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## Why do we study charm photoproduction in **UPCs?**

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- They provide constraints on the nuclear parton-distribution functions (nPDFs) of gluons in nuclei down to low x and on the emergence of the saturation regime.
- Ultra-peripheral heavy-ion collisions (UPCs) are characterized by a  $(10)^{2}$ clean experimental environment, without significant contamination from final-state interactions.
- Open-heavy flavor measurements provide access to a wide region of  $(x,Q^2)$  space.
- Since  $m_{c,b} > \Lambda_{QCD}$ , heavy-flavor measurements can be described by  $(1)^2$ perturbative QCD calculations down to  $p_T = 0$  GeV/c.



COIL

EE

 $\sim 11 \text{ m}$ 

 $\operatorname{HE}$ 

Beam line

HB

EB

Tracker

Interaction point

 $(100)^{2}$ 

## How our signal events look like?

- D<sup>0</sup> mesons can be produced in ultra-peripheral Pb-Pb collisions in scatterings of quasi-real photons emitted by one nucleus with **partons** from the other colliding nucleus
- Leading-order direct process: a photon from one nucleus scatters directly on a gluon of the other nucleus and creates a  $c - \overline{c}$  pair
  - The measurement includes contributions from:
    - Direct and resolved-photon mechanisms
    - Prompt (c  $\rightarrow$  D<sup>0</sup>) and non-prompt (b  $\rightarrow$  D<sup>0</sup>) events <sub>Pb</sub>
    - Decay channel:  $D^0 \rightarrow K^- \pi^+$  (and charge conj.)



#### New Level-1 trigger strategy for photoproduction with CMS:

- The CMS Level-1 (L1) trigger system uses custom hardware processes and fast information from the calorimeters to achieve the first significant data rate reduction
- **Challenge**: Very high rate compared to hadronic events  $\rightarrow$  need for a strong rate reduction and quick decision to maximize the collected statistics
- $\rightarrow$  **Solution**: Level-1 trigger that uses both ZDC and HCAL/ECAL information
- CMS has a unique coverage for photonuclear measurements:
- Tracker:  $|\eta| < 2.4$

- ECAL and HCAL:  $|\eta| < 3.0$
- HF calorimeter:  $3.0 < |\eta| < 5.2$
- ZDC: |η| > 8.3

#### Kinematic coverage and trigger choice:

- For high  $p_T D^0$  ( $p_T > 5$  GeV/c), L1 jet (8 GeV) trigger combined with ZDC exclusive OR condition (Xn0n or 0nXn)
- **For low p\_T D^0 (2 < p\_T < 5 GeV/c), jet** triggers become very inefficient  $\rightarrow$  L1 ZDCOR 1n: at least one of the ZDC signals under the 1 neutron threshold



 $D^0$  flight length

(decay length)

#### **Offline event selections**

**ZDC selection**: exactly one ZDC signal is above the 1n threshold (Xn0n or 0nXn, X>0)

ZDC

HCAL

 $\operatorname{HF}$ 

 $\sim 140~{\rm m}$ 

- **Rapidity gap** on the side of "empty" ZDC: No particle flow candidate is found above a given energy in the HF calorimeter
- Primary vertex selection and beam gas/machine induced background rejection

## **Offline D<sup>0</sup> selections**

- D<sup>0</sup> candidates reconstructed from oppositely charged pairs of high-purity tracks
- Rectangular cuts (validated with a BDT-based optimization), which exploit the following topological variables:



- 3D decay length, normalized by its uncertainty
- Pointing ( $\alpha$ ) and opening angles ( $\Delta \theta$ )
- Secondary vertex reconstruction probability

#### **Efficiency and trigger corrections:**

- Data-driven jet trigger correction in bins of  $D^0 p_T$  and y
- MC-based (Pythia 8) corrections in bins of  $D^0$  ( $p_T$ , y) used for event selection efficiency, D<sup>0</sup> acceptance, reconstruction and selection efficiency
- Pythia 8 distributions of generated-level D<sup>0</sup> reweighted according to FONLL-based predictions



flight

Secondary vertex

 $\mathbf{p}_{ ext{track }2}$ 

1.1.1.1.1.1

 $K^{-}$ 

#### Signal extraction and cross section calculations

- Signal yield extracted in intervals of  $D^0 p_T$  and y with unbinned fits to the  $D^0$  invariant mass distributions:
  - Double-Gaussian signal shape to model the distribution of signal D<sup>0</sup>
  - Single Gaussian for D<sup>0</sup> candidates with wrong mass hypothesis
  - Exponential shape for the combinatorial background
  - Crystal Ball function templates for  ${
    m D}^0 o {
    m K}^- {
    m K}^+$  and  ${
    m D}^0 o \pi^- \pi^+$  decay channels
- Corrected cross section:



# $\frac{d^2\sigma}{dp_{\rm T}dy}({\rm D}^0p_{\rm T},{\rm D}^0y) = \frac{1}{2}\frac{1}{\mathcal{LB}}\frac{N_{\rm D}^{\rm raw}}{\Delta p_{\rm T}\Delta y} \frac{1}{\epsilon_{\rm evt}\ \epsilon_{\rm trigger}\ P_{\rm prescale}\ \left(\alpha\ \epsilon_{\rm D}0\right)\ \epsilon_{\rm EM,pileup}$

Systematic uncertainties included for luminosity estimation, electromagnetic pileup correction, track reconstruction, event selection efficiencies, D<sup>0</sup> acceptance, selection and reconstruction efficiencies, yield extraction, trigger correction, MC reweighting

#### **Conclusions:**

- New experimental constraints obtained on nuclear matter with heavy-quark observables in a large region of x and Q<sup>2</sup>
- This measurement opens the way for a new experimental program exploiting fully-reconstructed heavy-flavour hadrons and heavy-flavour jets in ultraperipheral heavy-ion collisions at the LHC.

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