## Recent ATLAS results relevant for PDFs at low and high x, saturation in both pp and HI collisions

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on behalf of the ATLAS Collaboration

Diffraction and Low-x



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#### Introduction:

#### **PDF (Parton Distribution Function):**

- Distribution of quarks and gluons (partons) inside a proton.
- Typically used for calculations in proton-proton collisions.

#### nPDF (nuclear Parton Distribution Function):

- Distribution of partons inside a nucleus (which contains multiple protons and neutrons).
- Accounts for modifications due to the presence of multiple nucleons, such as shadowing, EMC effect, and anti-shadowing.

#### x/Q<sup>2</sup> Coverage:

- PDFs: Have extensive coverage over a wide range of x
   (Bjorken-x) and Q<sup>2</sup> (momentum transfer squared) values.
- nPDFs: The x/Q<sup>2</sup> coverage is less extensive, particularly in regions of very low and very high x, leading to uncertainties in certain calculations involving nuclei.



Observation of tt production in lepton+jets and dilepton channels in *p*+Pb collisions at  $\sqrt{s_{NN}}$  = 8.16 TeV with the ATLAS detector

Measurement of the centrality dependence of the dijet yield in *p*+Pb collisions at  $\sqrt{s_{NN}}$  = 8.16 TeV with the ATLAS detector

Photo-nuclear jet production in ultra-peripheral Pb+Pb collisions at  $\sqrt{s_{_{\rm NN}}}$ =5.02 TeV with the ATLAS detector

# Observation of tt production in lepton+jets and dilepton channels in *p*+Pb collisions at $\sqrt{s_{_{\rm NN}}}$ = 8.16 TeV with the ATLAS detector



arXiv:2405.05078v1
Submitted to JHEP

Event candidate for process in ttbar p+Pb

<u>Video</u>

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 $\sqrt{s_{_{NN}}} = 8.16 \text{ TeV}$ 



Run: 313100 Event: 168745611 2016-11-18 22:14:23

#### **Motivation:**

- In ATLAS observation of tt in *p*+Pb data individually in lepton+jets and dilepton channels **this talk** 
  - The first measurement using the **dilepton** channel in *p*+Pb collisions.
  - p+Pb data from 2016 with Integrated luminosity L = 165 nb<sup>-1</sup>
  - All the plots can be found <u>TOPQ-2023-32</u>
- Top quarks provide novel probes of nuclear modifications to parton distribution functions (nPDF) in a poorly constrained kinematic region.
- Bjorken- $x \sim 5 \cdot 10^{-3} 0.05$  and  $Q^2 \sim m_t^2 \sim 3 \cdot 10^4 \text{ GeV}^2$  ( $m_t$  stands for the top-quark mass)





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Simultaneous fit in 6 regions

4j1b1ℓ ( I+jets, 1b, ejets ), 4j1b1ℓ ( I+jets, 1b, mujets ), 4j2bincl1ℓ ( I+jets, >=2b, ejets ), j2bincl1ℓ ( I+jets, >=2b, mujets ), 2j1b2ℓ (dilepton, 1b ), 2j2bincl2ℓ (dilepton, >=2b ).

- The most signal events observed in the *ℓ*+jets
   2 b-jet regions.
- The best signal-to-background ratio found in the **dilepton 2 b-jet** region.

• The main background contributions:

W+jets (l+jets)
 Z+jets (dilepton)
 non-prompt and fake lepton background



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#### Fit results in signal regions:





• Best-fit value for the parameter we ultimately want to extract,

$$\mu_{Total} = 1.04 + 0.094 - 0.088$$

- ATLAS total uncertainty: ±9% (CMS total uncertainty: 18%)
- Relative statistical uncertainty: ~3.5%

Source	$\Delta \sigma_{t \bar{t}} / \sigma_{t \bar{t}}$	
	unc. up [%]	unc. down [%]
Jet energy scale	+4.6	-4.1
$t\bar{t}$ generator	+4.5	-4.0
Fake-lepton background	+3.1	-2.8
Background	+3.1	-2.6
Luminosity	+2.8	-2.5
Muon uncertainties	+2.3	-2.0
W+jets	+2.2	-2.0
<i>b</i> -tagging	+2.1	-1.9
Electron uncertainties	+1.8	-1.5
MC statistical uncertainties	+1.1	-1.0
Jet energy resolution	+0.4	-0.4
tī PDF	+0.1	-0.1
Systematic uncertainty	+8.3	-7.6

### **Observed significances:**

- **I+jets** observed significance : ~  $14\sigma$
- **Dilepton** observed significance : ~ 8 $\sigma$
- **Combined** observed significance: ~ 16 $\sigma$
- Observation of ttbar production in *p*+Pb in **combined** and also in **individual** I+jets and dilepton channels

## First observation in dilepton channel with *p*+Pb collisions

## Cross-Section & R<sub>pA</sub> Measurement:

• The measured inclusive tt cross section for the *p*+Pb collisions is:

$$\sigma_{t\bar{t}} = \mu_{t\bar{t}} \cdot A_{\text{Pb}} \cdot \sigma_{t\bar{t}}^{\text{th}} = 58.1 \pm 2.0 \text{ (stat.)}_{-4.4}^{+4.8} \text{ (syst.) nb}$$

• Nuclear modification factor, measured for the first time at the LHC, is defined as:

$$R_{pPb} = \sigma_{pPb} / (A_{Pb} \cdot \sigma_{pp}) = = 1.090 \pm 0.039 \text{ (stat.)}^{+0.094}$$
  
(syst.). -0.091

- The **most precise** tt cross-section measurement in HI collisions.
- The result is consistent with the cross section in pp collisions, extrapolated to
- $\sqrt{s} = 8.16$  TeV and scaled by  $A_{pb} = 208$ .





## Measurement of the centrality dependence of the dijet yield in *p*+Pb collisions at $\sqrt{s_{_{\rm NN}}}$ = 8.16 TeV with the ATLAS detector



- Measure the centrality dependence of dijet yield in proton-lead (*p*+Pb) collisions at 8.16 TeV using the ATLAS detector.
- Insight into modifications of parton distribution functions (PDFs) within a nuclear environment.
- Investigate potential scaling trends with the Bjorken x variable for partons from the proton in *p*+Pb collisions.
- Contribute to the understanding of small proton configurations and the role of nuclear modification in high-energy collisions.



#### Jets in 2016 *p*+Pb collisions:

- Dijets in *p*+Pb at 8.16 TeV allow to provide input to color transparency effects in p+A collisions.
- Anti-k<sub>t</sub> R = 0.4 calorimeter jets are considered in the analysis.
- Centrality is determined using  $\Sigma E_{T}$  in the Pb-going side of the Forward Calorimeter (FCal).

JER

JES

0.2

0.1

1.05

0.95

- Two centrality classes are studied: 0-10% - central events, 60-90% peripheral events.
- Jet energy scale (JES) and jet energy resolution (JER) show no significant centrality dependence.



#### **Kinematic domain:**



**X**<sub>p</sub> is the fraction of the proton's longitudinal momentum carried by the parton (quark or gluon) from the proton in the collision.

X<sub>Pb</sub> is the fraction of the lead nucleus' longitudinal momentum carried by the parton from the lead nucleus in the collision.

- It shows momentum fractions for p and Pb in different slices of  $y^*$  and  $y_b$ . It can cover a large  $x_p^- x_{Pb}^-$  space.
- The measurement probes the structure of both p and Pb over four orders of magnitude.
- The analysis is not directly carried out in parton system kinematic variables.

#### **Kinematic domain:**



• Work is underway to extract the cross section over the full phase space covered by this analysis.

Parton momentum fraction dependence:



### x<sub>p</sub> dependence

- A log-linear scaling as a function of  $\langle x_n \rangle$ .
- The strongest suppression in the proton valence region.
- The log-linear scaling disappears in the low-x<sub>n</sub> region.

### x<sub>pb</sub> dependence

- Log-linear decreasing trend with increasing suppression for each slice in y<sub>b</sub>.
- Strongest suppression observed at low-x<sub>pb</sub>.
- Values of R<sub>CP</sub> > 1 in the 10<sup>-2</sup> < x<sub>Pb</sub> < 2 · 10<sup>-1</sup> range (~anti-shadowing).

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Parton momentum fraction dependence:



- Left shows there is a strong suppression of jets produced in central vs peripheral collisions as a function of x<sub>p</sub>, right shows a strong suppression for low x<sub>pb</sub>
- There is no common trend but rather separate trends for different slices in  $y_{\rm h}$ .

## Photo-nuclear jet production in ultra-peripheral Pb+Pb collisions at $\sqrt{s}_{_{\rm NN}}$ =5.02 TeV with the ATLAS detector





- Ultra-Peripheral Collisions (UPCs) occur when the impact parameter is greater than the sum of the nuclear radii, leading to interactions via long-range electromagnetic forces.
- UPCs provide a unique way to probe the structure of nPDF through photon-hadron interactions.
- Goal to measure cross section of photo nuclear diet production
- The dijet pair allows us to measure differentially in the reconstructed kinematics of the hard scattering.
- The measurement is performed using a data set recorded in 2018 with an integrated luminosity of 1.72 nb<sup>-1</sup>



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#### **Selecting Photo-nuclear Jet Events:**

- Photo-nuclear jet events with cuts motivated by the particular event topology are selected
- **OnXn** requirement for nuclear breakup in exactly one ATLAS Zero-Degree Calorimeter (ZDC)

#### OnXn

- Indicates that there are no neutrons detected in one direction (denoted by "0n") and one or more neutrons detected in the opposite direction (denoted by "Xn").
- This condition helps identify photo-nuclear events, where a photon interacts with one nucleus, without causing the complete breakup of both nuclei.



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- The shape of the H<sub>T</sub> (the scalar sum of the transverse momenta of jets in an event) distribution agrees fairly well between data and MC.
- This distribution is also sensitive to differences in jet response between data and MC.
- The y<sub>iets</sub> distribution reflects the balance between energy on the photon and nucleus-going sides.



#### **Systematic Uncertainties:**

- The pseudorapidity gap selection (green) and sensitivity to the prior (cyan) uncertainties are subdominant
- The JES (magenta) and JER (blue) uncertainties are substantial but not dominant.
- Jet calibration sequence derived for high- $\mu$  data in a low- $\mu$  environment (orange) are dominant in most bins.



#### **UPC Dijet Cross-Sections:**

- At intermediate photon energies, we can access higher-x partons.
- Systematic uncertainties grow near the acceptance edge at high-x.

- Higher photon energy opens up the low-x shadowing region.
- Results are quite consistent with the theoretical model.

- The highest photon energy allows the most access to low *x*.
- Systematic control is more challenging near acceptance edges.



### Observation of tt production in lepton+jets and dilepton channels in p+Pb

- The first tt observation using the dilepton channel in p+Pb collisions at the LHC.
- The tt cross section is measured to be  $\sigma_{tf}$  = 57.9 ± 2.0 (stat.) <sup>+4.9</sup><sub>-4.5</sub> (syst.) nb.
- The most precise tt cross-section measurement in *p*+Pb collisions at the LHC.

## Measurement of the centrality dependence of the dijet yield in *p*+Pb

- New input to understand the suppression of the jet production in p+A collisions.
- The results are compatible with an interpretation in terms of color fluctuation effects.
- Next: extraction of the dijet cross-section to constrain nPDFs

### Photo-nuclear jet production in ultra-peripheral Pb+Pb

- Photo-nuclear jet production was measured by ATLAS in 5.02 TeV Pb+Pb collisions with 2018 data.
- Particle-Flow jets allow for the measurement to be extended even lower in jet  $p_T$  while maintaining good control over uncertainties.
- This data can add a wide range of kinematic coverage to existing nPDF constraints.

## Backup

$$p_{T,Avg} = \frac{p_{T,1} + p_{T,2}}{2}, \quad y_b = \frac{y_1^{c.m.} + y_2^{c.m.}}{2}, \text{ and}$$
  
 $y^* = \frac{|y_1^{c.m.} - y_2^{c.m.}|}{2},$ 

$$R_{\rm CP}(p_{T,{\rm Avg}}, y_b, y^*) = \frac{\frac{1}{\langle T_{AB}^{0\%-10\%} \rangle} \frac{1}{N_{\rm evt}^{0\%-10\%}} \frac{d^3 N_{\rm dijet}^{0\%-10\%}}{dp_{T,{\rm Avg}} dy_b dy^*}}{\frac{1}{\langle T_{AB}^{60\%-90\%} \rangle} \frac{1}{N_{\rm evt}^{60\%-90\%}} \frac{d^3 N_{\rm dijet}^{60\%-90\%}}{dp_{T,{\rm Avg}} dy_b dy^*}},$$

$$x_p = \frac{p_{T,1}e^{y_1^{\text{c.m.}}} + p_{T,2}e^{y_2^{\text{c.m.}}}}{\sqrt{s_{\text{NN}}}} \simeq \frac{2p_{T,\text{Avg}}}{\sqrt{s_{\text{NN}}}}e^{y_b}\cosh(y^*)$$

and

$$x_{\rm Pb} = \frac{p_{T,1}e^{-y_1^{\rm c.m.}} + p_{T,2}e^{-y_2^{\rm c.m.}}}{\sqrt{s_{\rm NN}}} \simeq \frac{2p_{T,{\rm Avg}}}{\sqrt{s_{\rm NN}}}e^{-y_b}\cosh(y^*),$$

- Large rapidity gaps on one side of the detector
- To veto  $\gamma \gamma \rightarrow \mathbf{q} \overline{\mathbf{q}}$ , also require  $\Delta \eta_A^{edge} < 3$ .
- At least two Particle-Flow jets with pT > 15 GeV



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Photon Energy  $0.004 < z_{v} < 0.008$ Photon Energy  $0.008 < z_{v} < 0.015$ Photon Energy  $0.015 < z_{v} < 0.027$ [hp/GeV] [/1016 [Jub [] [http://milling//mill [http://dull.com/ 101/ 101/ 101/ 35 < H<sub>2</sub> < 43 GeV 10<sup>15</sup> ATLAS Preliminary 35 < H<sub>T</sub> < 43 GeV ATLAS Preliminary 35 < H<sub>τ</sub> < 43 GeV</li> ATLAS Preliminary 10<sup>15</sup> 43 < H<sub>T</sub> < 53 GeV (×10<sup>-2</sup>) Pb+Pb 5.02 TeV. 1.72 nb<sup>-1</sup> 43 < H<sub>T</sub> < 53 GeV (×10<sup>-2</sup>) 43 < H<sub>T</sub> < 53 GeV (×10<sup>-2</sup>) Pb+Pb 5.02 TeV, 1.72 nb<sup>-1</sup> Pb+Pb 5.02 TeV. 1.72 nb 10<sup>1</sup> 53 < H\_ < 66 GeV (×10<sup>-4</sup>)  $0.004 < z_{\gamma} < 0.008$ 53 < H<sub>T</sub> < 66 GeV (×10<sup>-4</sup>) 53 < H<sub>T</sub> < 66 GeV (×10<sup>-4</sup>)  $0.008 < z_{\gamma} < 0.015$  $0.015 < z_{\gamma} < 0.027$ 66 < H<sub>T</sub> < 81 GeV (×10<sup>-6</sup>) UPC  $\gamma + A \rightarrow jets$ 66 < H\_ < 81 GeV (×10<sup>-6</sup>) UPC  $\gamma + A \rightarrow jets$ 66 < H<sub>T</sub> < 81 GeV (×10<sup>-6</sup>) UPC  $\gamma + A \rightarrow jets$ 81 < H<sub>2</sub> < 100 GeV (×10<sup>-8</sup>) anti-k. R=0.4 Jets 81 < H<sub>T</sub> < 100 GeV (×10<sup>-8</sup> anti-k. R=0.4 Jets 81 < H<sub>T</sub> < 100 GeV (×10<sup>-8</sup>) anti-k. R=0.4 Jets |<sup>2</sup><sup>10°</sup> O 100 < H<sub>T</sub> < 123 GeV (×10<sup>-10</sup>) 35 < M<sub>inte</sub> < 185 GeV Ġ, O 100 < H<sub>\*</sub> < 123 GeV (×10<sup>-10</sup>) 35 < M<sub>ints</sub> < 185 GeV R O 100 < H<sub>2</sub> < 123 GeV (×10<sup>-10</sup>) 35 < M<sub>iets</sub> < 185 GeV 123 < H<sub>7</sub> < 152 GeV (×10<sup>-12</sup>) 123 < H<sub>T</sub> < 152 GeV (×10<sup>-12</sup>) 123 < H<sub>1</sub> < 152 GeV (×10<sup>-12</sup>) ຼຍ ₹10<sup>6</sup> ູຢ\_<10<sup>6</sup> ຼ<sub>ອ</sub>ບ\_×10 ····\*··· د 40<sup>1</sup><sup>⊥</sup>10<sup>3⁺</sup> <sup>103\*</sup> الله 10<sup>3</sup> Hp 10<sup>3</sup> 10 10 10 10 10  $10^{-9}$ 10 10<sup>-12</sup>È 10<sup>-12</sup> ⊨ = = Pythia 8 γN → jets, 10<sup>-12</sup> = Pythia 8 γN → jets, : Pythia 8 γN → jets, ---nCTEQ PDFs with Pb photon flux nCTEQ PDFs with Pb photon flux nCTEQ PDFs with Pb photon flux 10<sup>-15</sup> 10-15 10<sup>-1</sup>  $10^{-2}$  $10^{-2}$  $10^{-2}$  $10^{-1}$ 10<sup>-1</sup> X, [70<sup>18</sup> € ∑ 10<sup>18</sup> 0 10<sup>18</sup> [np/ge/ 10<sup>18</sup> 10<sup>17</sup> 10<sup>17</sup>  $35 < H_T < 43 \text{ GeV}$ ATLAS Preliminary  $35 < H_T < 43 \text{ GeV}$ ATLAS Preliminary 35 < H\_ < 43 GeV ATLAS Preliminary 43 < H<sub>T</sub> < 53 GeV (×10<sup>-2</sup>) Pb+Pb 5.02 TeV, 1.72 nb<sup>-1</sup> 43 < H<sub>T</sub> < 53 GeV (×10<sup>-2</sup>) 43 < H<sub>T</sub> < 53 GeV (×10<sup>-2</sup>) Pb+Pb 5.02 TeV. 1.72 nb Pb+Pb 5.02 TeV, 1.72 nb1 /qn] 53 < H<sub>+</sub> < 66 GeV (×10<sup>-4</sup>)  $0.015 < x_{\rm A} < 0.200$ ▲ 53 < H<sub>T</sub> < 66 GeV (×10<sup>-4</sup>)  $0.015 < x_A < 0.200$ 53 < H<sub>T</sub> < 66 GeV (×10<sup>-4</sup>)  $0.015 < x_a < 0.200$ 66 < H\_ < 81 GeV (×10<sup>-6</sup>) UPC  $\gamma + A \rightarrow jets$ 66 < H<sub>T</sub> < 81 GeV (×10<sup>-6</sup>) UPC  $\gamma + A \rightarrow jets$ 66 < H<sub>y</sub> < 81 GeV (×10<sup>-6</sup> UPC  $\gamma + A \rightarrow jets$ 81 < H<sub>y</sub> < 100 GeV (×10<sup>-8</sup>) anti-k, R=0.4 Jets 81 < H<sub>2</sub> < 100 GeV (×10<sup>-8</sup>) anti-k, R=0.4 Jets 81 < H<sub>T</sub> < 100 GeV (×10<sup>-8</sup>) anti-k. R=0.4 Jets 35 < M<sub>jets</sub> < 185 GeV ď2 100 < H<sub>T</sub> < 123 GeV (×10<sup>-10</sup>) d2, 100 < H<sub>T</sub> < 123 GeV (×10<sup>-10</sup>) 35 < M<sub>iets</sub> < 185 GeV P 100 < H<sub>T</sub> < 123 GeV (×10<sup>-10</sup> 35 < M .... < 185 GeV 123 < H<sub>T</sub> < 152 GeV (×10<sup>-12</sup>) 123 < H<sub>v</sub> < 152 GeV (×10<sup>-12</sup> 123 < H<sub>T</sub> < 152 GeV (×10<sup>-12</sup> <sup>6</sup> م ్ల్<10<sup>6</sup> ຼຍ <u>₹</u>10<sup>6</sup> ⊥10 Нр 10<sup>-</sup> 10<sup>-</sup> 10<sup>-</sup> 10<sup>-</sup> 10<sup>-</sup> 년 10 북 10 10 10 10-6 10-6  $10^{-6}$ 10-5  $10^{-9}$ 10 10-12 10<sup>-12</sup>⊨ 10<sup>-12</sup> ⊨ : Pythia 8 γN → jets, = Pythia 8  $\gamma N \rightarrow jets$ . = := Pythia 8 γN → jets, ---nCTEQ PDFs with Pb photon flux nCTEQ PDFs with Pb photon flux nCTEQ PDFs with Pb photon flux  $10^{-3}$  $10^{-2}$  $10^{-3}$  $10^{-2}$  $10^{-3}$  $10^{-2}$ Zγ Zγ

Z<sub>v</sub> - Fraction of the photon's energy that is carried by the parton that interacts in the hard scattering.

 $Z_{\gamma}$ 

30 X - Fraction of the nuclear momentum carried by the parton from the nucleus participating in the scattering process. Santu Mondal (IEAP, CTU in Prague)