

# The real corrections to the Next-to-Leading Order Higgs Impact Factor at Physical Top Quark Mass

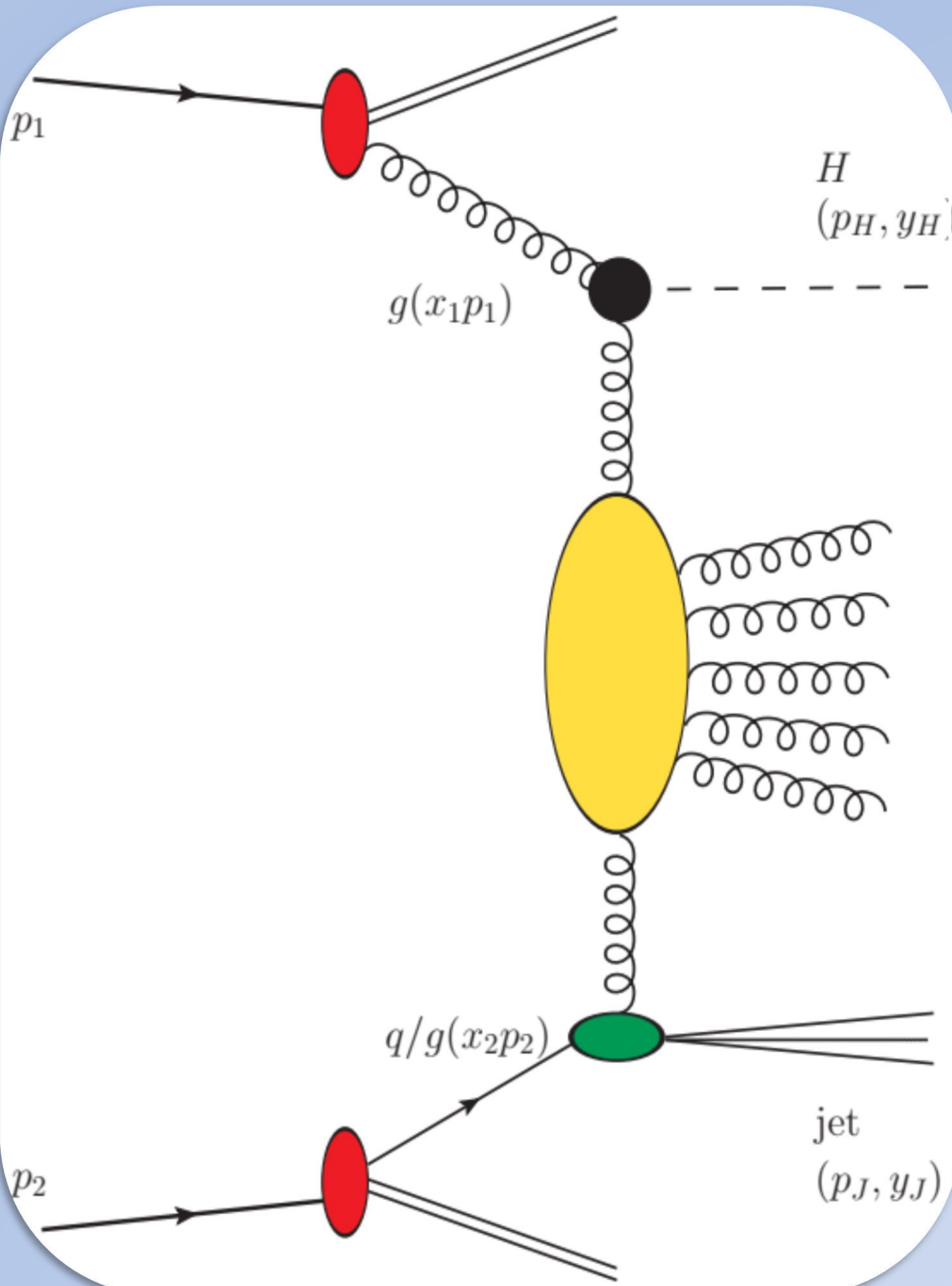
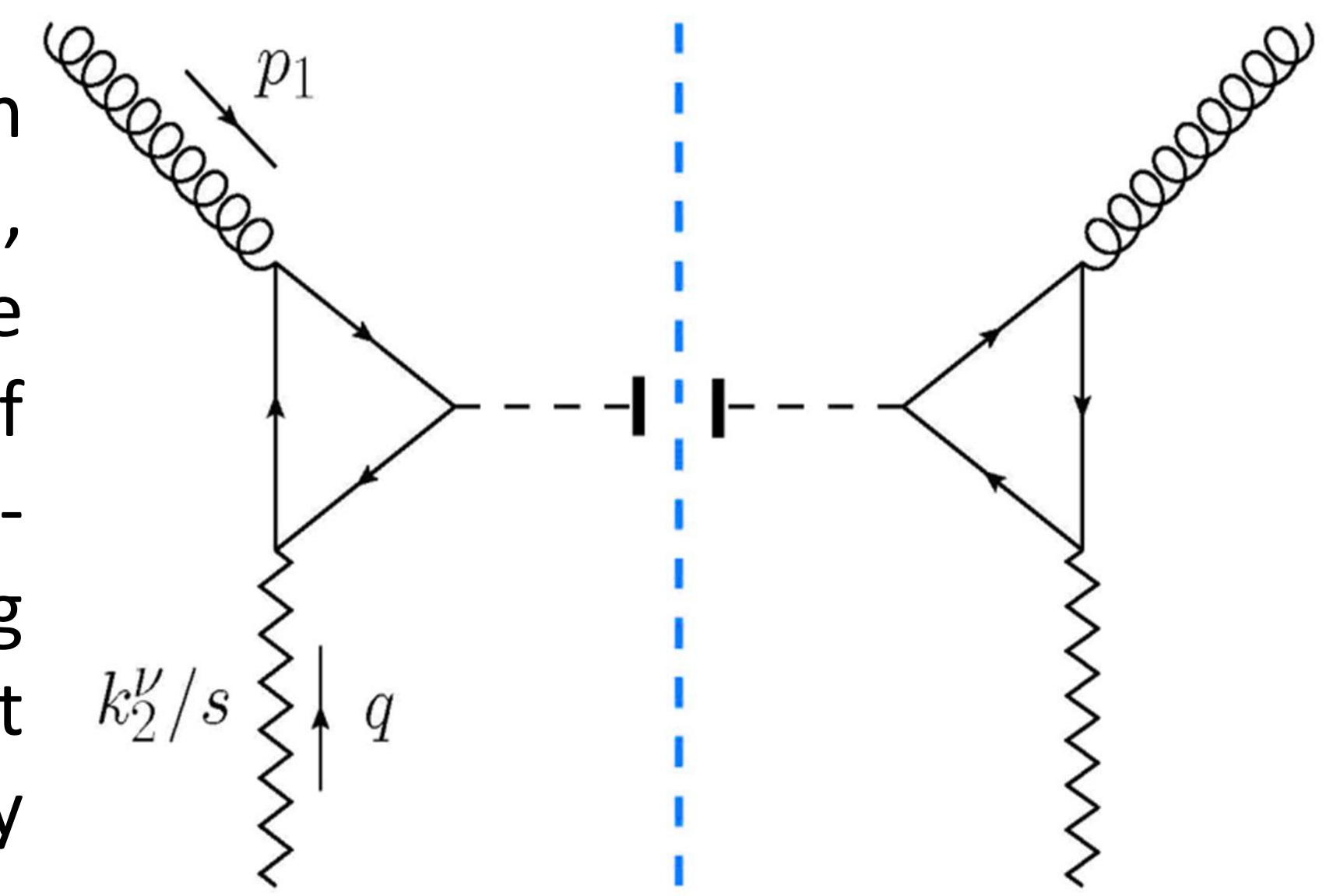
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## Overview

We calculate **real corrections** to the **forward Higgs boson impact factor** with full top-mass dependence. It will enable new predictions for high-energy processes such as: inclusive forward Higgs production, inclusive production of a forward Higgs with a backward object (jet or hadron) at the LHC and future hadron colliders.

The impact factor is a function that describes the transition between the initial and final state of a particle in its fragmentation region, capturing the specific dynamics of the interaction process. In the **BFKL approach**, differential cross sections take the peculiar form of a convolution, in transverse-momenta space, of two process-dependent *impact factors*, describing the transition of each colliding particle to a definite state, and a universal, process-independent *Green's function*, which encodes the resummation of energy logarithms.

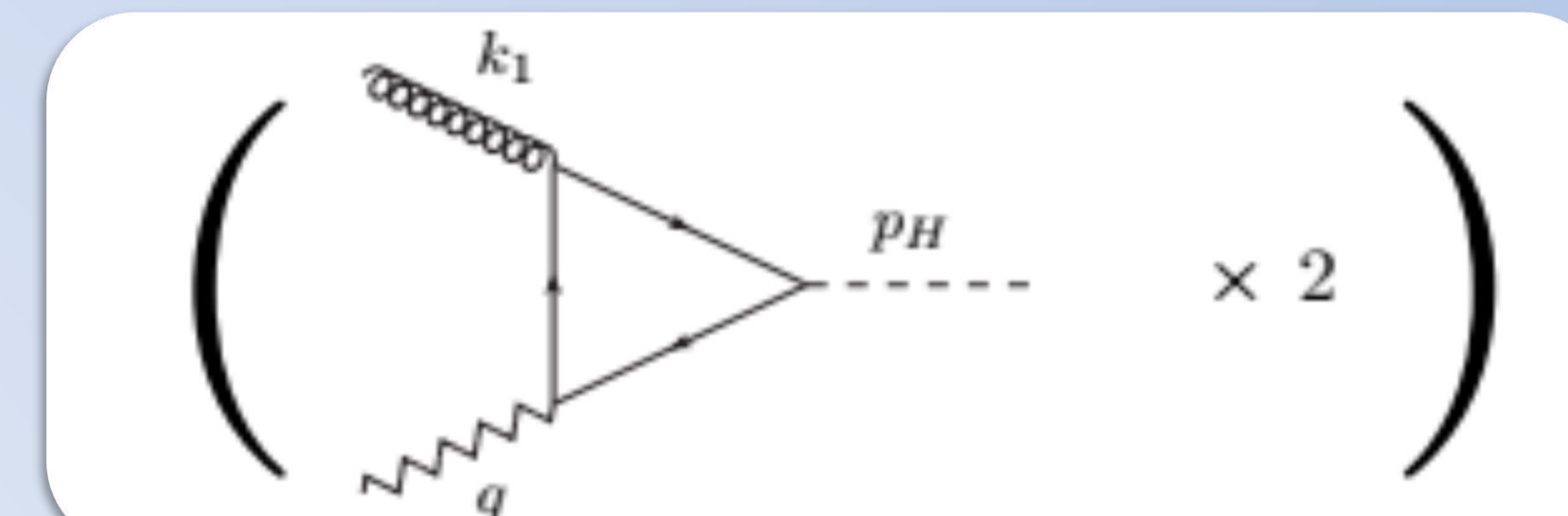


## LO Computation

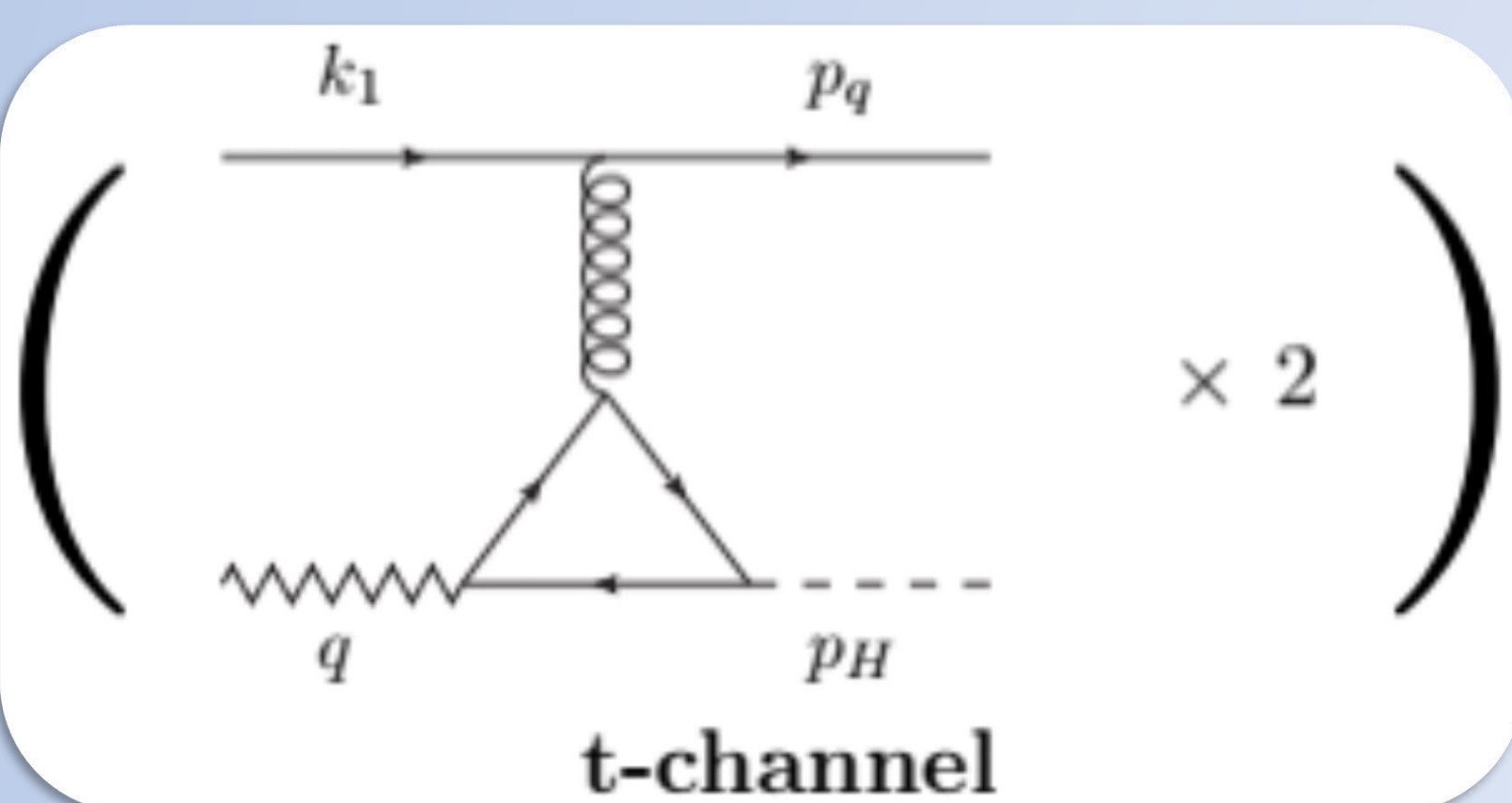
The proton-initiated impact factor reads

$$\frac{d\Phi_{PP}^{\{H\}^{(0)}}(\vec{q})}{dx_H d^2\vec{p}_H} = \int_{x_H}^1 \frac{dz_H}{z_H} f_g\left(\frac{x_H}{z_H}\right) \frac{d\Phi_{gg}^{\{H\}^{(0)}}(\vec{q})}{dz_H d^2\vec{p}_H} = \frac{|F_T(0, -\vec{q}^2, m_H^2)|^2 \vec{q}^2 f_g(x_H)}{8(1-\epsilon)\sqrt{N^2-1}} \delta^{(2)}(\vec{q} - \vec{p}_H),$$

where  $f_g(x_H)$  is the gluon parton distribution function.



The two triangular-like diagrams contributing to the Higgs impact factor at the LO.



The two triangular-like diagrams contributing to the quark-initiated contribution to the Higgs impact factor at the NLO.

## NLO Computation

The Next-to-Leading Order computation analyzes the impact factor for Higgs boson production with full top-mass dependence, addressing infrared, collinear, and rapidity divergences.

### Infrared Structure and Gauge Invariance

Analysis of infrared and collinear singularities in real impact factor corrections.

- **Collinear Singularity:** appears when  $\vec{p}_g = \vec{q} - \vec{p}_H \rightarrow 0$
- **Soft Singularity:** occurs when  $z_H \rightarrow 1$ , indicating gluon energy loss

The complete result for the infrared limit of the impact factor is

$$\frac{d\Phi_{gg}^{\{Hg\}}(z_H, \vec{p}_H, \vec{q}; s_0)}{dz_H d^2\vec{p}_H} \Big|_{\vec{q} \sim \vec{p}_H} \xrightarrow{s_0 \rightarrow 0} \frac{g^2 |F_T(0, -\vec{p}_H^2, m_H^2)|^2 N}{4(1-\epsilon)\sqrt{N^2-1} (2\pi)^{D-1}} \left( \frac{\vec{p}_H^2}{(\vec{q} - \vec{p}_H)^2} \frac{z_H}{1-z_H} \theta(s_\Lambda - s_{gR}) + \frac{1}{(\vec{q} - \vec{p}_H)^2} \left[ z_H(1-z_H)\vec{p}_H^2 + 2(1-\epsilon) \frac{1-z_H(\vec{p}_H \cdot (\vec{q} - \vec{p}_H))^2}{(\vec{q} - \vec{p}_H)^2} \right] \right).$$

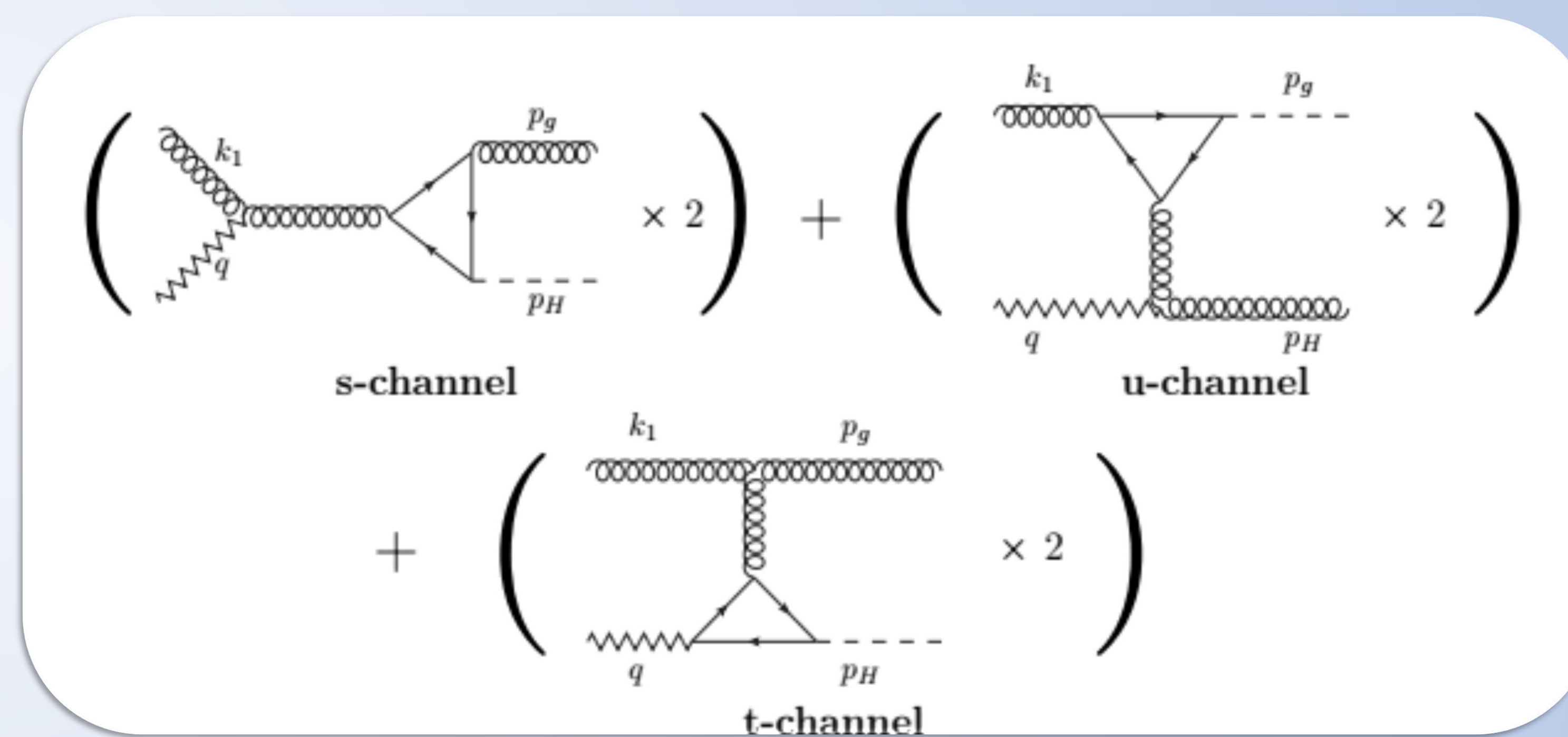
This is the proper structure of divergences and moreover the full impact factor vanishes when the transverse momentum of the Reggeon approaches zero, in agreement with the gauge invariance.

### Rapidity Divergence within the BFKL factorization framework

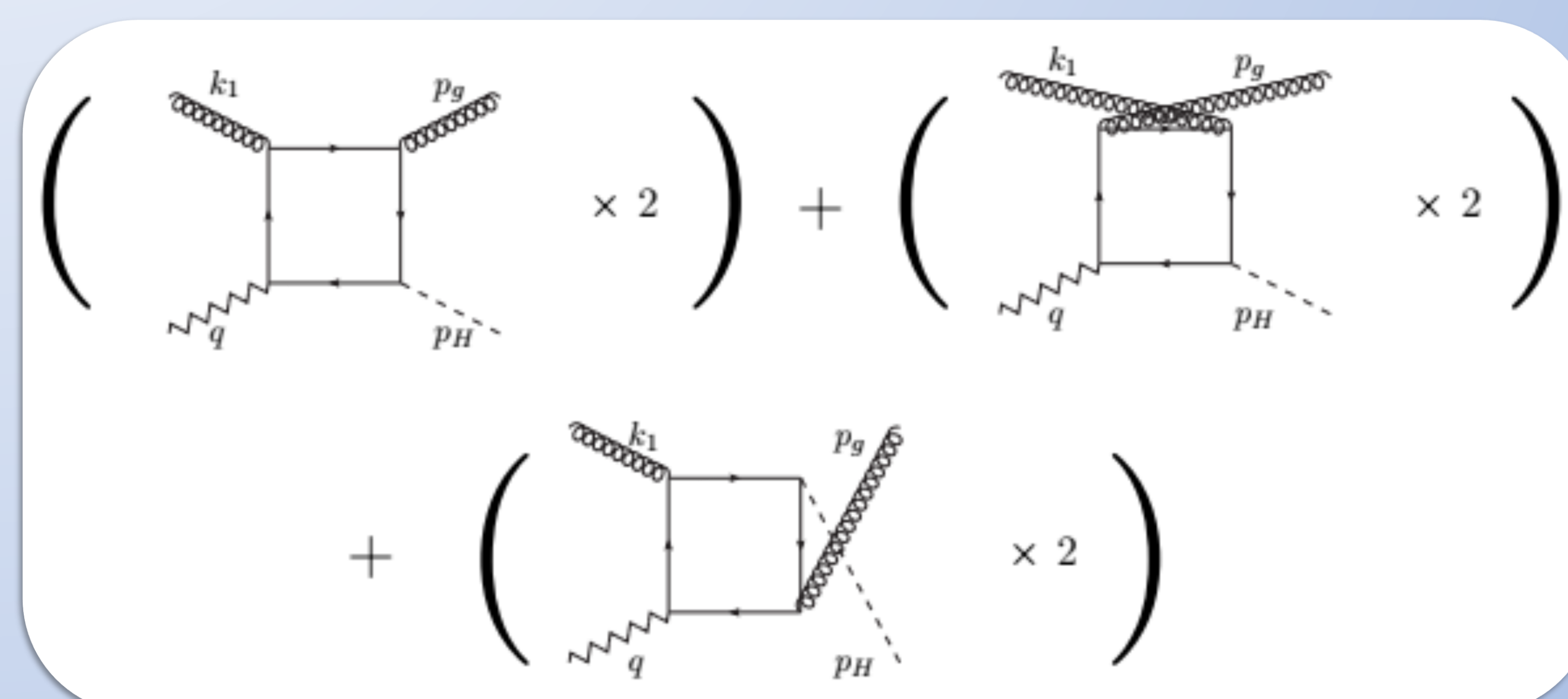
The high-rapidity limit of the impact factor is

$$\frac{d\Phi_{gg}^{\{Hg\}}(z_H, \vec{p}_H, \vec{q}; s_0)}{dz_H d^2\vec{p}_H} \Big|_{z_H \rightarrow 1} = \frac{g^2 |F_T(0, -\vec{p}_H^2, m_H^2)|^2 N}{4(1-\epsilon)\sqrt{N^2-1} (2\pi)^{D-1}} \frac{\vec{q}^2}{(\vec{q} - \vec{p}_H)^2} \frac{1}{(1-z_H)} \theta\left(s_\Lambda - \frac{(\vec{q} - \vec{p}_H)^2}{1-z_H}\right).$$

The rapidity divergence is cancelled through the BFKL counter-term, aligning with theoretical predictions.



The six triangular-like diagrams contributing to the gluon-initiated contribution to the Higgs impact factor at the NLO.



The six box-like diagrams contributing to the gluon initiated contribution to the Higgs impact factor at the NLO.

## Further information

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