

# Quarkonium production at the LHC: experiment and theory overview

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A. Andronic - University of Münster

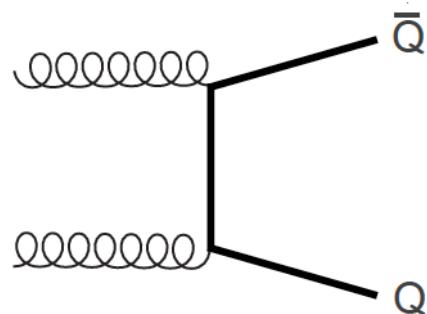


- Measurements and models in pp collisions (brief)
- Data and models in Pb-Pb collisions (compared to RHIC data)

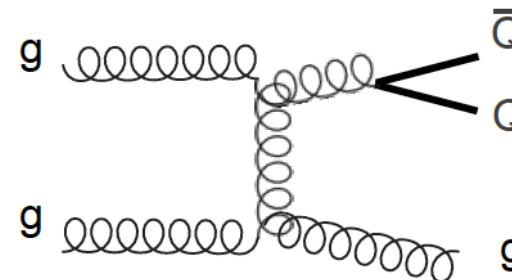
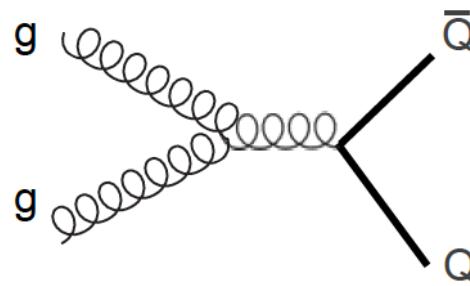
# Heavy-quark and quarkonium production

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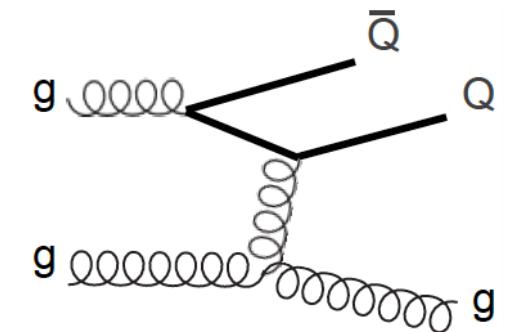
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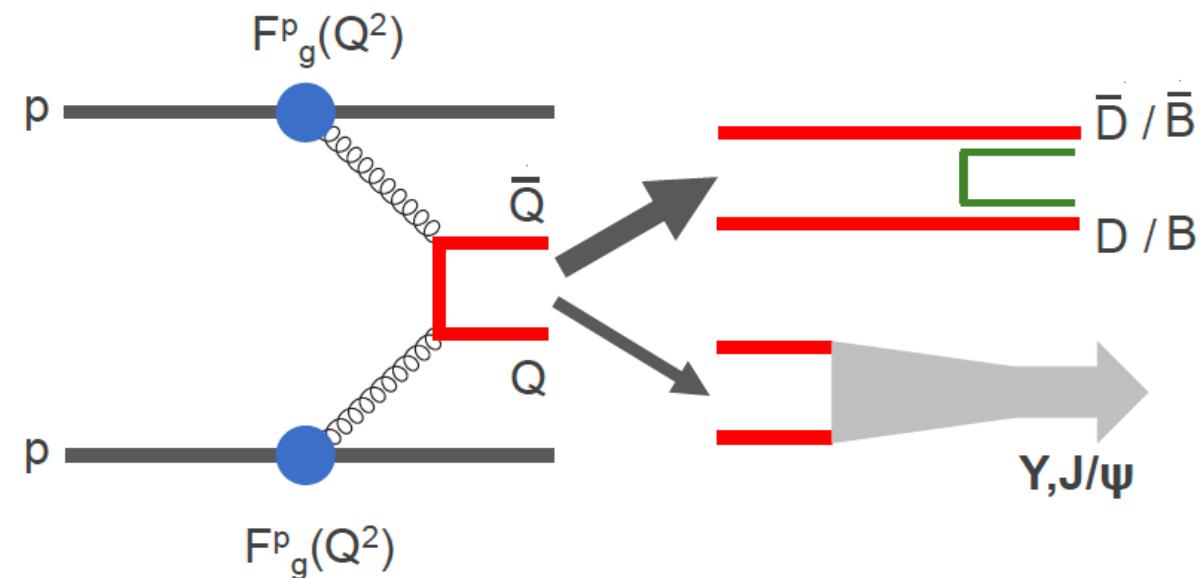
flavor production



gluon splitting



flavor excitation



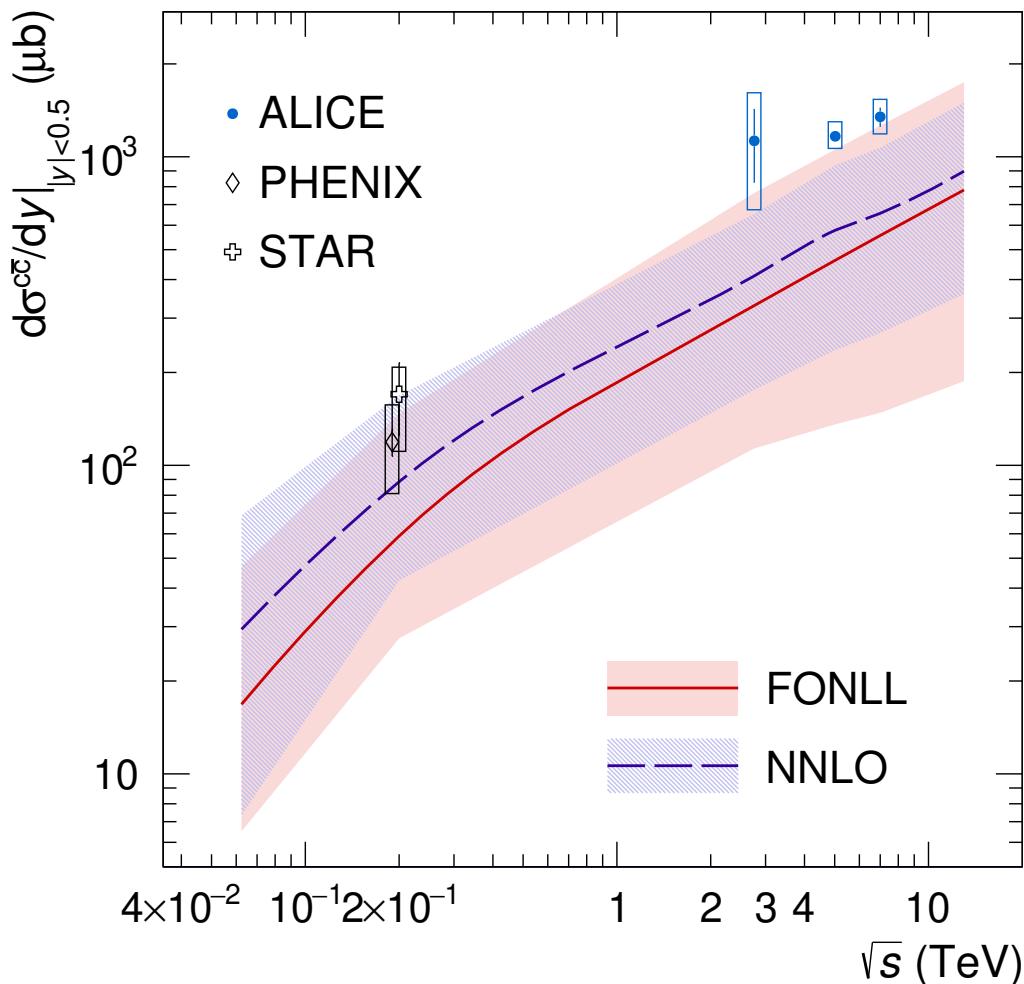
open heavy flavor

quarkonium

# Charm data in pp collisions

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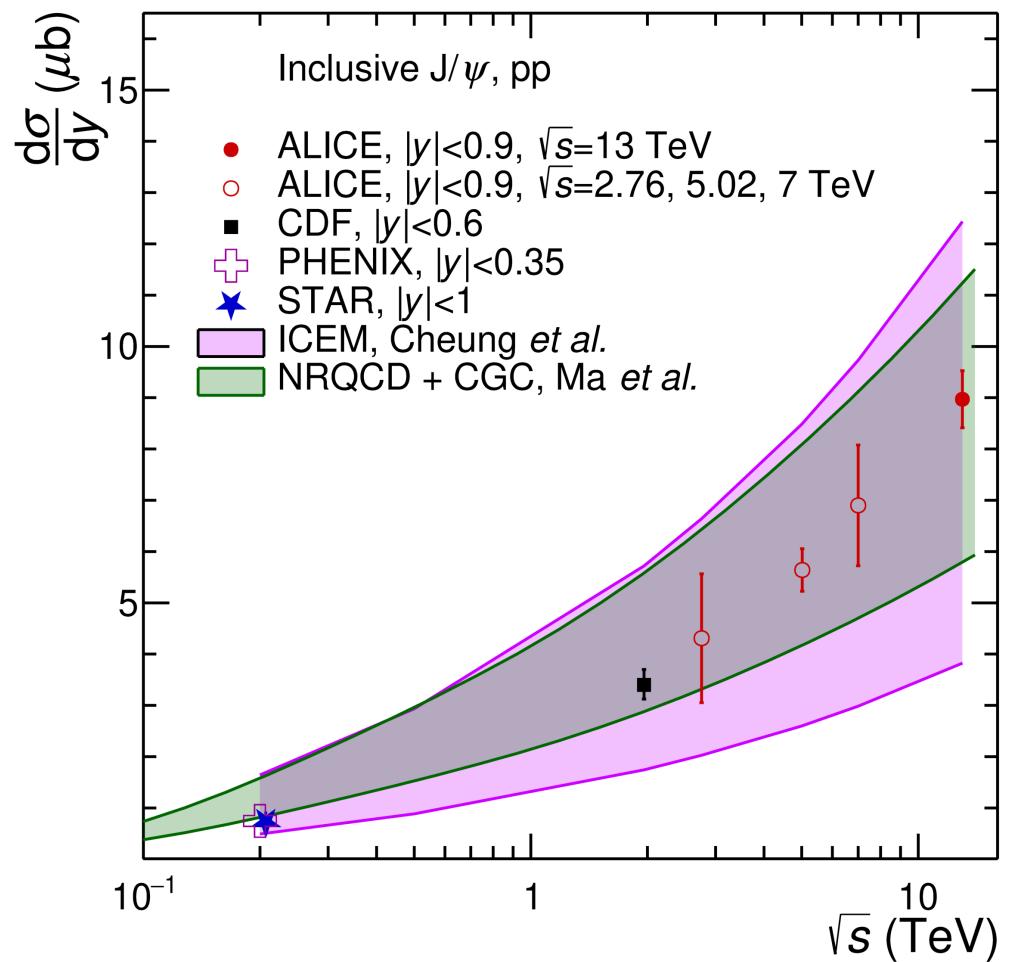


ALI-PUB-500755

ALICE, PRD 105 (2022) L011103

FONLL, NNLO: perturbative QCD calculations

abundant charm(onium) production at the LHC; charmonium/charm  $\sim 1\%$



ALICE, Eur. Phys. J. C 81 (2021) 1121

# Quarkonium models

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Non-relativistic QCD (NRQCD):

as  $Q, \bar{Q}$  move slowly (in quarkonium c.m.s.,  $v$ ), exploits hierarchy of scales  
 $m_Q \gg m_Q v, m_Q v^2, \Lambda_{QCD}$  ( $\simeq 200$  MeV)

sum over all states (singlet, octet) of the product of  $\sigma_{Q,\bar{Q}}$  (pQCD) times the long-distance matrix elements that describe the non-perturbative formation process (derived from fits of data or calculated in lattice QCD)

Color Evaporation Model, CEM:  $\sigma_{CEM}^{pp \rightarrow H+X} = k_H \cdot \sigma_{Q,\bar{Q}}$  (quarkonium  $H$ )

Color Singlet Model, GSM:  $\sigma_{CSM}^{pp \rightarrow H+X} \sim \sigma_{Q,\bar{Q}}(v=0) |R_H(0)|^2$  ( $S$  states)

$R_H(0)$  is the radial wavefunction at origin for hadron (quarkonium)  $H$

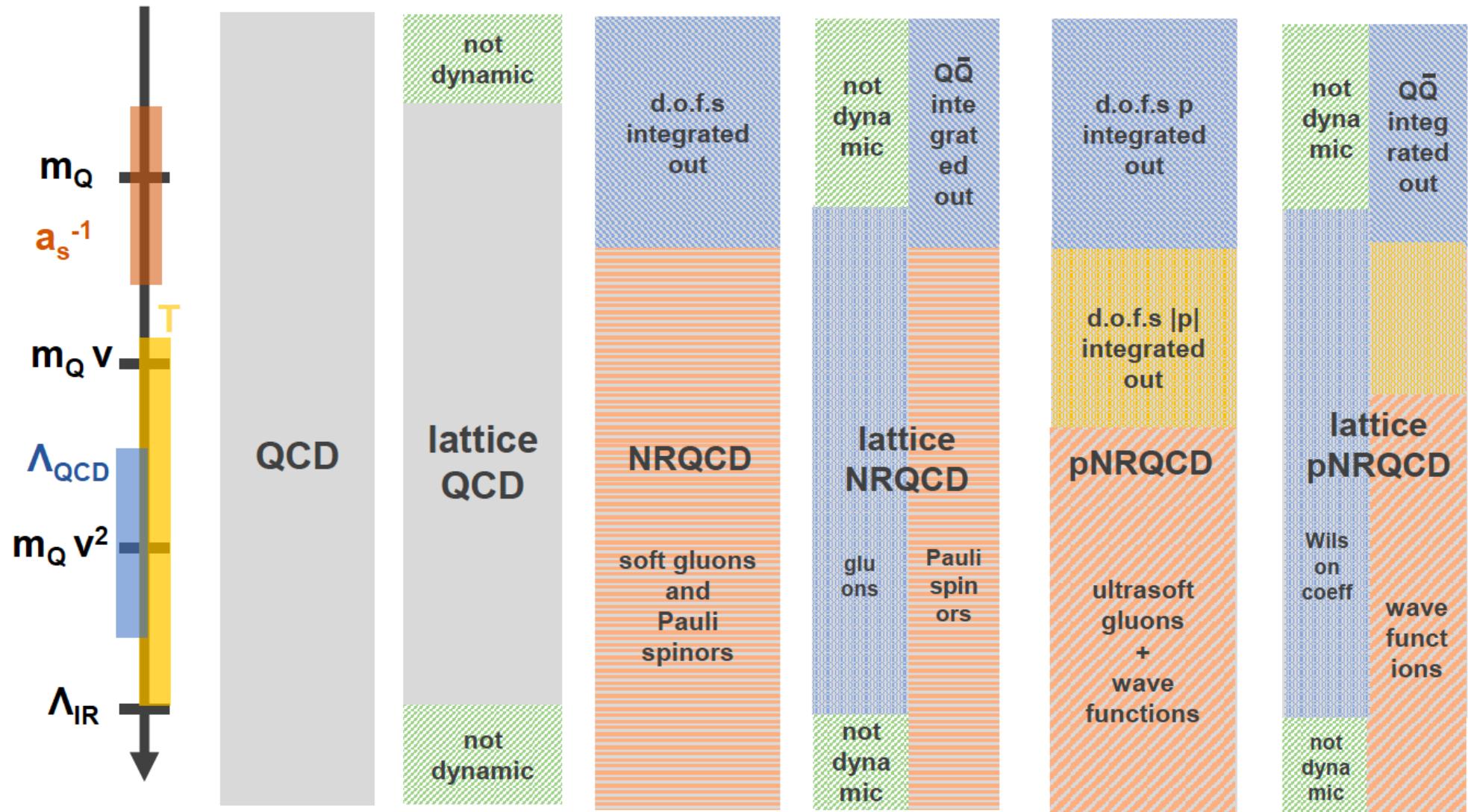
Bodwin, [arXiv:1012.4215](https://arxiv.org/abs/1012.4215); Soto, [arXiv:1101.2392](https://arxiv.org/abs/1101.2392);

# Quarkonium modelling

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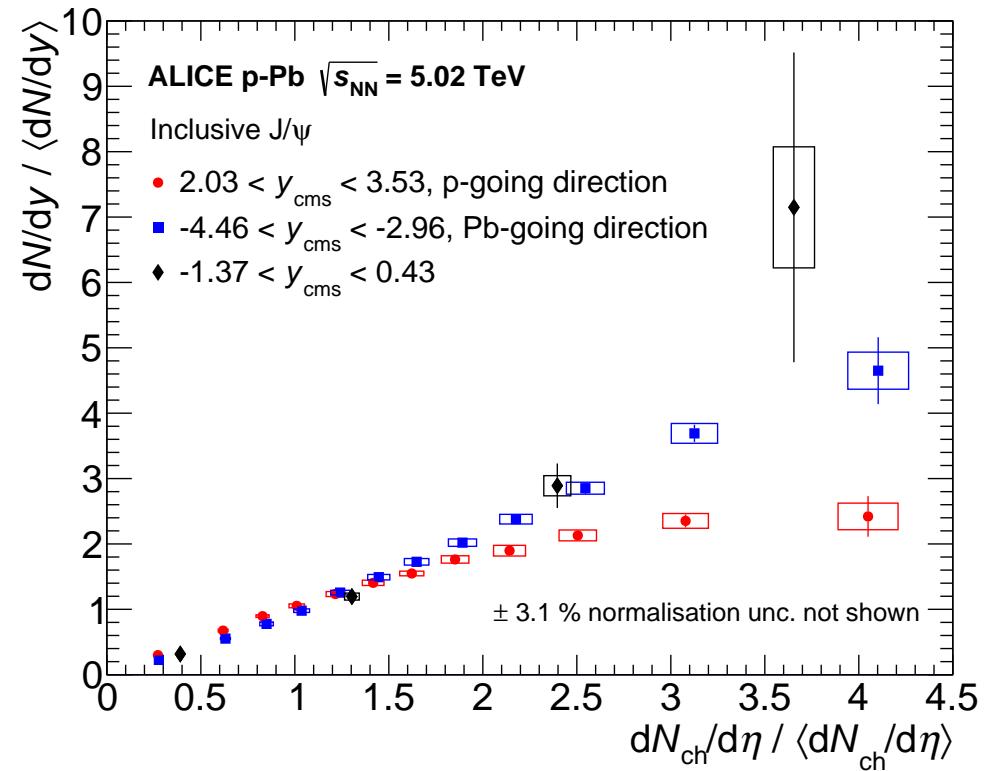
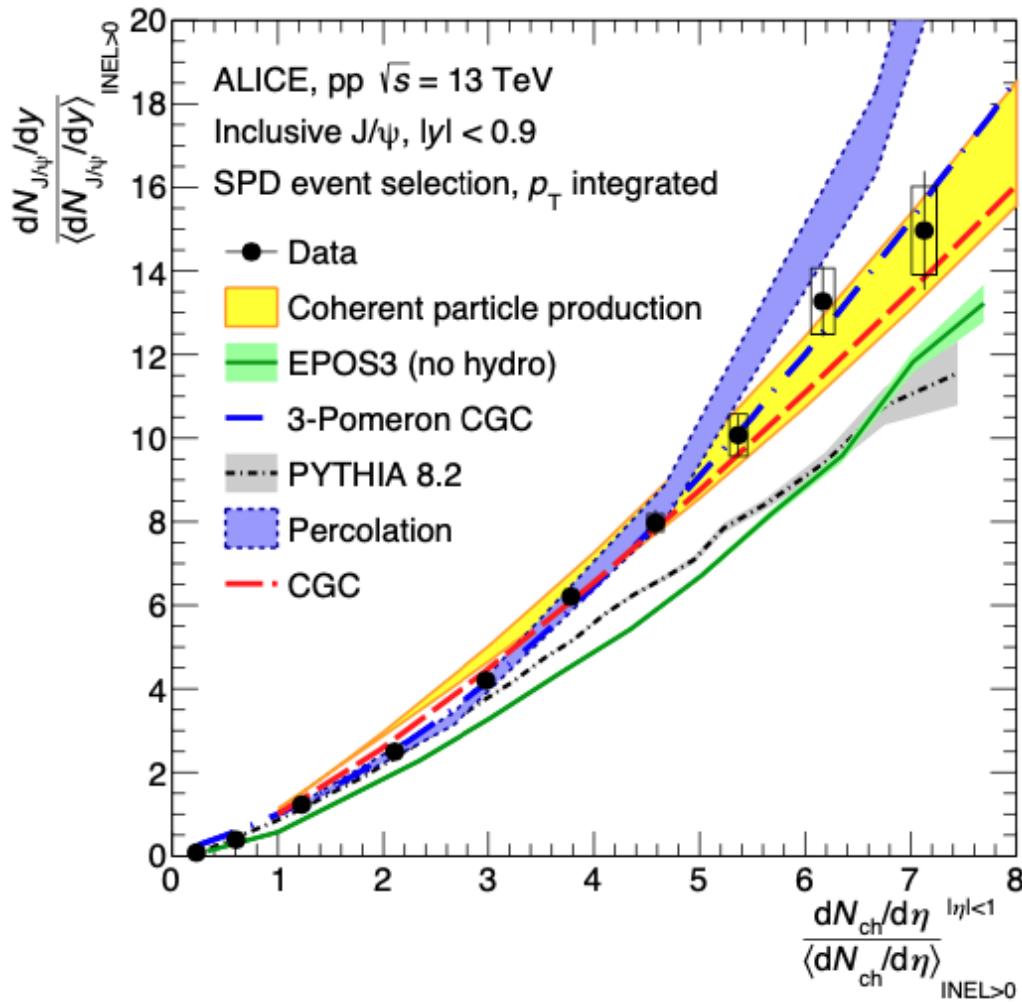
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...is a rich theory field



# $J/\psi$ vs. event activity

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What is trivial here? (1st diagonal?)

"Energy cost" similar for MB event ( $\Delta y=1$ ) and  $J/\psi$  ( $\simeq 6$  GeV)

*Similar behaviour measured for D mesons;  $\Upsilon$ ; charged part.  $p_T=5-6$  GeV/c*

# Charmonium and deconfined matter

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the original idea: Matsui & Satz, [Phys.Lett. B 178 \(1986\) 178](#)

"If high energy heavy-ion collisions lead to the formation of a hot quark-gluon-plasma, then color screening prevents  $c\bar{c}$  binding in the deconfined interior of the interaction region."

Refinements: "sequential suppression":

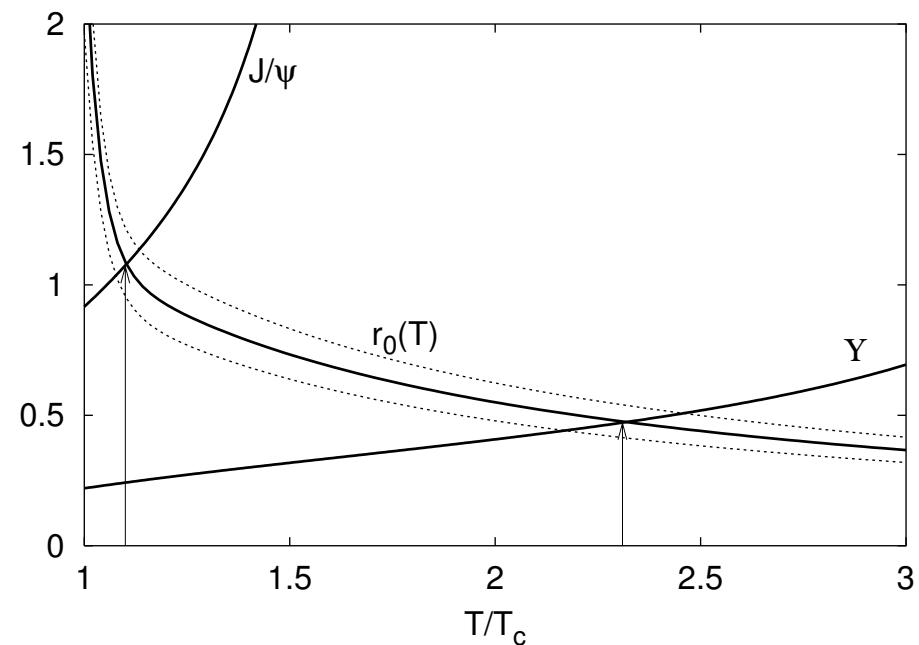
Digal et al., [PRD 64 \(2001\) 75](#)

no  $q\bar{q}$  bound state if

$$r_{q\bar{q}}(T) > r_0(T) \simeq \frac{1}{g(T)T}$$

$r_0$  Debye length in QGP

$\Rightarrow q\bar{q}$  "thermometer" of QGP



Thermal picture ( $n_{partons} = 5.2T^3$  for 3 flavors)

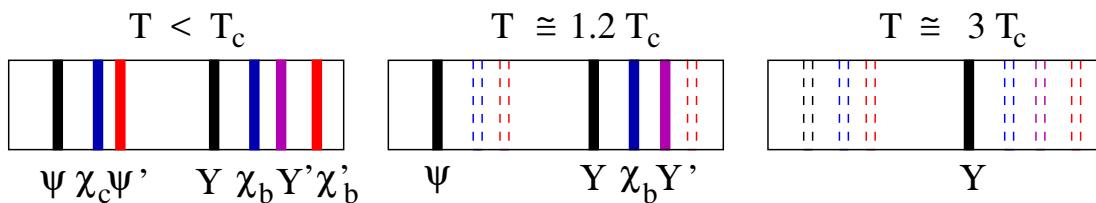
for  $T=500$  MeV:  $n_p \simeq 84/\text{fm}^3$ , mean separation  $\bar{r}=0.2 \text{ fm} < r_{J/\psi}$

# Quarkonium and the QGP - “sequential suppression”

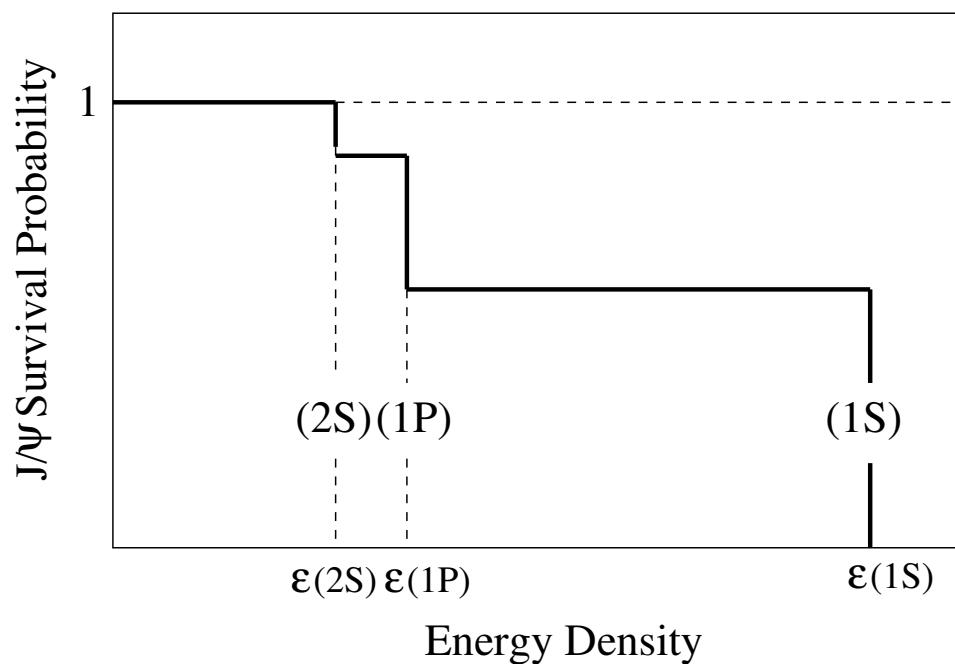
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Digal et al., PRD 64 (2001) 75; H.Satz, arXiv:1310.1209



quantitative: lattice QCD  
(next talk, by P. Petreczky)



← assumed feed-down to  $J/\psi$ :  
10% from  $\psi(2S)$  and 30% from  $\chi_c$

Experimentally:

$$R_{AA}^{J/\psi} = \frac{dN_{J/\psi}^{AA}/dy}{N_{coll} \cdot dN_{J/\psi}^{pp}/dy} \quad \text{vs. } N_{part}$$

or  $N_{ch}$

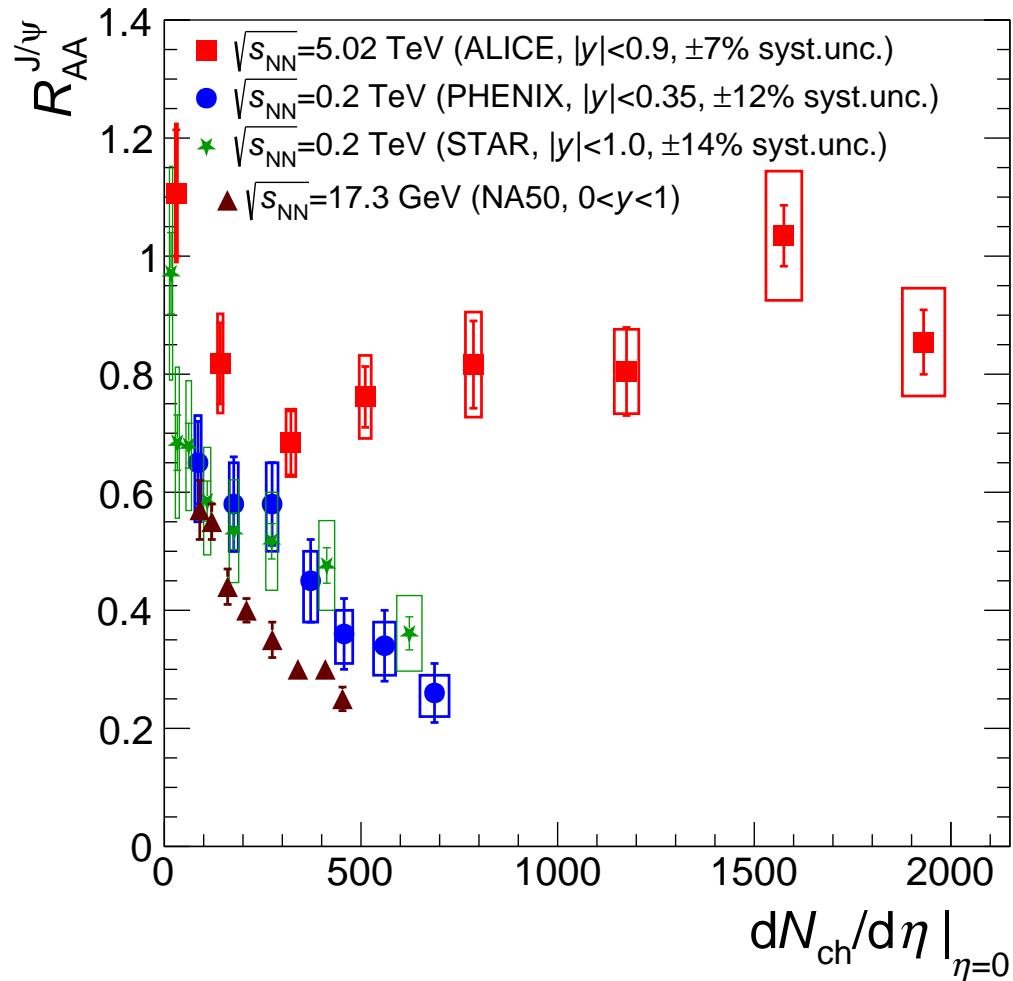
*quarkonium as thermometer*

# Charmonium data: SPS, RHIC, LHC

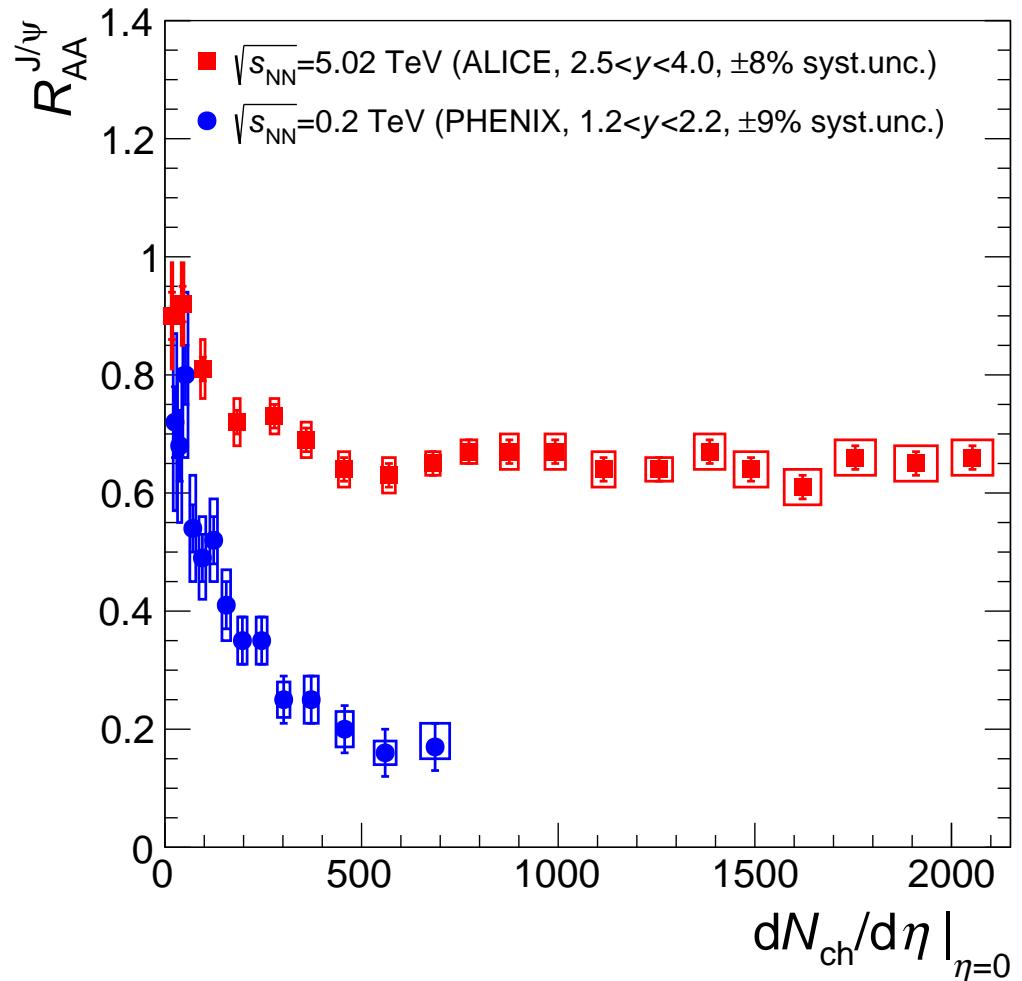
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midrapidity



forward rapidity



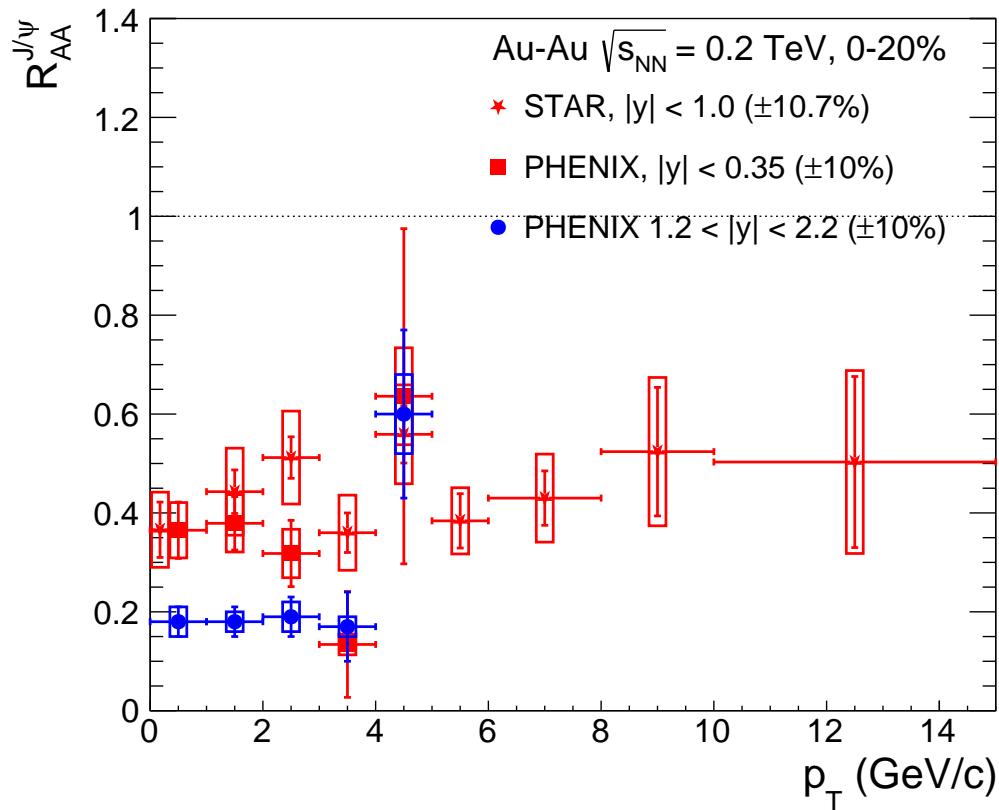
suppression at lower energies, but not obviously a sequential one  
at the LHC another effect offsets the dissociation

# Charmonium data: RHIC and LHC

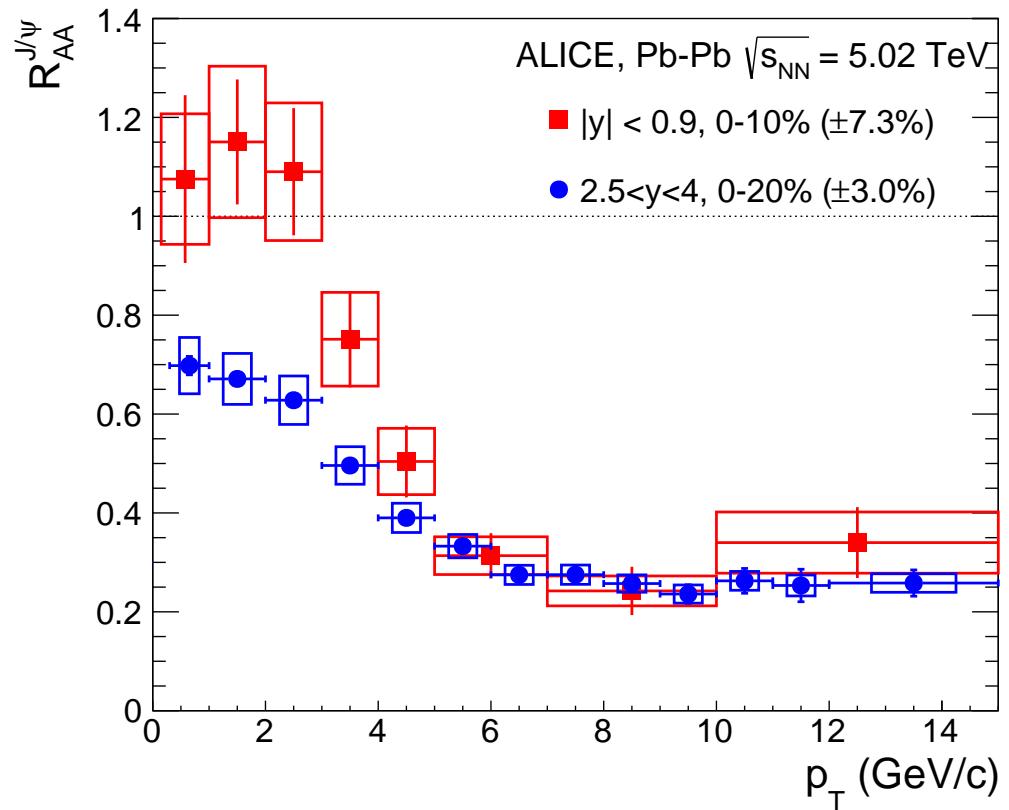
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RHIC



LHC

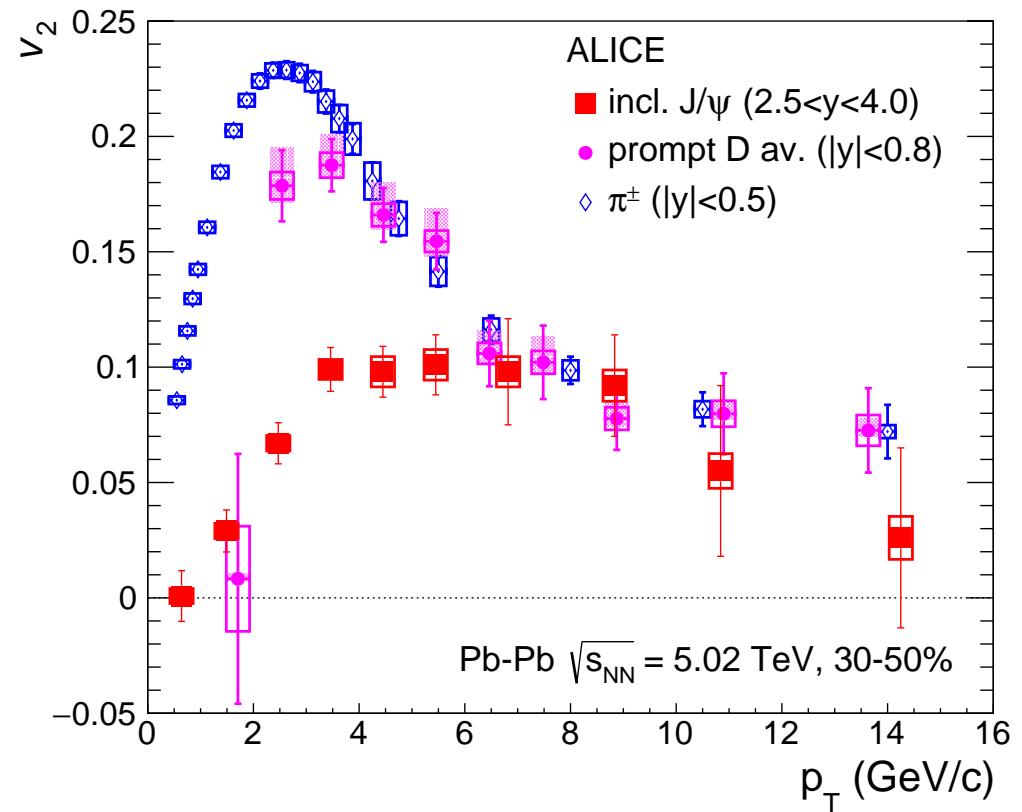
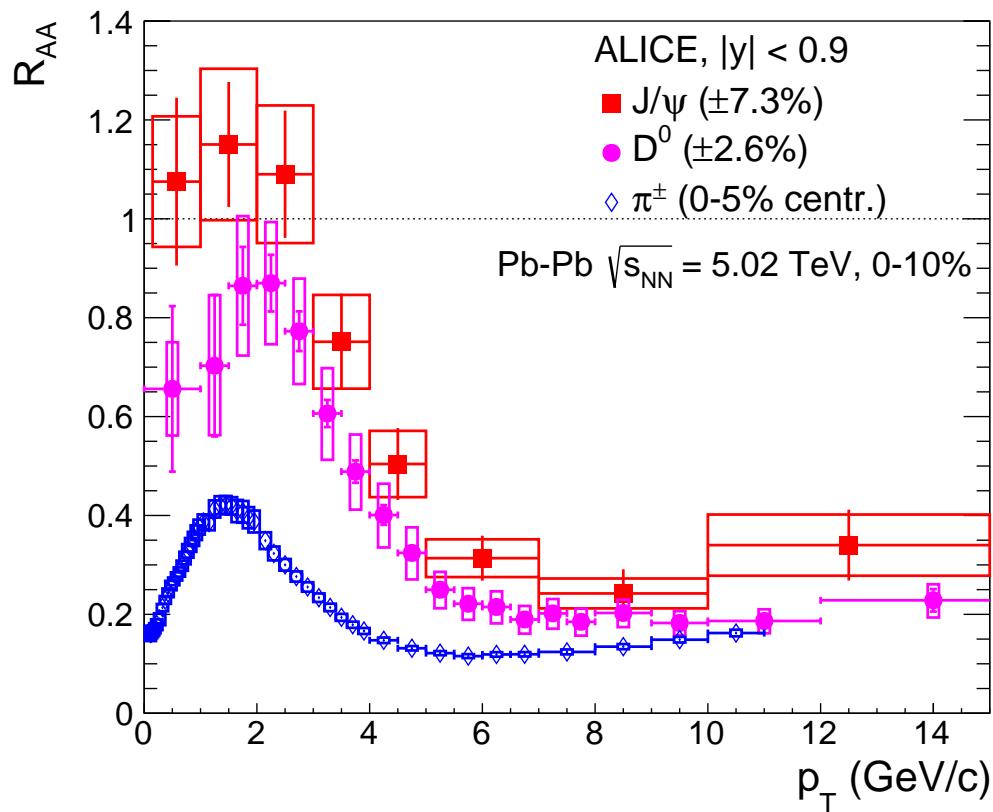


production at low  $p_T$  is enhanced at the LHC compared to RHIC  
...stronger at midrapidity

# Charmonium data at the LHC: in perspective

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ALICE, [arXiv:2303.13361](https://arxiv.org/abs/2303.13361), [arXiv:2110.09420](https://arxiv.org/abs/2110.09420), [arXiv:1802.09145](https://arxiv.org/abs/1802.09145)

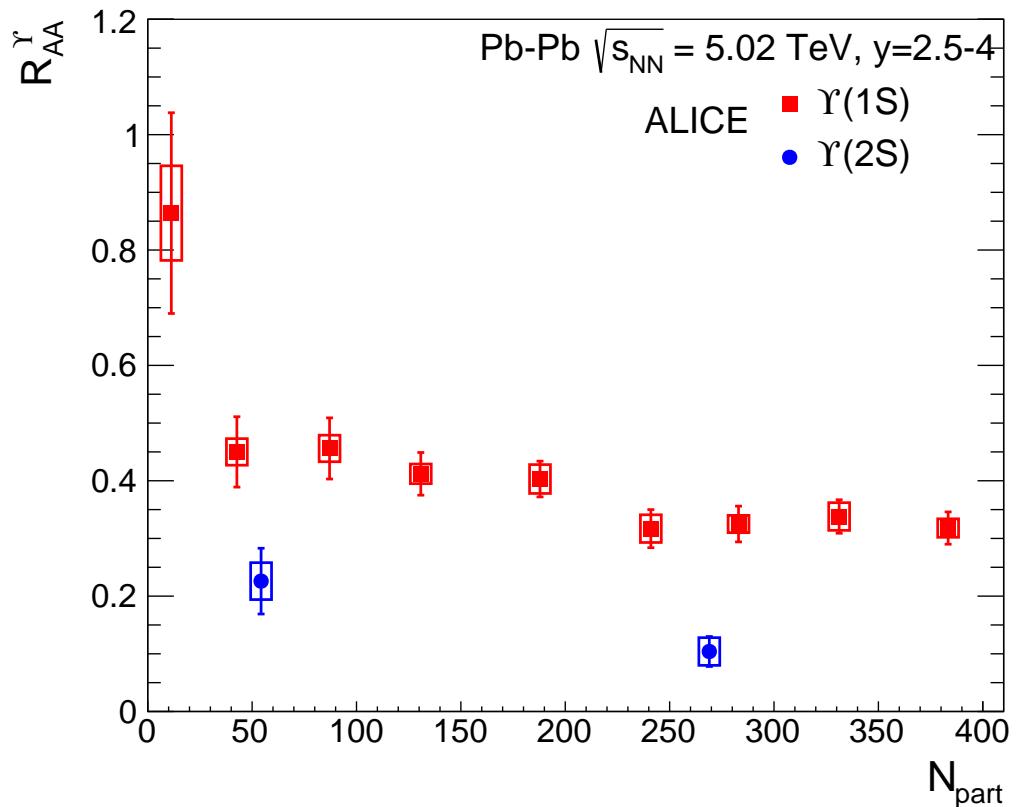
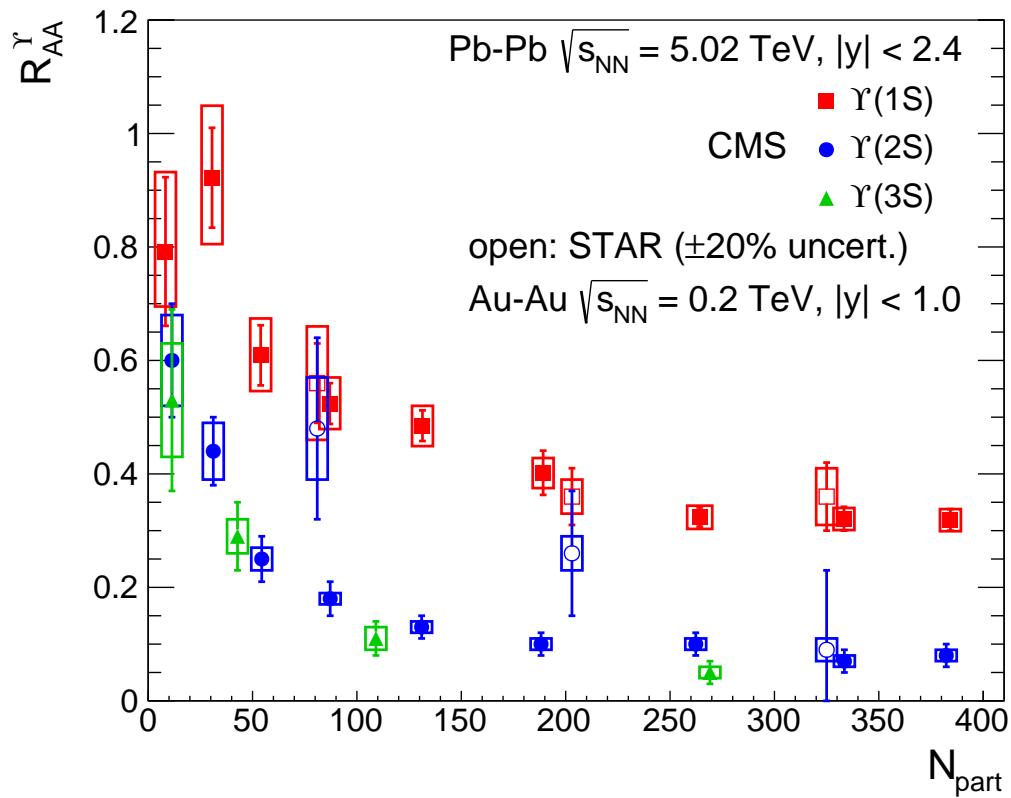
a very clear hierarchy at low/intermediate  $p_T$

charm production is a pQCD process, pion production largely a thermal process

# Bottomonium data: RHIC and LHC

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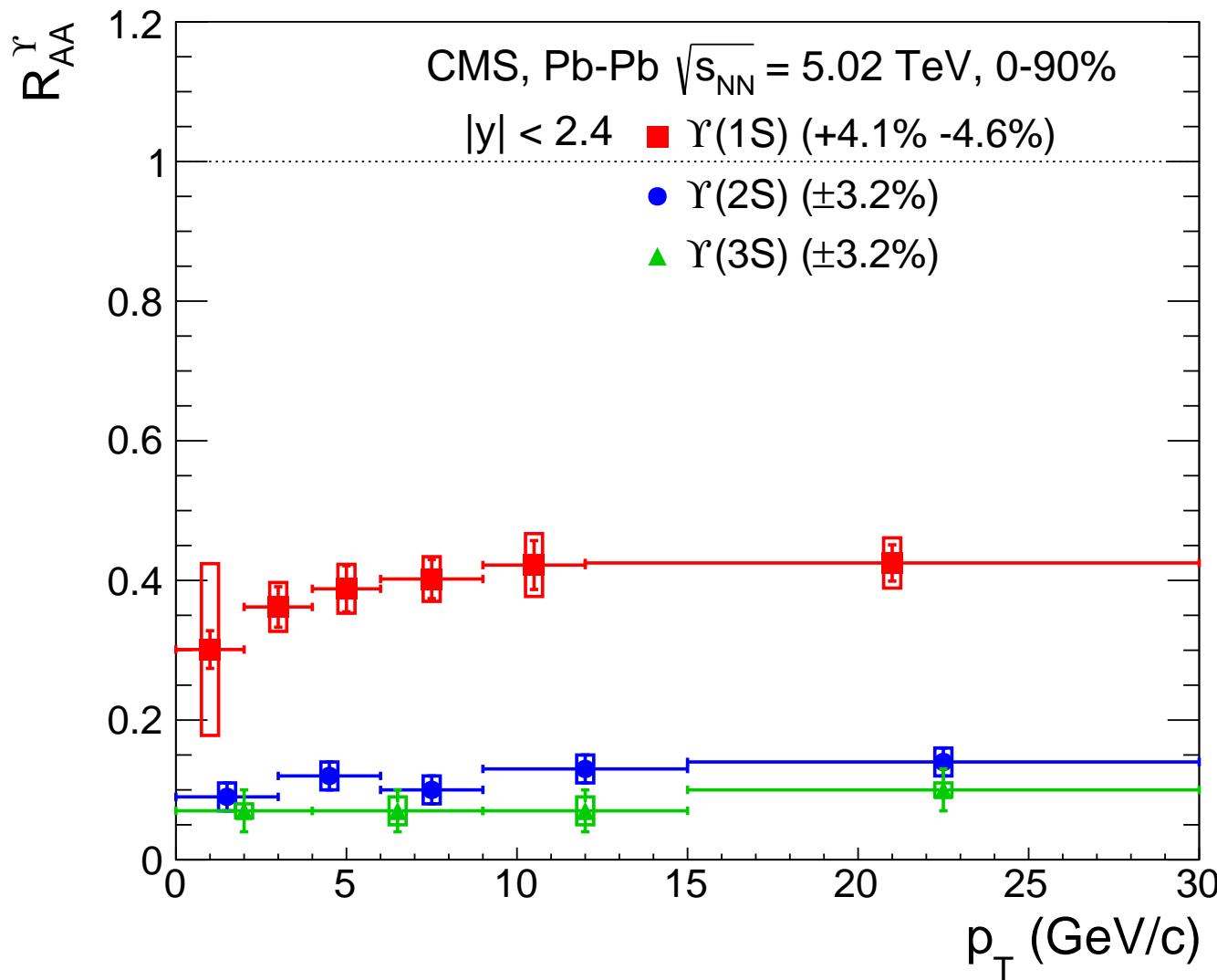
significant suppression, hierarchy between states (sequential?)

not significantly different at LHC and RHIC (except perhaps  $\gamma(2S)$ )

# Bottomonium data at the LHC

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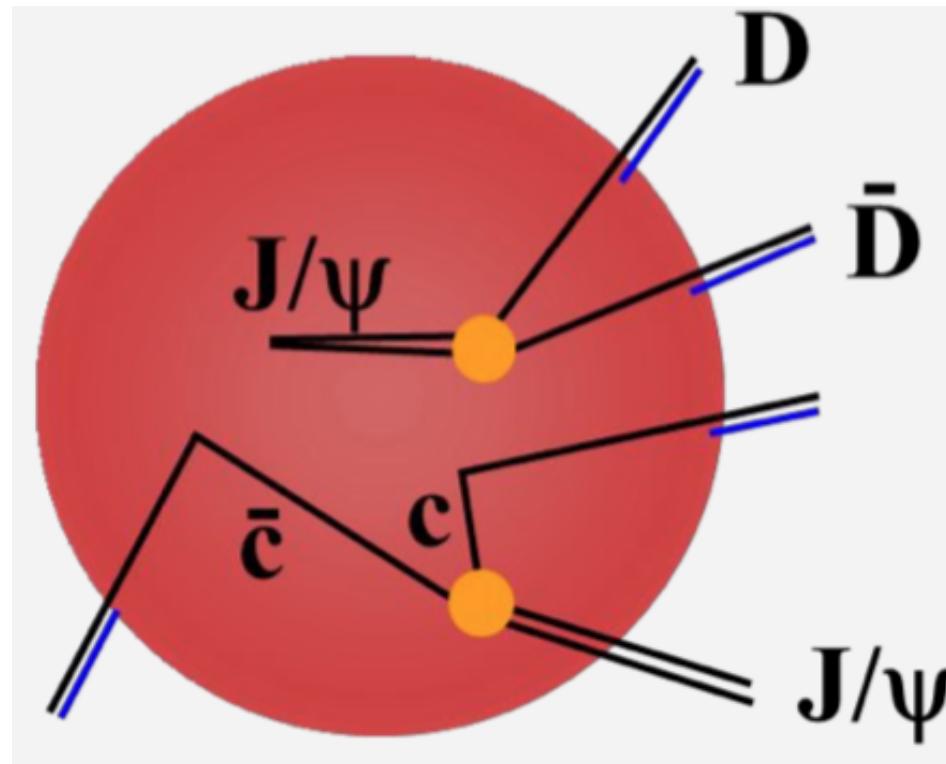


weak  $p_T$  dependence (dissociation expected to be stronger at low  $p_T$ )

# There is dissociation and there is (re)generation

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This picture is implemented in models in 2 basic ways:

- 1) statistical hadronization: full screening, generation at hadronization  
( $c, \bar{c}$  produced in initial scattering and fully thermalized in QGP)
- 2) transport: continuous destruction and (re)generation, also from different  $c, \bar{c}$   
(time evolution of  $T$  constrained by other measurements)

# The statistical hadronization model

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Braun-Munzinger, Stachel, PLB 490 (2000) 196; NPA 789 (2006) 334, PLB 652 (2007) 259

all charm quarks are produced in primary hard collisions ( $t_{c\bar{c}} \sim 1/2m_c \simeq 0.1 \text{ fm}/c$ )

...survive and thermalize in QGP (thermal, but not chemical equilibrium)

charmed hadrons are formed at chemical freeze-out together with all hadrons

“generation” ...no  $J/\psi$  survival in QGP (full screening)

(if supported by data)  $J/\psi$  loses status as “thermometer” of QGP

...and gains status as a powerful observable for the QCD phase boundary

Predicts  $p_T$  spectra too: hydrodynamics (MUSIC) (input for  $\beta_T$  in blast-wave formula)

[JHEP 07 \(2021\) 035, arXiv:2308.14821](#)

# SHM for charm (SHMc)

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$pQCD$  production, "throw in":  $N_{c\bar{c}} = 9.6 \rightarrow g_c = 30.1$  ( $I_1/I_0 = 0.974$ )

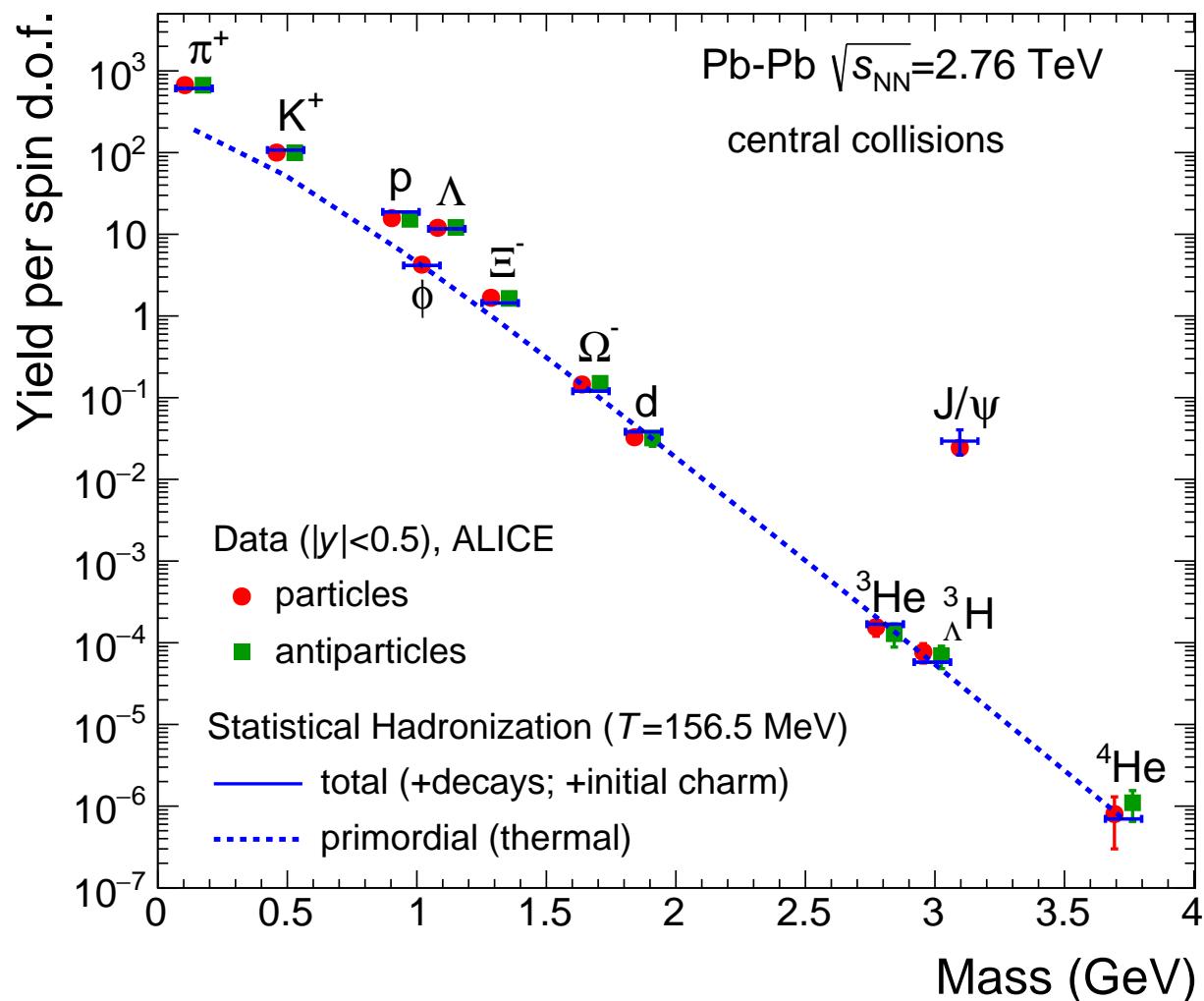
LHC, central collisions

assume:

- full thermalization of  $c, \bar{c}$   
("mobility" in  $V \simeq 4000 \text{ fm}^3$ )
- full color screening  
(Matsui-Satz)

Braun-Munzinger, Stachel, [PLB 490 \(2000\) 196](#)

Model predicts all charm  
chemistry ( $\psi(2S), X(3872)$ )



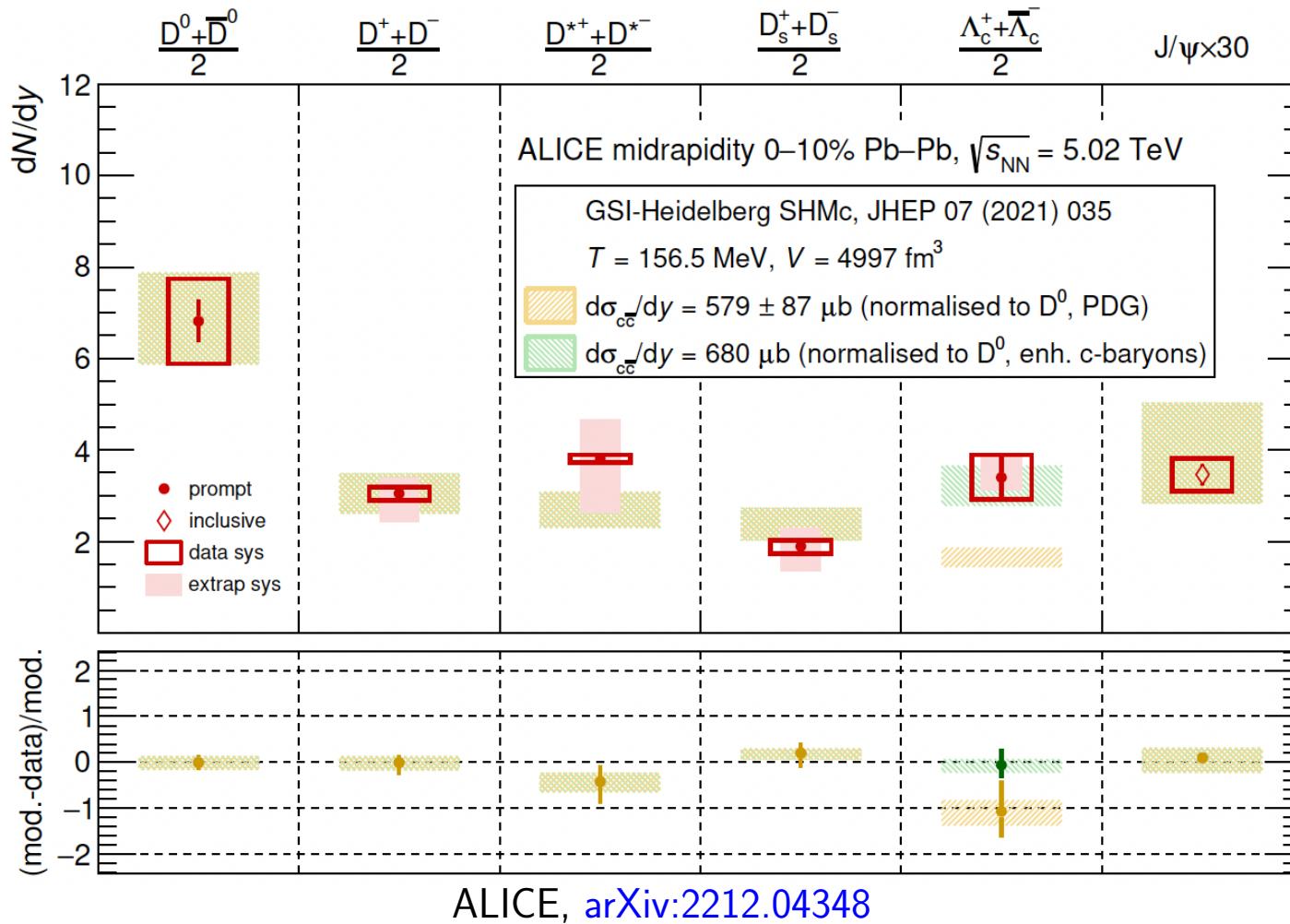
$\pi, K^\pm, K^0$  from charm included in the thermal fit  
(0.7%, 2.9%, 3.1% for  $T=156.5 \text{ MeV}$ )

[PLB 797 \(2019\) 134836](#)

# Charm data and SHMc model

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Enh. c-baryons: *tripled* the excited charm-baryon states, and  $d\sigma_{c\bar{c}}/dy$ : +19%

RQM: He,Rapp, PLB 795 (2019) 117; LQCD, Bazavov et al., PLB 737 (2014) 210

leaves the mesonic sector unaffected, for the commensurately larger  $\sigma_{c\bar{c}}$

# Transport models

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implement screening picture with space-time evolution of the fireball (hydro-like)  
continuous destruction and “(re)generation” (“recombination”)  
Boltzmann equation (loss and gain terms)

Thews et al., PRC 63 (2001) 054905 ...

“TAMU”, PLB 664 (2008) 253, NPA 859 (2011) 114, EPJA 48 (2012) 72

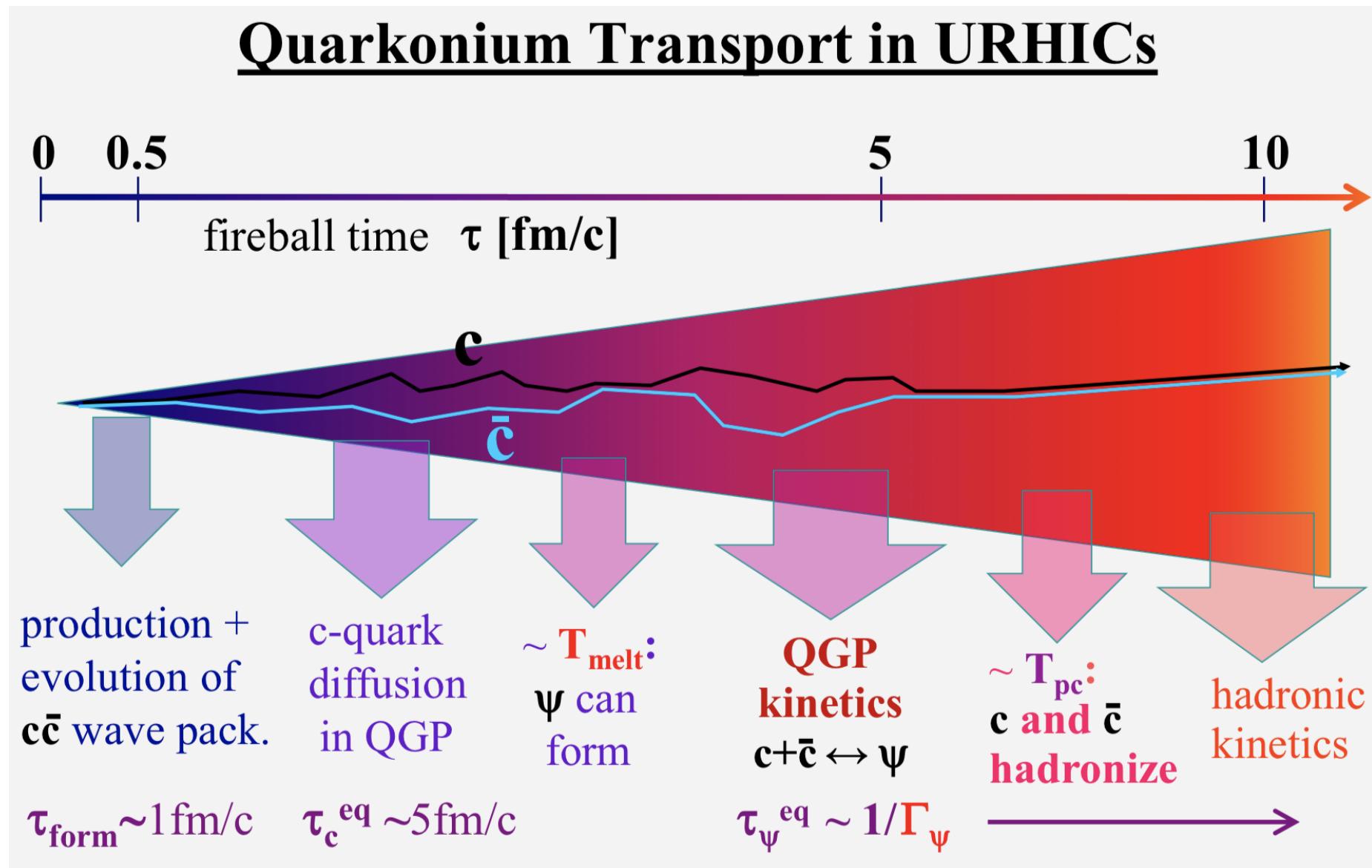
“Tsinghua”, PLB 607 (2005) 107, PLB 678 (2009) 72, arXiv:1401.5845

Predicts  $R_{AA}$ ,  $v_2(p_T)$  (TAMU describes D mesons too)

# Transport models - schematics

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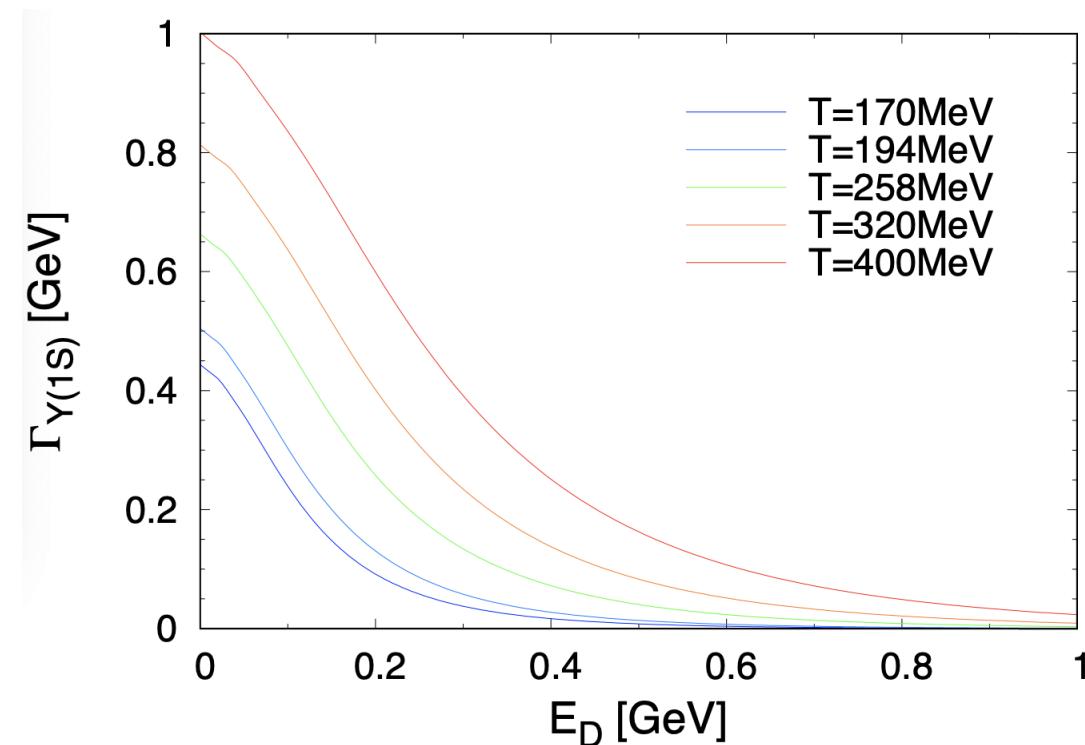
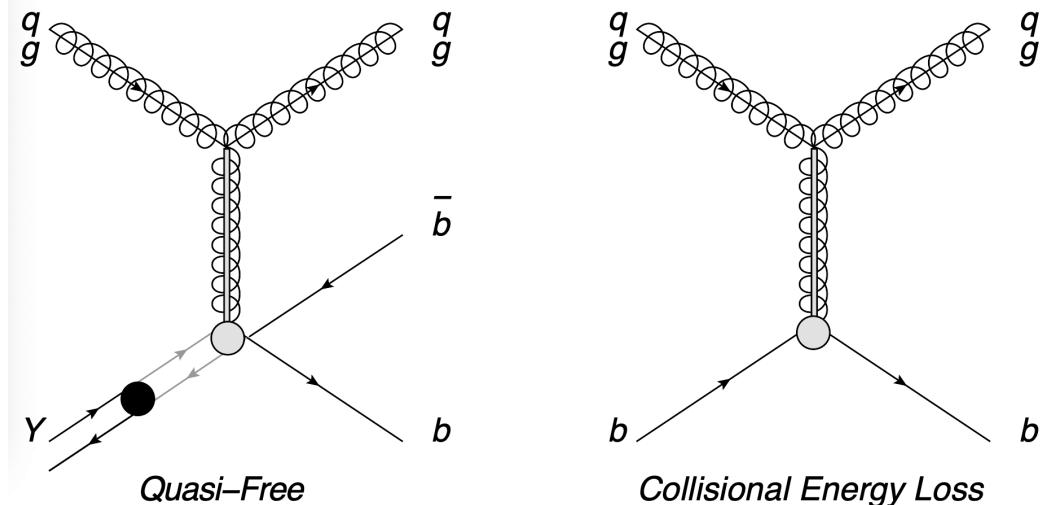


# Transport models - the essentials

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$$\frac{dN_\gamma}{d\tau} = -\Gamma(T(\tau)) [N_\gamma(\tau) - N_\gamma^{eq}(T(\tau))]$$



Du, Liu, Rapp, [Phys. Lett. B 796 \(2019\) 20](#)

Quasi-free = inelastic (dissociation) = imaginary part of HQ potential

$K$  factor: enhancement over perturbative results ( $K = 5$  in  $\Gamma_{Y(1S)}$  above)

$N^{eq}$ : SHM generation

## 2 other models

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- “Comovers” model (Santiago):  
( invented to describe suppression in the (final-state) hadronic medium at SPS  
 $J/\psi + \pi \rightarrow D + \bar{D}$  )  
at LHC: gluo-dissociation and a (re)generation component (dominant for  $J/\psi$ )  
(the Boltzmann equation of the transport model is also in the comoving system)

E. Ferreiro et al., PLB 731 (2014) 57

- Hydrodynamic model (Kent State Univ.)  
hydro gives energy density vs. space-time  
suppression probability vs.  $\varepsilon$  gives  $R_{AA}$  (of  $Y$ )

Strickland, Bazow, NPA 879 (2012) 25

## ...and one more: the quantum approach

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Open quantum systems description + Effective NRQCD theory

quantum evolution of the  $b\bar{b}$  pair in QGP (heat bath; hydrodynamics)

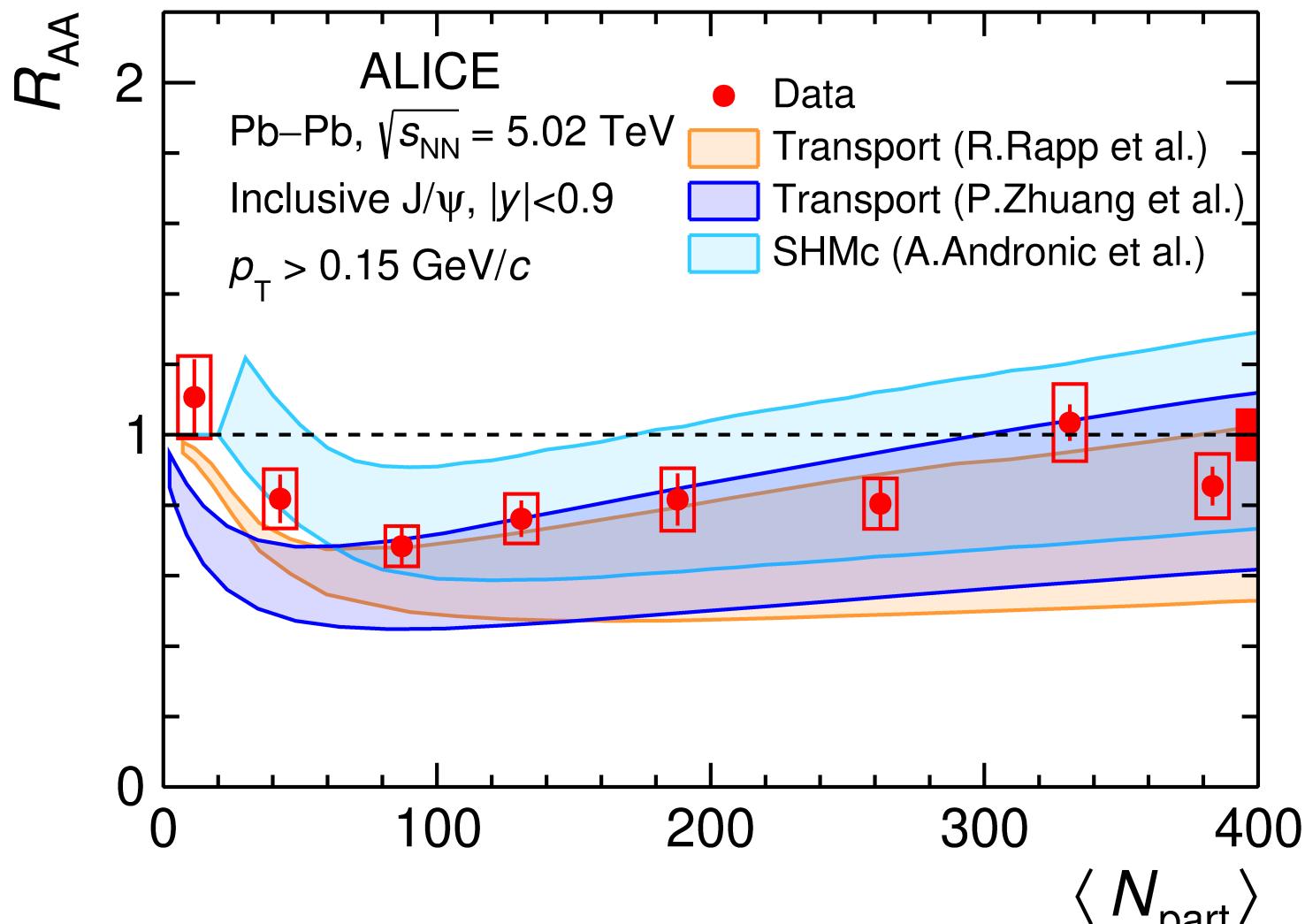
Lindblad equation (accuracy: at NLO in the binding energy over temperature)

Strickland, Thapa, PRD 108 (2023) 014031

# SHMc vs. transport models

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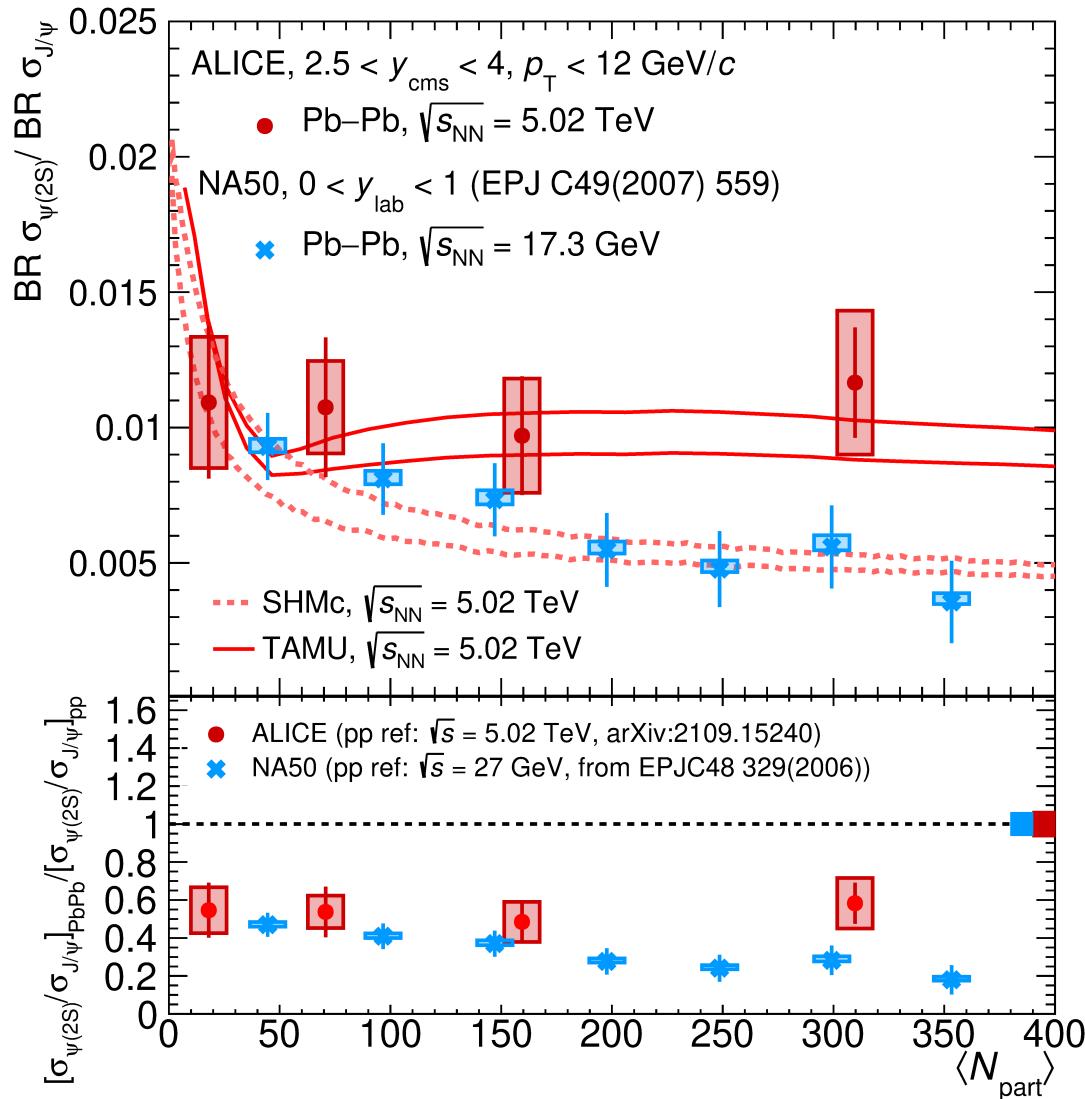
SHMc:  $d\sigma_{c\bar{c}}/dy$  via normalization to  $D^0$  in Pb-Pb 0-10%, ALICE, [arXiv:2110.09420](https://arxiv.org/abs/2110.09420)

$dN/dy = 6.82 \pm 1.03$  ( $|y| < 0.5$ ; FONLL for  $y=2.5-4$ ; assuming hadronization fractions in data as in SHMc)

# $\psi(2S)/J/\psi$ at the LHC (and SPS)

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ALI-PUB-528400

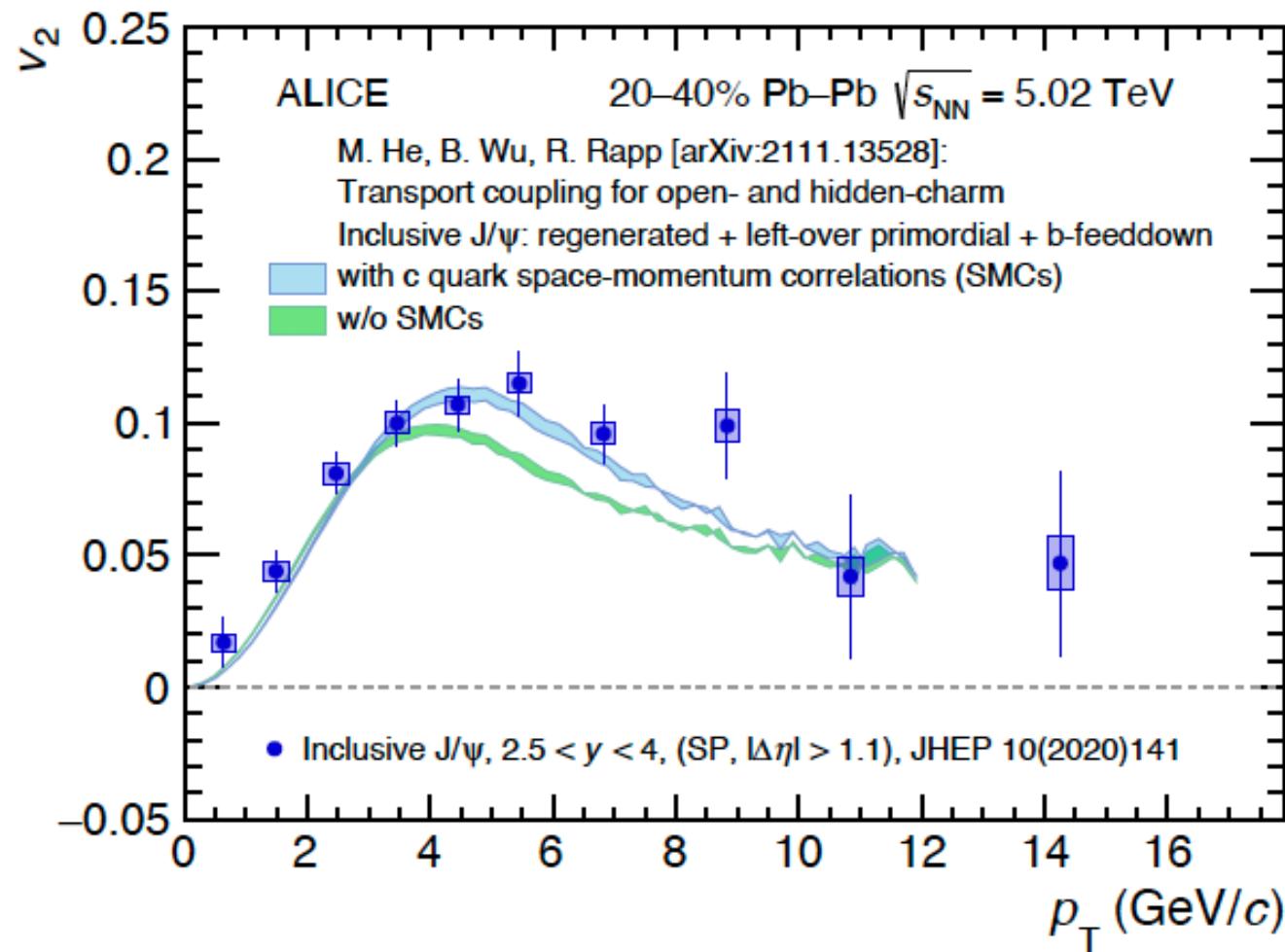
ALICE, arXiv:2210.08893

In SHMc uncertainty only due to nuclear-corona  
( $\sigma_{c\bar{c}}$  cancels out completely)

# $\text{J}/\psi$ elliptic flow: data and transport model

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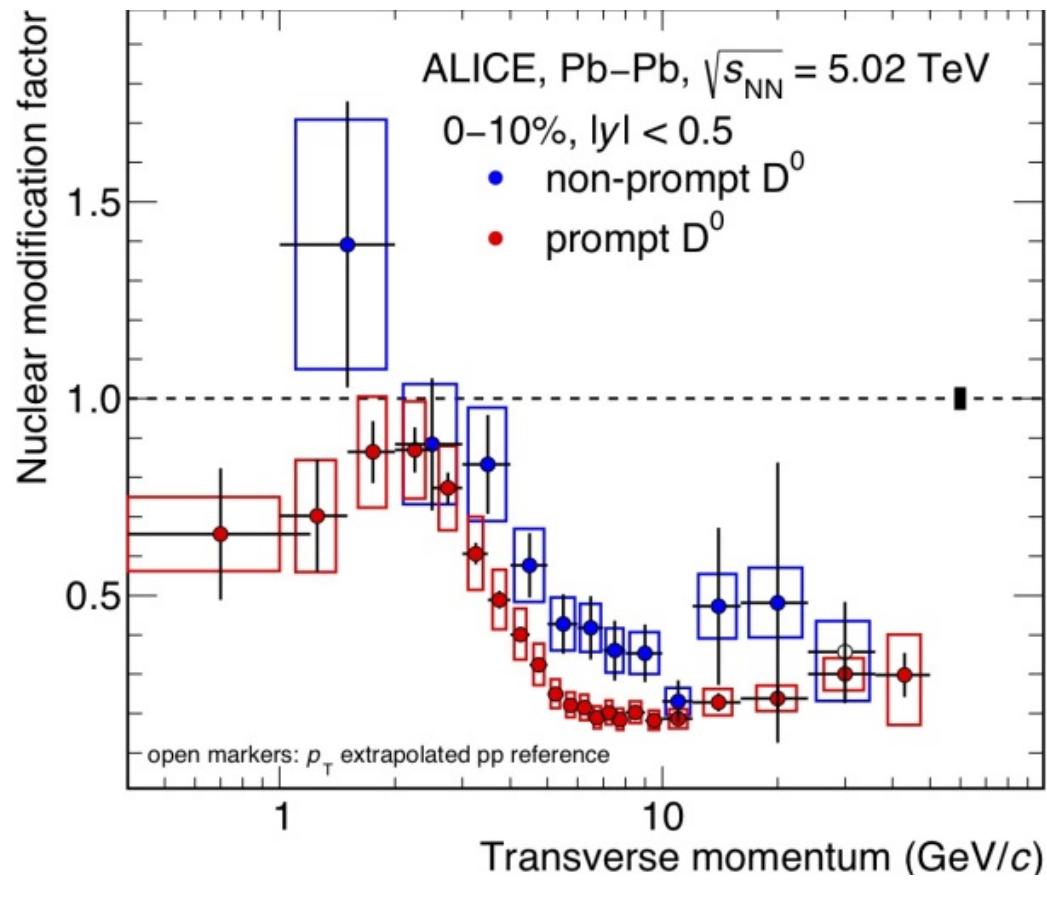
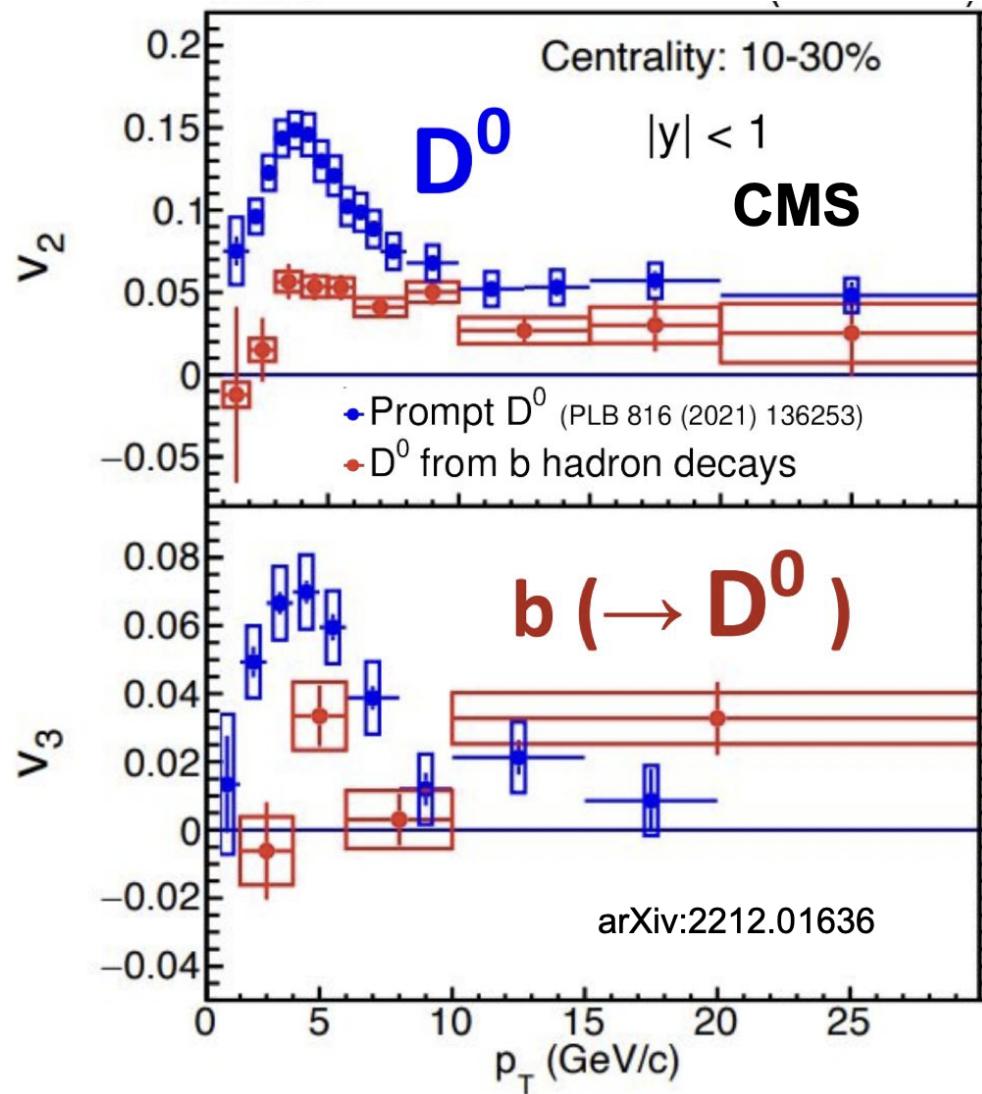
ALICE, [arXiv:2212.04348](https://arxiv.org/abs/2212.04348)

a very good description of data by the TAMU model

# Beauty quark thermalization?

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ALICE, arXiv:2202.00815, ATLAS

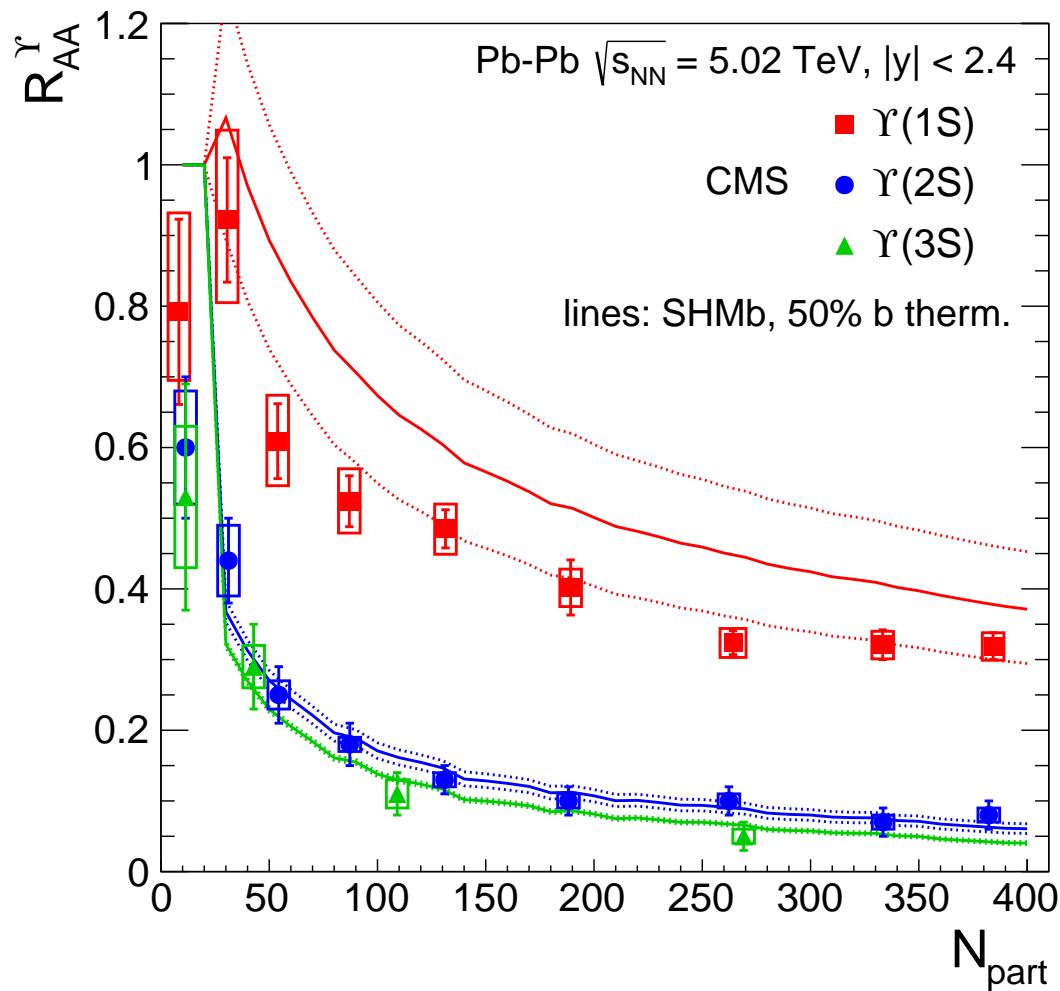
definitely strong flow but clearly less strong than for charm (CMS, QM'22, HIN-21-003)

...and a strong coupling with the medium (less energy loss than charm,  $p_T \simeq 10$  GeV/c)

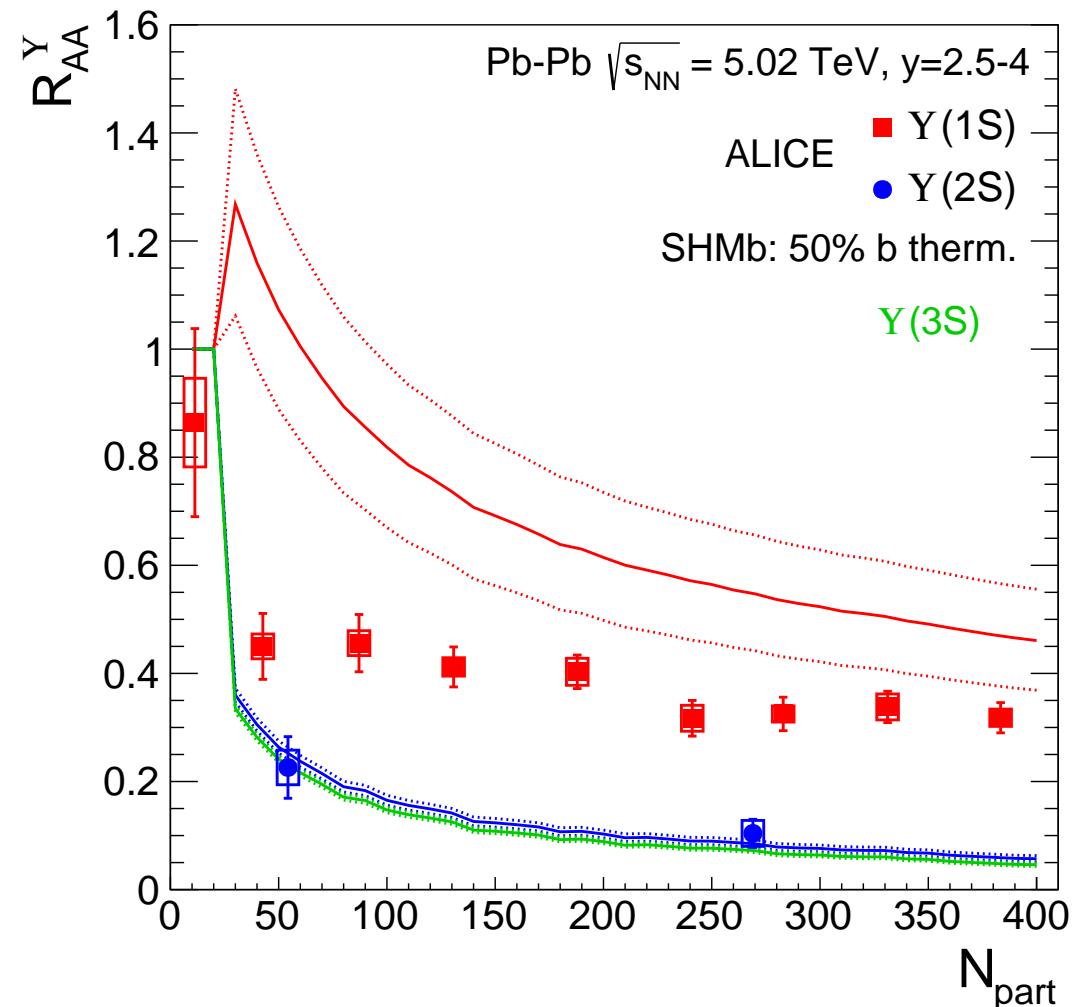
# $R_{AA}$ , 50% $b\bar{b}$ thermalized

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CMS, PRL 120 (2018) 142301



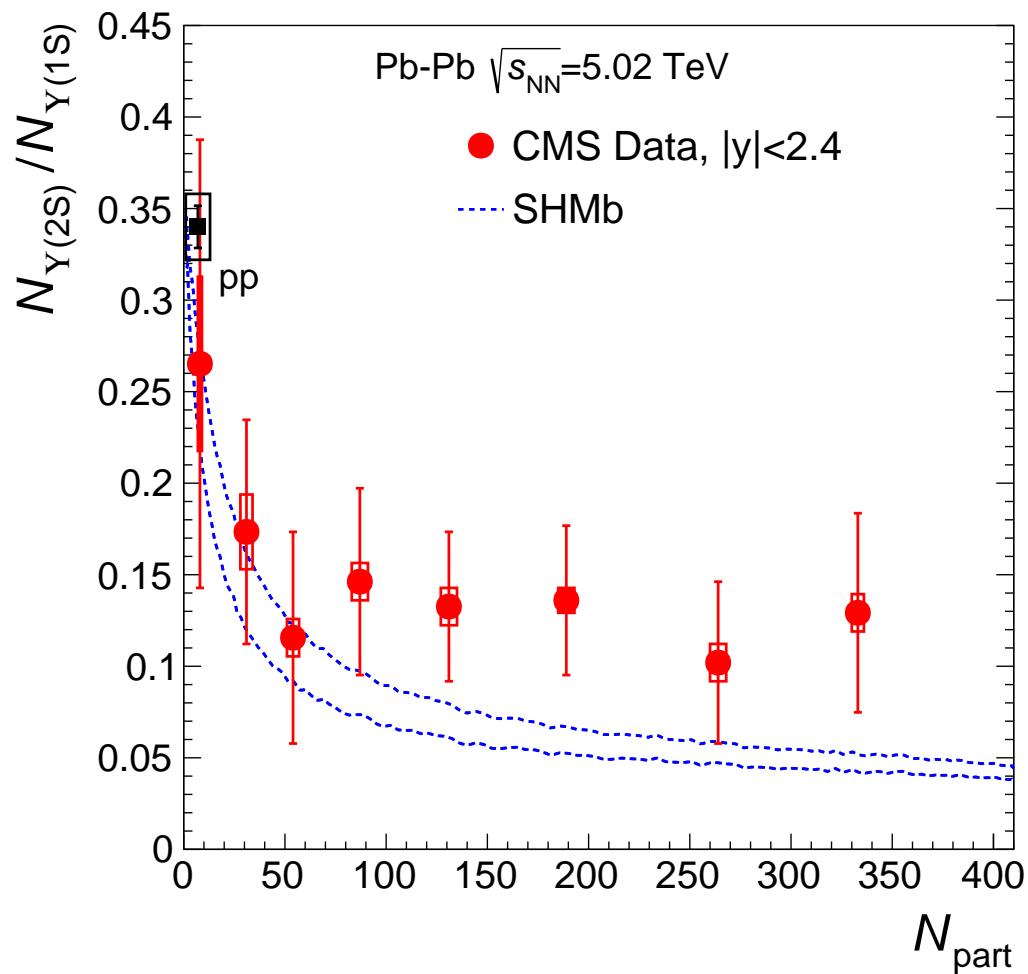
ALICE, PLB 822 (2021) 136579

*What does non-thermalized beauty produce? (no room for it in SHMb)*

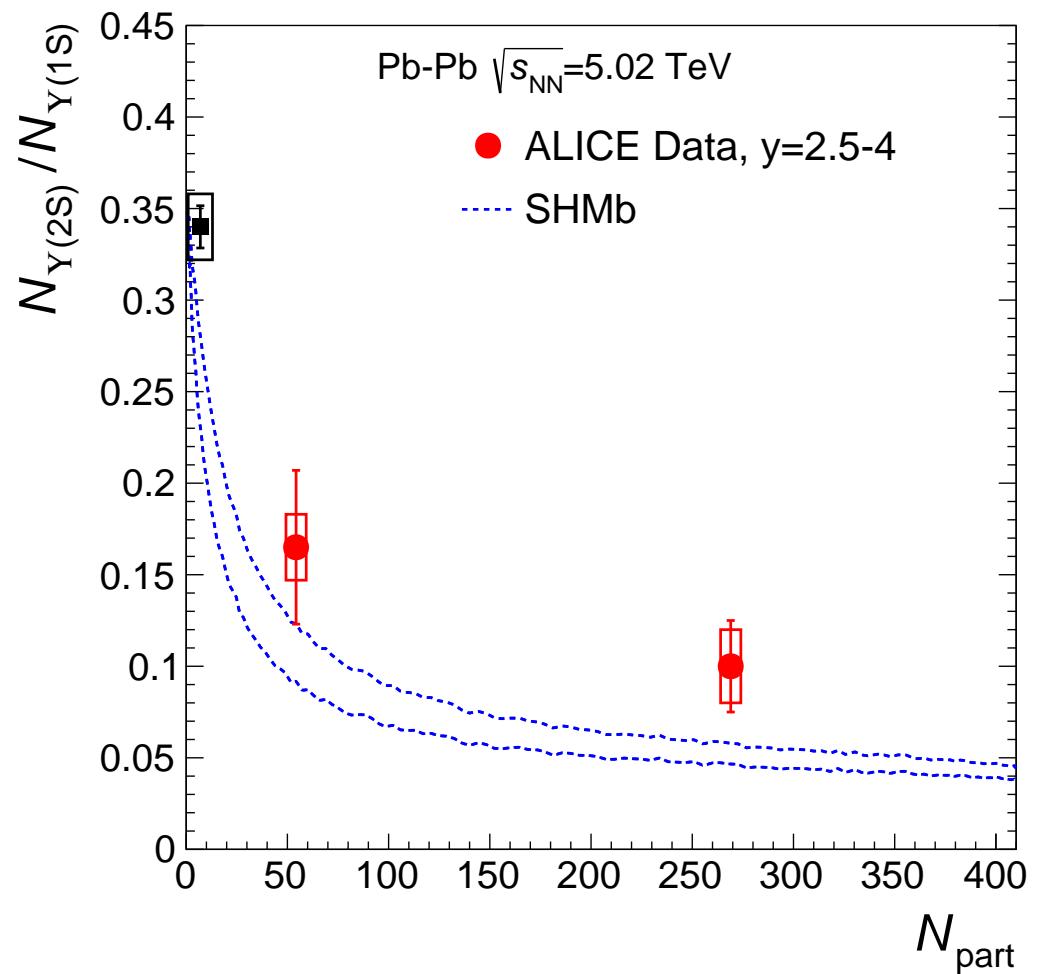
# $\Upsilon(2S)/\Upsilon(1S)$ ratio (100% b thermalization)

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CMS, PRL 120 (2018) 142301



ALICE, PLB 822 (2021) 136579

ALICE pp:  $\Upsilon(2S)/\Upsilon(1S) = 0.5 \pm 0.1$ , arXiv:2109.15240

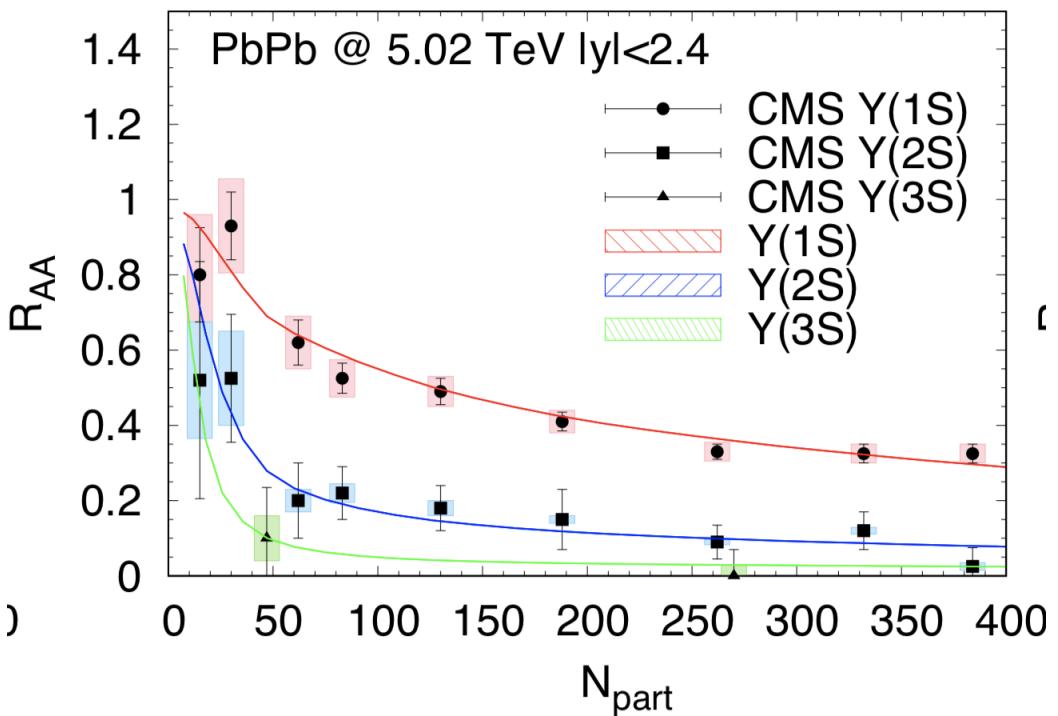
SHMb uncert.: nuclear-corona (fraction)

# $\Upsilon$ description in the transport model

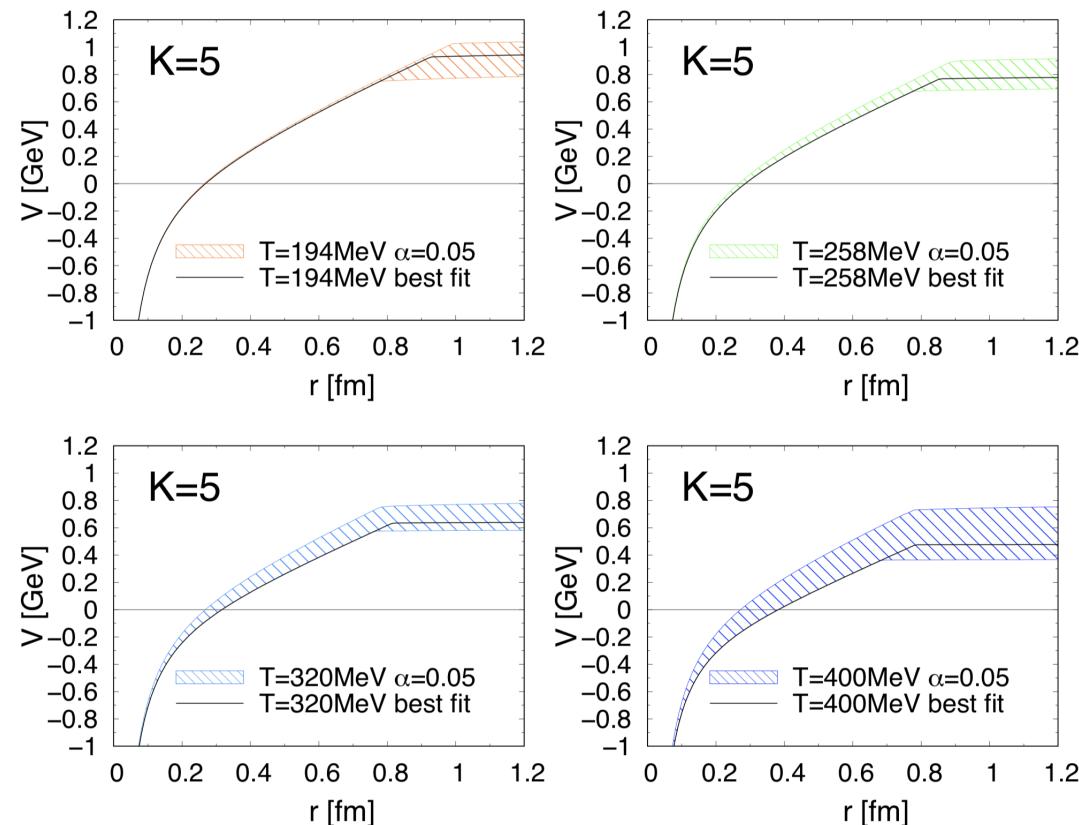
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...is very good; allows extraction of in-medium (Cornell) potential



(re)generation important for  $\Upsilon(2S)$



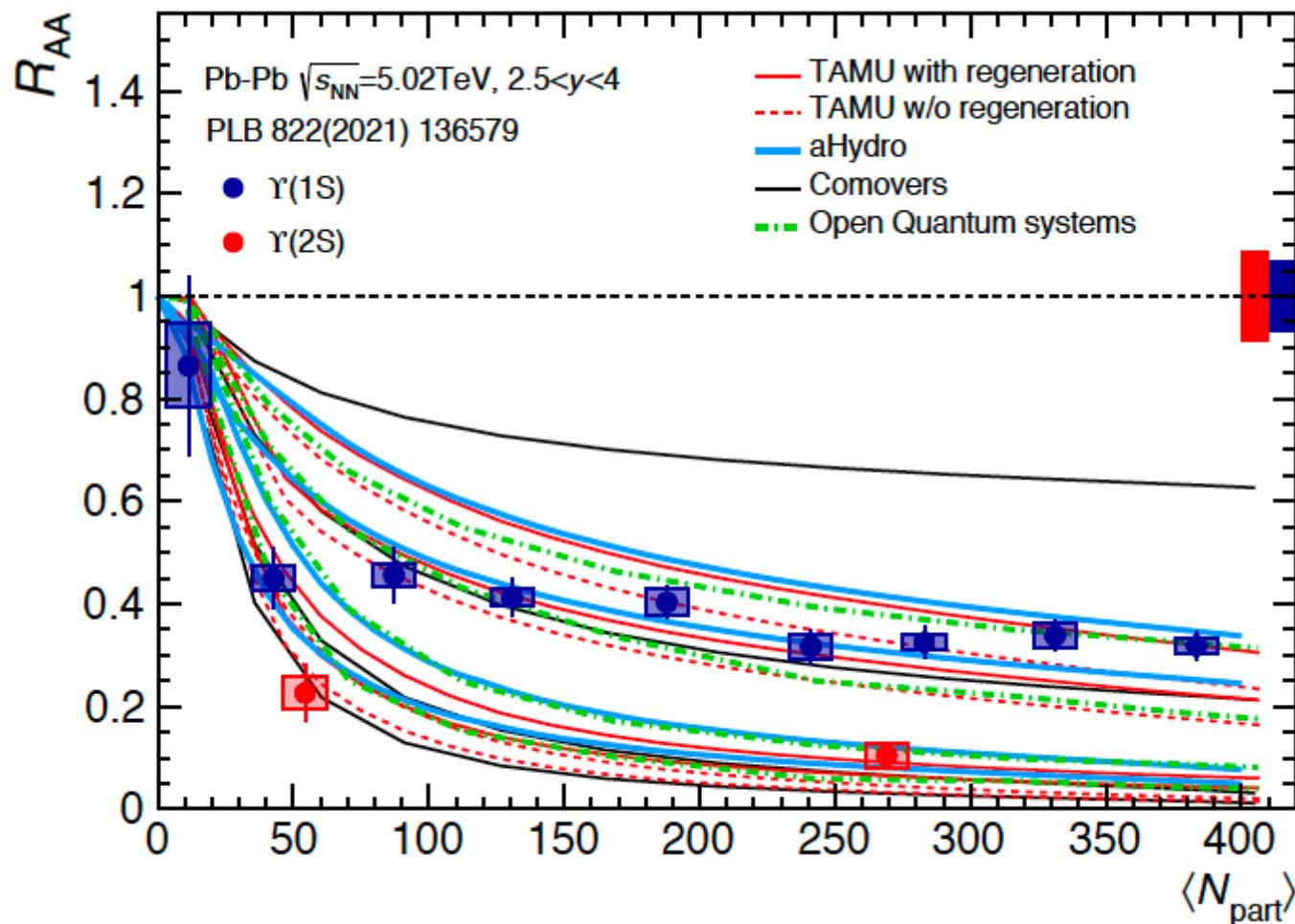
Transport Model (TAMU), Du, Liu, Rapp, Phys. Lett. B 796 (2019) 20

*Substantial remnants of the long-range color confining force in QGP*

# $\Upsilon$ suppression data and models

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ALICE, [arXiv:2212.04348](https://arxiv.org/abs/2212.04348)

All models (except perhaps Comovers ...large uncert.) reproduce the data well  
TAMU: Regeneration important for  $\Upsilon(2S)$

# Summary / Conclusions: charm

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- Charm quarks seem to thermalize very effectively (close to 100%) in QGP
- SHM: all charmonium *and open charm states* are generated exclusively at hadronization (chemical freeze-out) ...full color screening

The model is very successful in reproducing the  $J/\psi$  and open charm data  
*A handle for hadronization  $T$  with a mass scale well above  $T$*

- Transport models: continuous  $J/\psi$  destruction and (re)generation in QGP  
(only up to 2/3 of the  $J/\psi$  yield (LHC, central collisions) originates from deconfined  $c$  and  $\bar{c}$  quarks)

*Discriminating the two pictures implies providing an answer to fundamental questions related to the fate of hadrons in a hot deconfined medium.*

A precision ( $\pm 10\%$ ) measurement of  $d\sigma_{c\bar{c}}/dy$  in Pb-Pb (Au-Au) collisions needed for a stringent test  
(within reach with the upgraded detectors at the LHC and RHIC)

# Summary / Conclusions: beauty

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- Full beauty thermalization seems not realized in nature
  - ...with 30-50% of beauty quarks fully thermalized SHM can explain the  $\Upsilon$  data
  - ...but this fraction is (significantly) dependent on the b-hadron spectrum

What does non/partially-thermalized beauty produce?

no  $\Upsilon$  because strong coupling with the medium destroys the  $b-\bar{b}$  correlation?

- Transport and Hydro models are successfully reproducing  $\Upsilon$  suppression
  - Transport: regeneration important for  $\Upsilon(2S)$  (at the LHC)
- Quantum approaches on strong rise

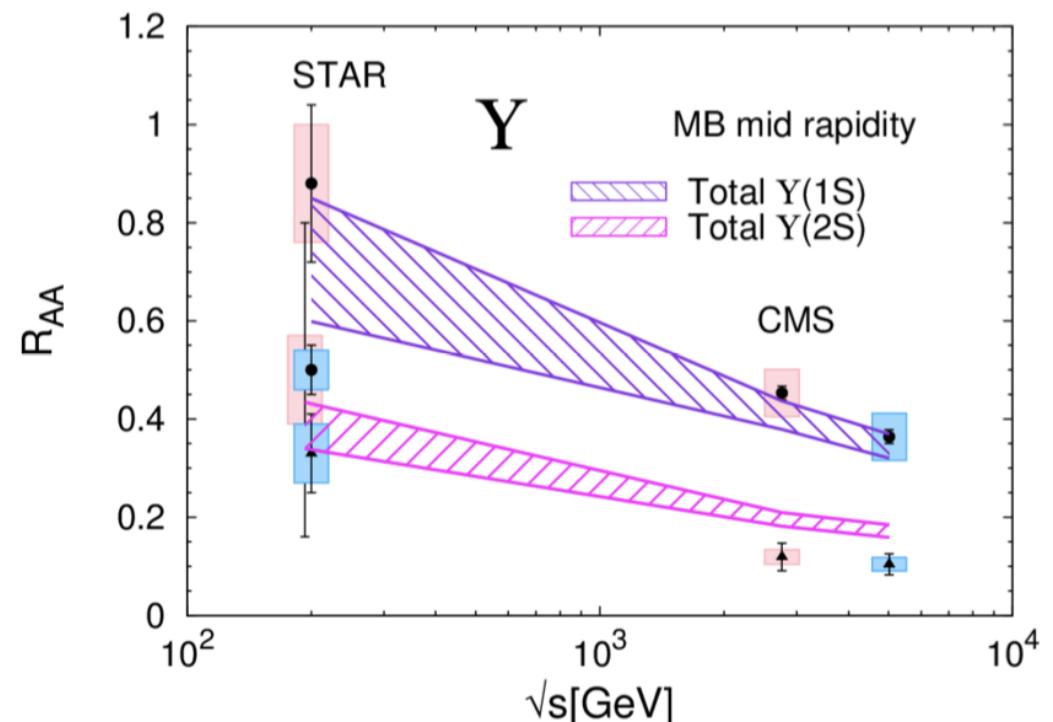
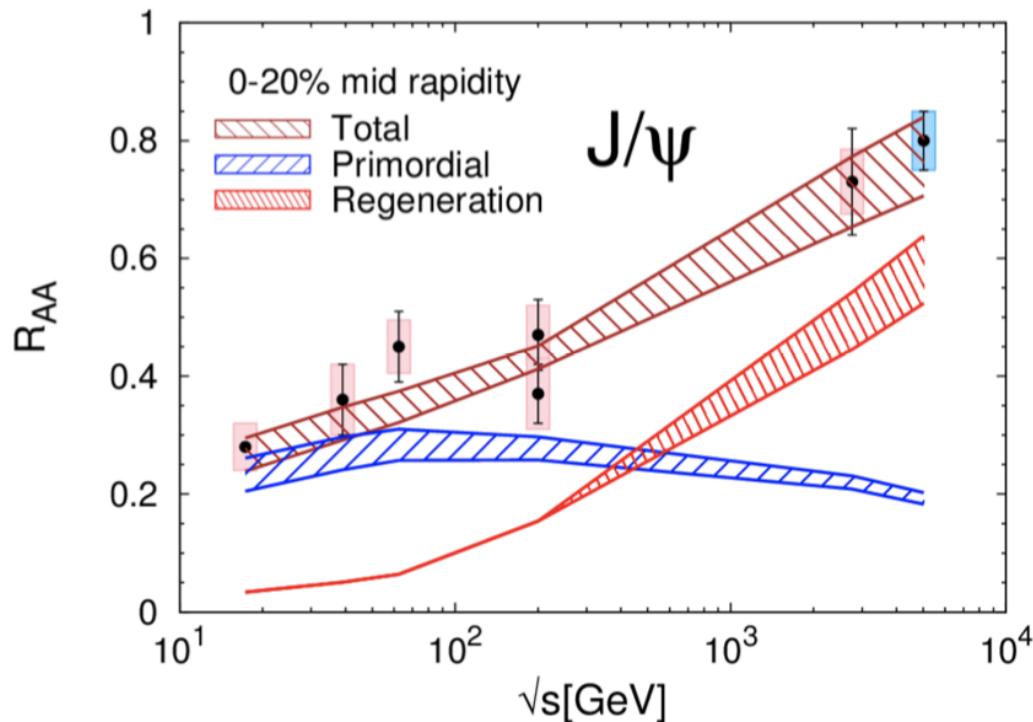
## Additional material

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# Fractions primordial, (re)generated - energy dependence

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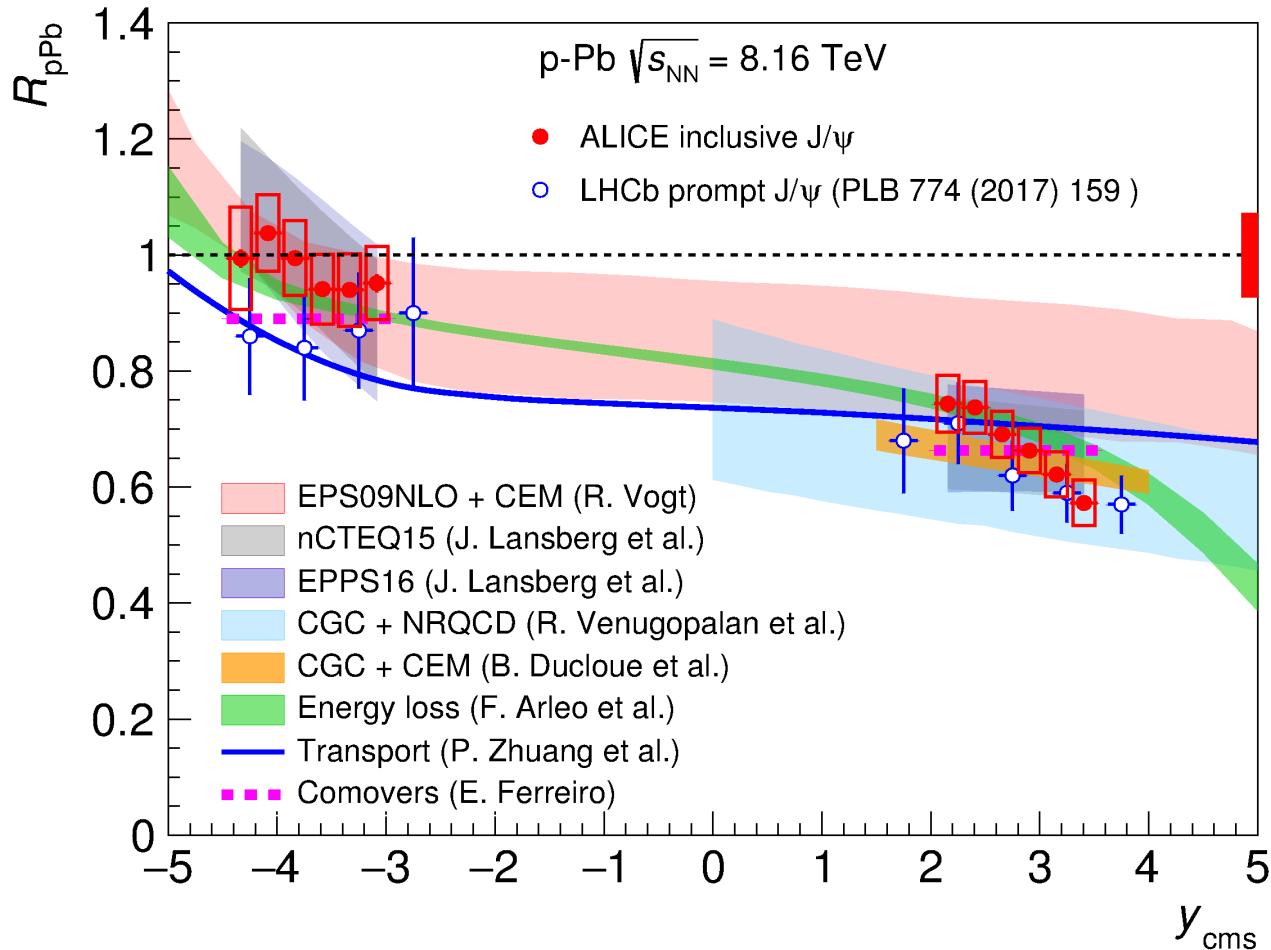


Rapp, Du, [arXiv:1704.07923](https://arxiv.org/abs/1704.07923)

# $\text{J}/\psi$ production in p–Pb collisions

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$$R_{\text{pPb}} = \frac{dN_{\text{pPb}}/dp_{\text{T}}dy}{\langle N_{\text{coll}}^{\text{pPb}} \rangle \cdot dN_{\text{pp}}/dp_{\text{T}}dy}$$

$$\langle N_{\text{coll}}^{\text{pPb}} \rangle \simeq 7$$

Shadowing describes data  
(shadowing uncert. are large)

Color Glass Condensate also  
successful

ALICE, [JHEP 07 \(2018\) 160](#)

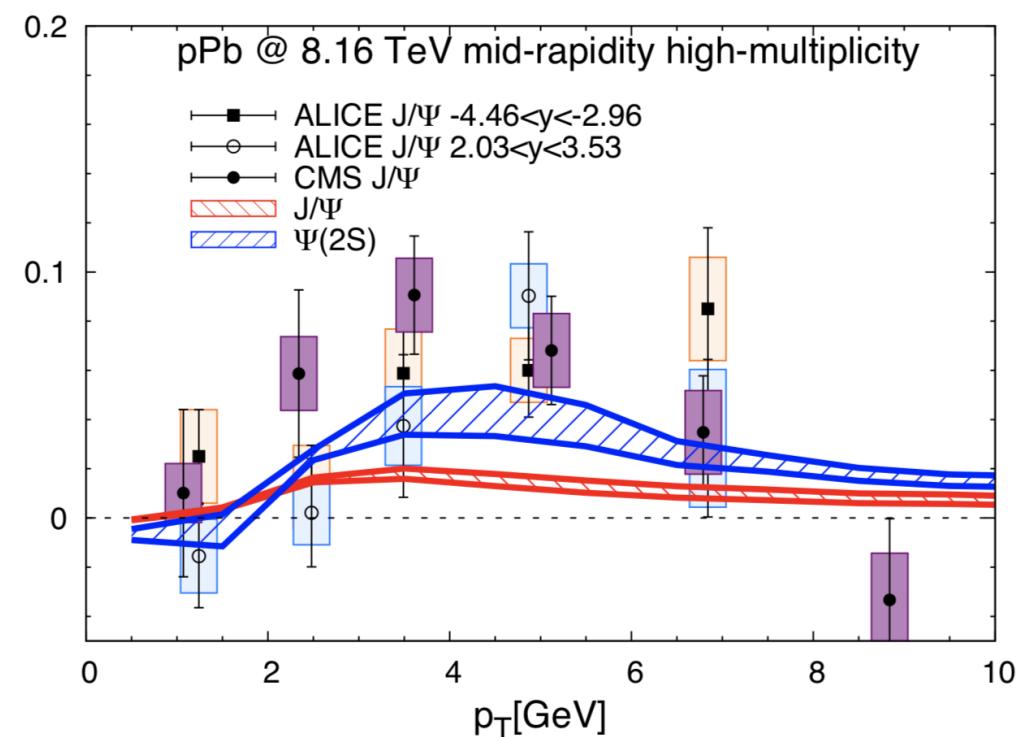
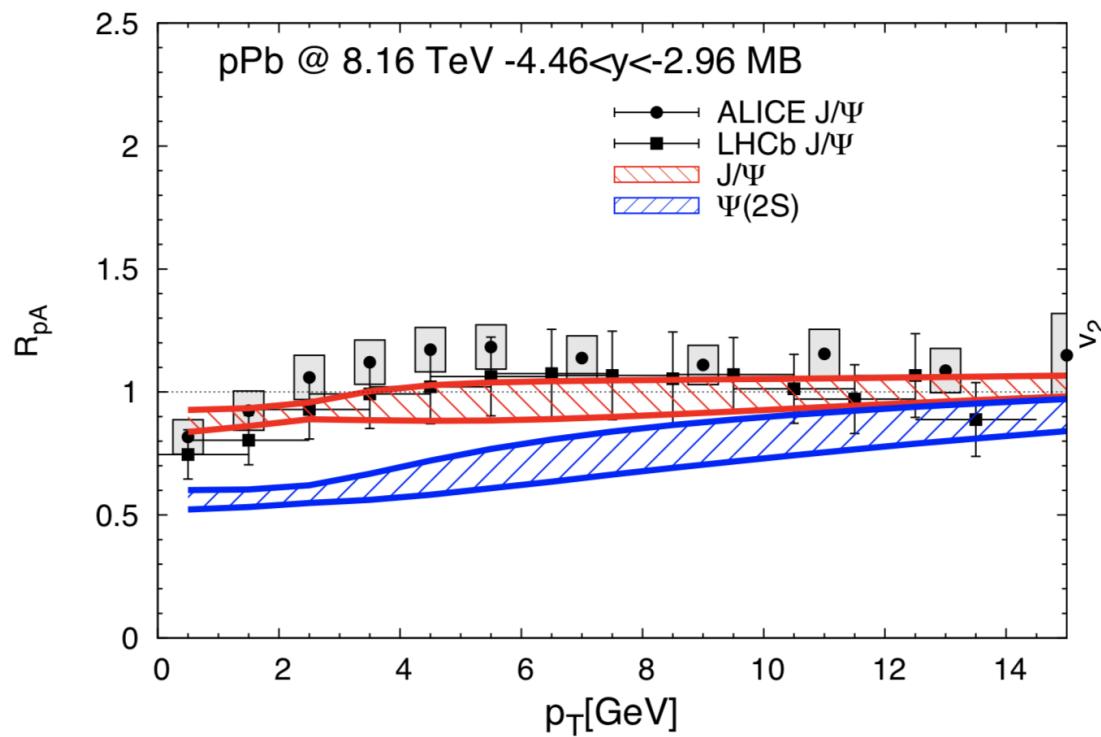
LHCb, [PLB 774 \(2017\) 159](#)

Seen also with Run 1 data (5.02 TeV): ALICE, [JHEP 02 \(2014\) 073](#), [06 \(2015\) 55](#)

# $J/\psi$ production in p–Pb collisions

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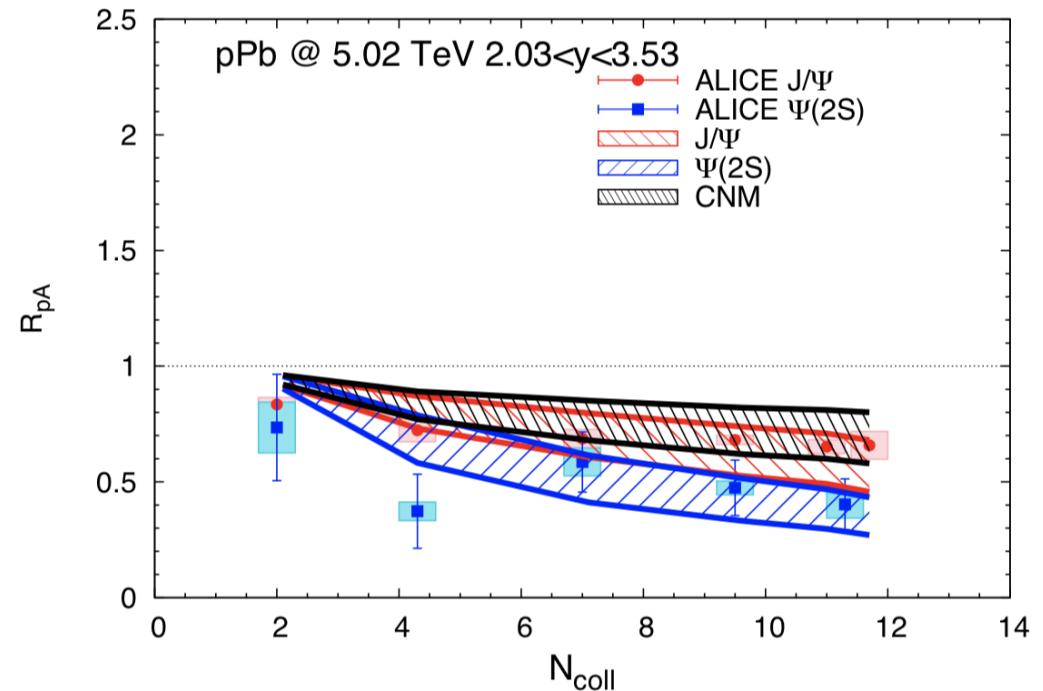
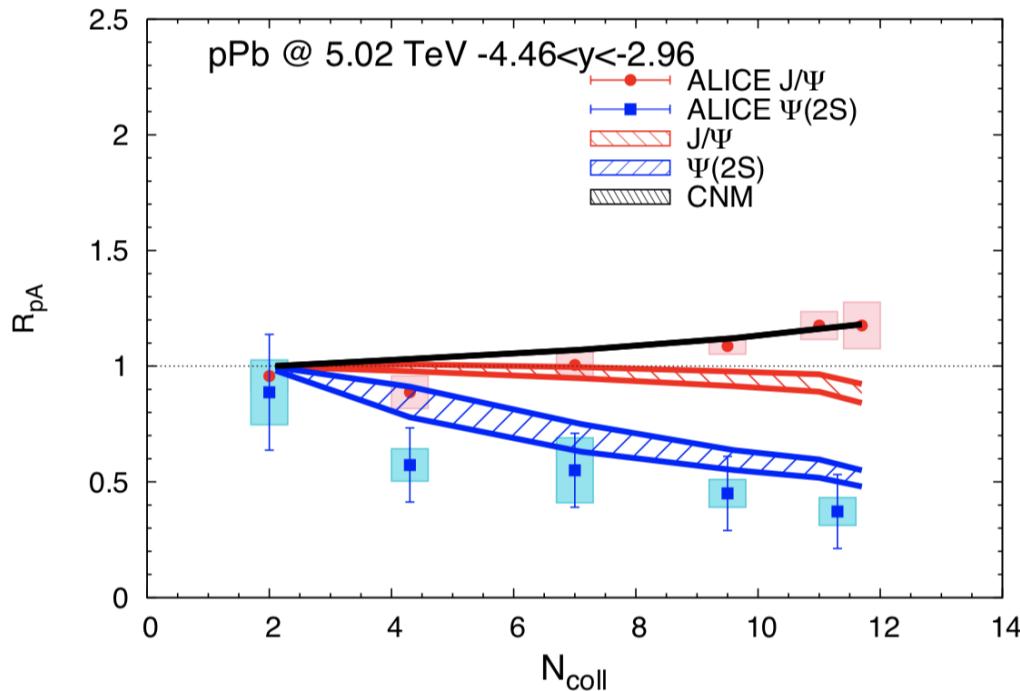
Bands: Transport Model, TAMU Du, Rapp, [JHEP 1903 \(2019\) 015](#)

*Need experimentally (in reach for Run 3,4): better precision; also  $v_3$ ; separate  $B$  component;  $v_2$  of  $\psi(2S)$  ?*

# $J/\psi$ and $\psi(2S)$ production in p–Pb collisions

A. Andronic

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Du, Rapp, JHEP 1903 (2019) 015