

# Study of $\langle p_T \rangle$ scaling with $m/n_q$ in relativistic heavy-ion collisions

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# Outline

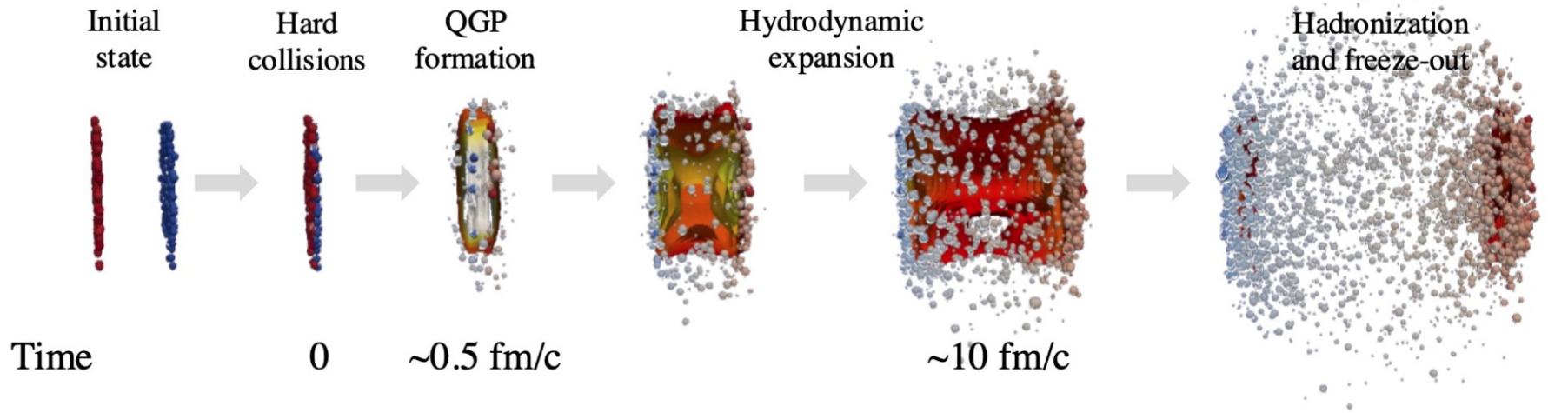
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Motivation and data used

Results

Conclusions

# Heavy-ion collision evolution



Visualization by J.E. Bernhard, arXiv:1804.06469

- thermal (kinetic) “freeze-out” (FO) --> kinetic FO temperature and average transverse flow velocity
- the collective transverse expansion of the system --> entirely generated and develops throughout the entire evolution of the system created in the collision
- the shape of the transverse momentum distributions of the identified charged hadrons --> sensitive to the dynamics of the nucleus-nucleus collisions.

# Particle source

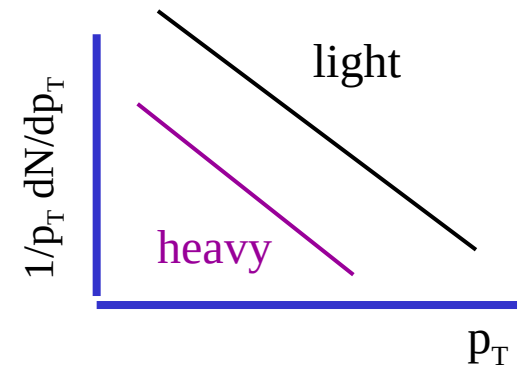
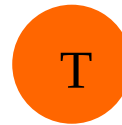
Final state spectra → system properties at thermal freeze-out

Thermal source → spectrum slope reflects the temperature of the fireball

$$\frac{dN}{m_T dm_T} \sim e^{-m_T/T_{\text{slope}}}$$

$$\text{where } m_T \equiv \sqrt{p_T^2 + m^2}$$

purely thermal source



# Particle source

Final state spectra → system properties at thermal freeze-out

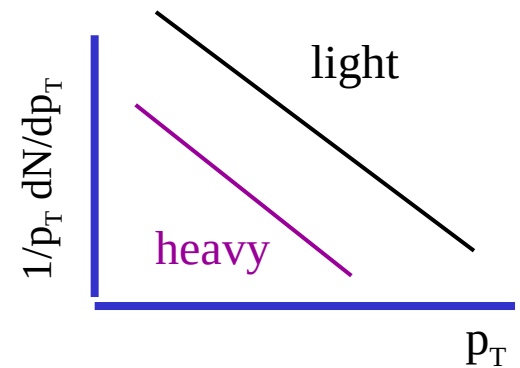
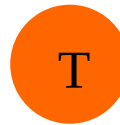
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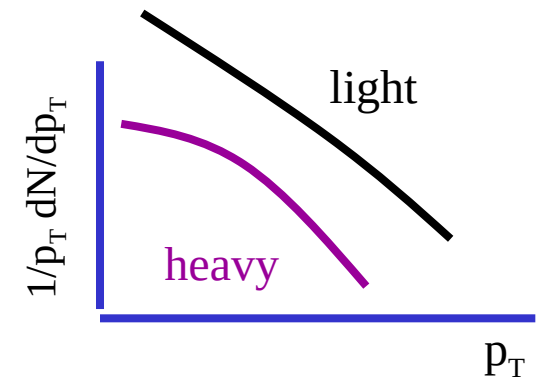
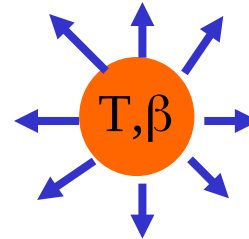
$$\text{where } m_T \equiv \sqrt{p_T^2 + m^2}$$

Data → shape is different in p-p and A-A → stronger effect for heavier particles

purely thermal source



explosive source



Flow → collective motion of particles (due to high pressure arising from compression and heating of nuclear matter) superimposed on thermal motion

# Data

STAR-BES Au-Au data at  $\sqrt{s_{NN}} = 7.7, 11, 19.6, 27$  and  $39$  GeV

The average transverse momentum:

$$\langle p_T \rangle = \frac{\int_0^\infty p_T (2\pi p_T) f(p_T) dp_T}{\int_0^\infty (2\pi p_T) f(p_T) dp_T}.$$

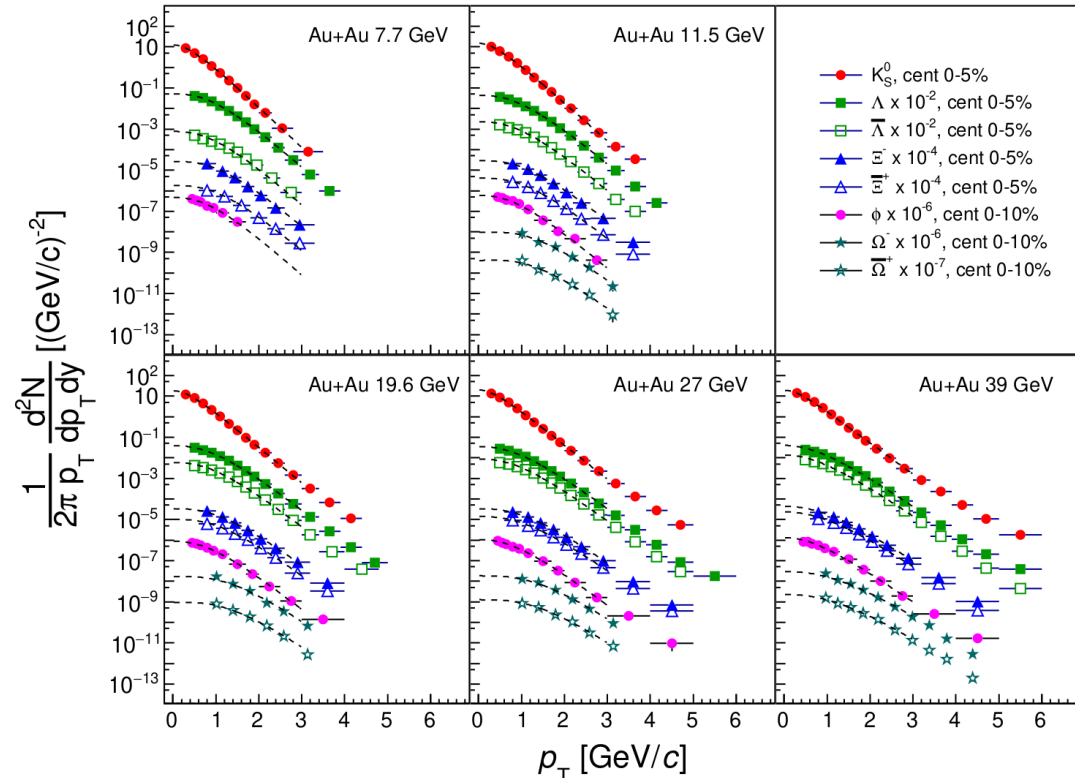
For  $K_S^0, \phi, \Omega \rightarrow$  an exponential function:

$$\frac{d^2 N}{2\pi p_T dp_T dy} = A \exp\left(-\frac{m_T - m}{T}\right)$$

For  $\Lambda, \Xi \rightarrow$  Boltzmann function:

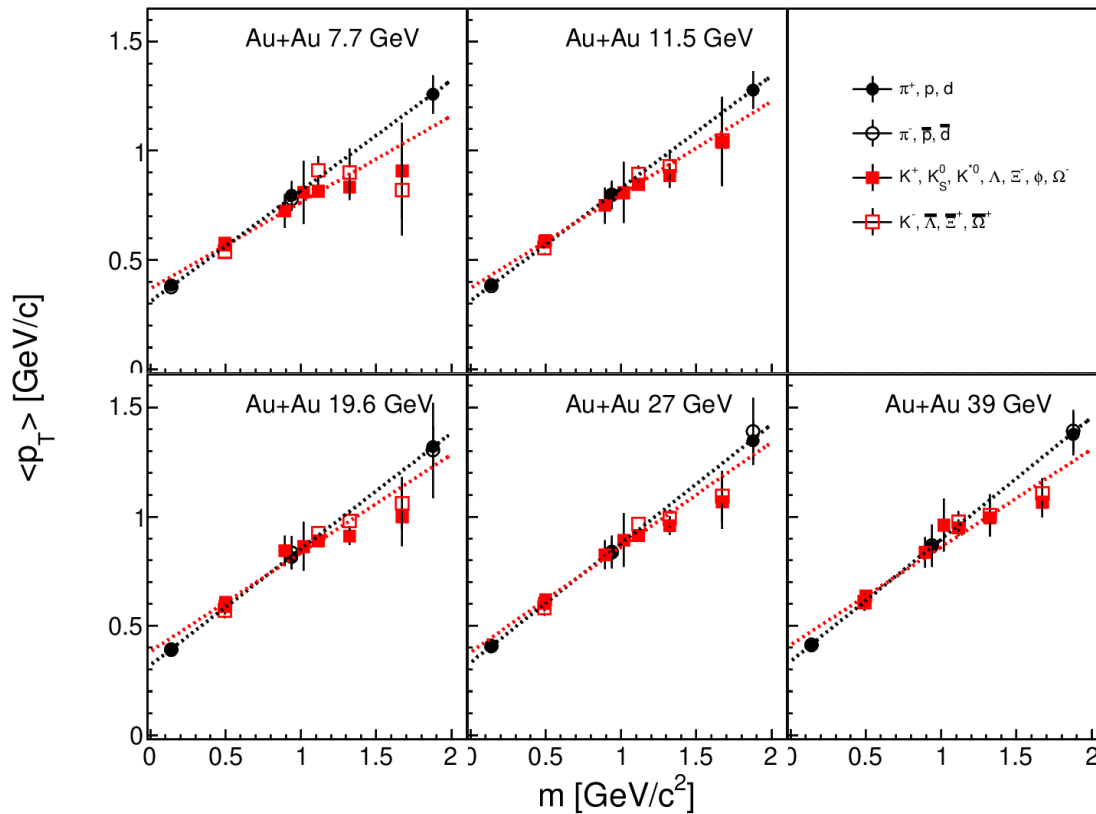
$$\frac{d^2 N}{2\pi p_T dp_T dy} = A m_T \exp\left(-\frac{m_T - m}{T}\right)$$

Fit range:  $0 < p_T < 2$  GeV/c (for  $K_S^0, \phi, \Lambda$ ) and  $0 < p_T < 2.6$  GeV/c ( $\Xi, \Omega$ )



J. Adam et al., STAR coll., Phys.Rev.C 102 (2020),034909;  
L. Adamczyk et al., Phys. Rev. C 93, 21903 (2016)

# Mean transverse momentum



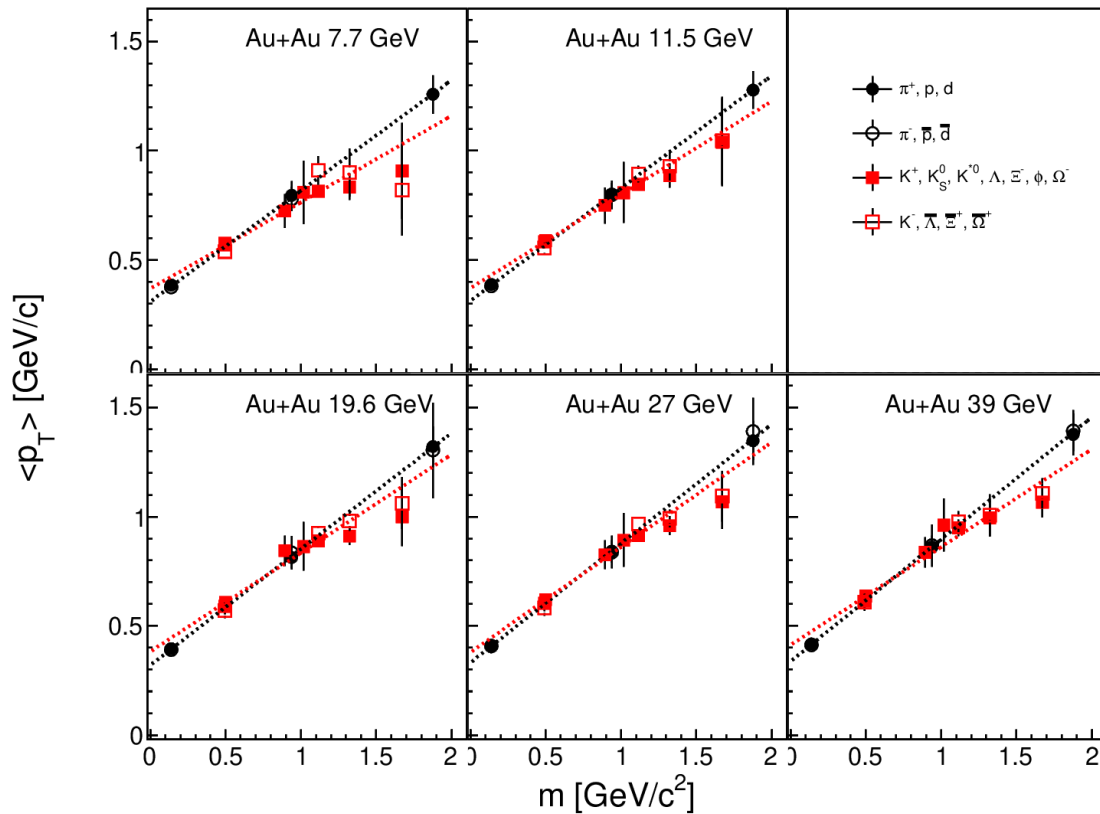
$\langle p_T \rangle$  reflects the slopes of the  $p_T$  spectra

$\langle p_T \rangle$  increases with the mass for all RHIC-BES energies

central Au-Au collisions  
(0-5%/0-10% centrality class)

J. Adam et al., STAR coll., Phys.Rev.C 102 (2020),034909; L. Adamczyk et al., Phys. Rev. C 96, 044904 (2017); L. Adamczyk et al., Phys. Rev. C 93, 21903 (2016); J. Adams et al, STAR coll., Phys. Rev. C 99, 064905 (2019), M. S. Abdallah et al., STAR coll., Phys. Rev. C 107 (2023) 34907

# Mean transverse momentum

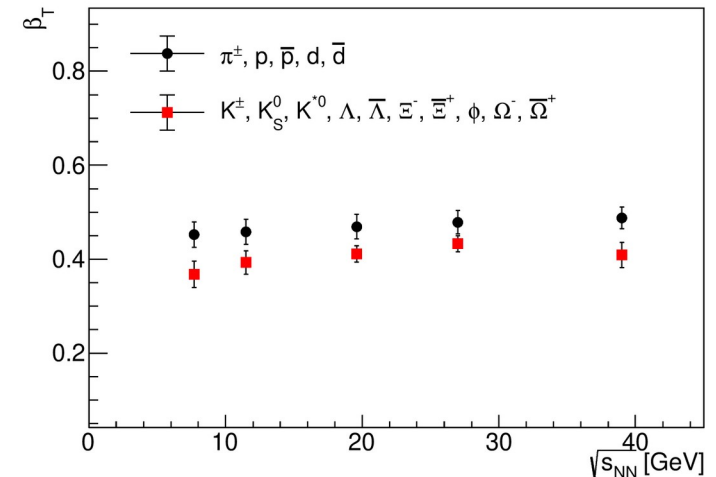


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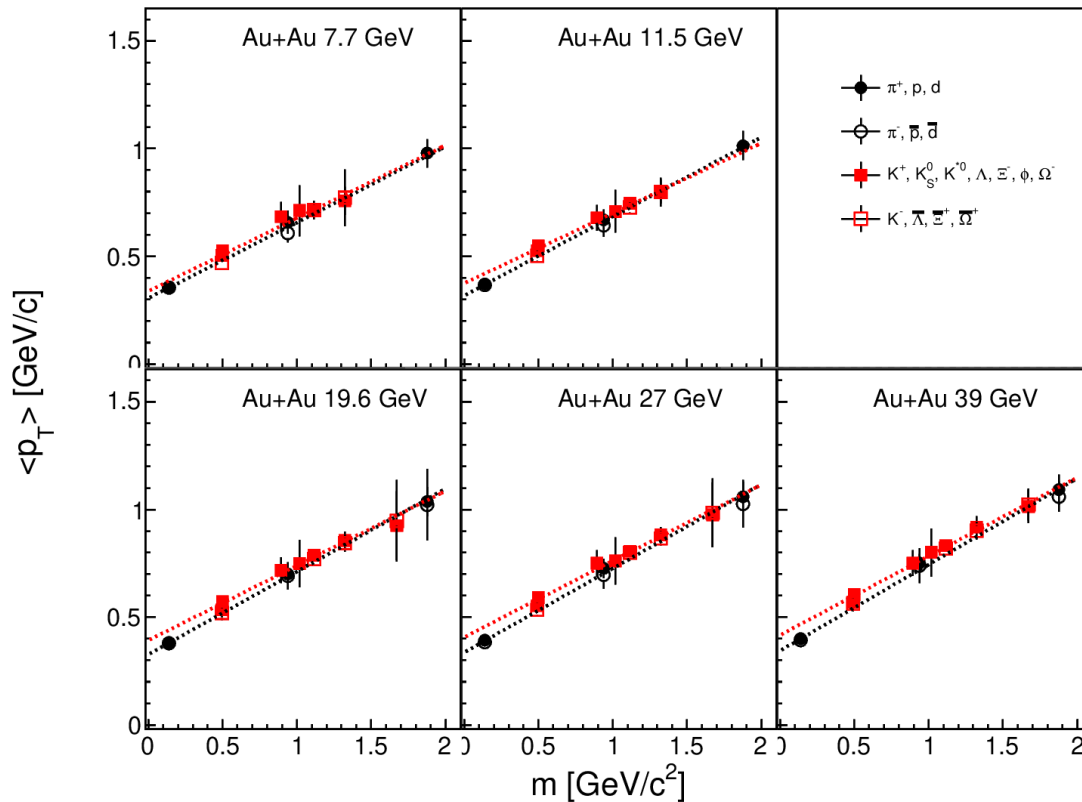
$$\langle p_T \rangle = C' + m \beta_T' = C' + m_0 \gamma_T' \beta_T'$$



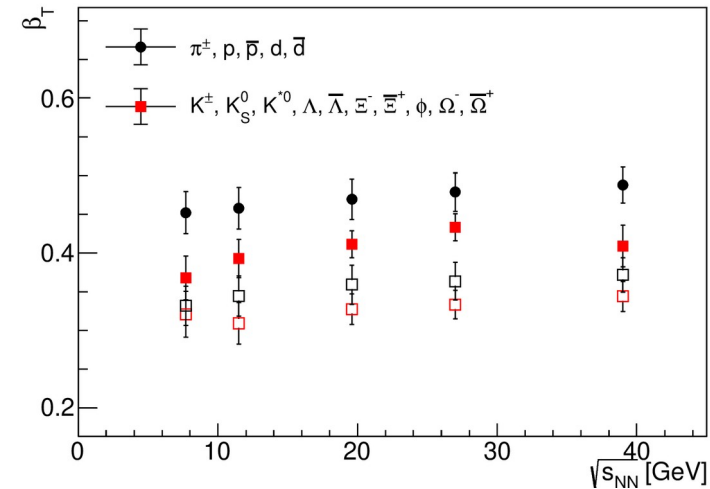
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# Mean transverse momentum

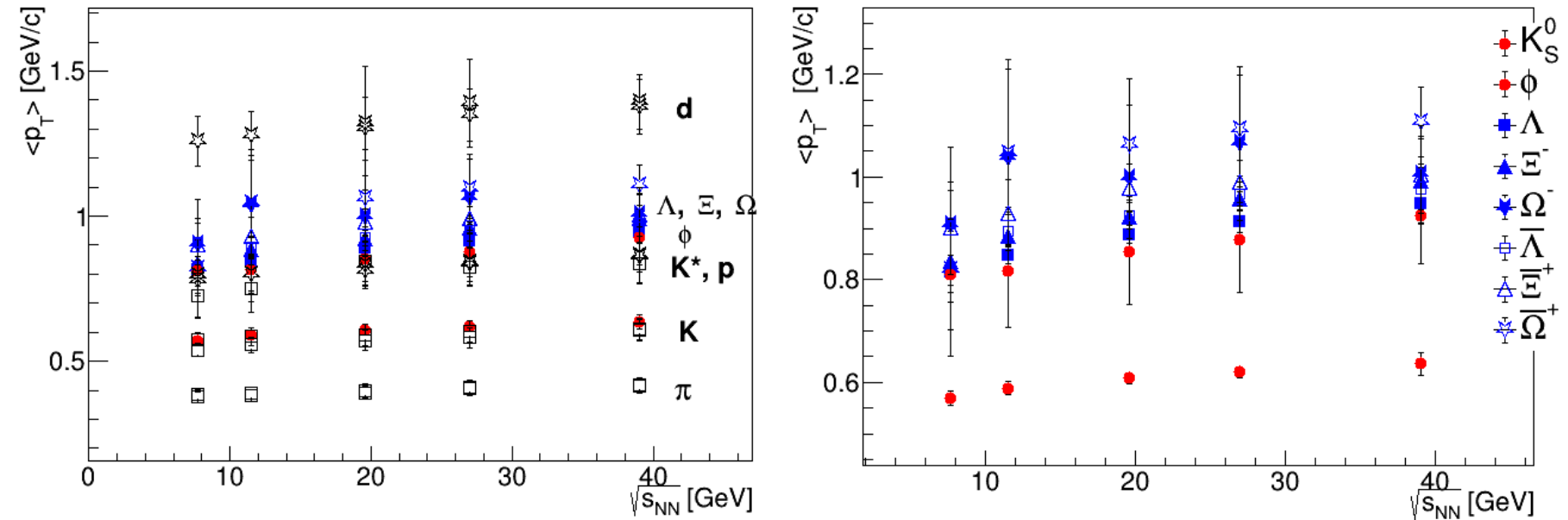


$\langle \beta_T \rangle \rightarrow$  smaller in peripheral collisions  $\rightarrow$  weaker transverse collective flow



peripheral Au-Au collisions  
(40-50%/40-60% centrality class)

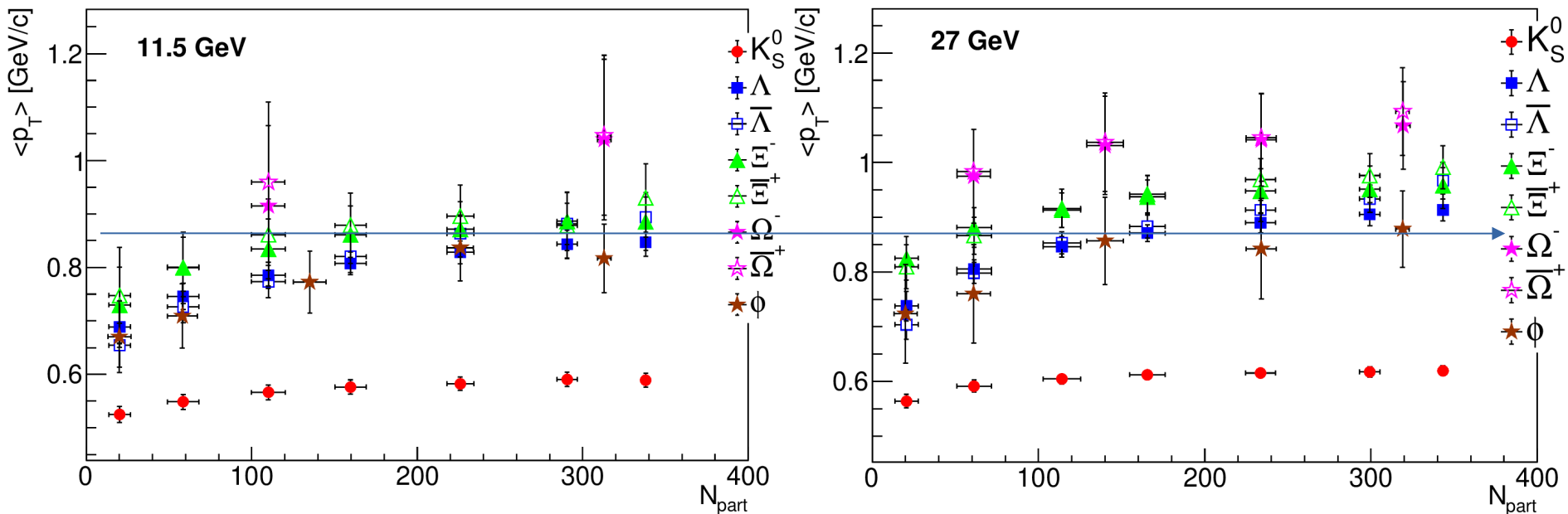
# $\langle p_T \rangle$ vs energy



- mass dependence of  $\langle p_T \rangle$  --> transverse collective flow --> larger  $p_T$  kick for particles with higher mass

- very weak energy dependence

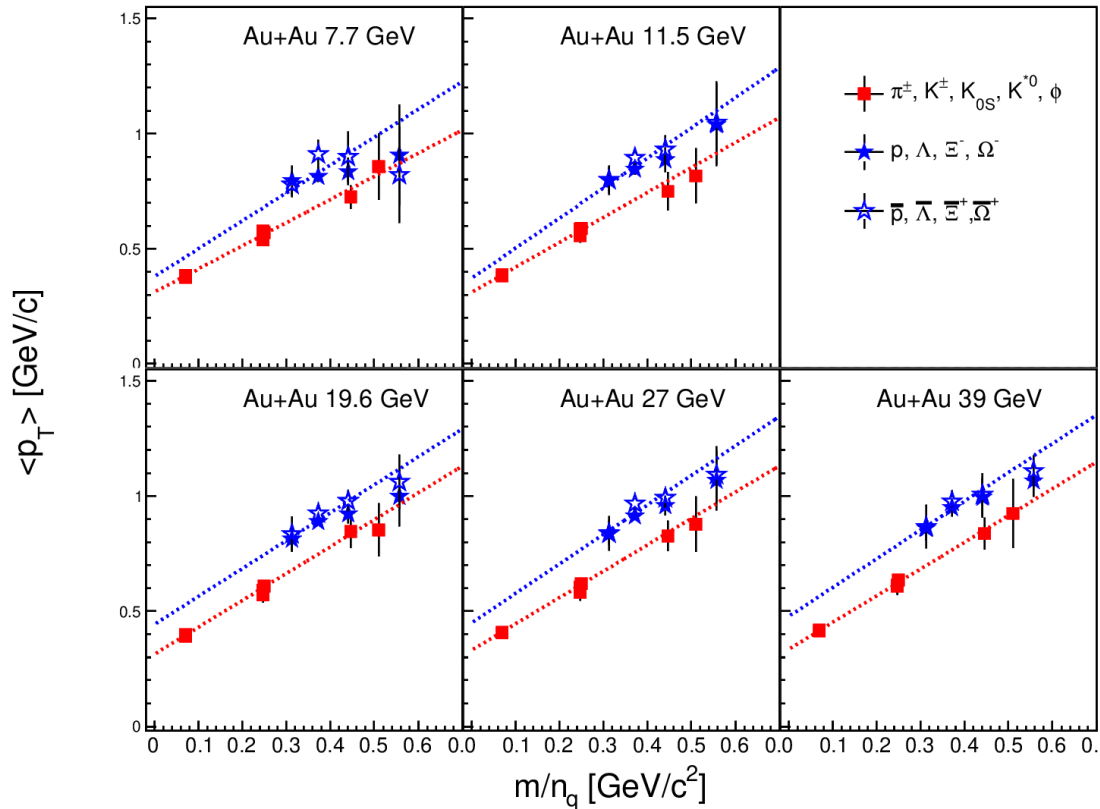
# $\langle p_T \rangle$ vs centrality



-  $\langle p_T \rangle$  increases with  $N_{part}$  for all energies --> gradual development of the transverse collective motion with increasing medium volume

-  $\langle p_T \rangle$  for  $K_S^0$  increases slightly with centrality, while for particles with higher mass, the increase is stronger --> contribution of the collective flow proportional with the particle mass

# $\langle p_T \rangle$ vs $m/n_q$



0-5%/0-10%

- 0-5%:  $\pi, K^{+/-}, K^0_S, p, \Lambda, \Xi$

- 0-10%:  $\phi$  (all energies),  $\Omega, K^{*0}$  (11.5-39 GeV)

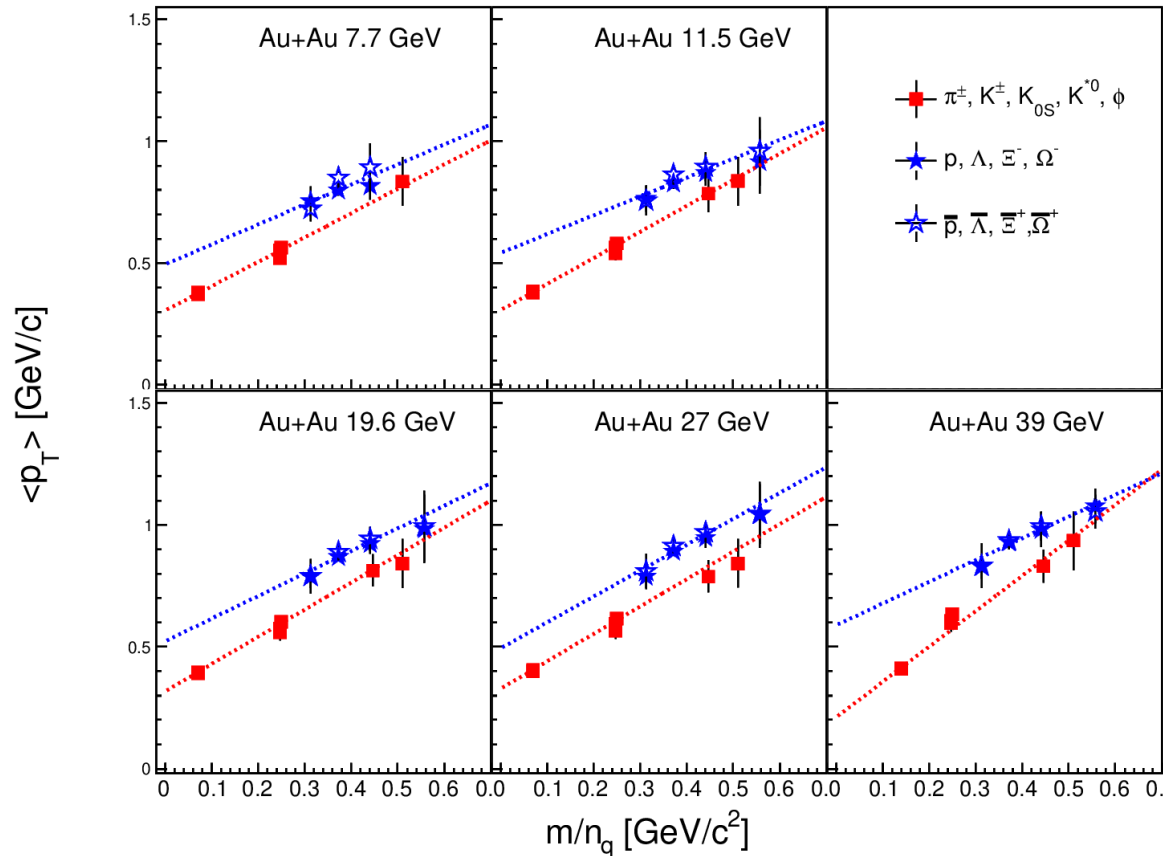
- 0-20%:  $K^{*0}$  (7.7 GeV)

- 0-60%:  $\Omega$  (7.7 GeV)

-  $\langle p_T \rangle$  of baryons as a function of reduced mass increases with a different slope compared with  $\langle p_T \rangle$  of mesons

- for all energies, the average  $p_T$  for baryons is larger than for mesons

# $\langle p_T \rangle$ vs $m/n_q$

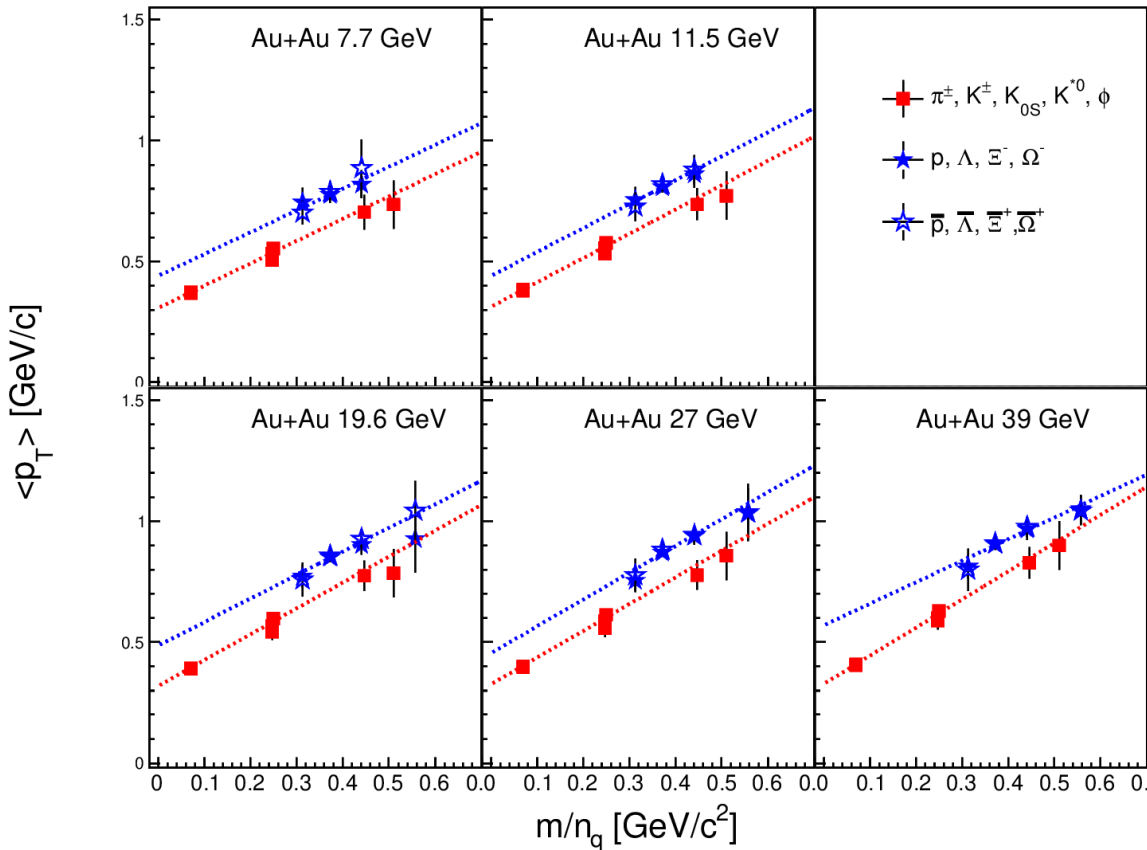


10-20%

- at 7.7 GeV, no  $\Omega$  data in this centrality class

- at 11.5 GeV, 10-60% centrality class for  $\Omega$

# $\langle p_T \rangle$ vs $m/n_q$

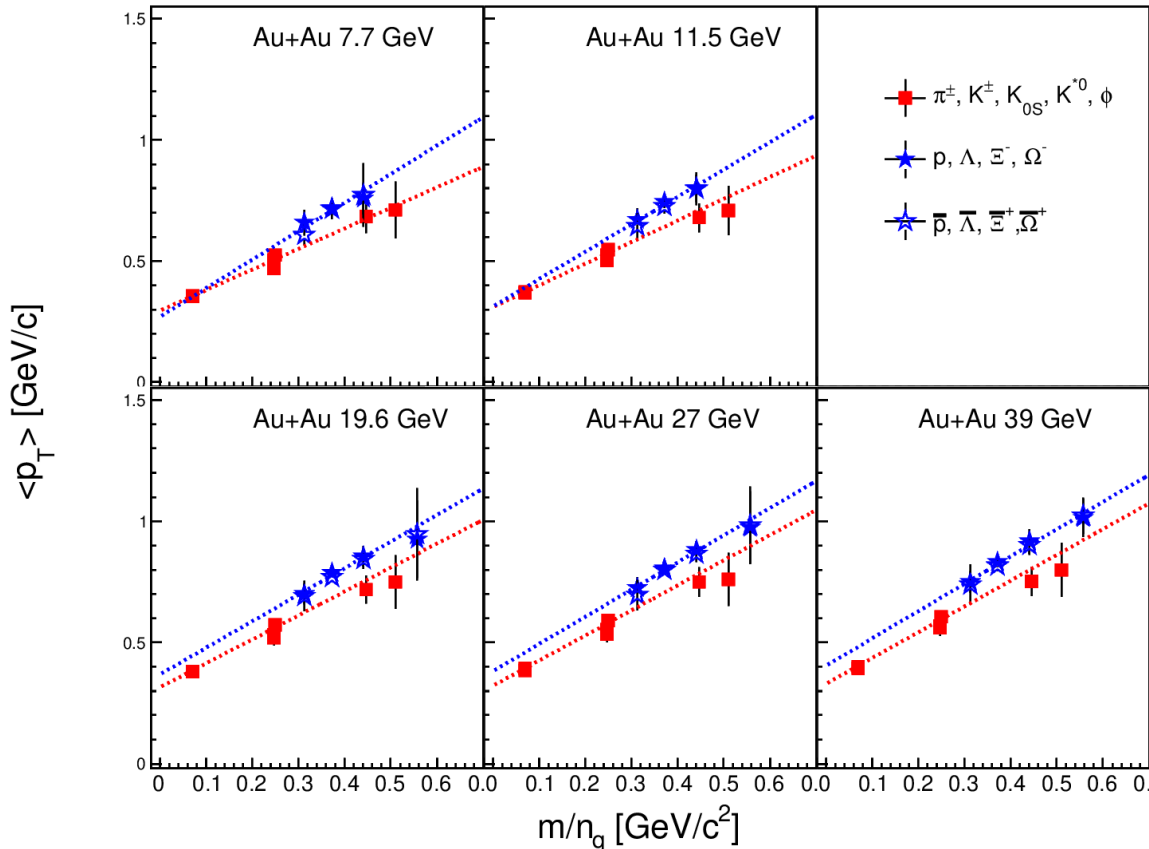


20-30%/20-40%

- 20-30%:  $\pi, K^{+/-}, K_S^0, K^{*0}, \phi, p, \Lambda, \Xi$

- 20-40%:  $K^{*0}$  (7.7 GeV),  $\Omega$  (19.6-39 GeV)

# $\langle p_T \rangle$ vs $m/n_q$



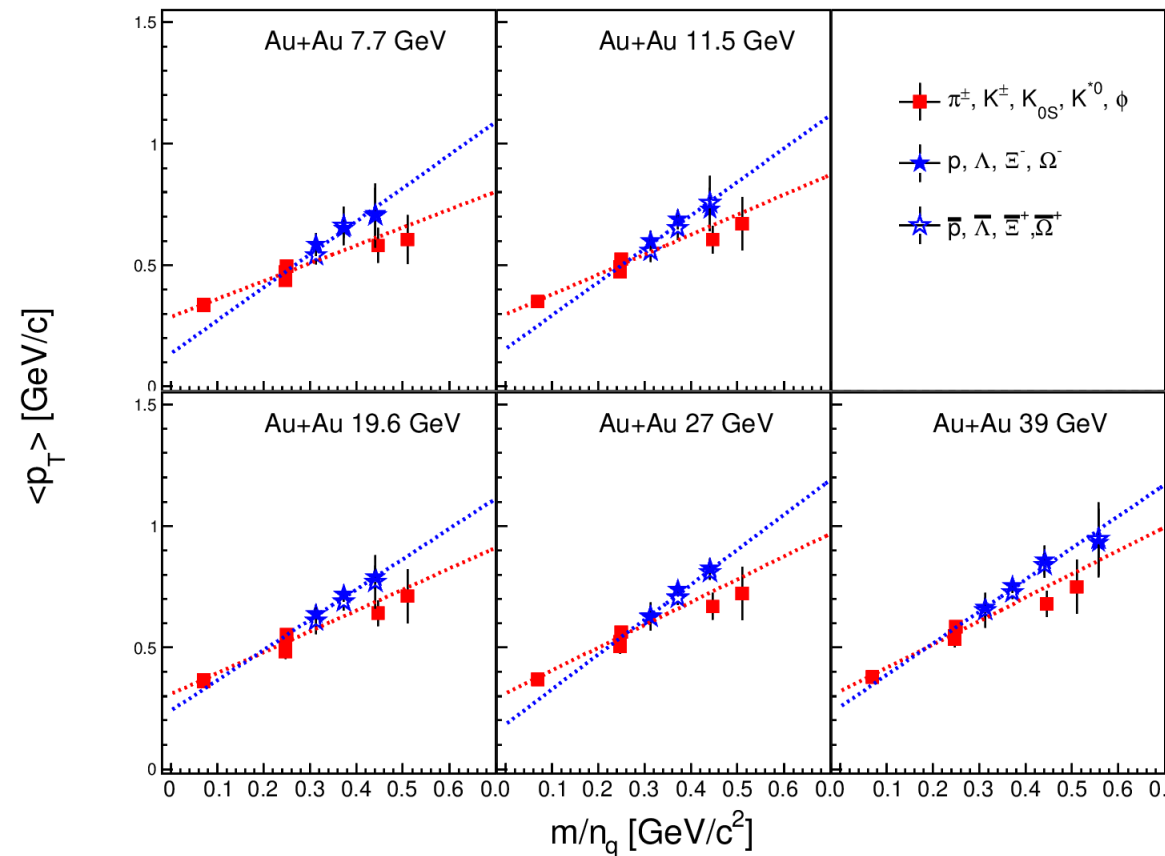
40-50%/40-60%

- 40-50%:  $\pi, K^{+/-}, p$

- 40-60%:  $K_{0S}^0, \phi, \Lambda, \Xi, K^{*0}$

- 40-60%:  $\Omega$  (19.6-39 GeV)

# $\langle p_T \rangle$ vs $m/n_q$



60-70%/60-80%

- 60-70%:  $\pi$ ,  $K^{+/-}$ ,  $p$

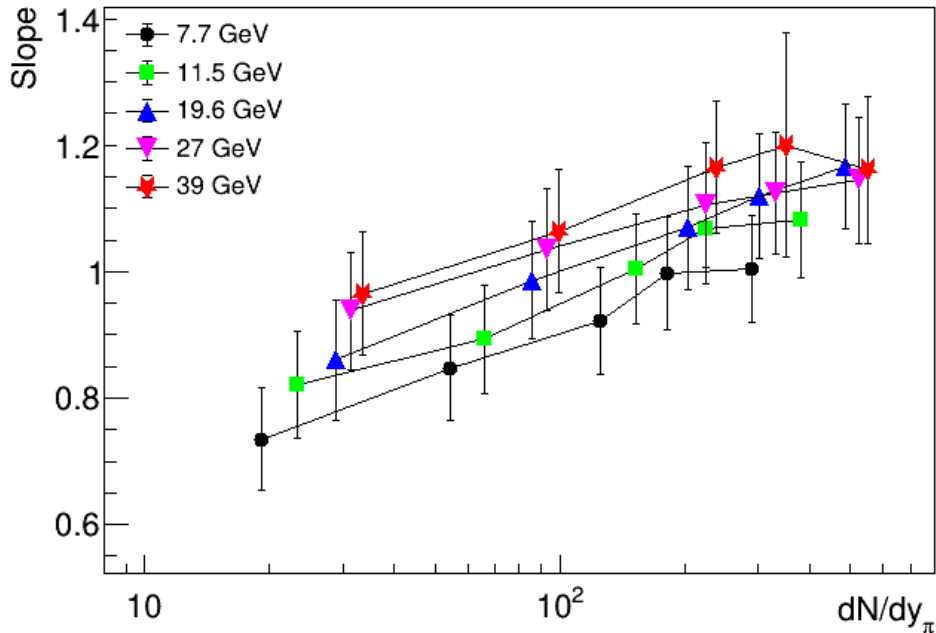
- 60-80%:  $K_S^0$ ,  $\phi$ ,  $\Lambda$ ,  $\Xi$ ,  $K^{*0}$

- 60-80%:  $\Omega$  (39 GeV)

- the differences in  $\langle p_T \rangle$  values between the baryons and mesons are smaller in peripheral collisions



# Meson slopes



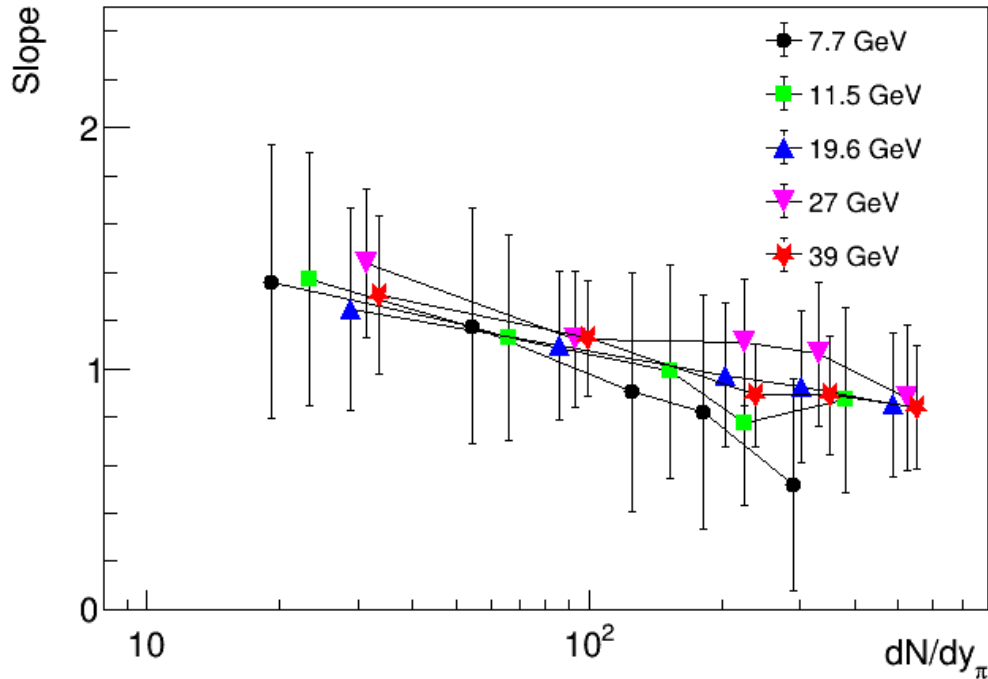
- the slopes of  $\langle p_T \rangle = f(m/n_q)$  for mesons extracted from linear fits

- pion multiplicities are taken from L. Adamczyk et al., STAR Collaboration, Phys. Rev. C 96, 044904 (2017)

$$dN/dy_\pi = 1.5 \cdot (dN/dy_{\pi^+} + dN/dy_{\pi^-})$$

- for all energies, the slope increases with pion multiplicity from peripheral to semi-central collisions, while for central collisions seems to saturate

# Baryon slopes

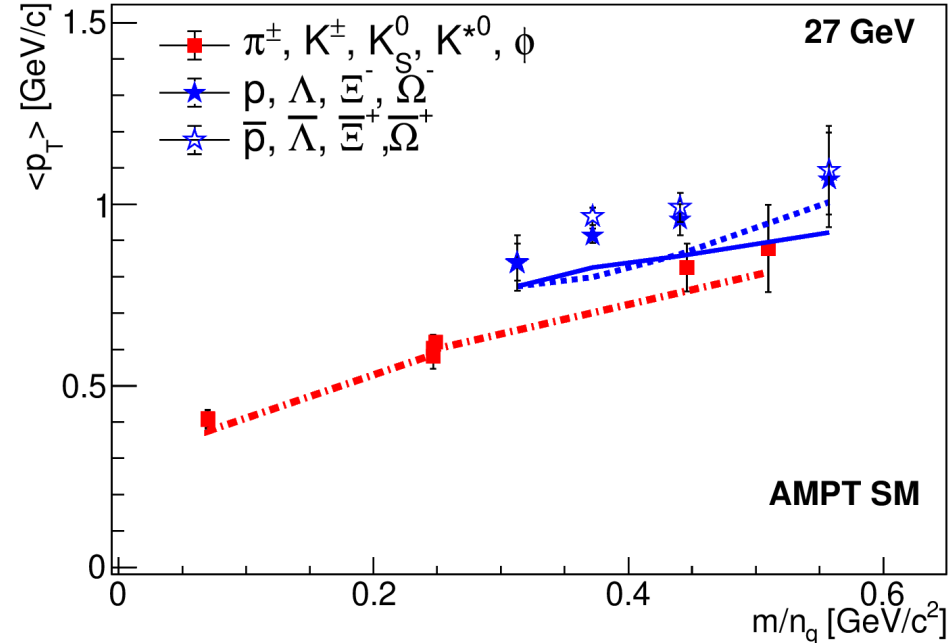
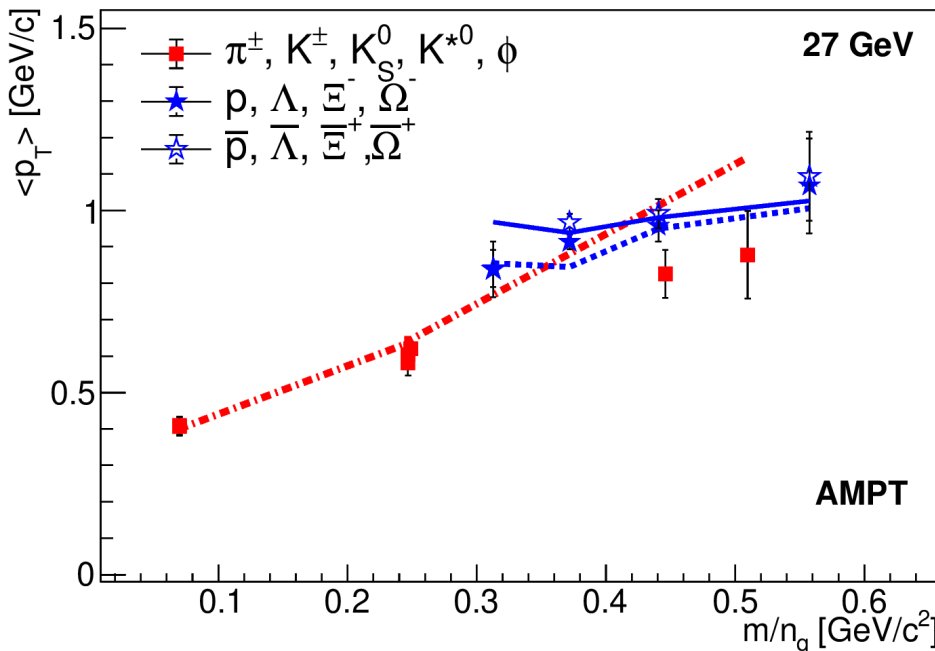


- for baryons, the  $\langle p_T \rangle = f(m/n_q)$  dependence is fitted with a 1st degree polynomial

- pion multiplicities are taken from L. Adamczyk et al., STAR Collaboration, Phys. Rev. C 96, 044904 (2017)

- different behaviour for baryons compared to meson slopes

# AMPT model comparison – central collisions



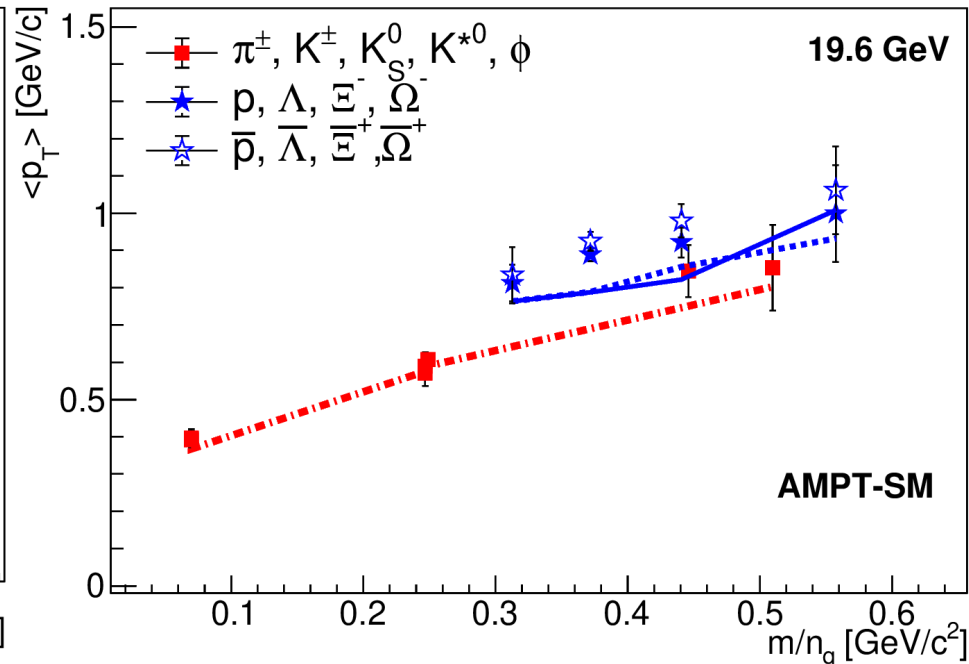
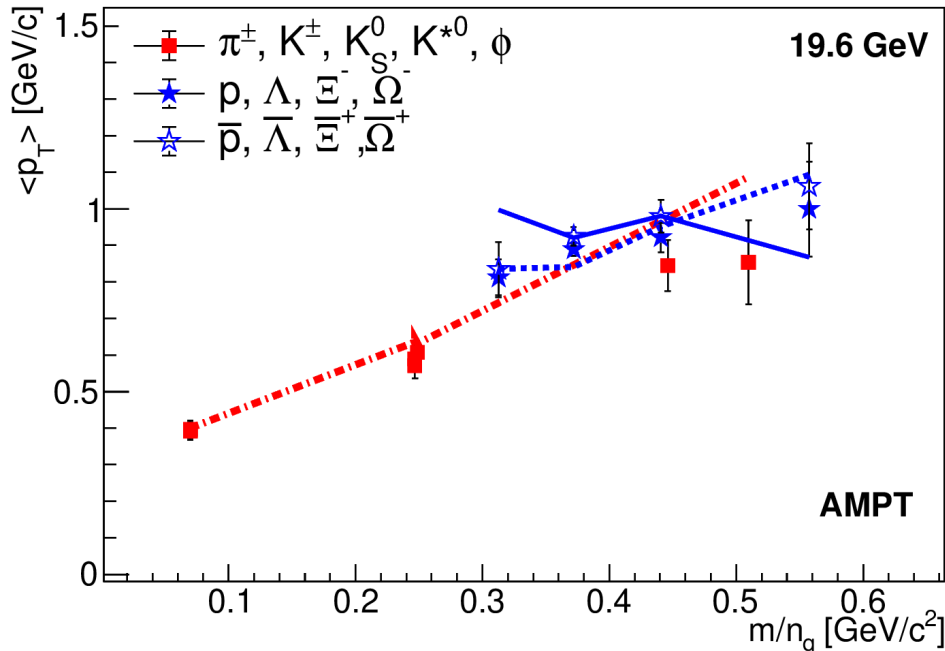
- 0-5%/0-10% centrality

- A Multi-Phase Transport (AMPT) --> a Monte Carlo transport model for heavy-ion collisions at relativistic energies --> includes both initial partonic and final hadronic interactions, and the transition between these two phases of matter.

- AMPT-default describes better the (anti)baryon  $\langle p_T \rangle$ ; but overestimates the  $\phi$   $\langle p_T \rangle$

- AMPT-SM underestimates the  $\langle p_T \rangle$  for baryons and for  $\phi$

# Model comparison – central collisions

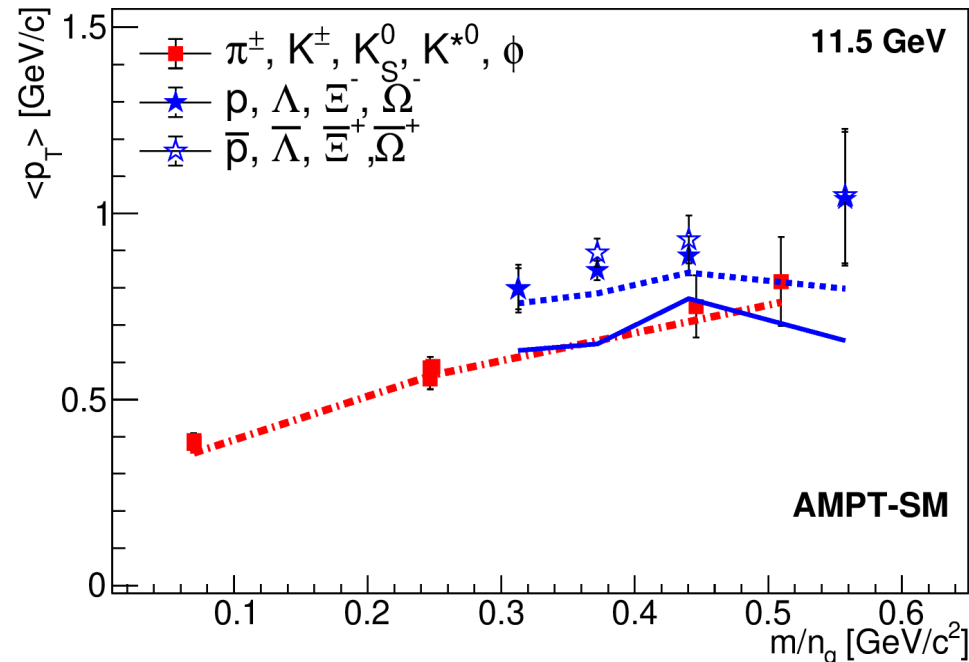
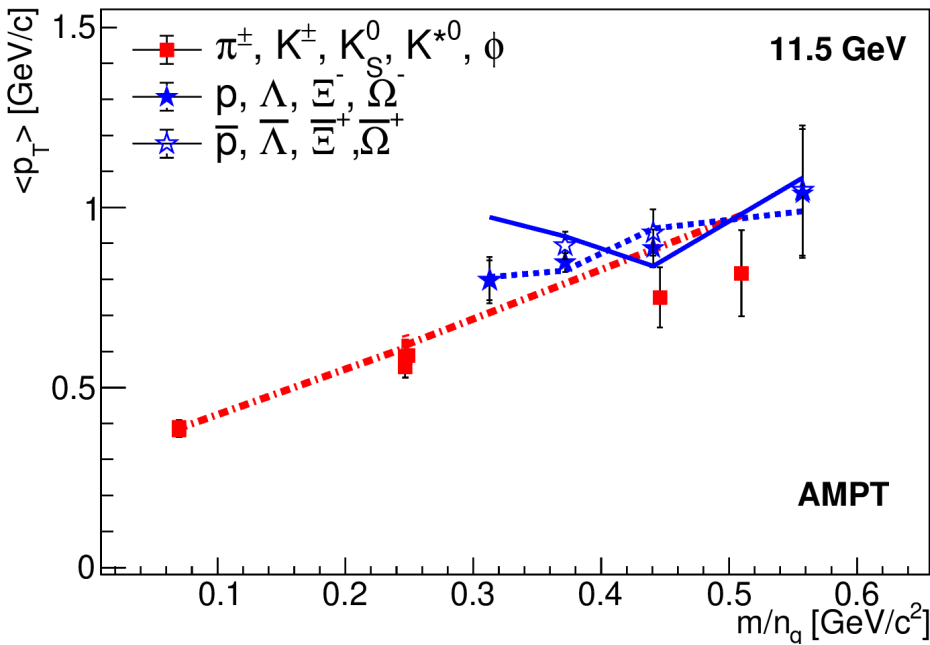


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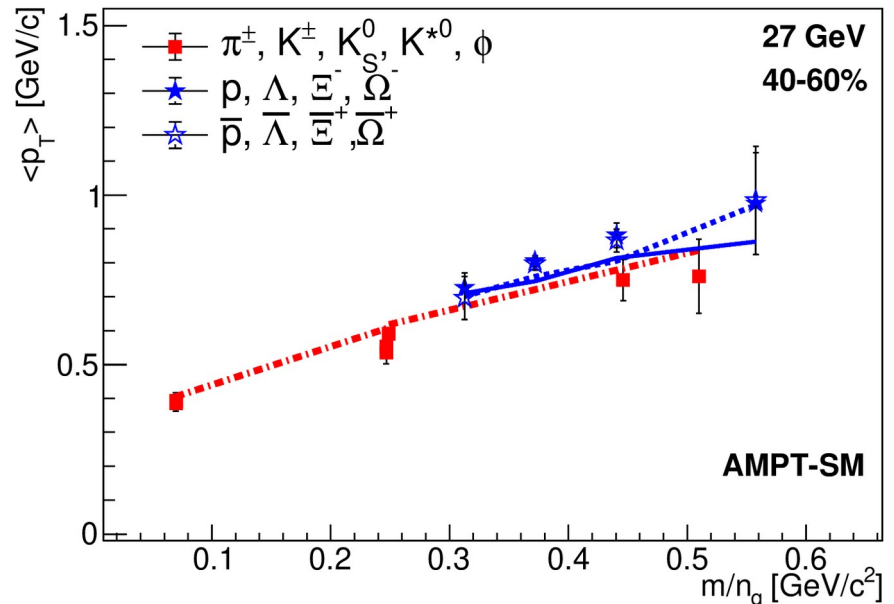
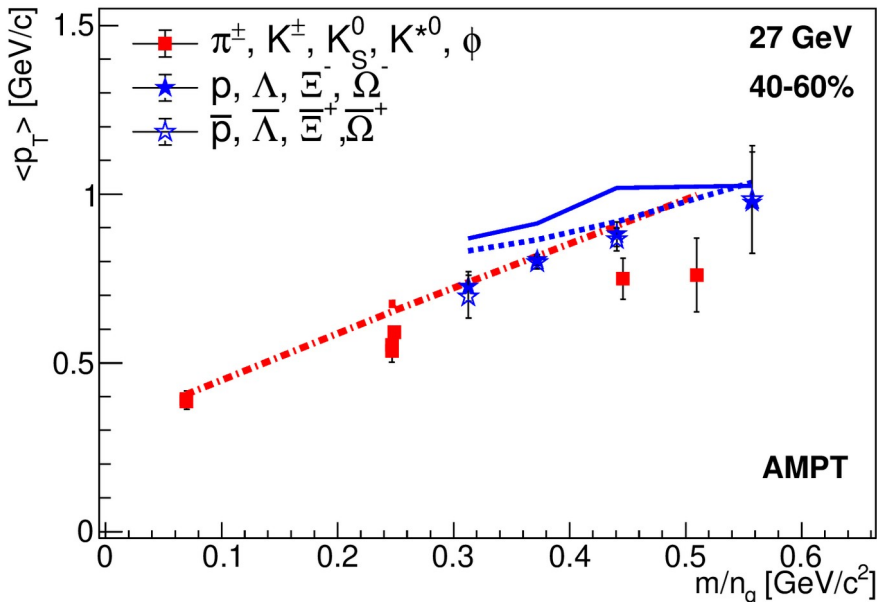


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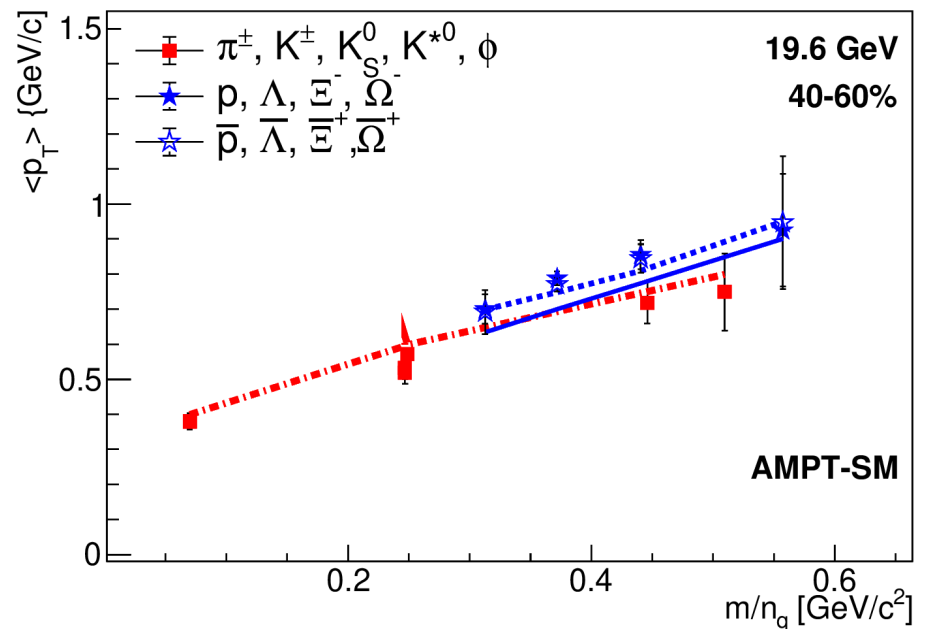
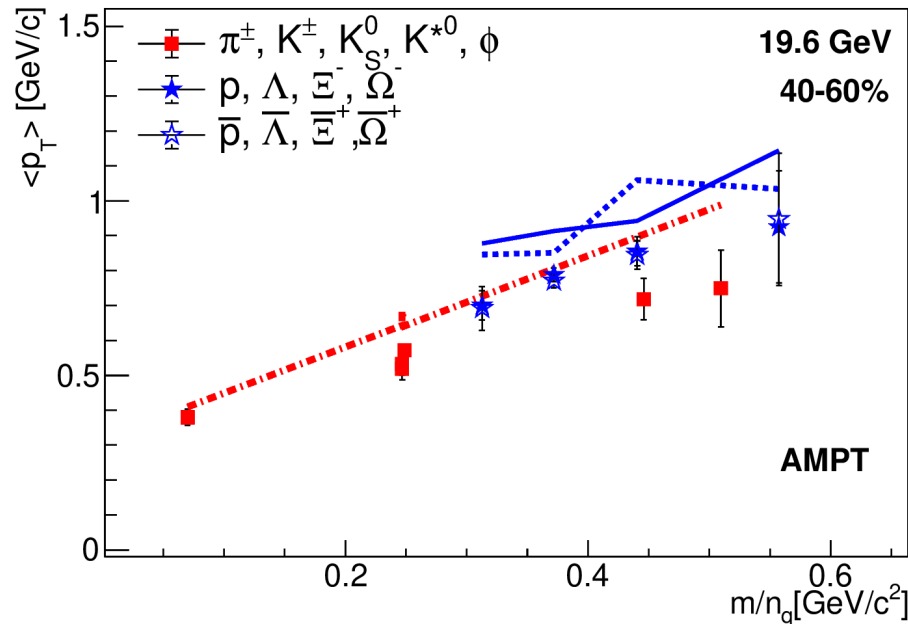
- AMPT-SM describes  $\langle p_T \rangle$  for mesons; smaller values of baryon  $\langle p_T \rangle$  compared to antibaryons

# Model comparison – peripheral collisions



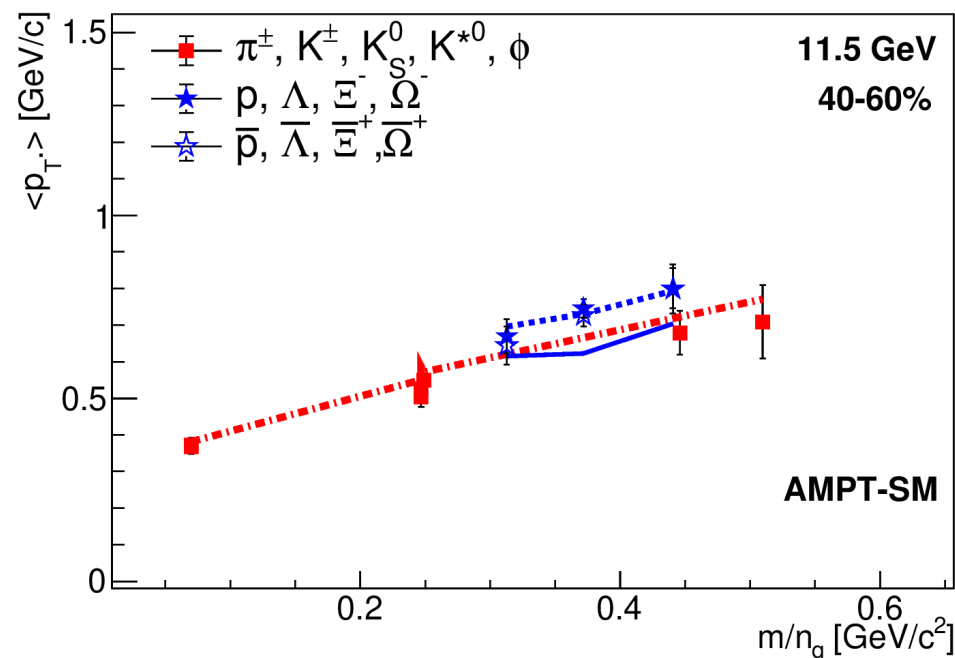
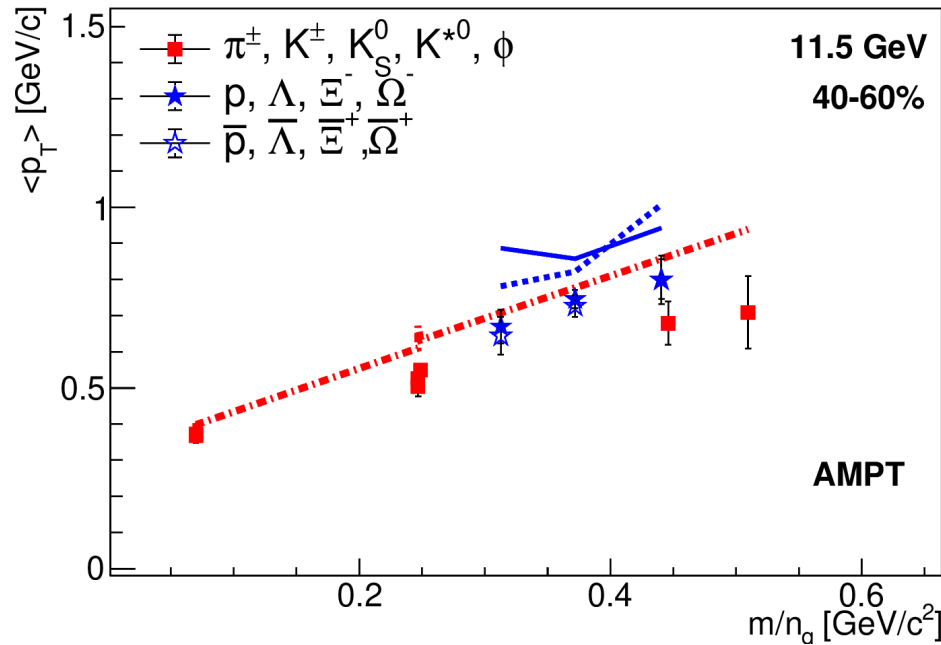
- 40-05%/40-60% centrality
- AMPT-default overestimates the (anti)baryon and meson  $\langle p_T \rangle$
- AMPT-SM describes better  $\langle p_T \rangle$  for mesons and baryons

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- AMPT-SM describes better  $\langle p_T \rangle$  for mesons and baryons; the  $\langle p_T \rangle$  for antibaryons is larger than that of baryons



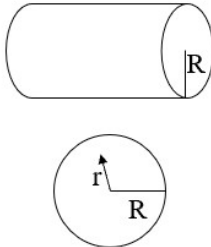
# Blast-wave model

→ A cylindrical expanding fireball in local thermal equilibrium, in which the particles are locally thermalized at a kinetic freeze-out temperature and are moving with a common transverse collective flow velocity (E. Schnedermann et. al, PRC48 (1993) 2462).

→ The  $p_T$  spectrum of produced particles described by:

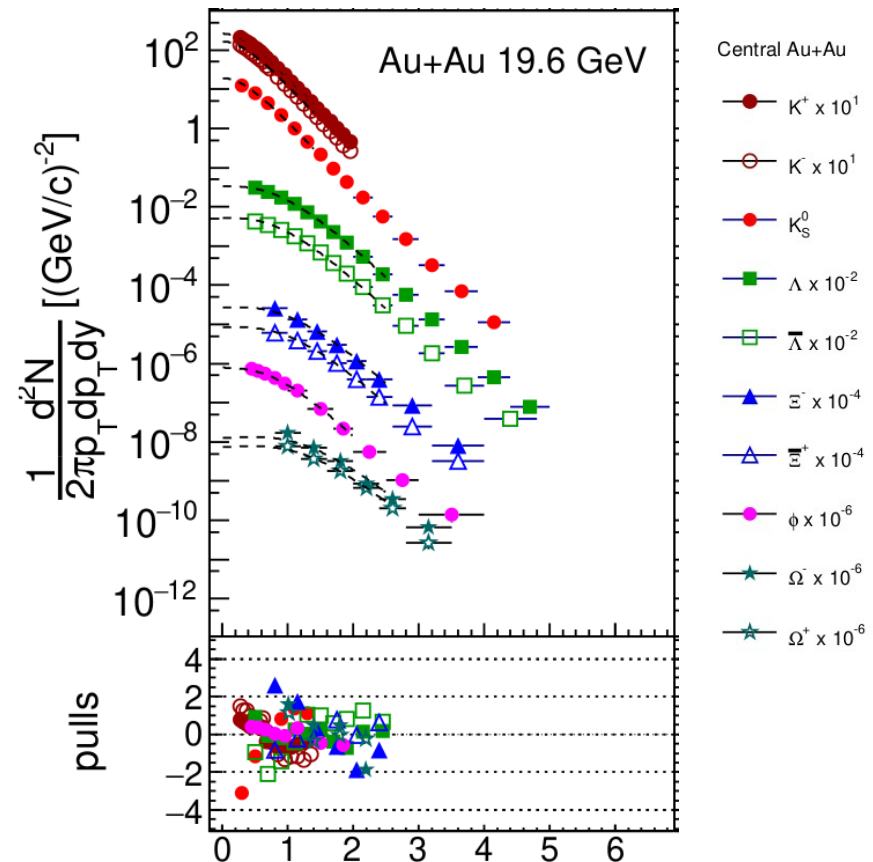
$$\frac{dN}{p_T dp_T} \propto \int_0^R r dr m_T I_0 \left( \frac{p_T \sinh \rho(r)}{T_{\text{kin}}} \right) \times K_1 \left( \frac{m_T \cosh \rho(r)}{T_{\text{kin}}} \right),$$

→ The radial flow velocity profile parametrized as:

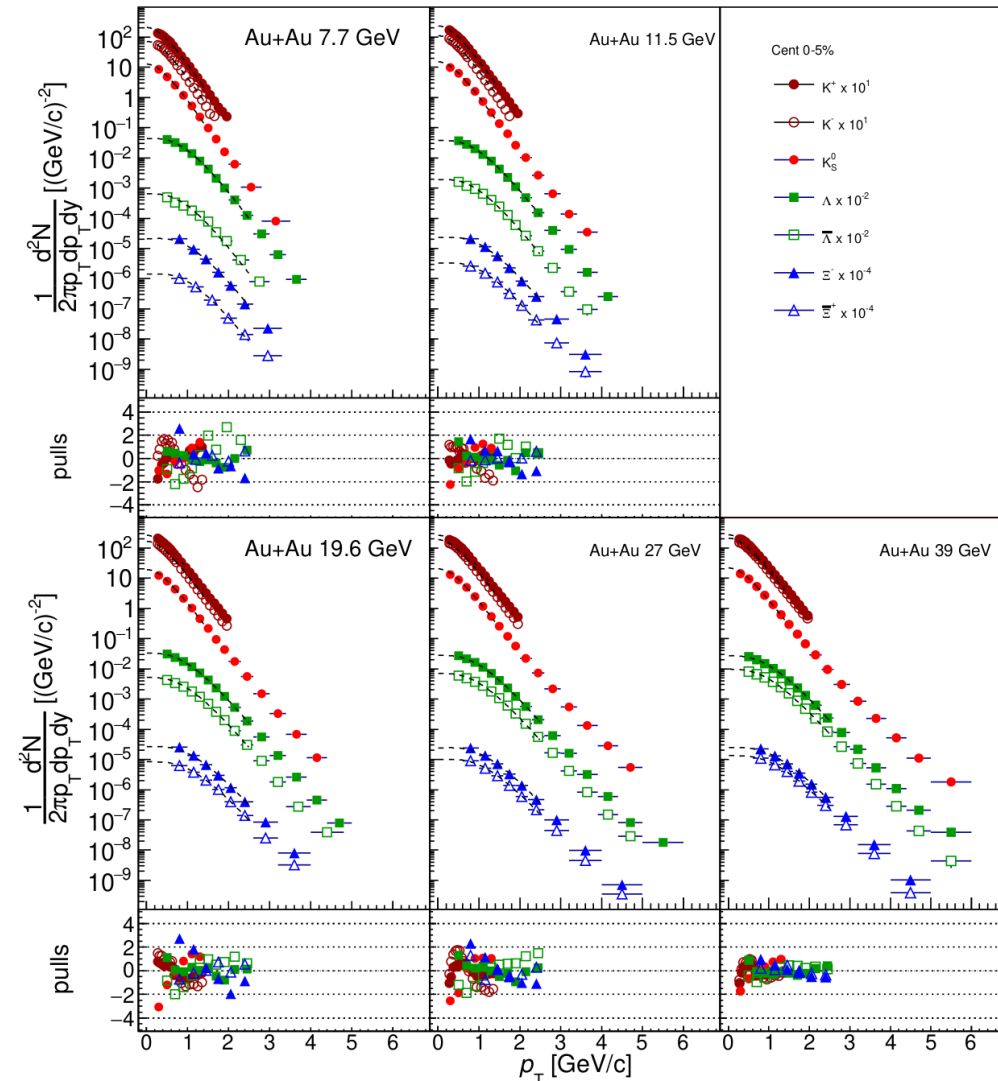
$$\beta_T(r) = \beta_s \left( \frac{r}{R} \right)^n$$


→ The average transverse radial flow velocity is:

$$\langle \beta_T \rangle = \frac{2}{n+2} \beta_s$$



# Blast-wave analysis of strange hadron spectra



--> BW fits on the  $p_T$  spectra of strange hadrons in Au+Au collisions at  $\sqrt{s_{NN}} = 7.7, 11.5, 19.6, 27, 39$  GeV from the STAR experiment

--> standard BW analyses -->  $\pi^+, \pi^-, K^+, K^-, p$  and anti- $p$  spectra fits

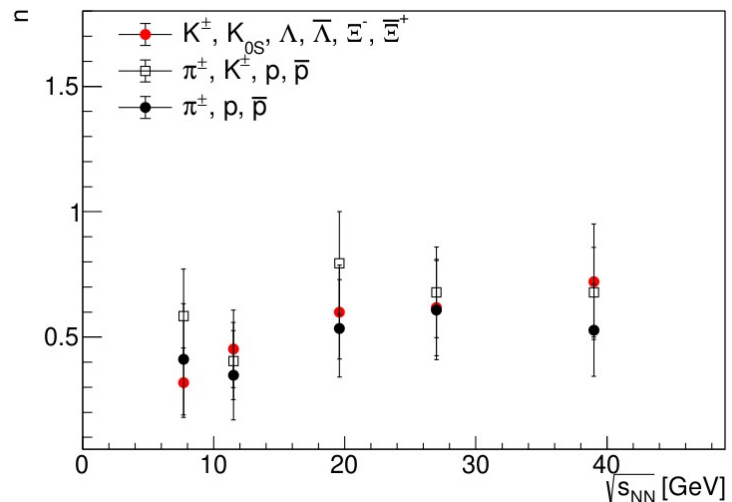
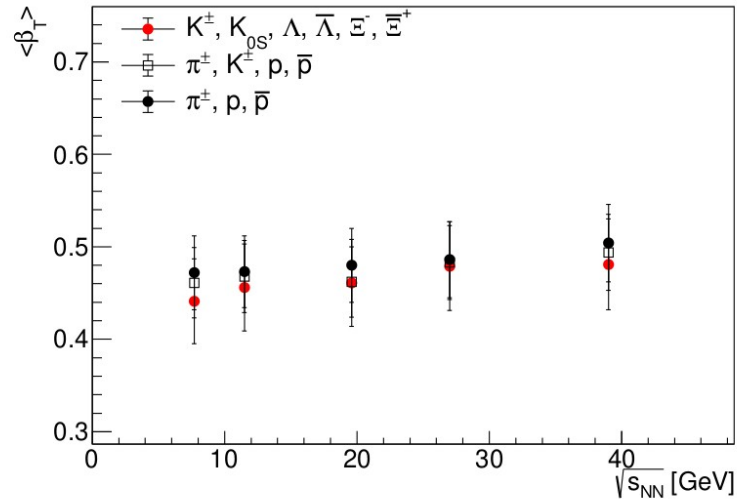
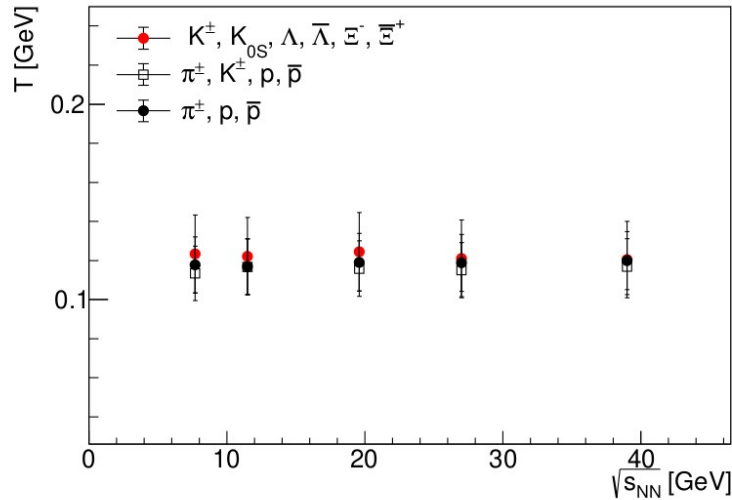
--> fit ranges: for kaons  $p_T < 1.4$  GeV/c; for strange baryons  $p_T < 2.5$  GeV/c

--> pulls distributions calculated as:

(data – fit value)/data error

J. Adam et al., STAR coll., Phys.Rev.C 102 (2020),034909;  
L. Adamczyk et al., Phys. Rev. C 96, 044904 (2017); L. Adamczyk et al., Phys. Rev. C 93, 21903 (2016)

# Energy dependence



--> in most central Au-Au collisions, similar  $T$  and  $\langle\beta_T\rangle$  within errors in this energy range for strange, non-strange and bulk hadrons --> indicate common freeze-out conditions for the strange/non-strange hadrons

--> very weak energy dependence of  $T$  and  $\langle\beta_T\rangle$  parameters

-->  $n$  increases slowly with energy

# Conclusions

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- in Au+Au collisions at RHIC-BES energies,  $\langle p_T \rangle$  increases with particle mass indicating the presence of the radial flow in the system  $\rightarrow$  two different linear trends can be observed for non-strange hadrons and strange hadrons separately.
- $\langle p_T \rangle$  of each studied particle species increases with  $N_{\text{part}}$  at all RHIC-BES energies, indicating a stronger collective motion in more central collisions
- weak energy dependence for the analyzed energy range
- different behaviour of  $\langle p_T \rangle$  as a function of  $m/n_q$  for mesons and baryons
- for central collisions, the default AMPT is generally better than the SM version to describe  $\langle p_T \rangle$  for mesons and baryons, while peripheral collisions  $\langle p_T \rangle$  are better described by AMPT-SM
- the values of strange particles freeze-out  $T$  and  $\langle \beta_T \rangle$  are similar within errors with the corresponding parameter values extracted from BGBW fits on bulk and non-strange particles  $p_T$  spectra.