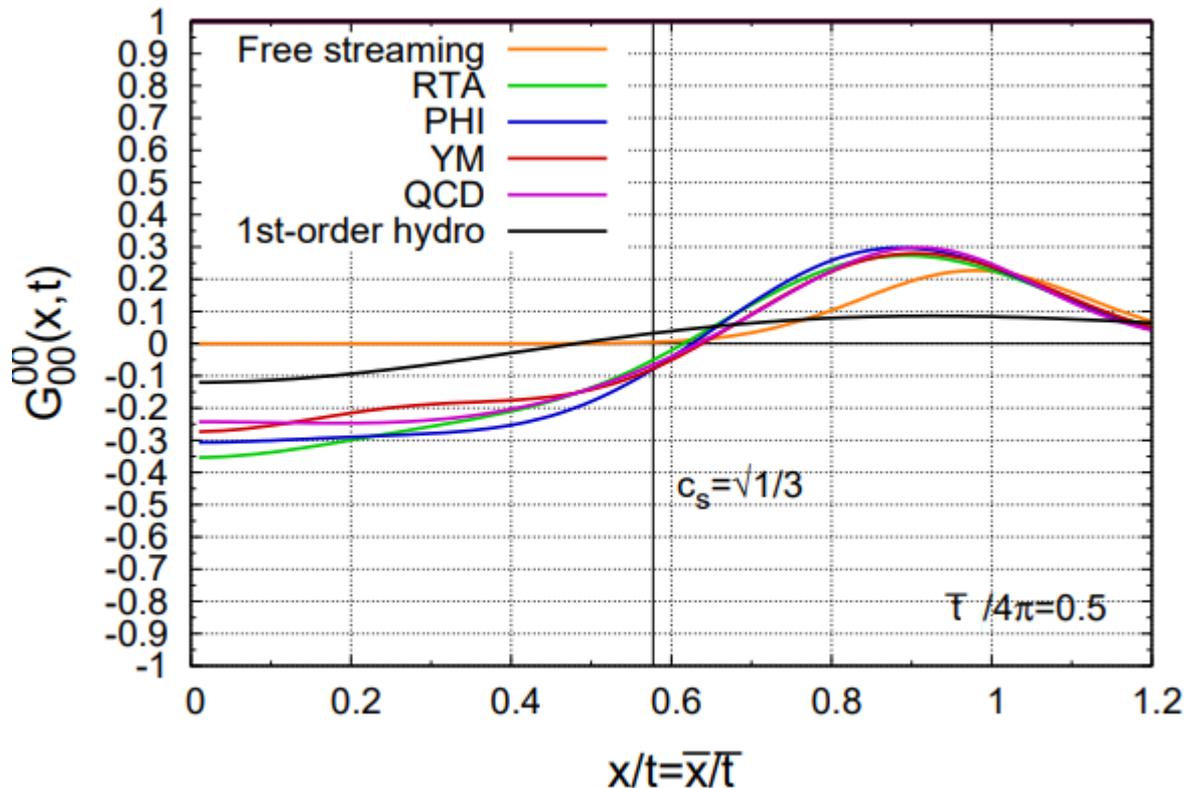


- **Sound mode** dominates at small k (**non-hydro mode** dominates at large k)
- Universality of residue at some degree

Response in position space

Response in position among kinetic theories

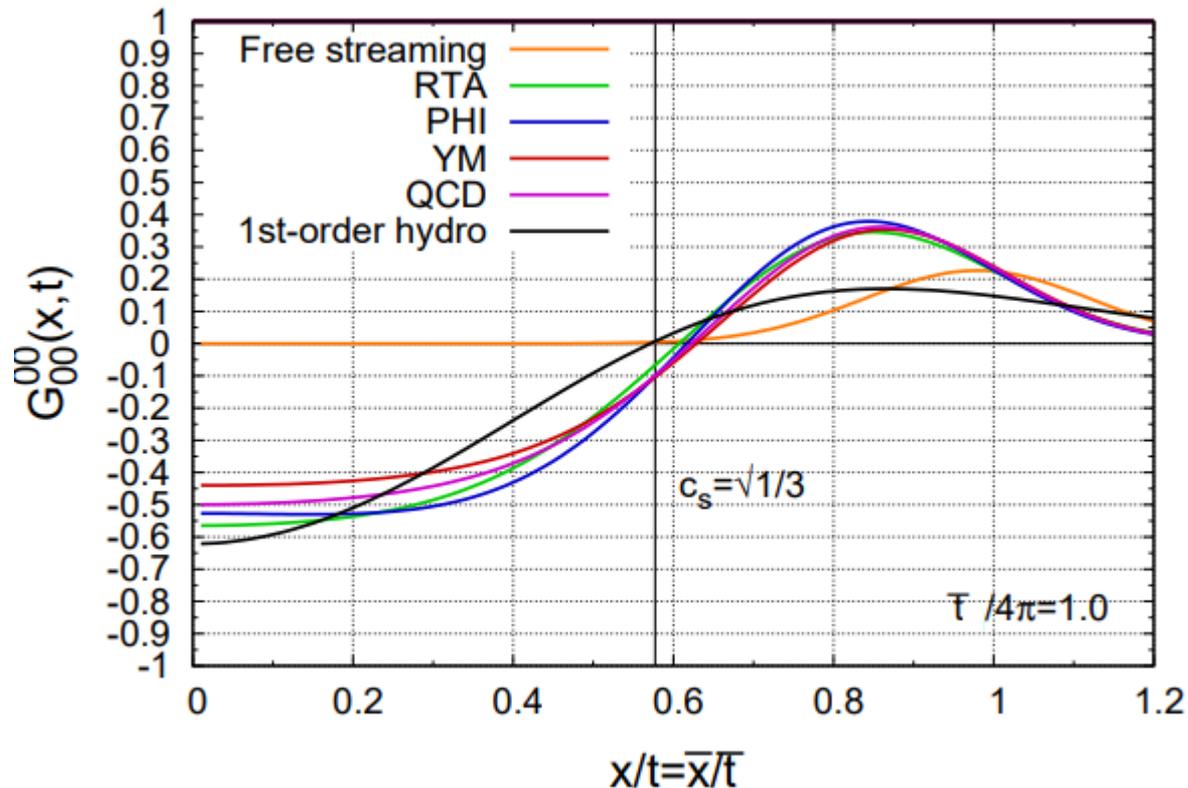
- Universality of position response



Response in position space

Response in position among kinetic theories

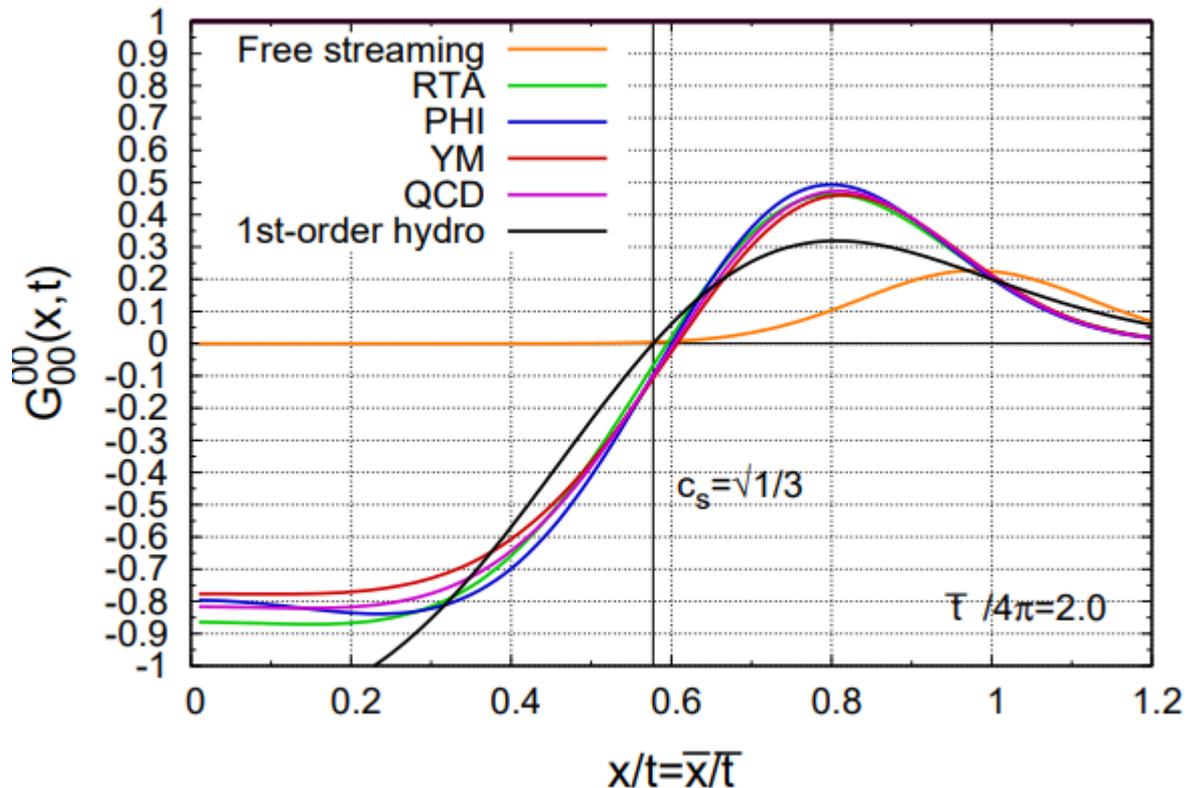
- Universality of position response



Response in position space

Response in position among kinetic theories

- Kinetic theories get closer to hydrodynamics at later time



Summary

What are expected and verified

- Kinetic theories **converge to hydrodynamics** at small k
- Linear response of kinetic theories can be described by **sound**+**non-hydro** modes

What are strikingly new

- Universality of **energy-momentum response functions** among kinetic theories
(Even at early time and large k which are in the non-equilibrium region)

What are future perspectives

- Possible to push hydrodynamic description to non-equilibrium regime?

Thanks!