



# Centrality dependent Levy HBT at PHENIX

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NEW TRENDS IN HIGH ENERGY AND  
LOW-X PHYSICS

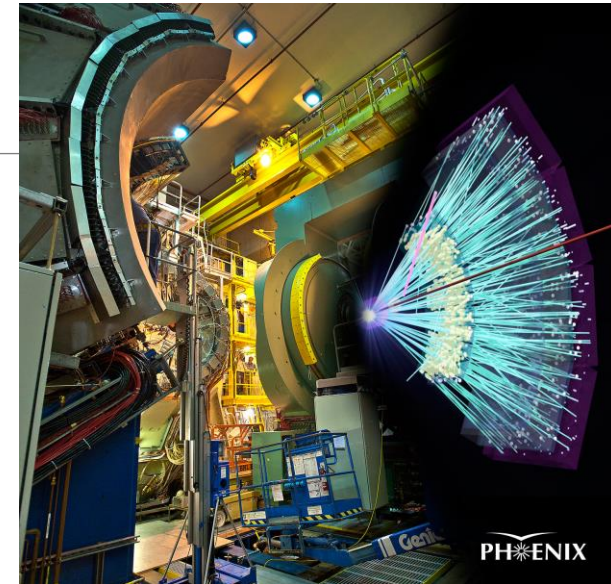


# The PHENIX and the BES

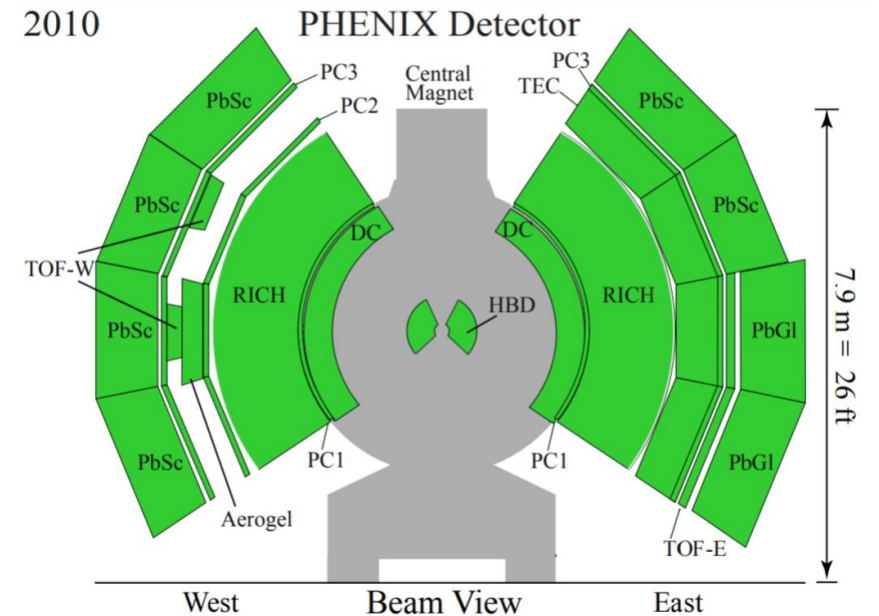
Collision energies: 7.7 to 200 GeV

20-400 MeV in  $\mu_B$ , 140-170 MeV in  $T$

This talk: 200 GeV Au+Au



$\sqrt{s_{NN}}$ [GeV]											
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200	<input checked="" type="checkbox"/>	<input checked="" type="checkbox"/>	<input checked="" type="checkbox"/>	<input checked="" type="checkbox"/>	<input checked="" type="checkbox"/>	<input checked="" type="checkbox"/>	<input checked="" type="checkbox"/>	<input checked="" type="checkbox"/>	<input checked="" type="checkbox"/>	<input checked="" type="checkbox"/>	<input checked="" type="checkbox"/>
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7.7									<input checked="" type="checkbox"/>		



# Femtoscscopy – general remarks

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Originates from radio astronomy

- Hanbury-Brown and Twiss observed intensity correlation
- In high energy physics, Goldhaber, Goldhaber, Lee and Pais

Technique to access the spatio-temporal structure of the particle emitting source

$$C_2(p_1, p_2) = \frac{N_2(p_1, p_2)}{N_1(p_1)N_1(p_2)}$$

where we can use the Yano-Koonin formula to relate the mom. dists. to the source:

$$N_2(p_1, p_2) = \int dx_1 dx_2 S(x_1, p_1) S(x_2, p_2) |\Psi_2(x_1, x_2)|^2$$

$S$ : source function,  $\Psi_2$  two-particle wavefunction

# Femtoscscopy – two approaches

Assume the source shape:  **$S \sim \text{Gaussian}$**

Measure in a clean environment, e. g. in  $pp$

Learn about the final state interactions hidden in the **wave function**

Program in ALICE:

$p - K, p - p, p - \Lambda, \Lambda - \Lambda, p - \Xi, p - \Omega,$

$p - \Sigma, p - \phi, N - \Sigma, N - \Lambda$

Assume the **wave function**: free planewave

$$|\Psi_2|^2 = 1 + \cos((p_1 - p_2)x)$$

Not too realistic: Coulomb (and strong) FSI

What is the interacting wave function?

$$\Psi_2 \sim \frac{\Gamma(1+i\eta)}{e^{\frac{\pi\eta}{2}}} [e^{ikr} F(-i\eta, 1, i(kr - \mathbf{k}\mathbf{r}))] \\ +\mathbf{r} \rightarrow -\mathbf{r}$$

(more complicated with strong interaction)

Learn about the **source size** and **shape**

# Final state interactions

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Like-charged pions → Coulomb correction

Strong final state interaction may play a role

Effect of the resonances: core-halo model

- Long-lived resonances contribute to the halo
- In-medium mass modifications could cause specific  $m_T$  dependence

Partially coherent particle production (core-halo model)

Aharonov-Bohm like effect: the hadron gas acts as a background field, the correlated bosons paths are the closed loop

# Levy parametrization of the $C_2$

Generalized Gaussian – Levy distribution

$$\mathcal{L}(\alpha, R, r) = \frac{1}{(2\pi)^3} \int d^3q e^{iqr} e^{-\frac{1}{2}|qR|^\alpha}$$

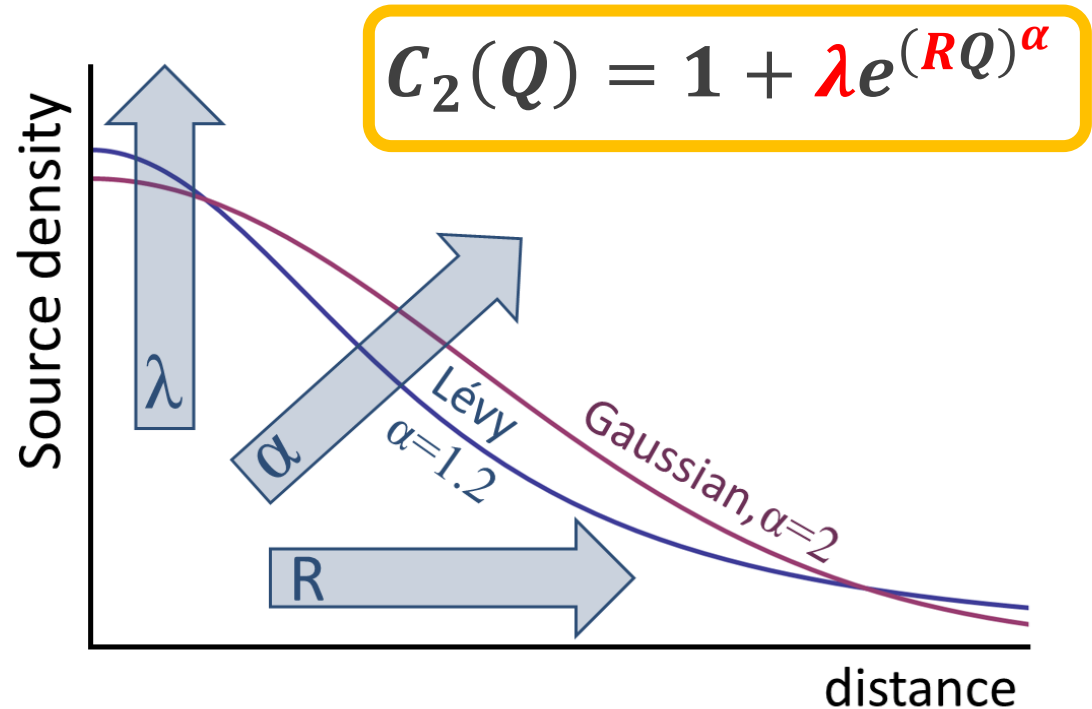
$\alpha = 2$ : Gaussian,  $\alpha = 1$ : Cauchy,  $0 < \alpha \leq 2$ : Levy

Assume the source to be Levy!

$\lambda(K)$ : core-halo parameter

$R(K)$ : Levy-scale parameter

$\alpha(K)$ : Levy index of stability



# Physics in the parameters

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Possible interpretations of the  $\lambda$ :

1. Specific  $m_T$  suppression linked to in-medium mass modification of  $\eta'$
2. Measuring two- and three particle correlations could shed light on partially coherent particle production (see core-halo model):

$$f_c(K) = \frac{N_c(K)}{N(K)} \quad \text{and} \quad p_c(K) = \frac{N_c^p(K)}{N_c(K)}$$

$$\lambda_2 = f_c^2 [(1 - p_c)^2 + 2p_c(1 - p_c)]$$
$$\lambda_3 = 2f_c^3 [(1 - p_c)^3 + 3p_c(1 - p_c)^2] + 3f_c^2 [(1 - p_c)^2 + 2p_c(1 - p_c)]$$

$$\kappa_3 = \frac{(\lambda_3 - 3\lambda_2)}{2\sqrt{\lambda_2^3}}$$

Independent from  $f_c$



# Physics in the parameters

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## Possible interpretation of the $R$ :

- Important:  $R_{Levy} \neq R_{Gauss}$
- Is it related to the size? Check hydro-like scaling:  $\frac{1}{R^2} = A m_T + B$
- Seen in Gaussian parametrizations

## Possible interpretation of the $\alpha$ :

- Surprising similarity with the critical exponent of the spatial correlation in 3D

$$\text{spatial corr.} \sim r^{-1-\eta} \quad \text{symm. Levy dist.} \sim r^{-1-\alpha}$$

- Sudden change in  $\alpha$  could be a sign for critical behavior
- Could be the sign of anomalous diffusion or jets

MOTIVATION TO MEASURE HBT VERY PRECISELY



# Centrality dependent HBT analysis from PHENIX

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Au+Au @ 200 GeV from Run 10,  $\pi^+\pi^+ + \pi^-\pi^-$

$\alpha, R, \lambda, \frac{1}{R^2}, \frac{1}{\hat{R}}, \frac{\lambda}{\lambda_{max}}$  in 6 cent bin (0-10% ... 50-60%) and 24  $m_T$  bins

1D variable  $Q = |q_{LCMS}|$  (instead of  $q_{inv} = |q_{PCMS}|$ )

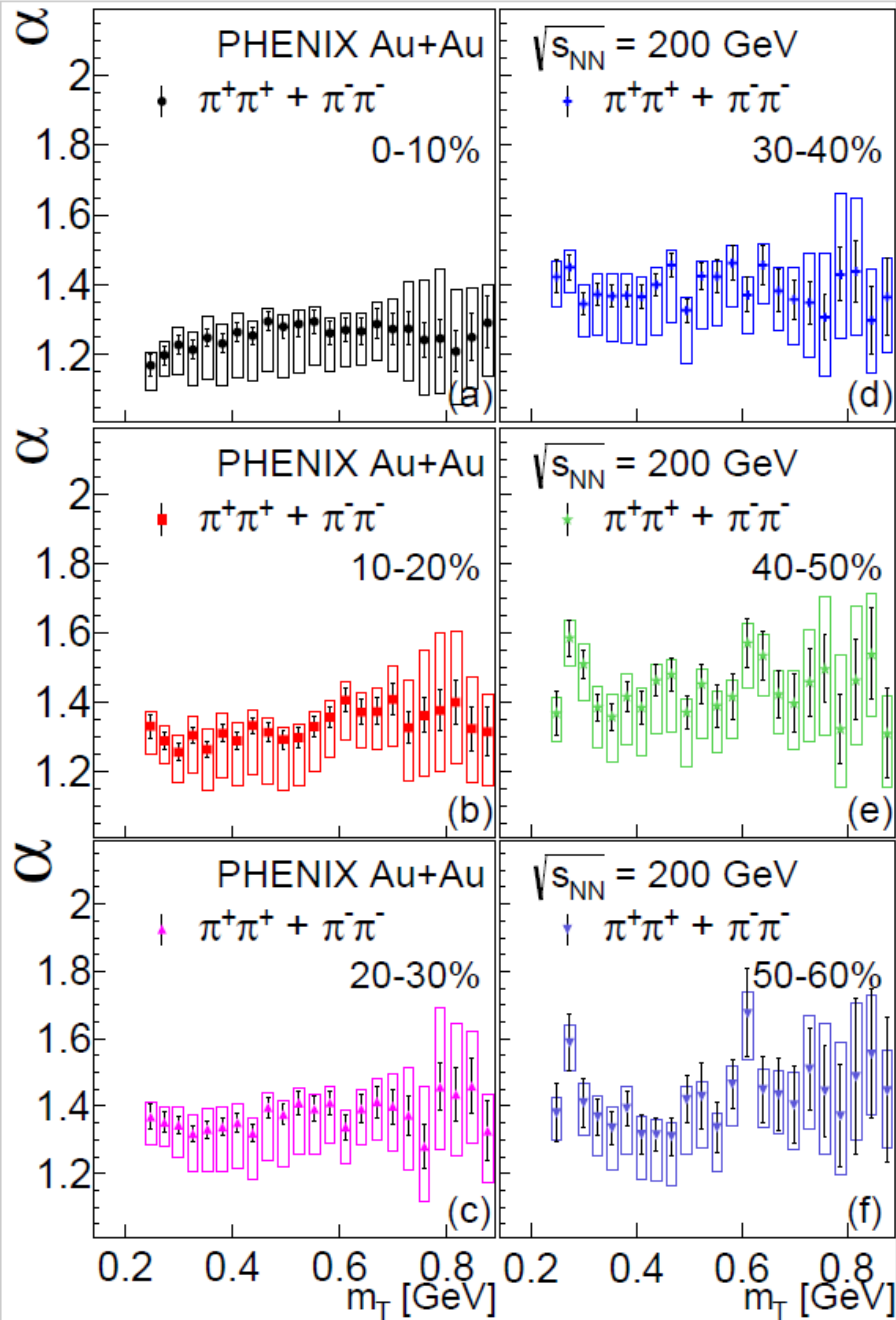
Fit function incorporates CC FSI (weighting for var. change)

Costume track and pair cuts to obtain clean sample

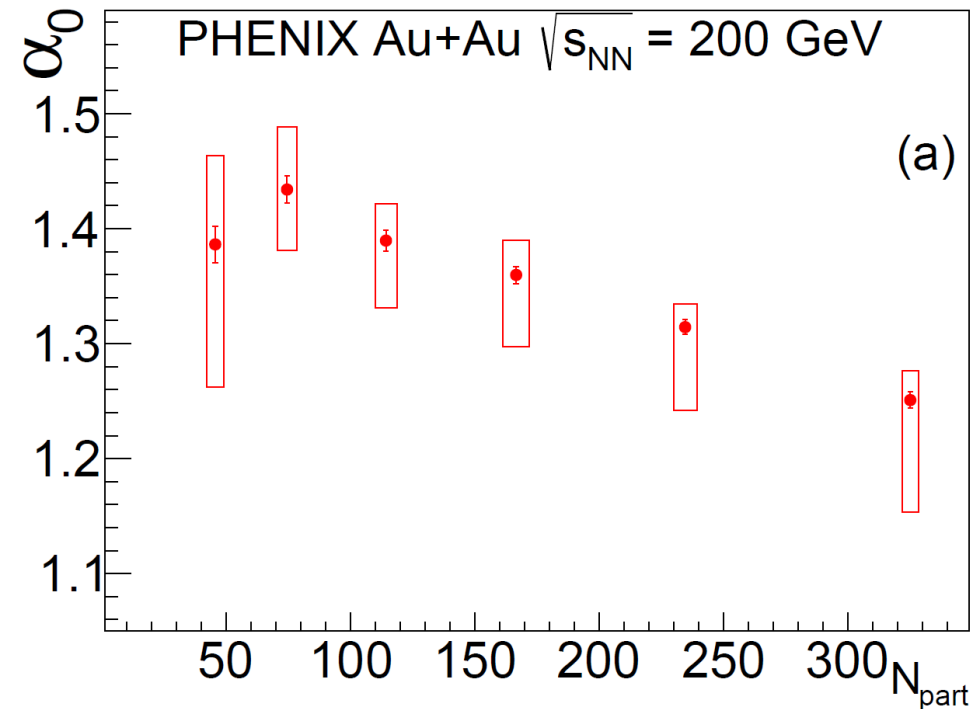
Submitted to PRC, available at [arXiv:2407.08586](https://arxiv.org/abs/2407.08586)

Let's see the results!

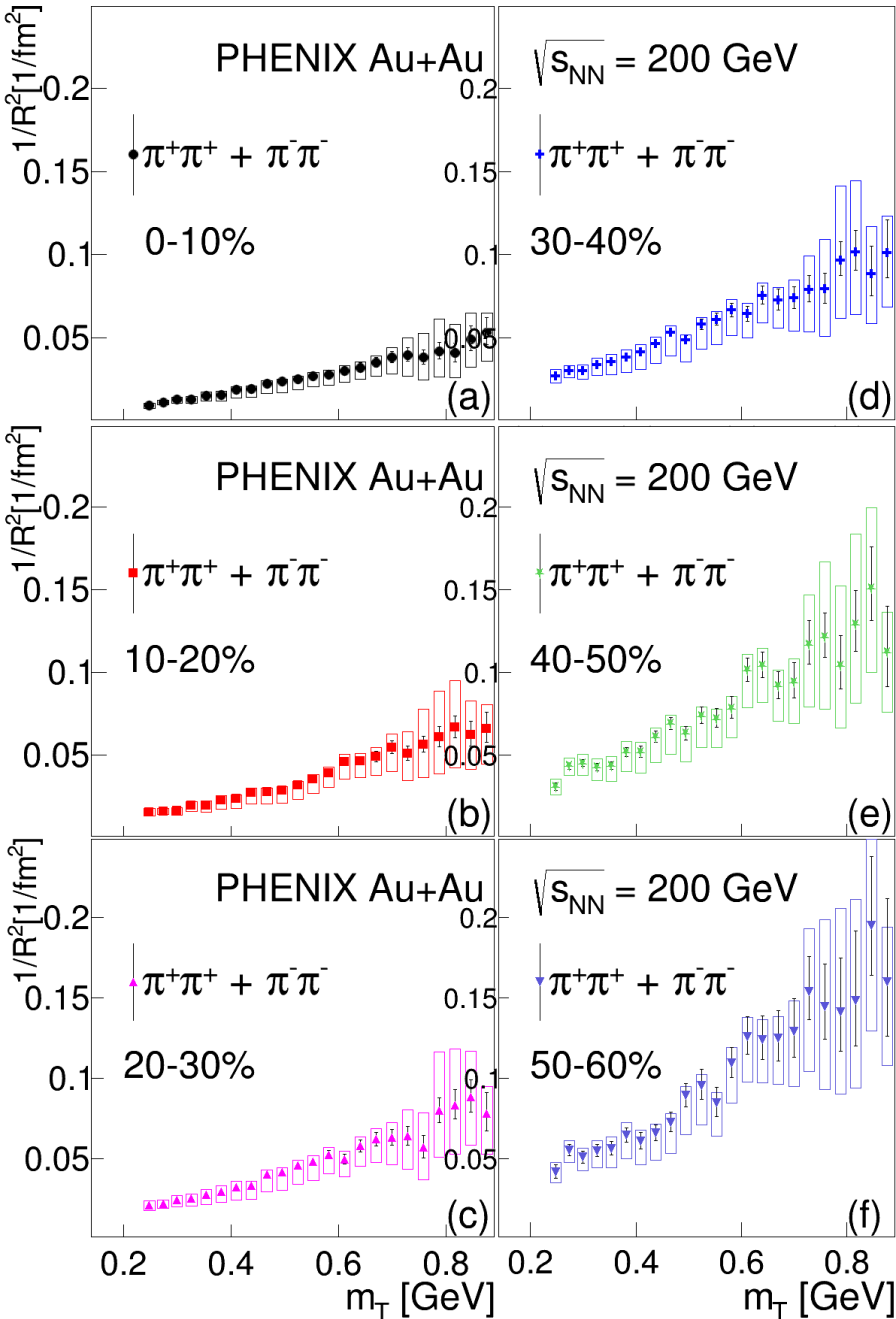
$$\alpha(m_T, N_{part})$$



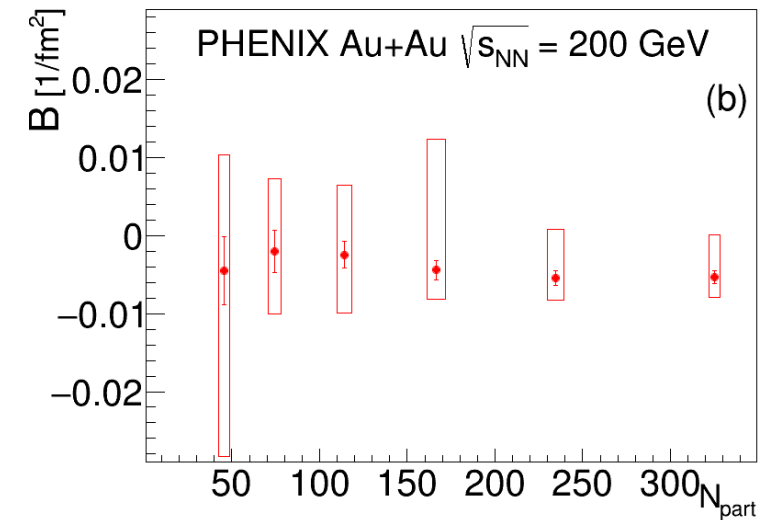
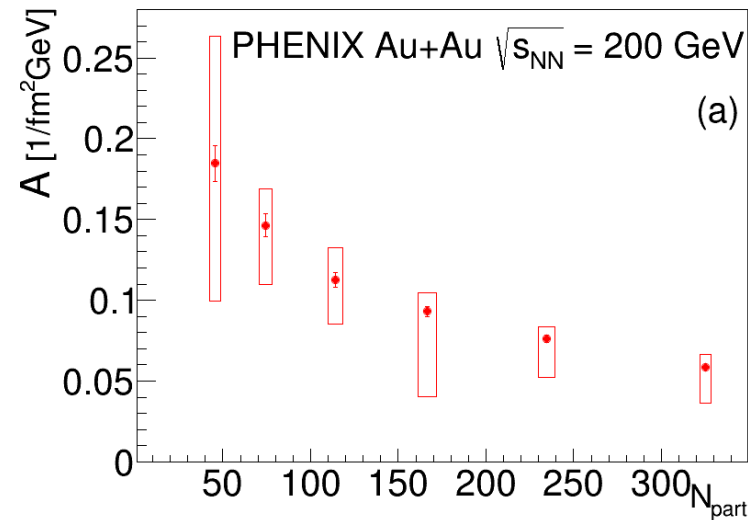
- Does not depend on  $m_T$ , does depend on  $N_{part}$
- $N_{part}$  dep. has model selection power
- Anomalous diffusion, QCD jets, resonances???



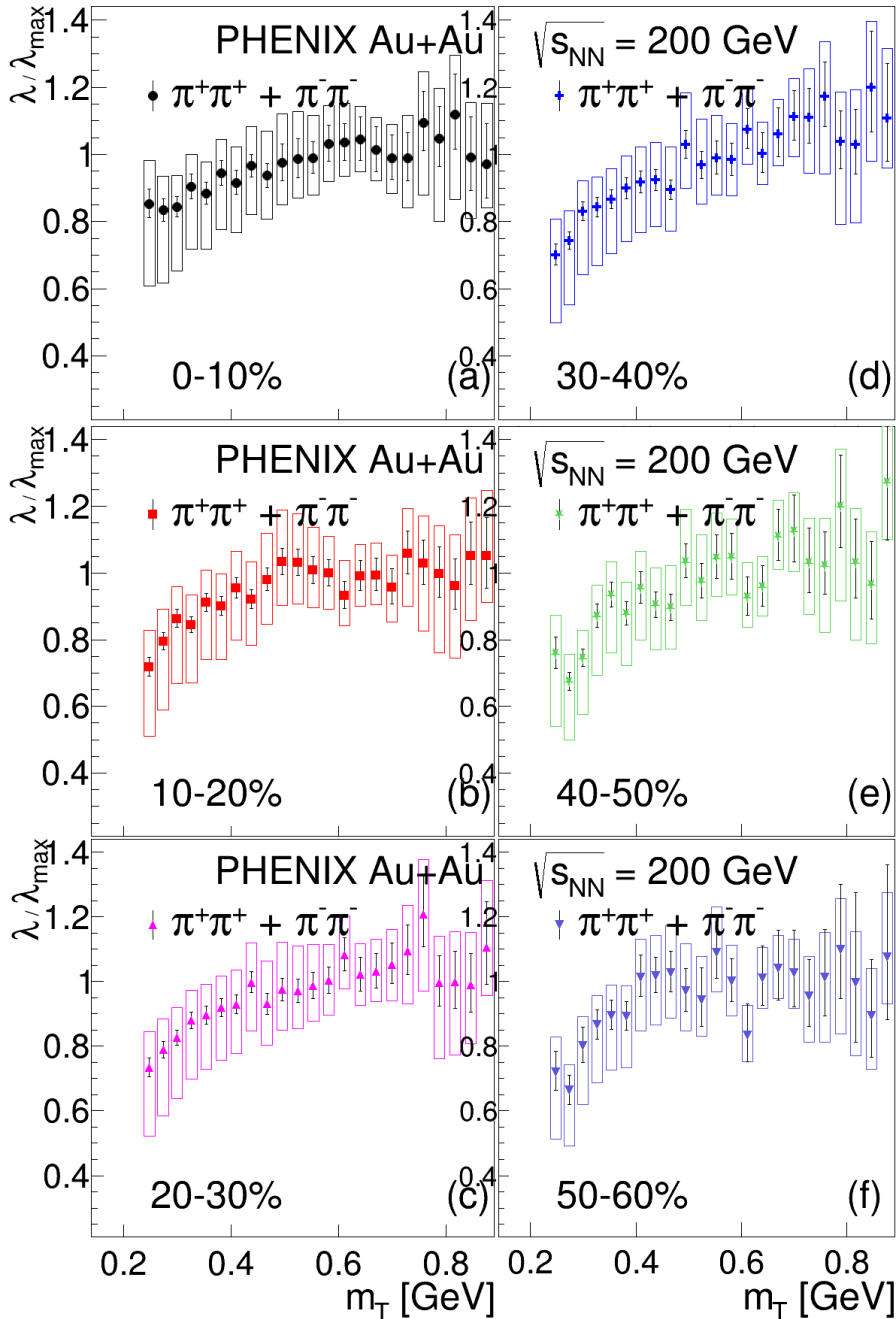
# $R(m_T, N_{part})$



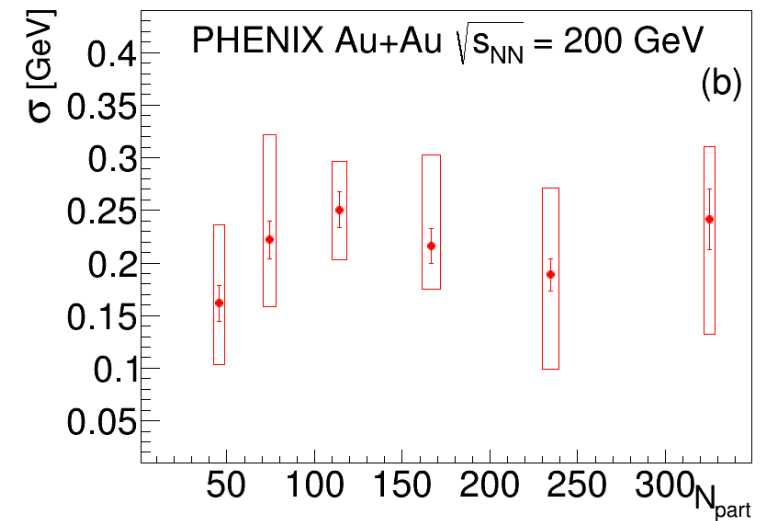
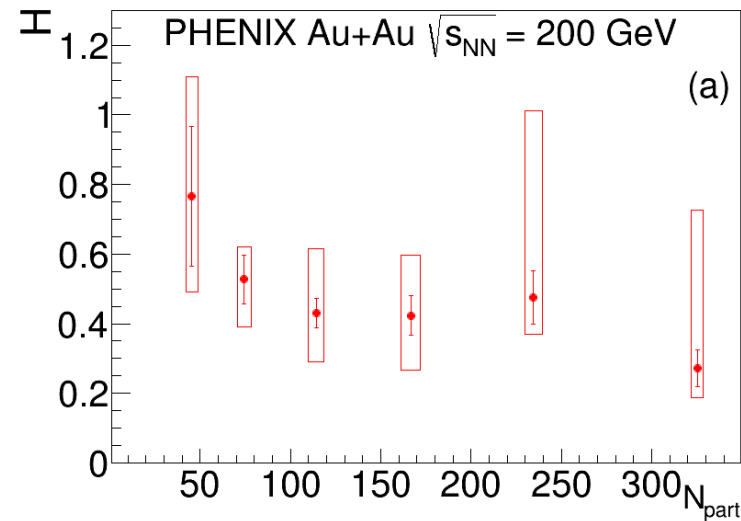
- Hydro scaling but not RMS, i.e.,  $R_{Gauss} \neq R_{Levy}$
- Centrality ordering, monotonic behavior
- Related to the size?



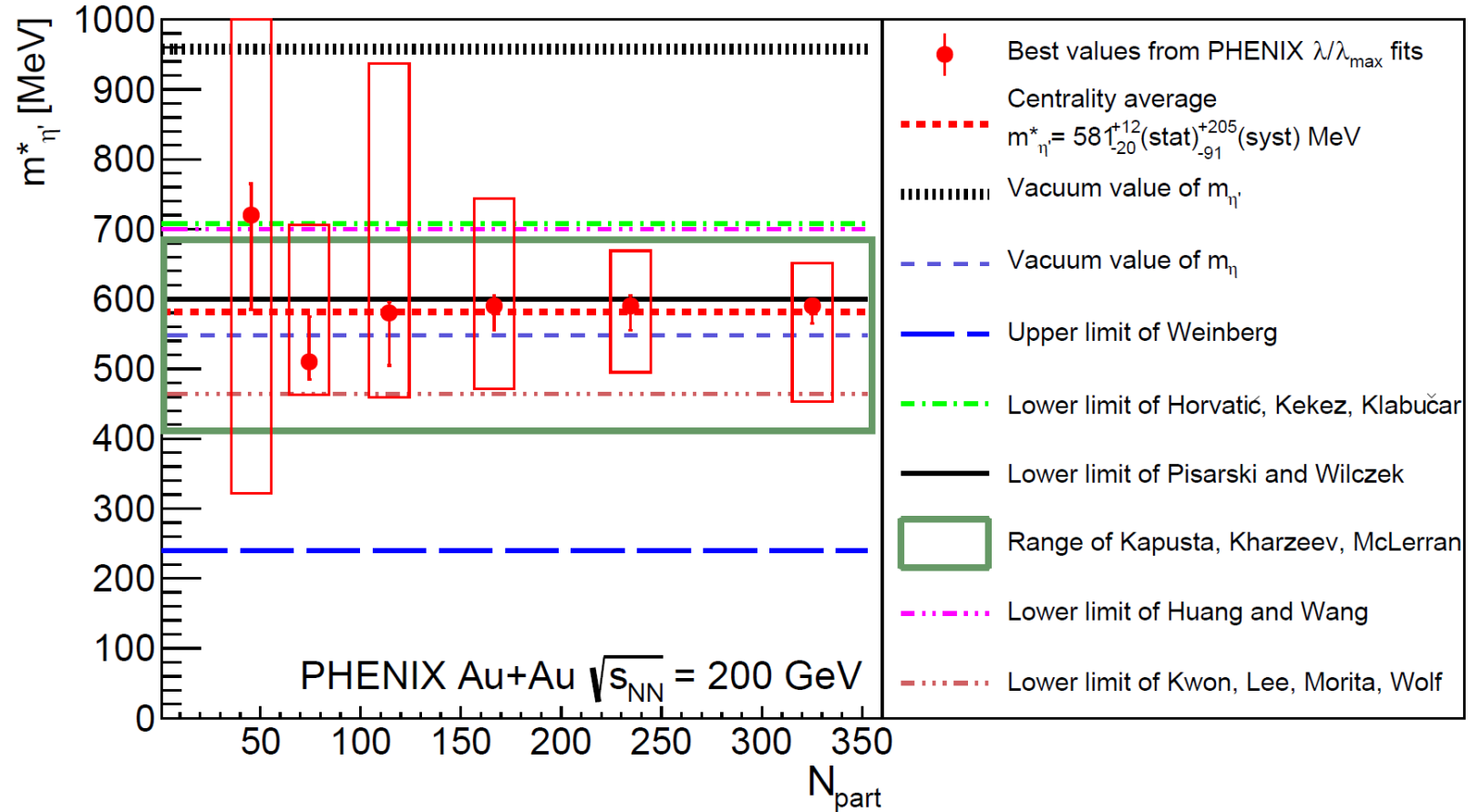
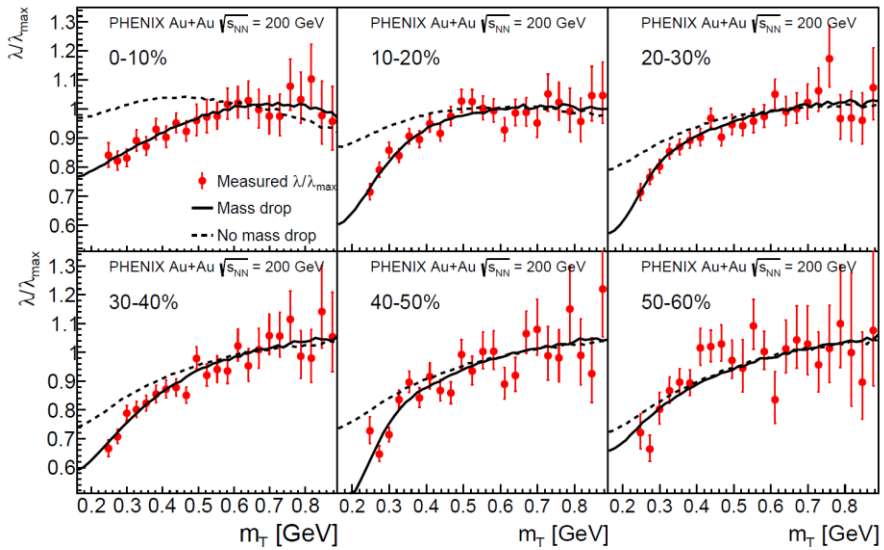
$$\lambda(m_T, N_{part})$$



- Suppression on every centrality
- Centrality independent!
- Sign of the  $\eta'$  in medium mass modification?
- Let's compare the results to Monte Carlo simulations

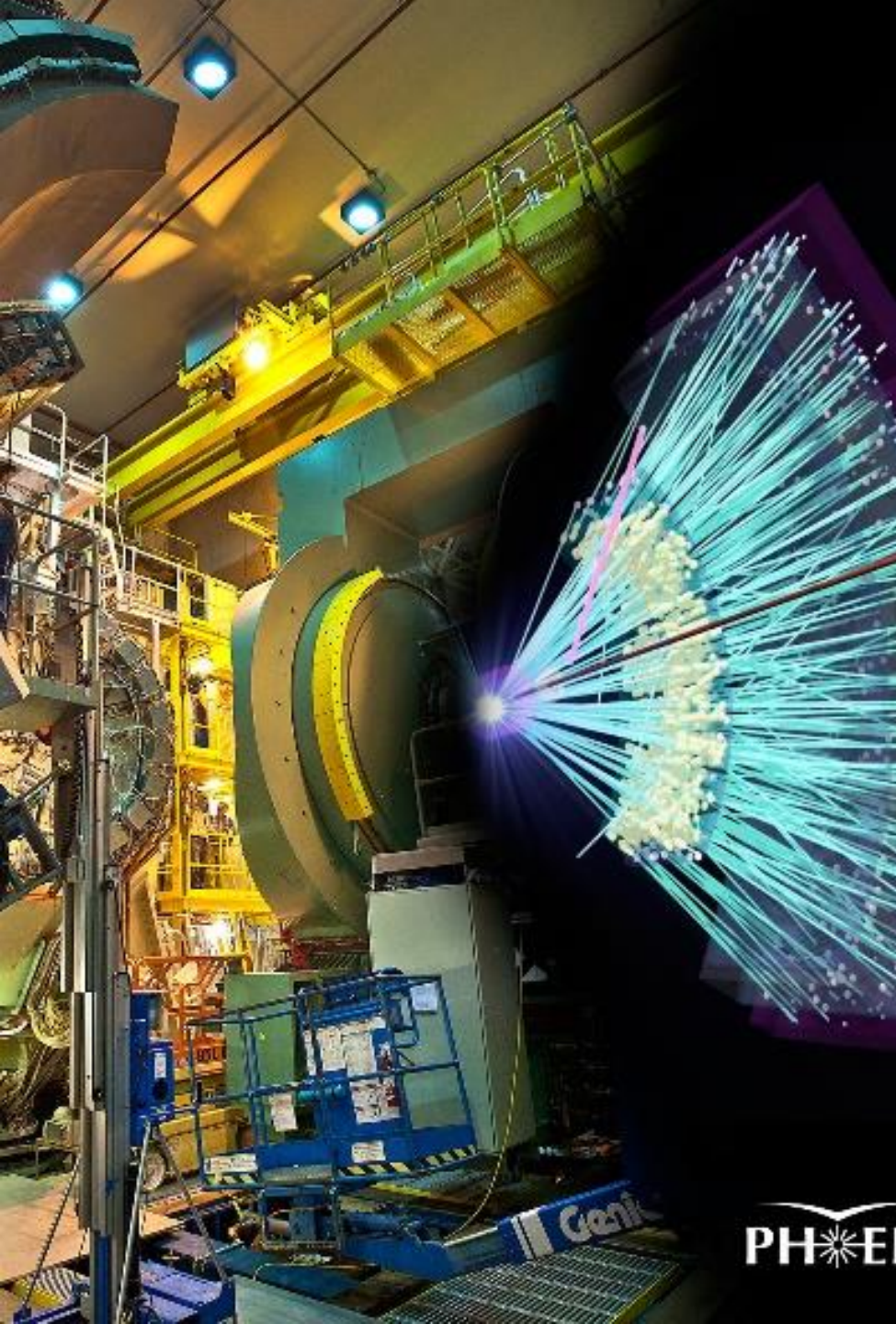


# Sign of $\eta'$ mass modification?



- Calculate  $\lambda$  from SHAREv3 for resonance production
- Resonance chain decay included
- In medium  $\eta'$  mass modification is compatible with the results in each centrality class





# Summary and outlook

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Precise measurement of BEC requires Levy

$$1 < \alpha < 2$$

Levy scale  $R$  exhibits hydro scaling  $\rightarrow$  size?

Strength parameter indicates in medium mass modification of  $\eta'$  meson

Paper is on the way [arXiv:2407.08586](https://arxiv.org/abs/2407.08586)

THANK YOU FOR YOUR ATTENTION!

# ZIMÁNYI SCHOOL 2024



**L. Kassák: Image architecture**

## 24th ZIMÁNYI SCHOOL WINTER WORKSHOP ON HEAVY ION PHYSICS

**December 2-6, 2024**

**Budapest, Hungary**



**József Zimányi (1931 - 2006)**

<http://zimanyischool.kfki.hu/24/>



# References

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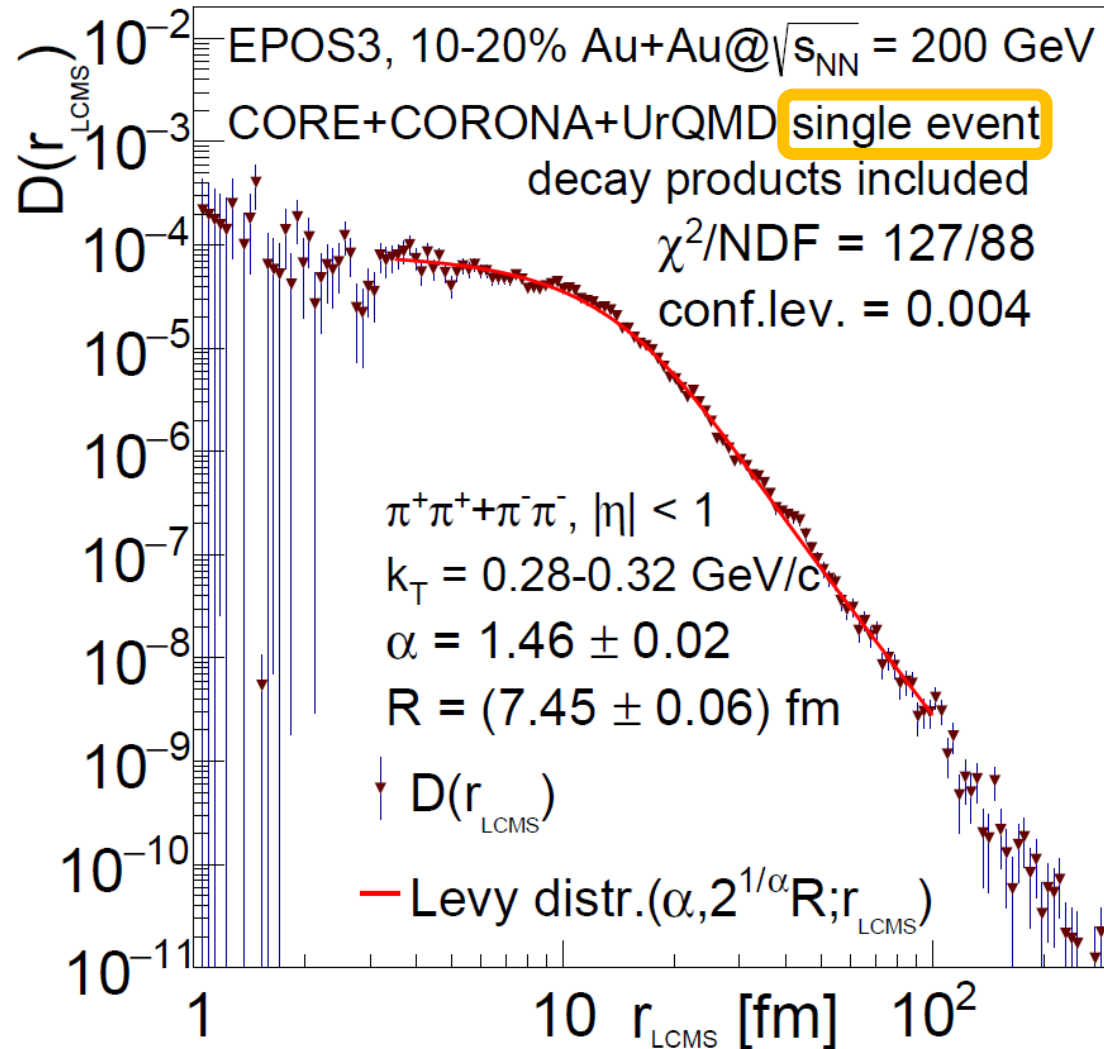
[Universe 4 \(2018\) 2, 31](#)

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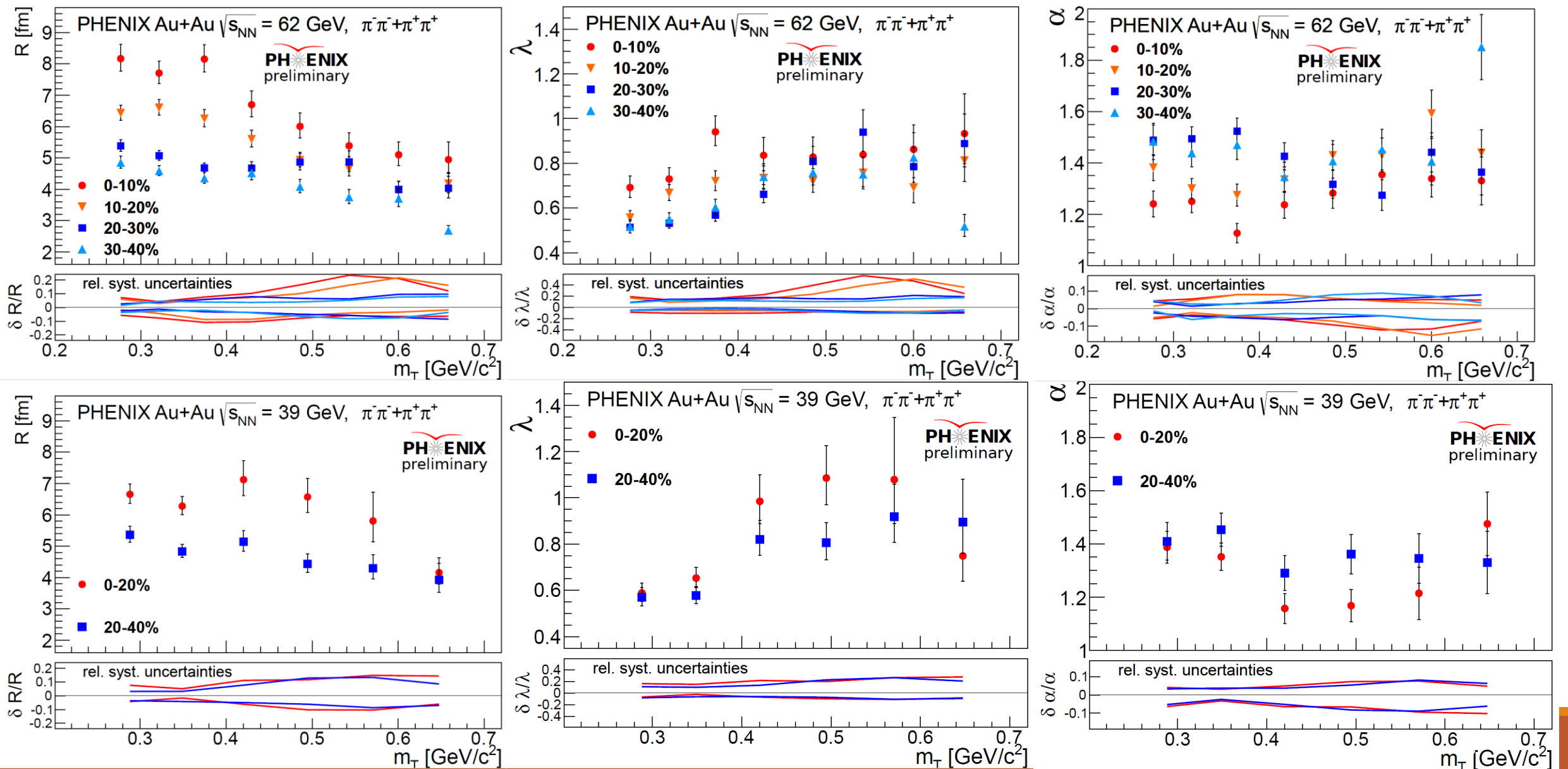
[https://arxiv.org/abs/2407.08586](#)

# EPOS simulation – event-by-event correlation

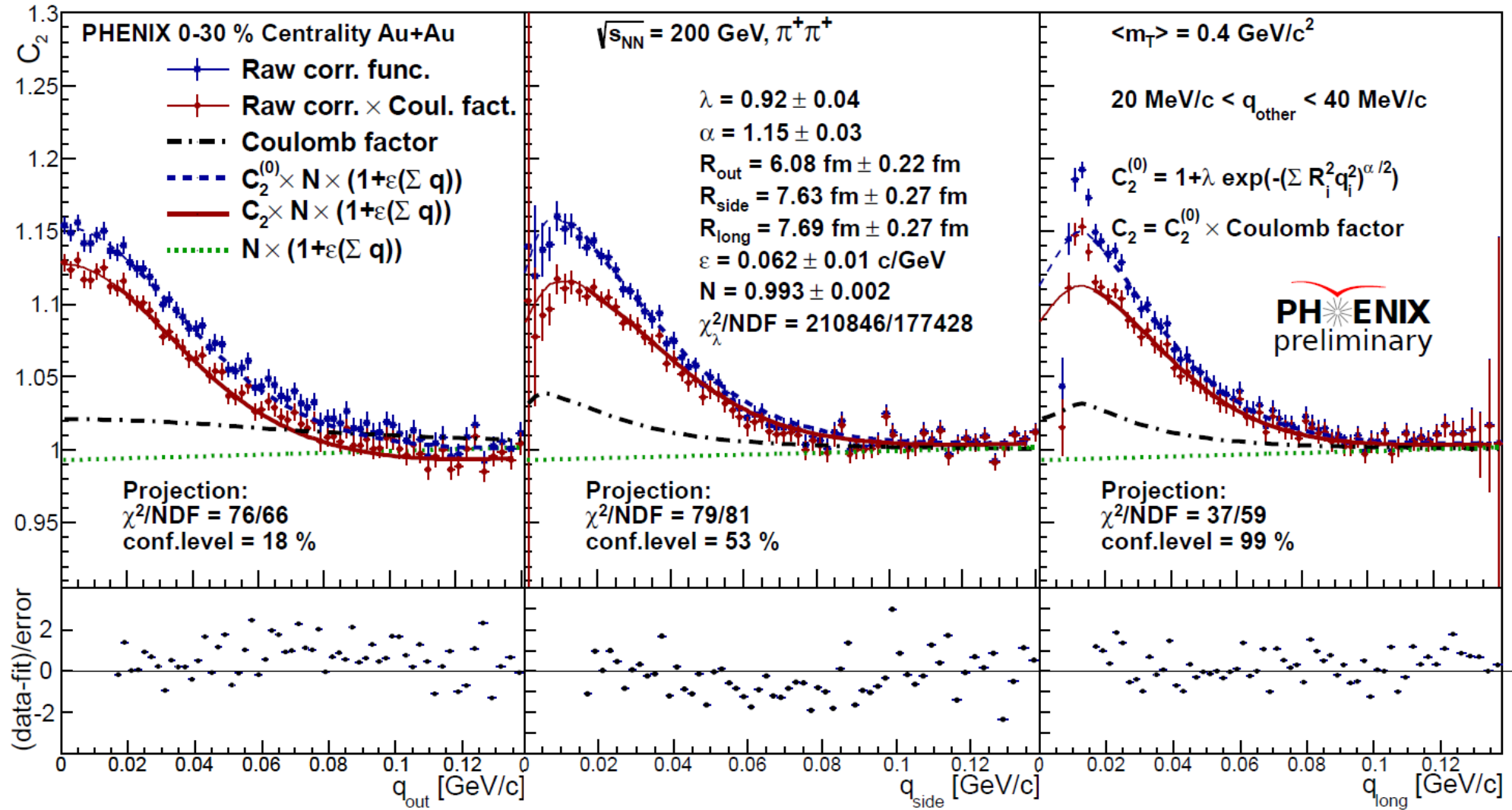


- Core-halo picture is included
- UrQMD for the hadronic cascade
- Levy gives the good description
- It is a single event!
- This analysis support that the origin of the Levy shape could be explained only with the experimental averaging
- This analysis also support the role of the resonances, i.e., the anomalous diffusion
- With this confidence let's look at other experiments

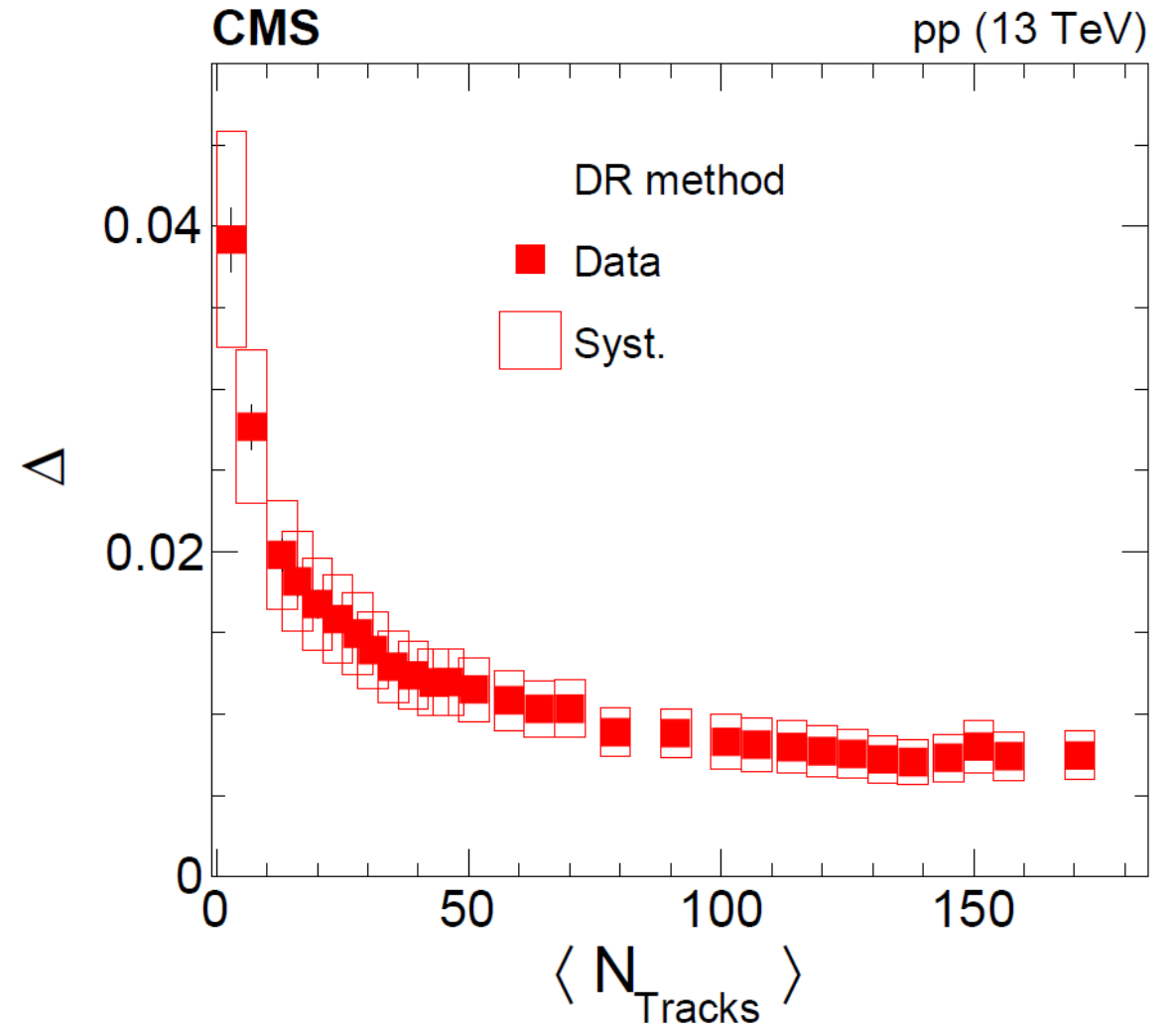
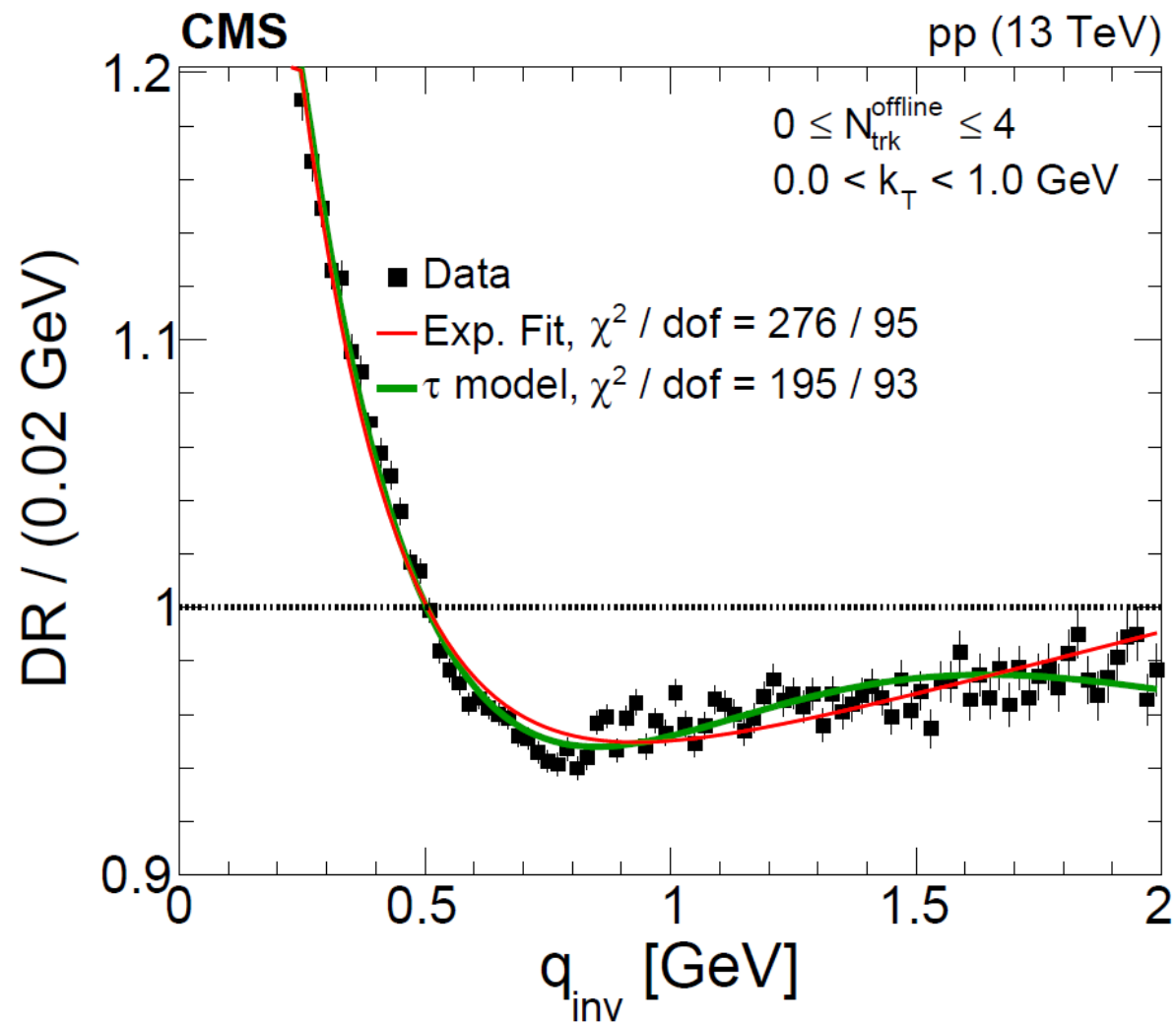
# Backup slides



# Backup slides



# Backup slides



# Backup slides

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$$Q = \sqrt{(p_{1x} - p_{2x})^2 + (p_{1y} - p_{2y})^2 + q_{\text{long,LCMS}}^2},$$

where

$$q_{\text{long,LCMS}}^2 = \frac{4(p_{1z}E_2 - p_{2z}E_1)^2}{(E_1 + E_2)^2 - (p_{1z} + p_{2z})^2}$$

# Femtoscscopy – the core-halo model

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Usually pions, kaons, protons are measured

Resonance contributions are considerable: core-halo model

$$S(x, p) = \sqrt{\lambda} S_{core}(x, p) + (1 - \sqrt{\lambda}) S_{halo}^{Rh}(x, p)$$

Let's introduce the pair source function as

$$D_{AB}(x, p) = \int d^3R S_A\left(R + \frac{x}{2}, p\right) S_B\left(R - \frac{x}{2}, p\right)$$

With this the pair source function in the core-halo model:

$$D(x, p) = \lambda D_{cc}(x, p) + \underbrace{2\sqrt{\lambda}(1 - \sqrt{\lambda}) D_{ch}(x, p) + (1 - \sqrt{\lambda})^2 D_{hh}(x, p)}_{\text{Notation: } D_{(h)}/(1 - \lambda)}$$



# Femtoscscopy – general form

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With  $K = 0.5(p_1 + p_2)$  and  $Q = p_1 - p_2$ ! Also assume that  $p_1 \approx p_2$

$$C_2(Q, K) \approx \lambda \int d^3r D_{cc}(r, K) \left| \Psi_2^{(Q)}(r) \right|^2 + (1 - \lambda) \int d^3r D_{(h)}(r, K) \left| \Psi_2^{(Q)}(r) \right|^2$$

If we take the  $R_h \rightarrow \infty$  limit the Bowler-Sinyukov formula is given:

$$C_2(Q, K) \approx 1 - \lambda + \lambda \int d^3r D_{cc}(r, K) \left| \Psi_2^{(Q)}(r) \right|^2$$

The simple planewave case (i.e. no FSI):

$$C_2^{(0)}(Q, K) = 1 + \lambda \frac{\tilde{D}_c(Q, K)}{\tilde{D}_c(Q = 0, K)}$$

# On the 3D variable of the correlation function

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$$C_2(Q, K) \approx 1 - \lambda + \lambda \int d^3r D_{cc}(r, K) \left| \Psi_2^{(Q)}(r) \right|^2$$

The  $K$  dependence is much smoother than the  $Q$  dependence

Use the  $Q$  as a variable and measure the  $K$  dep. of the params.

$$Q \cdot K = (p_1 - p_2)(p_1 + p_2) = p_1^2 - p_2^2 = 0 \rightarrow Q_0 = \vec{Q} \frac{\vec{K}}{K_0}$$

$C_2(Q)$  can be transformed to  $C_2(\vec{Q})$

Go to LCM system where  $\vec{Q} = (Q_{out}, Q_{side}, Q_{long})$

# On the **1D** variable of the correlation function

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What about in 1D? Could be necessary due to the lack of statistics

Usual choice:  $q_{inv} = \sqrt{-Q^\mu Q_\mu}$ , arguable choice!

$$q_{inv} = (1 - \beta_t^2)Q_{out}^2 + Q_{side}^2 + Q_{long}^2$$

But  $q_{inv}$  could be very small even if  $Q_{out}^2 \approx Q_{side}^2 \approx Q_{long}^2 \neq 0$

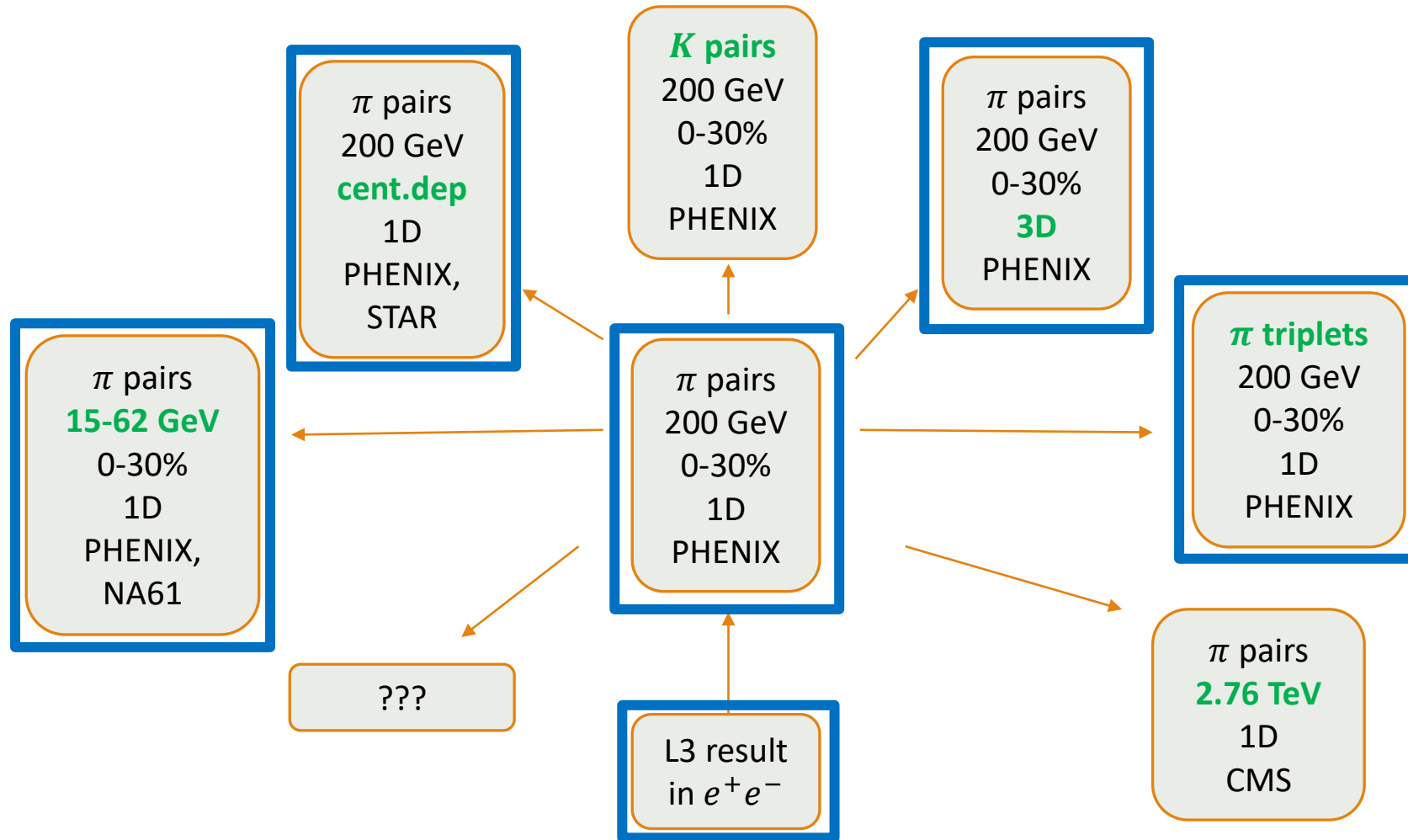
It is also known that the source approximately spherical at RHIC

$$q_{inv} = |q_{PCMS}| \Rightarrow Q = |q_{LCMS}|$$

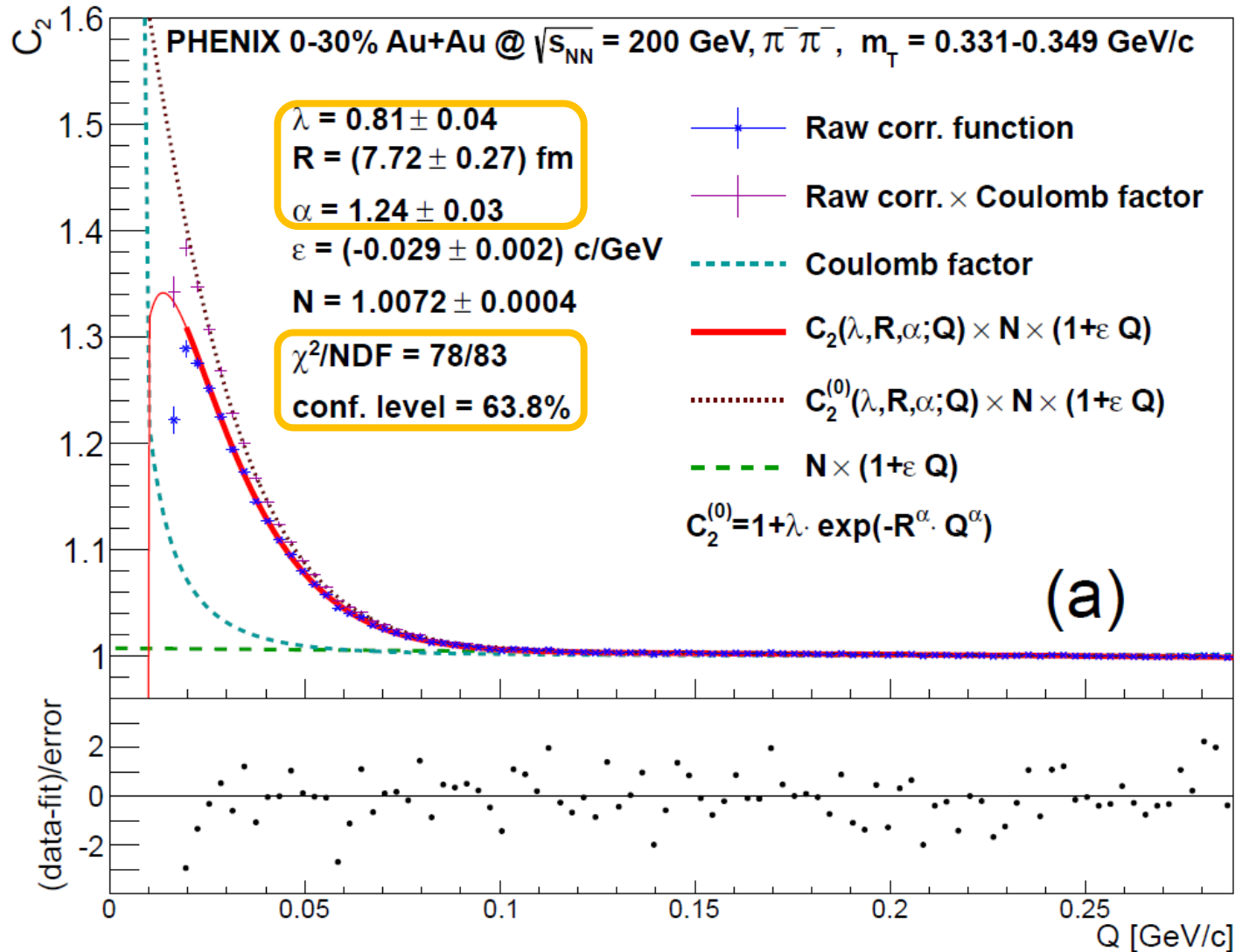
1D variable!

Here, sphericity preserved, so  $Q$  independent of the direction of  $q_{LCMS}$

# The tree of the Levy analyses

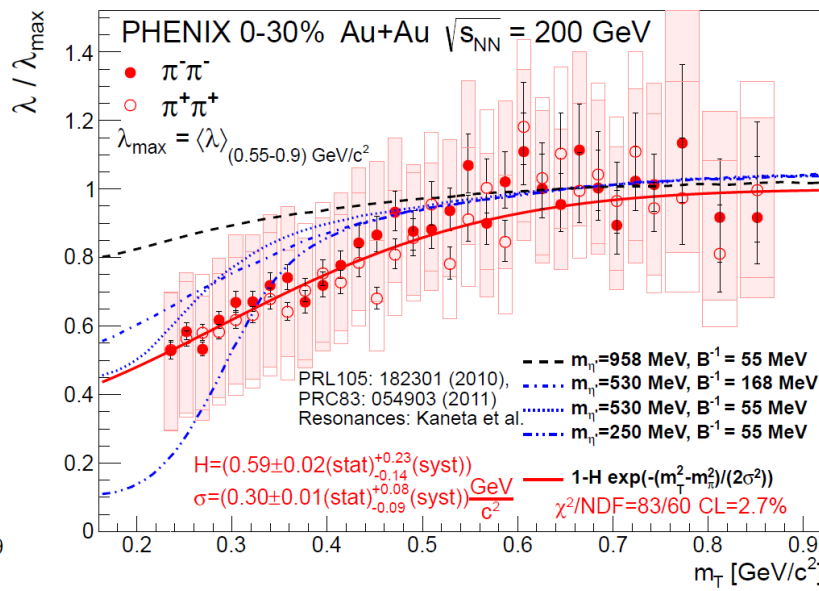
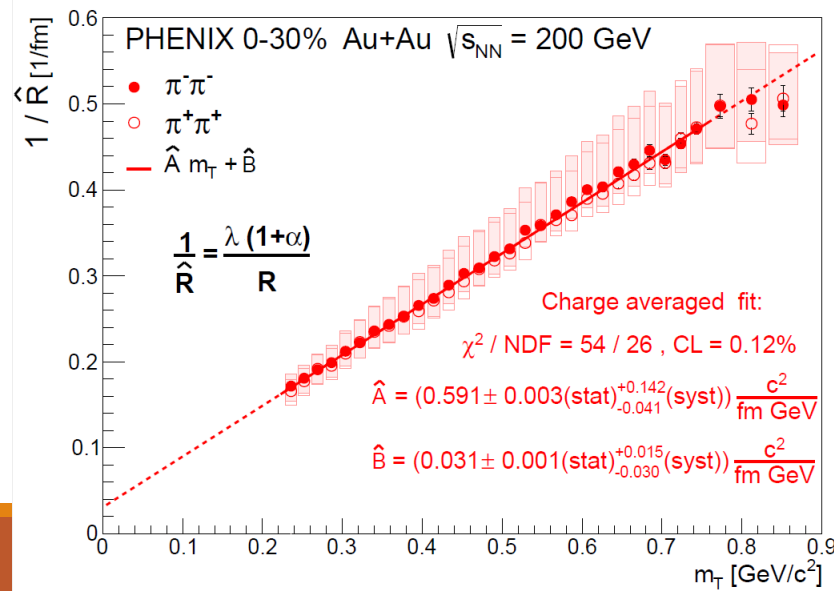
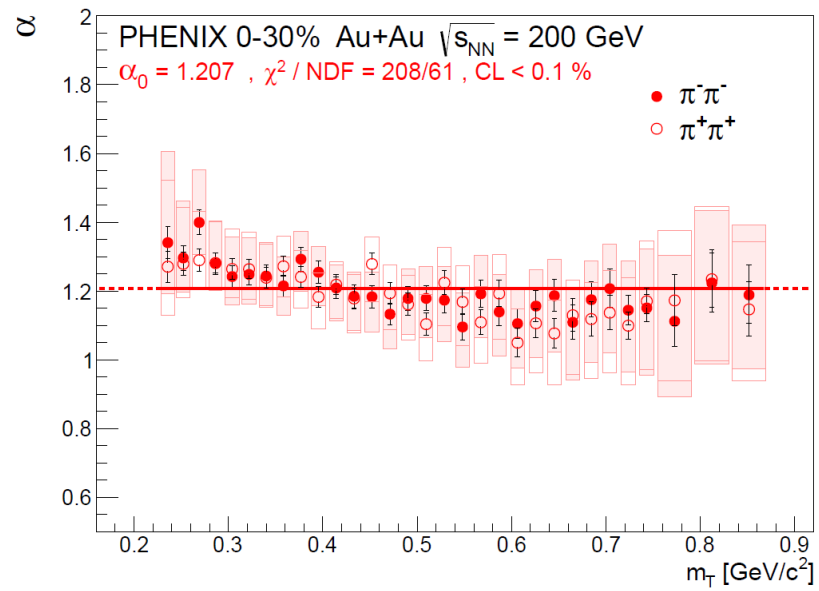
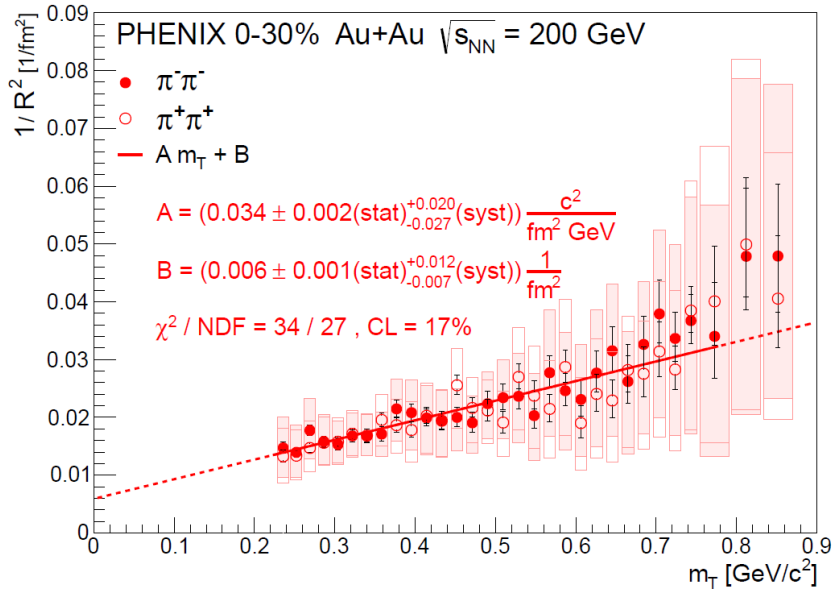


# The first results – PHENIX 0-30% Au+Au



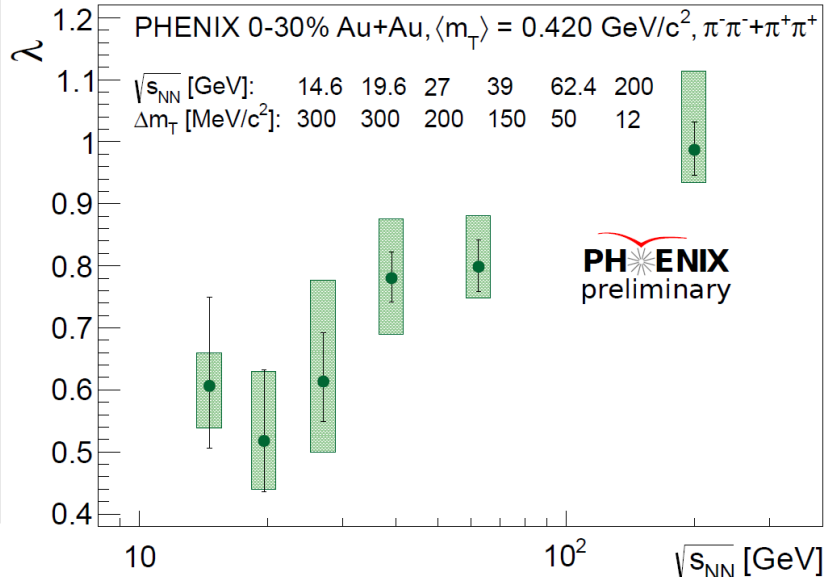
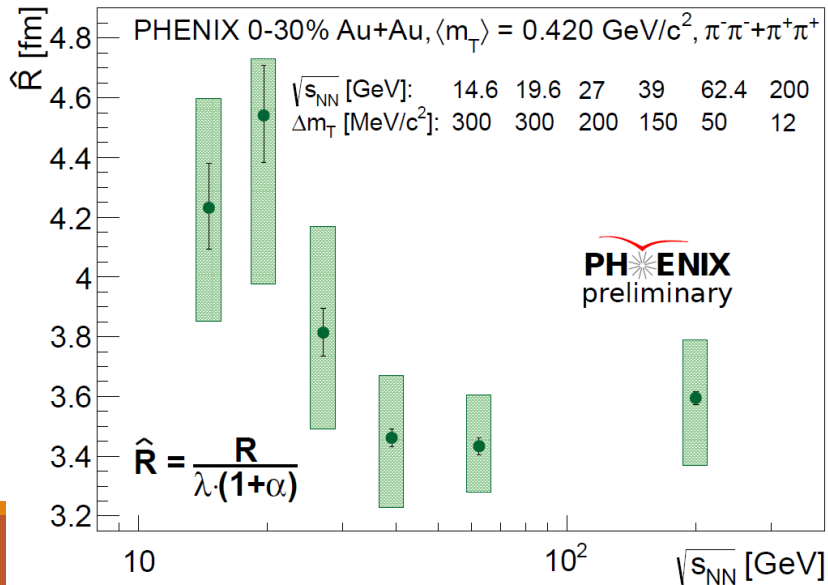
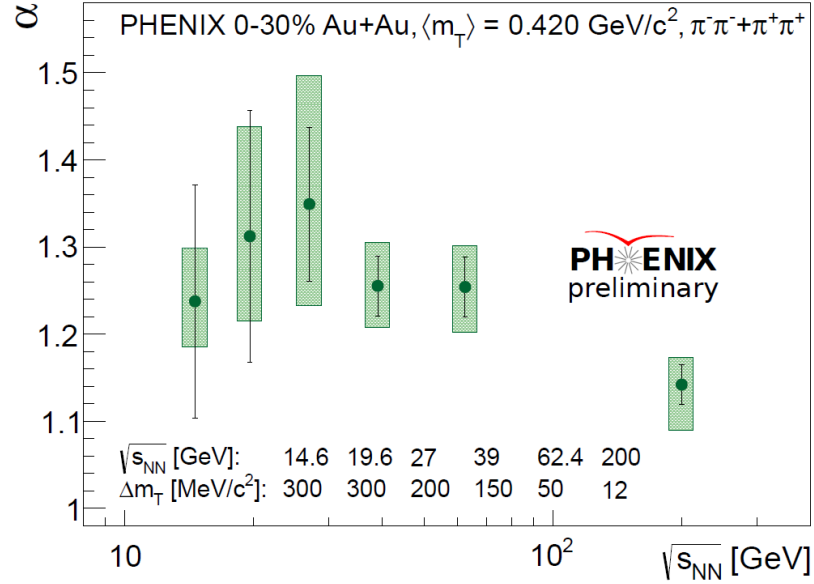
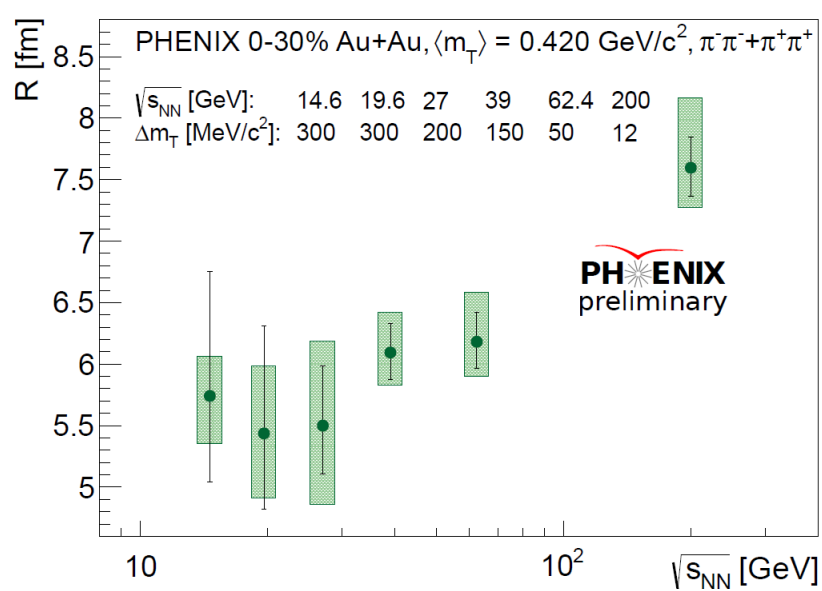
- Measured correlation function in 31  $m_T$  bin with 0-30% cent.
- Coulomb correction incorporated into the fit function
- $\alpha \neq 2$  nor  $\alpha \neq 1$
- The fits are acceptable in terms of confidence level and  $\chi^2/NDF$
- Gaussian parametrization cannot describe the data

# The first results – PHENIX 0-30%, Au+Au



- $R$  exhibits hydro scaling
- $1 < \alpha < 2, \langle \alpha \rangle \approx 1.2$
- $\lambda(m_T)$  suppressed which compatible with modified  $\eta'$  mass in the medium (compared with a resonance model)
- New scaling parameter
  - Interpretation?
- Interpretation of  $\alpha$  ?
- Let's see the  $N_{part}$  and  $\sqrt{s_{NN}}$  dependence

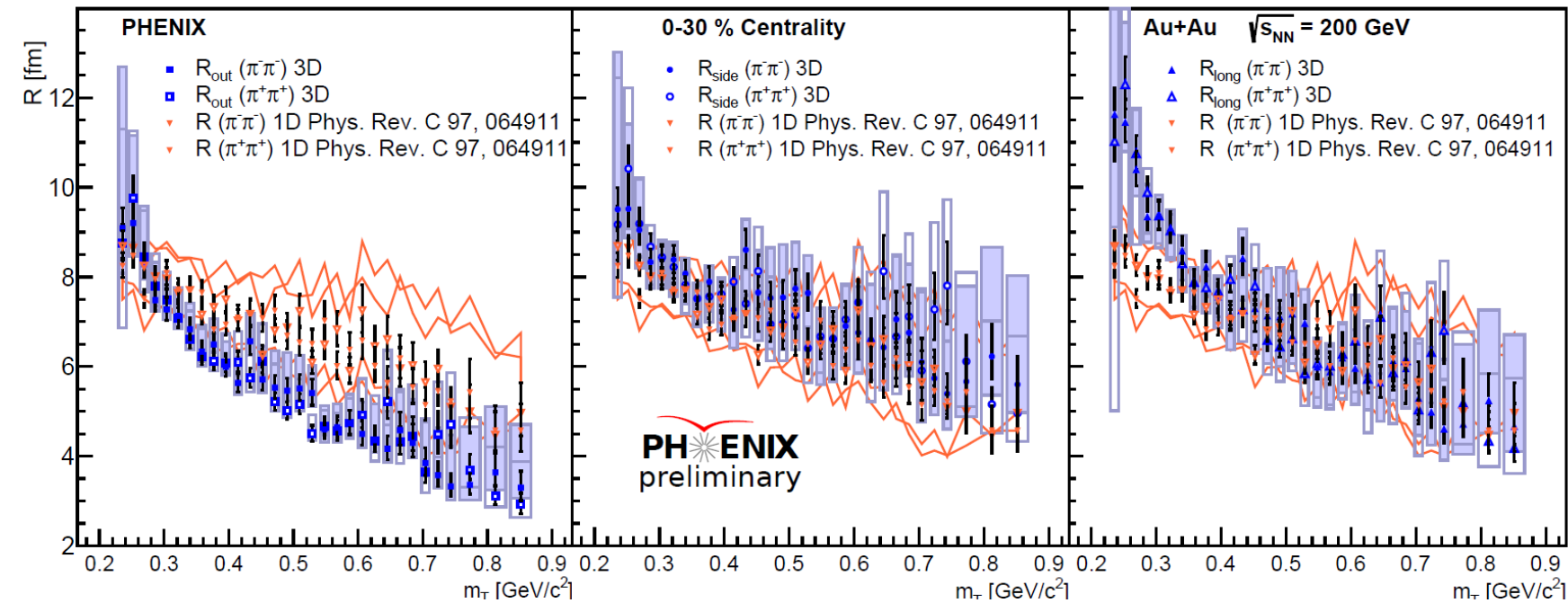
# $\sqrt{s_{NN}}$ dependence – PHENIX Au+Au



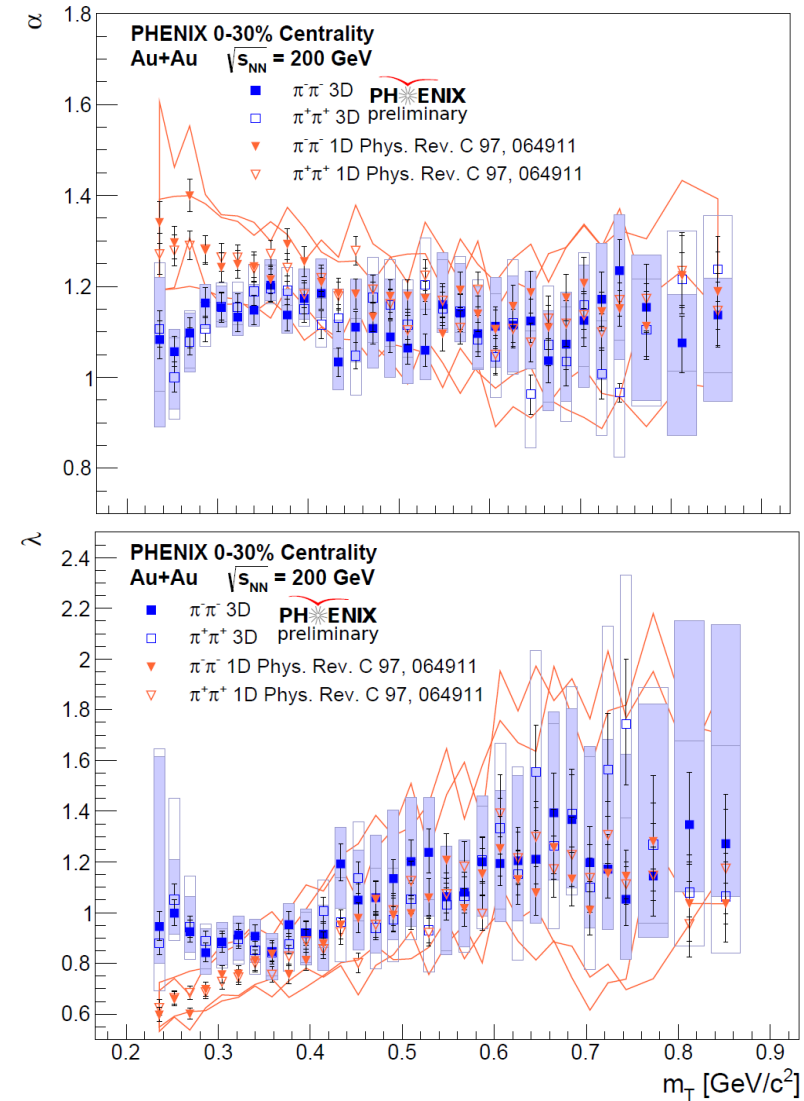
- Integrated in  $m_T$  due to the lack of statistics
- $\alpha$  does not really depend on  $\sqrt{s_{NN}}$
- Non-monotonic behavior of  $\hat{R}$  observed
  - Interpretation?
- For  $\sqrt{s_{NN}} \geq 39 \text{ GeV}$  there are  $m_T$  dependent analysis but the trends are not clear



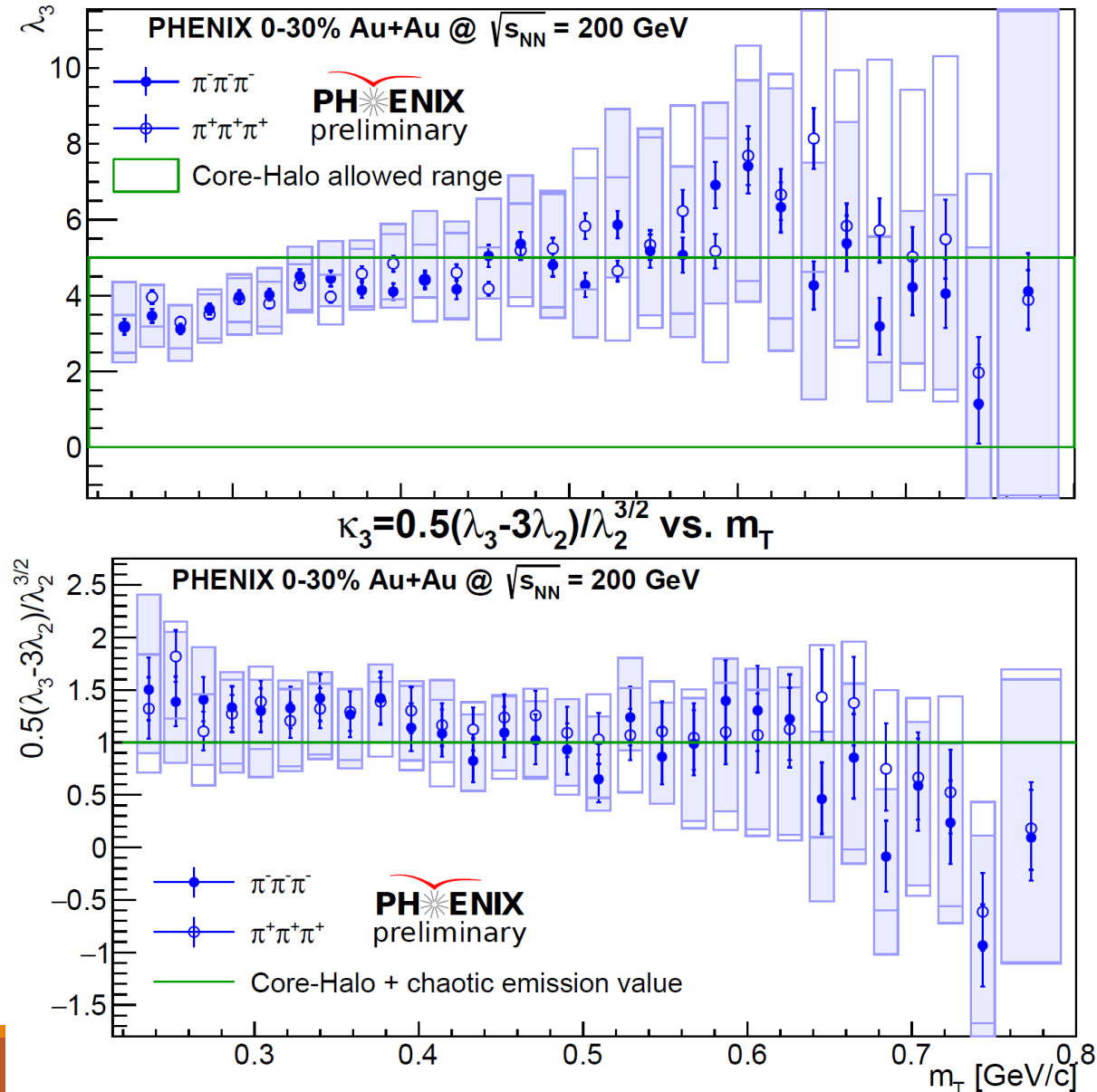
# 3D correlation – PHENIX 0-30% Au+Au



- 3D measurement gives very similar results compared to 1D
- The source appears to be spherical
- $\lambda$  suppression is there in 3D too, with small discrepancy
- Preliminary data!



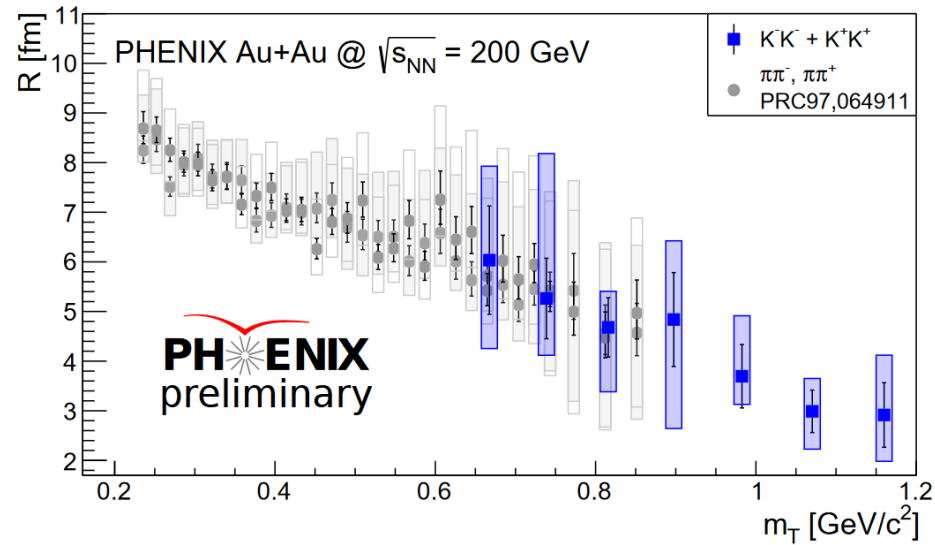
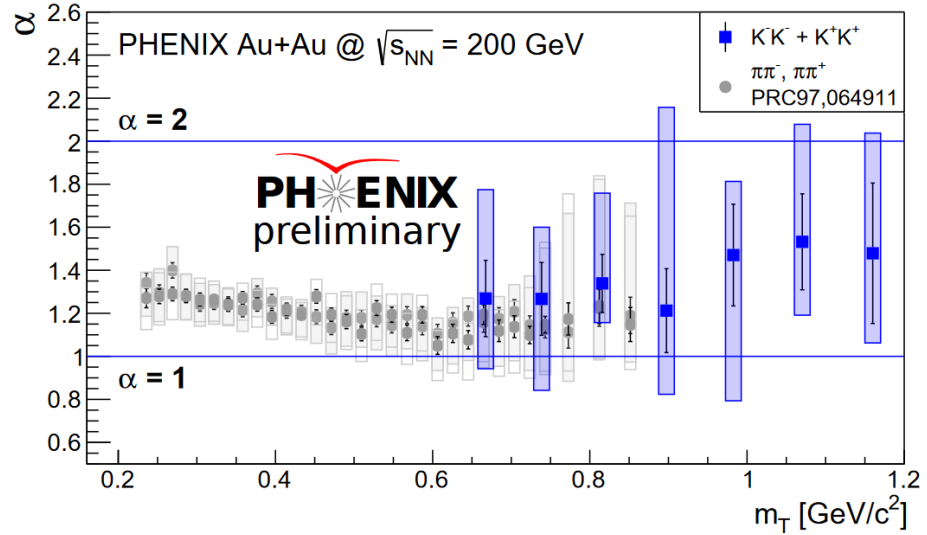
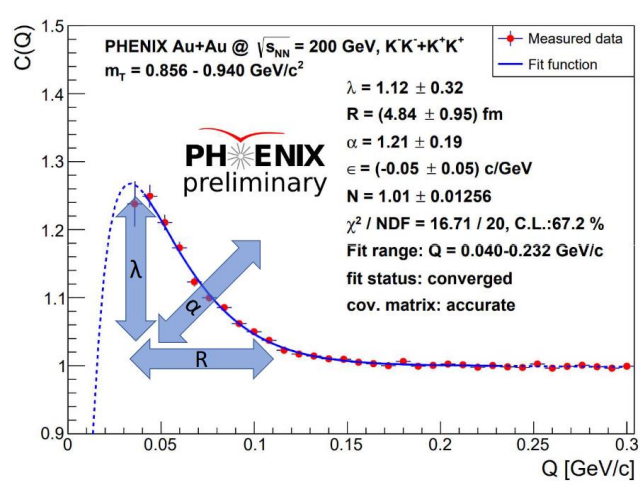
# 3 particle correlation – PHENIX 0-30% Au+Au



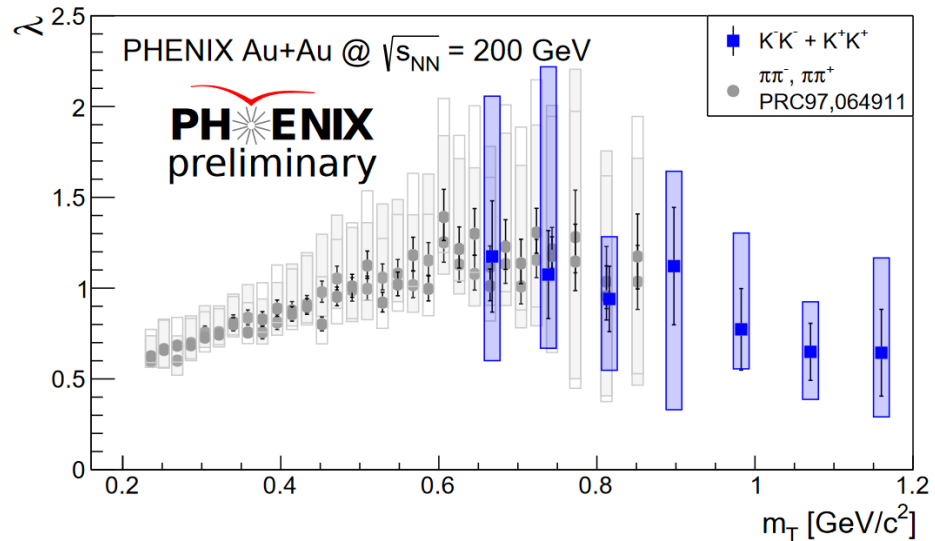
$$\kappa_3 = \frac{(\lambda_3 - 3\lambda_2)}{2\sqrt{\lambda_2^3}}$$

- From the definition:
  - No coherence:  $p_c = 0 \Rightarrow \kappa = 1$
  - Coherence:  $p_c > 0 \Rightarrow \kappa < 1$
- The source seems to be chaotic

# PHENIX @ 200 GeV – kaon correlation in Au+Au



- $\alpha_K \approx \alpha_\pi$  underlying Levy process?
- $\lambda$  exhibits decreasing trends – unidentified hadrons
- R supports its geometrical interpretation as before
- Preliminary results



# Partial conclusions and critiques

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Gaussian parametrization clearly not acceptable in terms of  $\chi^2/NDF$  and CL

Levy gives satisfactory description of the measured 1D data at RHIC BES 1 energies in Au+Au collisions

$1 < \alpha < 2$ , doesn't depend on  $m_T$  strongly but centrality dependent

Why? Two main explanation besides the aforementioned:

- We use 1D variable which has an influence. In 3D it would be Gaussian!
- We measure the average of many Gaussian correlation functions with different width so the average is not Gaussian