

Status of Flavor Physics and Anomalies

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Flavor Physics



Promising Indirect Probes of New Physics

Probe more generic new physics




- ▶ Test bedrock assumptions of particle physics

Lorentz invariance; CPT invariance; ...

($\Lambda \gtrsim M_{\text{Planck}} \sim 10^{19} \text{ GeV}$)

Reach to higher new physics scales



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► **Test bedrock assumptions of particle physics**

Lorentz invariance; CPT invariance; ...

($\Lambda \gtrsim M_{\text{Planck}} \sim 10^{19} \text{ GeV}$)

► **Test (approximate) accidental symmetries of the SM**

Baryon Number: e.g. proton decay

($\Lambda \sim \Lambda_{\text{GUT}} \sim 10^{16} \text{ GeV}$)

Lepton Number: e.g. neutrinoless double beta decay

($\Lambda \sim \Lambda_{\text{see-saw}} \sim 10^{12} \text{ GeV}$)

Flavor: e.g. flavor changing neutral currents

($\Lambda \sim 10^3 - 10^8 \text{ GeV}$)

CP: e.g. electric dipole moments

($\Lambda \sim 10^3 - 10^8 \text{ GeV}$)

Reach to higher new physics scales

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CP: e.g. electric dipole moments
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- ▶ **Test “ordinary” Standard Model processes**

Higgs precision program; Electroweak precision observables; muon anomalous magnetic moment; ...
($\Lambda \sim 10^3 \text{ GeV}$)

Reach to higher new physics scales

The Standard Model and Beyond

The diagram illustrates the Standard Model Lagrangian, \mathcal{L}_{SM} , with several callouts pointing to specific terms or features:

- CC problem**: Callout pointing to the Λ^4 term.
- Hierarchy problem**: Callout pointing to the $\Lambda^2 H^2$ term.
- Vacuum stability?**: Callout pointing to the λH^4 term.
- Strong CP problem**: Callout pointing to the $F_{\mu\nu} \tilde{F}^{\mu\nu}$ term.
- SM flavor puzzle**: Callout pointing to the $+ Y H \bar{\Psi} \Psi$ term.

$$\mathcal{L}_{\text{SM}} \sim \Lambda^4 + \Lambda^2 H^2 + \lambda H^4 + \bar{\Psi} \not{D} \Psi + (D_\mu H)^2 + (F_{\mu\nu})^2 + F_{\mu\nu} \tilde{F}^{\mu\nu} + Y H \bar{\Psi} \Psi$$

The Standard Model and Beyond

$$\mathcal{L}_{\text{SM}} \sim \Lambda^4 + \Lambda^2 H^2 + \lambda H^4$$
$$+ \bar{\Psi} \not{D} \Psi + (D_\mu H)^2 + (F_{\mu\nu})^2 + F_{\mu\nu} \tilde{F}^{\mu\nu}$$
$$+ Y H \bar{\Psi} \Psi + \frac{1}{\Lambda} (LH)^2 + \frac{1}{\Lambda^2} \sum_i \mathcal{O}_i^{\text{dim6}} + \dots$$

CC problem

Hierarchy problem

Vacuum stability?

Strong CP problem

SM flavor puzzle

Neutrino masses

Flavorful new physics?

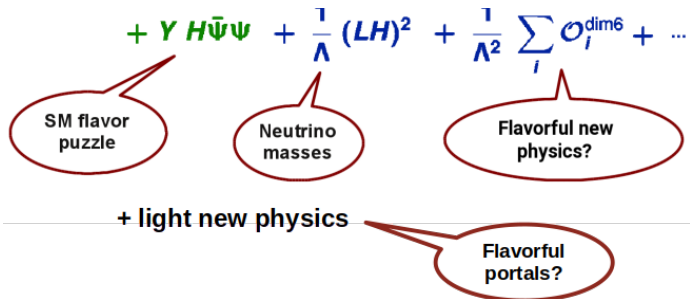
The Standard Model and Beyond

The diagram shows the Standard Model Lagrangian \mathcal{L}_{SM} with several callouts pointing to specific terms or parts of the equation:

- CC problem** points to the Λ^4 term.
- Hierarchy problem** points to the $\Lambda^2 H^2$ term.
- Vacuum stability?** points to the λH^4 term.
- Strong CP problem** points to the $F_{\mu\nu} \tilde{F}^{\mu\nu}$ term.
- SM flavor puzzle** points to the $Y H \bar{\Psi} \Psi$ term.
- Neutrino masses** points to the $\frac{1}{\Lambda} (LH)^2$ term.
- Flavorful new physics?** points to the $\frac{1}{\Lambda^2} \sum_i \mathcal{O}_i^{\text{dim6}}$ term.
- Flavorful portals?** points to the **+ light new physics** text below the Lagrangian.

$$\begin{aligned} \mathcal{L}_{\text{SM}} \sim & \Lambda^4 + \Lambda^2 H^2 + \lambda H^4 \\ & + \bar{\Psi} \not{D} \Psi + (D_\mu H)^2 + (F_{\mu\nu})^2 + F_{\mu\nu} \tilde{F}^{\mu\nu} \\ & + Y H \bar{\Psi} \Psi + \frac{1}{\Lambda} (LH)^2 + \frac{1}{\Lambda^2} \sum_i \mathcal{O}_i^{\text{dim6}} + \dots \\ & + \text{light new physics} \end{aligned}$$

Two Basic Flavor Questions



Q1: What is the origin of the hierarchical flavor structure of the SM?

Q2: Are there new sources of flavor violation beyond the SM?

The Standard Model Flavor Puzzle

Why are there **three flavors** of quarks and leptons?



What is the origin of the hierarchies in the **fermion spectrum**?

What is the origin of the hierarchies in the **quark mixing**?

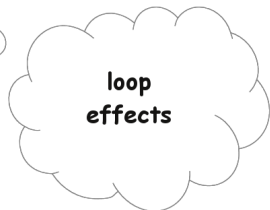
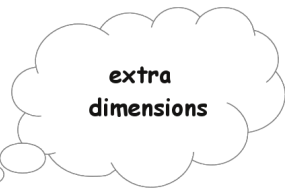
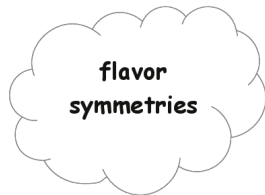
Is **lepton mixing** anarchic?

Addressing the SM Flavor Puzzle



Addressing the SM Flavor Puzzle

recent years have seen
a revival of symmetry based
flavor model building



revisiting horizontal symmetries;
flavor deconstruction;
modular flavor symmetries;
generalized symmetries; ...



... Cordova et al. 2212.13193; Davighi et al. 2201.07245, 2212.06163, 2303.1520, 2305.16280;
Fernandez Navarro et al. 2305.07690, 2311.05683; Cornella et al. 2306.08026; Asadi et al.
2308.01340; Chen et al. 2312.09255; Greljo et al. 2309.11547, 2406.01696, 2406.02687;
... many many more ...

Flavor Violation in SMEFT

1 : X^3		2 : H^6		3 : $H^4 D^2$		5 : $\psi^2 H^3 + \text{h.c.}$	
Q_G	$f^{ABC} G_{\mu\nu}^A G_{\rho\sigma}^B G_{\rho\mu}^C$	Q_H	$(H^\dagger H)^3$	$Q_{H\Box}$	$(H^\dagger H)\Box(H^\dagger H)$	Q_{eH}	$(H^\dagger H)(\bar{l}_p e_r H)$
$Q_{\bar{G}}$	$f^{ABC} \bar{G}_{\mu\nu}^A G_{\rho\sigma}^B G_{\rho\mu}^C$			Q_{HD}	$(H^\dagger D_\mu H)^\dagger (H^\dagger D_\mu H)$	Q_{uH}	$(H^\dagger H)(\bar{q}_p u_r \tilde{H})$
Q_W	$\epsilon^{IJK} W_{\mu\nu}^I W_{\rho\sigma}^J W_{\rho\mu}^K$					Q_{dH}	$(H^\dagger H)(\bar{q}_p d_r \tilde{H})$
$Q_{\bar{W}}$	$\epsilon^{IJK} \bar{W}_{\mu\nu}^I W_{\rho\sigma}^J W_{\rho\mu}^K$						
4 : $X^2 H^2$		6 : $\psi^2 XH + \text{h.c.}$		7 : $\psi^2 H^2 D$			
Q_{HG}	$H^\dagger H G_{\mu\nu}^A G^{A\mu\nu}$	Q_{eW}	$(\bar{l}_p \sigma^{\mu\nu} e_r) \tau^I H W_{\mu\nu}^I$	$Q_{H1}^{(1)}$	$(H^\dagger i \overleftrightarrow{D}_\mu H)(\bar{l}_p \gamma^\mu l_r)$		
$Q_{H\bar{G}}$	$H^\dagger H \bar{G}_{\mu\nu}^A G^{A\mu\nu}$	Q_{eB}	$(\bar{l}_p \sigma^{\mu\nu} e_r) H B_{\mu\nu}$	$Q_{H1}^{(3)}$	$(H^\dagger i \overleftrightarrow{D}_\mu^\dagger H)(\bar{l}_p \tau^I \gamma^\mu l_r)$		
Q_{HW}	$H^\dagger H W_{\mu\nu}^I W^{I\mu\nu}$	Q_{uG}	$(\bar{q}_p \sigma^{\mu\nu} T^A u_r) \tilde{H} G_{\mu\nu}^A$	Q_{He}	$(H^\dagger i \overleftrightarrow{D}_\mu H)(\bar{e}_p \gamma^\mu e_r)$		
$Q_{H\bar{W}}$	$H^\dagger H \bar{W}_{\mu\nu}^I W^{I\mu\nu}$	Q_{uW}	$(\bar{q}_p \sigma^{\mu\nu} u_r) \tau^I \tilde{H} W_{\mu\nu}^I$	$Q_{Hq}^{(1)}$	$(H^\dagger i \overleftrightarrow{D}_\mu H)(\bar{q}_p \gamma^\mu q_r)$		
Q_{HB}	$H^\dagger H B_{\mu\nu} B^{\mu\nu}$	Q_{uB}	$(\bar{q}_p \sigma^{\mu\nu} u_r) \tilde{H} B_{\mu\nu}$	$Q_{Hq}^{(3)}$	$(H^\dagger i \overleftrightarrow{D}_\mu^\dagger H)(\bar{q}_p \tau^I \gamma^\mu q_r)$		
$Q_{H\bar{B}}$	$H^\dagger H \bar{B}_{\mu\nu} B^{\mu\nu}$	Q_{dG}	$(\bar{q}_p \sigma^{\mu\nu} T^A d_r) H G_{\mu\nu}^A$	Q_{Hu}	$(H^\dagger i \overleftrightarrow{D}_\mu H)(\bar{u}_p \gamma^\mu u_r)$		
Q_{HWB}	$H^\dagger \tau^I H W_{\mu\nu}^I B^{\mu\nu}$	Q_{dW}	$(\bar{q}_p \sigma^{\mu\nu} d_r) \tau^I H W_{\mu\nu}^I$	Q_{Hd}	$(H^\dagger i \overleftrightarrow{D}_\mu H)(\bar{d}_p \gamma^\mu d_r)$		
$Q_{H\bar{W}B}$	$H^\dagger \tau^I H \bar{W}_{\mu\nu}^I B^{\mu\nu}$	Q_{dB}	$(\bar{q}_p \sigma^{\mu\nu} d_r) H B_{\mu\nu}$	$Q_{Hud} + \text{h.c.}$	$i(\tilde{H}^\dagger D_\mu H)(\bar{u}_p \gamma^\mu d_r)$		
8 : $(\bar{L}L)(\bar{L}L)$		8 : $(\bar{R}R)(\bar{R}R)$		8 : $(\bar{L}L)(\bar{R}R)$			
Q_{ll}	$(\bar{l}_p \gamma_\mu l_r)(\bar{l}_s \gamma^\mu l_t)$	Q_{ee}	$(\bar{e}_p \gamma_\mu e_r)(\bar{e}_s \gamma^\mu e_t)$	Q_{le}	$(\bar{l}_p \gamma_\mu l_r)(\bar{e}_s \gamma^\mu e_t)$		
$Q_{qq}^{(1)}$	$(\bar{q}_p \gamma_\mu q_r)(\bar{q}_s \gamma^\mu q_t)$	Q_{uu}	$(\bar{u}_p \gamma_\mu u_r)(\bar{u}_s \gamma^\mu u_t)$	Q_{lu}	$(\bar{l}_p \gamma_\mu l_r)(\bar{u}_s \gamma^\mu u_t)$		
$Q_{qq}^{(3)}$	$(\bar{q}_p \gamma_\mu \tau^I q_r)(\bar{q}_s \gamma^\mu \tau^I q_t)$	Q_{dd}	$(\bar{d}_p \gamma_\mu d_r)(\bar{d}_s \gamma^\mu d_t)$	Q_{ld}	$(\bar{l}_p \gamma_\mu l_r)(\bar{d}_s \gamma^\mu d_t)$		
$Q_{lq}^{(1)}$	$(\bar{l}_p \gamma_\mu l_r)(\bar{q}_s \gamma^\mu q_t)$	Q_{eu}	$(\bar{e}_p \gamma_\mu e_r)(\bar{u}_s \gamma^\mu u_t)$	Q_{qe}	$(\bar{q}_p \gamma_\mu q_r)(\bar{e}_s \gamma^\mu e_t)$		
$Q_{lq}^{(3)}$	$(\bar{l}_p \gamma_\mu \tau^I l_r)(\bar{q}_s \gamma^\mu \tau^I q_t)$	Q_{ed}	$(\bar{e}_p \gamma_\mu e_r)(\bar{d}_s \gamma^\mu d_t)$	$Q_{qu}^{(1)}$	$(\bar{q}_p \gamma_\mu q_r)(\bar{u}_s \gamma^\mu u_t)$		
		$Q_{ud}^{(1)}$	$(\bar{u}_p \gamma_\mu u_r)(\bar{d}_s \gamma^\mu d_t)$	$Q_{qu}^{(8)}$	$(\bar{q}_p \gamma_\mu T^A q_r)(\bar{u}_s \gamma^\mu T^A u_t)$		
		$Q_{ud}^{(8)}$	$(\bar{u}_p \gamma_\mu T^A u_r)(\bar{d}_s \gamma^\mu T^A d_t)$	$Q_{qd}^{(1)}$	$(\bar{q}_p \gamma_\mu q_r)(\bar{d}_s \gamma^\mu d_t)$		
				$Q_{qd}^{(8)}$	$(\bar{q}_p \gamma_\mu T^A q_r)(\bar{d}_s \gamma^\mu T^A d_t)$		
8 : $(\bar{L}R)(\bar{R}L) + \text{h.c.}$		8 : $(\bar{L}R)(\bar{L}R) + \text{h.c.}$					
Q_{ledq}	$(\bar{l}_p^i e_r)(\bar{d}_s q_{tj})$	$Q_{quqd}^{(1)}$	$(\bar{q}_p^i u_r) \epsilon_{jk} (\bar{q}_s^k d_t)$				
		$Q_{quqd}^{(8)}$	$(\bar{q}_p^i T^A u_r) \epsilon_{jk} (\bar{q}_s^k T^A d_t)$				
		$Q_{lequ}^{(1)}$	$(\bar{l}_p^i e_r) \epsilon_{jk} (\bar{q}_s^k u_t)$				
		$Q_{lequ}^{(3)}$	$(\bar{l}_p^i \sigma_{\mu\nu} e_r) \epsilon_{jk} (\bar{q}_s^k \sigma^{\mu\nu} u_t)$				

2499 baryon number conserving dim. 6 operators

Grzadkowski et al. 1008.4884,

Alonso et al 1312.2014

Flavor Violation in SMEFT

1 : X^3		2 : H^6		3 : $H^4 D^2$		5 : $\psi^2 H^3 + \text{h.c.}$	
Q_G	$f^{ABC} G_{\mu\nu}^A G_{\nu\rho}^B G_{\rho\mu}^C$	Q_H	$(H^\dagger H)^3$	$Q_{H\Box}$	$(H^\dagger H)\Box(H^\dagger H)$	Q_{eH}	$(H^\dagger H)(\bar{l}_p e_r H)$
$Q_{\tilde{G}}$	$f^{ABC} \tilde{G}_{\mu\nu}^A \tilde{G}_{\nu\rho}^B \tilde{G}_{\rho\mu}^C$			Q_{HD}	$(H^\dagger D_\mu H)^\dagger (H^\dagger D_\mu H)$	Q_{uH}	$(H^\dagger H)(\bar{u}_p u_r \tilde{H})$
Q_W	$e^{JKL} W_{\mu\nu}^J W_{\nu\rho}^K W_{\rho\mu}^L$					Q_{dH}	$(H^\dagger H)(\bar{d}_p d_r \tilde{H})$
$Q_{\tilde{W}}$	$e^{JKL} \tilde{W}_{\mu\nu}^J \tilde{W}_{\nu\rho}^K \tilde{W}_{\rho\mu}^L$						

4 : $X^2 H^2$		6 : $\psi^2 X H + \text{h.c.}$		7 : $\psi^2 H^2 D$	
Q_{HC}	$H^\dagger H G_{\mu\nu}^A G^{A\mu\nu}$	Q_{eW}	$(\bar{l}_p \sigma^{\mu\nu} e_r) \tau^I H W_{\mu\nu}^I$	$Q_{H1}^{(1)}$	$(H^\dagger i \overleftrightarrow{D}_\mu H)(\bar{l}_p \gamma^\mu l_r)$
Q_{HC}	$H^\dagger H \tilde{G}_{\mu\nu}^A G^{A\mu\nu}$	Q_{uH}	$(\bar{l}_p \sigma^{\mu\nu} e_r) H B_{\mu\nu}$	$Q_{H1}^{(2)}$	$(H^\dagger i \overleftrightarrow{D}_\mu^I H)(\bar{l}_p \tau^I \gamma^\mu l_r)$
Q_{HW}	$H^\dagger H W_{\mu\nu}^I W^{I\mu\nu}$	Q_{uG}	$(\bar{q}_p \sigma^{\mu\nu} t^A u_r) \tilde{H} G_{\mu\nu}^A$	Q_{H2}	$(H^\dagger i \overleftrightarrow{D}_\mu H)(\bar{q}_p \gamma^\mu q_r)$
$Q_{H\tilde{W}}$	$H^\dagger H \tilde{W}_{\mu\nu}^I W^{I\mu\nu}$	Q_{uW}	$(\bar{q}_p \sigma^{\mu\nu} u_r) \tau^I \tilde{H} W_{\mu\nu}^I$	$Q_{H3}^{(1)}$	$(H^\dagger i \overleftrightarrow{D}_\mu H)(\bar{q}_p \gamma^\mu q_r)$
Q_{HB}	$H^\dagger H B_{\mu\nu} B^{\mu\nu}$	Q_{uB}	$(\bar{q}_p \sigma^{\mu\nu} u_r) \tilde{H} B_{\mu\nu}$	$Q_{H3}^{(2)}$	$(H^\dagger i \overleftrightarrow{D}_\mu^I H)(\bar{q}_p \tau^I \gamma^\mu q_r)$
$Q_{H\tilde{B}}$	$H^\dagger H \tilde{B}_{\mu\nu} B^{\mu\nu}$	Q_{dG}	$(\bar{q}_p \sigma^{\mu\nu} T^A d_r) H G_{\mu\nu}^A$	Q_{H4}	$(H^\dagger i \overleftrightarrow{D}_\mu H)(\bar{q}_p \gamma^\mu u_r)$
Q_{HWB}	$H^\dagger \tau^I H W_{\mu\nu}^I B^{\mu\nu}$	Q_{dW}	$(\bar{q}_p \sigma^{\mu\nu} d_r) \tau^I H W_{\mu\nu}^I$	Q_{H5}	$(H^\dagger i \overleftrightarrow{D}_\mu H)(\bar{d}_p \gamma^\mu d_r)$
$Q_{H\tilde{W}d}$	$H^\dagger \tau^I H \tilde{W}_{\mu\nu}^I B^{\mu\nu}$	Q_{dH}	$(\bar{q}_p \sigma^{\mu\nu} d_r) \tau^I H B_{\mu\nu}$	$Q_{H6d} + \text{h.c.}$	$i(\tilde{H}^\dagger D_\mu H)(\bar{q}_p \gamma^\mu d_r)$

2499 baryon number conserving dim. 6 operators

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Alonso et al 1312.2014

4 fermion interactions

8 : $(\bar{L}L)(\bar{L}L)$		8 : $(\bar{R}R)(\bar{R}R)$		8 : $(\bar{L}L)(\bar{R}R)$	
Q_{ll}	$(\bar{l}_p \gamma_\mu l_r)(\bar{l}_s \gamma^\mu l_t)$	Q_{ee}	$(\bar{e}_p \gamma_\mu e_r)(\bar{e}_s \gamma^\mu e_t)$	Q_{le}	$(\bar{l}_p \gamma_\mu l_r)(\bar{e}_s \gamma^\mu e_t)$
$Q_{ll}^{(1)}$	$(\bar{q}_p \gamma_\mu q_r)(\bar{q}_s \gamma^\mu q_t)$	Q_{uu}	$(\bar{u}_p \gamma_\mu u_r)(\bar{u}_s \gamma^\mu u_t)$	Q_{lu}	$(\bar{l}_p \gamma_\mu l_r)(\bar{u}_s \gamma^\mu u_t)$
$Q_{ll}^{(2)}$	$(\bar{q}_p \gamma_\mu \tau^I q_r)(\bar{q}_s \gamma^\mu \tau^I q_t)$	Q_{dd}	$(\bar{d}_p \gamma_\mu d_r)(\bar{d}_s \gamma^\mu d_t)$	Q_{ld}	$(\bar{l}_p \gamma_\mu l_r)(\bar{d}_s \gamma^\mu d_t)$
$Q_{ll}^{(3)}$	$(\bar{l}_p \gamma_\mu \tau^I q_r)(\bar{q}_s \gamma^\mu q_t)$	Q_{eu}	$(\bar{e}_p \gamma_\mu e_r)(\bar{u}_s \gamma^\mu u_t)$	Q_{qe}	$(\bar{q}_p \gamma_\mu q_r)(\bar{e}_s \gamma^\mu e_t)$
$Q_{ll}^{(4)}$	$(\bar{l}_p \gamma_\mu \tau^I l_r)(\bar{q}_s \gamma^\mu \tau^I q_t)$	Q_{ed}	$(\bar{e}_p \gamma_\mu e_r)(\bar{d}_s \gamma^\mu d_t)$	$Q_{qu}^{(1)}$	$(\bar{q}_p \gamma_\mu q_r)(\bar{u}_s \gamma^\mu u_t)$
		$Q_{ud}^{(1)}$	$(\bar{u}_p \gamma_\mu u_r)(\bar{d}_s \gamma^\mu d_t)$	$Q_{qu}^{(8)}$	$(\bar{q}_p \gamma_\mu T^A q_r)(\bar{u}_s \gamma^\mu T^A u_t)$
		$Q_{ud}^{(2)}$	$(\bar{u}_p \gamma_\mu T^A u_r)(\bar{d}_s \gamma^\mu T^A d_t)$	$Q_{qd}^{(1)}$	$(\bar{q}_p \gamma_\mu q_r)(\bar{d}_s \gamma^\mu d_t)$
				$Q_{qd}^{(8)}$	$(\bar{q}_p \gamma_\mu T^A q_r)(\bar{d}_s \gamma^\mu T^A d_t)$

8 : $(\bar{L}R)(\bar{R}L) + \text{h.c.}$		8 : $(\bar{L}R)(\bar{L}R) + \text{h.c.}$	
Q_{ledq}	$(\bar{l}_p^i e_r)(\bar{d}_s q_t^i)$	$Q_{quqd}^{(1)}$	$(\bar{q}_p^i u_r) e_{jk} (\bar{q}_s^j d_t^k)$
		$Q_{quqd}^{(8)}$	$(\bar{q}_p^i T^A u_r) e_{jk} (\bar{q}_s^j T^A d_t^k)$
		$Q_{lequ}^{(1)}$	$(\bar{l}_p^i e_r) e_{jk} (\bar{q}_s^j u_t^k)$
		$Q_{lequ}^{(2)}$	$(\bar{l}_p^i \sigma_{\mu\nu} e_r) e_{jk} (\bar{q}_s^j \sigma^{\mu\nu} u_t^k)$

Flavor Violation in SMEFT

1 : X^3	2 : H^6	3 : $H^4 D^2$	5 : $\psi^2 H^3 + \text{h.c.}$
Q_G	$f^{ABC} G_{\mu\nu}^A G_{\nu\rho}^B G_{\rho\mu}^C$	$Q_{H\Box}$	Q_{eH}
$Q_{\tilde{G}}$	$f^{ABC} \tilde{G}_{\mu\nu}^A \tilde{G}_{\nu\rho}^B \tilde{G}_{\rho\mu}^C$	Q_{HD}	Q_{uH}
Q_W	$e^{JKL} W_{\mu\nu}^J W_{\nu\rho}^K W_{\rho\mu}^L$		Q_{dH}
$Q_{\tilde{W}}$	$e^{JKL} \tilde{W}_{\mu\nu}^J \tilde{W}_{\nu\rho}^K \tilde{W}_{\rho\mu}^L$		

4 : $X^2 H^2$	6 : $\psi^2 X H + \text{h.c.}$	7 : $\psi^2 H^2 D$
Q_{HC}	Q_{eW}	$Q_{H1}^{(1)}$
Q_{HC}	Q_{uH}	$Q_{H1}^{(2)}$
Q_{HW}	Q_{uG}	Q_{H2}
$Q_{H\tilde{W}}$	Q_{uW}	$Q_{Hq}^{(1)}$
Q_{HR}	Q_{uR}	$Q_{Hq}^{(2)}$
$Q_{H\tilde{B}}$	Q_{dG}	Q_{Hu}
Q_{HWB}	Q_{dW}	Q_{Hd}
$Q_{H\tilde{W}d}$	Q_{dD}	$Q_{Hud} + \text{h.c.}$

2499 baryon number conserving dim. 6 operators

Grzadkowski et al. 1008.4884,

Alonso et al 1312.2014

4 fermion interactions

dipole transitions

8 : $(\bar{L}L)(\bar{L}L)$	8 : $(\bar{R}R)(\bar{R}R)$	8 : $(\bar{L}L)(\bar{R}R)$
Q_{ll}	Q_{ee}	Q_{le}
$Q_{ll}^{(1)}$	Q_{uu}	Q_{lu}
$Q_{ll}^{(2)}$	Q_{dd}	Q_{ld}
$Q_{ll}^{(3)}$	Q_{eu}	Q_{qe}
$Q_{ll}^{(4)}$	Q_{ed}	$Q_{qe}^{(1)}$
	$Q_{ud}^{(1)}$	$Q_{qu}^{(8)}$
	$Q_{ud}^{(2)}$	$Q_{qd}^{(1)}$
		$Q_{qd}^{(8)}$

8 : $(\bar{L}R)(\bar{R}L) + \text{h.c.}$

8 : $(\bar{L}R)(\bar{L}R) + \text{h.c.}$

Q_{leq_1}	$Q_{qq_1}^{(1)}$	$Q_{qq_1}^{(8)}$
Q_{leq_2}	$Q_{qq_2}^{(1)}$	$Q_{qq_2}^{(8)}$
	$Q_{leq_3}^{(1)}$	$Q_{leq_3}^{(8)}$
	$Q_{leq_3}^{(2)}$	

Flavor Violation in SMEFT

1 : X^3	2 : H^6	3 : $H^4 D^2$	5 : $\psi^2 H^3 + \text{h.c.}$
Q_G	$f^{ABC} G_{\mu\nu}^A G_{\nu\rho}^B G_{\rho\mu}^C$	$Q_{H\Box}$	Q_{eH}
$Q_{\tilde{G}}$	$f^{ABC} \tilde{G}_{\mu\nu}^A \tilde{G}_{\nu\rho}^B \tilde{G}_{\rho\mu}^C$	Q_{HD}	Q_{uH}
Q_W	$e^{JKL} W_{\mu\nu}^J W_{\nu\rho}^K W_{\rho\mu}^L$		Q_{dH}
$Q_{\tilde{W}}$	$e^{JKL} \tilde{W}_{\mu\nu}^J \tilde{W}_{\nu\rho}^K \tilde{W}_{\rho\mu}^L$		

4 : $X^2 H^2$	6 : $\psi^2 X H + \text{h.c.}$	7 : $\psi^2 H^2 D$
Q_{HC}	Q_{eW}	$Q_{H1}^{(1)}$
Q_{HC}	Q_{uH}	$Q_{H1}^{(2)}$
Q_{HW}	Q_{uG}	Q_{H2}
$Q_{H\tilde{W}}$	Q_{uW}	$Q_{Hq}^{(1)}$
Q_{HR}	Q_{uR}	$Q_{Hq}^{(2)}$
$Q_{H\tilde{B}}$	Q_{dG}	Q_{Hu}
Q_{HWB}	Q_{dW}	Q_{Hd}
$Q_{H\tilde{W}d}$	Q_{dD}	$Q_{Hed} + \text{h.c.}$

8 : $(\bar{L}L)(\bar{L}L)$	8 : $(\bar{R}R)(\bar{R}R)$	8 : $(\bar{L}L)(\bar{R}R)$
Q_{ll}	Q_{ee}	Q_{le}
$Q_{ll}^{(1)}$	Q_{uu}	Q_{lu}
$Q_{ll}^{(2)}$	Q_{dd}	Q_{ld}
$Q_{ll}^{(3)}$	Q_{eu}	Q_{ue}
$Q_{ll}^{(4)}$	Q_{ed}	$Q_{ue}^{(1)}$
	$Q_{ud}^{(1)}$	$Q_{ue}^{(2)}$
	$Q_{ud}^{(2)}$	$Q_{ue}^{(3)}$
		$Q_{ue}^{(4)}$

8 : $(\bar{L}R)(\bar{R}L) + \text{h.c.}$	8 : $(\bar{L}R)(\bar{L}R) + \text{h.c.}$
Q_{leq1}	$Q_{qq1}^{(1)}$
	$Q_{qq1}^{(2)}$
	$Q_{leq2}^{(1)}$
	$Q_{leq2}^{(2)}$

2499 baryon number conserving dim. 6 operators

Grzadkowski et al. 1008.4884,

Alonso et al 1312.2014

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Q_G	$f^{ABC} G_{\mu\nu}^A G_{\nu\rho}^B G_{\rho\mu}^C$	Q_H	$(H^\dagger H)^3$	$Q_{H\Box}$	$(H^\dagger H)\Box(H^\dagger H)$	Q_{eH}	$(H^\dagger H)(\bar{l}_p e_r H)$
$Q_{\tilde{G}}$	$f^{ABC} \tilde{G}_{\mu\nu}^A G_{\nu\rho}^B G_{\rho\mu}^C$			Q_{HD}	$(H^\dagger D_\mu H)^\dagger (H^\dagger D_\mu H)$	Q_{uH}	$(H^\dagger H)(\bar{u}_p u_r \tilde{H})$
Q_W	$e^{JKL} W_{\mu\nu}^J W_{\nu\rho}^K W_{\rho\mu}^L$					Q_{dH}	$(H^\dagger H)(\bar{d}_p d_r H)$
$Q_{\tilde{W}}$	$e^{JKL} \tilde{W}_{\mu\nu}^J W_{\nu\rho}^K W_{\rho\mu}^L$						

4 : $X^2 H^2$		6 : $\psi^2 X H + \text{h.c.}$		7 : $\psi^2 H^2 D$	
Q_{HC}	$H^\dagger H G_{\mu\nu}^A G^{A\mu\nu}$	Q_{eW}	$(\bar{l}_p \sigma^{\mu\nu} e_r) \tau^I H W_{\mu\nu}^I$	$Q_{H1}^{(1)}$	$(H^\dagger i \overleftrightarrow{D}_\mu H)(\bar{l}_p \gamma^\mu l_r)$
Q_{HC}	$H^\dagger H \tilde{G}_{\mu\nu}^A G^{A\mu\nu}$	Q_{uH}	$(\bar{l}_p \sigma^{\mu\nu} e_r) H B_{\mu\nu}$	$Q_{H1}^{(2)}$	$(H^\dagger i \overleftrightarrow{D}_\mu^I H)(\bar{l}_p \tau^I \gamma^\mu l_r)$
Q_{HW}	$H^\dagger H W_{\mu\nu}^I W^{I\mu\nu}$	Q_{uG}	$(\bar{q}_p \sigma^{\mu\nu} t^A u_r) \tilde{H} G_{\mu\nu}^A$	Q_{H2}	$(H^\dagger i \overleftrightarrow{D}_\mu H)(\bar{q}_p \gamma^\mu e_r)$
$Q_{H\tilde{W}}$	$H^\dagger H \tilde{W}_{\mu\nu}^I W^{I\mu\nu}$	Q_{uW}	$(\bar{q}_p \sigma^{\mu\nu} u_r) \tau^I \tilde{H} W_{\mu\nu}^I$	$Q_{H3}^{(1)}$	$(H^\dagger i \overleftrightarrow{D}_\mu H)(\bar{q}_p \gamma^\mu q_r)$
Q_{HB}	$H^\dagger H B_{\mu\nu} B^{\mu\nu}$	Q_{uB}	$(\bar{q}_p \sigma^{\mu\nu} u_r) \tilde{H} B_{\mu\nu}$	$Q_{H3}^{(2)}$	$(H^\dagger i \overleftrightarrow{D}_\mu^I H)(\bar{q}_p \tau^I \gamma^\mu q_r)$
$Q_{H\tilde{B}}$	$H^\dagger H \tilde{B}_{\mu\nu} B^{\mu\nu}$	Q_{dG}	$(\bar{q}_p \sigma^{\mu\nu} T^A d_r) H G_{\mu\nu}^A$	Q_{H4}	$(H^\dagger i \overleftrightarrow{D}_\mu H)(\bar{q}_p \gamma^\mu u_r)$
Q_{HWB}	$H^\dagger \tau^I H W_{\mu\nu}^I B^{\mu\nu}$	Q_{dW}	$(\bar{q}_p \sigma^{\mu\nu} d_r) \tau^I H W_{\mu\nu}^I$	Q_{H5}	$(H^\dagger i \overleftrightarrow{D}_\mu H)(\bar{d}_p \gamma^\mu d_r)$
$Q_{H\tilde{W}B}$	$H^\dagger \tau^I H \tilde{W}_{\mu\nu}^I B^{\mu\nu}$	Q_{dB}	$(\bar{q}_p \sigma^{\mu\nu} d_r) H B_{\mu\nu}$	Q_{H6+d}	$i(\tilde{H}^\dagger D_\mu H)(\bar{q}_p \gamma^\mu d_r)$

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$Q_{ll}^{(1)}$	$(\bar{q}_p \gamma_\mu q_r)(\bar{q}_s \gamma^\mu q_t)$	Q_{uu}	$(\bar{u}_p \gamma_\mu u_r)(\bar{u}_s \gamma^\mu u_t)$	Q_{lu}	$(\bar{l}_p \gamma_\mu l_r)(\bar{u}_s \gamma^\mu u_t)$
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$Q_{ll}^{(4)}$	$(\bar{l}_p \gamma_\mu \tau^I l_r)(\bar{q}_s \gamma^\mu \tau^I q_t)$	Q_{ed}	$(\bar{e}_p \gamma_\mu e_r)(\bar{d}_s \gamma^\mu d_t)$	$Q_{qu}^{(1)}$	$(\bar{q}_p \gamma_\mu q_r)(\bar{u}_s \gamma^\mu u_t)$
		$Q_{ud}^{(1)}$	$(\bar{u}_p \gamma_\mu u_r)(\bar{d}_s \gamma^\mu d_t)$	$Q_{qu}^{(8)}$	$(\bar{q}_p \gamma_\mu T^A q_r)(\bar{u}_s \gamma^\mu T^A u_t)$
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8 : $(\bar{L}R)(\bar{R}L) + \text{h.c.}$		8 : $(\bar{L}R)(\bar{L}R) + \text{h.c.}$	
Q_{leq_1}	$(\bar{l}_p^c e_r)(\bar{d}_s q_t)$	$Q_{qq_1}^{(1)}$	$(\bar{q}_p^c)_{i,j,k} (q_s^k d_t)$
		$Q_{qq_1}^{(8)}$	$(\bar{q}_p^c T^A)_{i,j,k} (q_s^k T^A d_t)$
		$Q_{leq_2}^{(1)}$	$(\bar{l}_p^c e_r)_{i,j} (\bar{q}_s^i u_t)$
		$Q_{leq_2}^{(3)}$	$(\bar{l}_p^c \sigma_{\mu\nu} e_r)_{i,j,k} (\bar{q}_s^i \sigma^{\mu\nu} u_t)$

2499 baryon number conserving dim. 6 operators

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Q_{HG}	$H^\dagger H \tilde{G}_{\mu\nu}^A G^{A\mu\nu}$	Q_{uH}	$(\bar{l}_p \sigma^{\mu\nu} e_r) H B_{\mu\nu}$	$Q_{H1}^{(2)}$	$(H^\dagger i \overleftrightarrow{D}_\mu^I H)(\bar{l}_p \tau^I \gamma^\mu l_r)$
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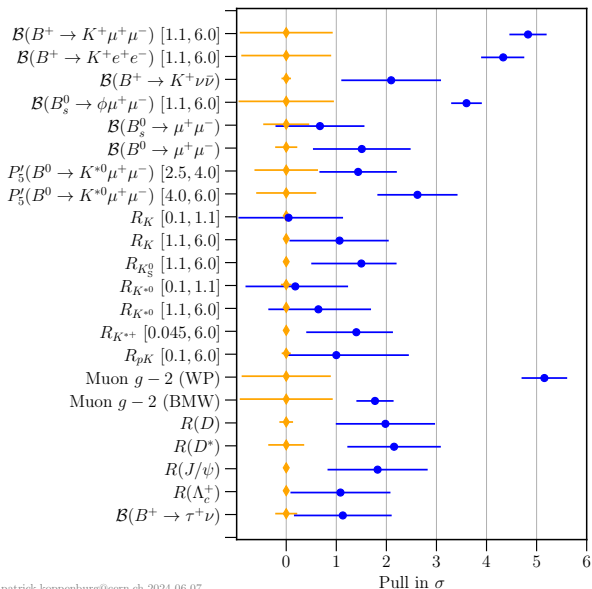
Higgs penguins

“Leave no stone unturned”
= probe as many operators as possible

8 : $(\bar{L}L)(\bar{L}L)$		8 : $(\bar{R}R)(\bar{R}R)$		8 : $(\bar{L}L)(\bar{R}R)$	
Q_{ll}	$(\bar{l}_p \gamma_\mu l_r)(\bar{l}_s \gamma^\mu l_t)$	Q_{ee}	$(\bar{e}_p \gamma_\mu e_r)(\bar{e}_s \gamma^\mu e_t)$	Q_{le}	$(\bar{l}_p \gamma_\mu l_r)(\bar{e}_s \gamma^\mu e_t)$
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		$Q_{leq_2}^{(1)}$	$(\bar{l}_p^c e_r) e_j k (\bar{q}_s^c u_t)$
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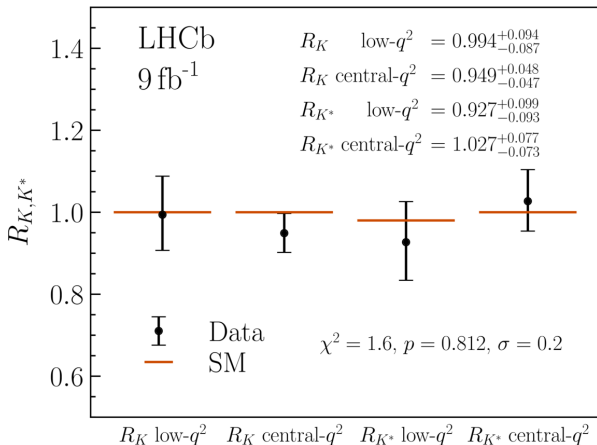
Guidance from Anomalies?



patrick.koppenburg@cern.ch 2024-06-07

Lepton Flavor Universality Tests in $b \rightarrow s\ell\ell$

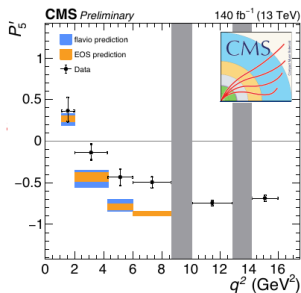
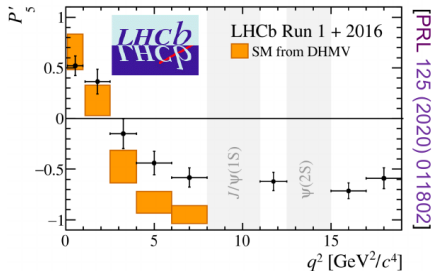
LHCb 2212.09152, 2212.09153



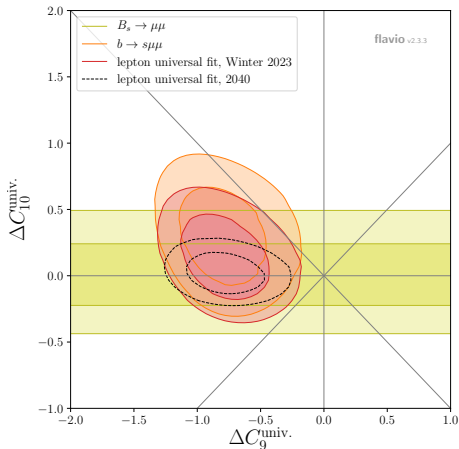
R_K and R_{K^*} are consistent with SM expectations at the $\sim 5\%$ level

New Physics in $b \rightarrow s\mu\mu$?

Many other experimental results on $b \rightarrow s\mu\mu$ don't agree well with SM predictions.
"Anomalies" both in branching ratios and angular distributions (P'_5).



Fits of $b \rightarrow s\ell\ell$ Data to Lepton Universal New Physics



WA, Gadam, Profumo 2306.15017

(also Greljo et al. 2212.10497; Ciuchini et al. 2212.10516;

Alguero et al. 2304.07330; Guadagnoli et al. 2308.00034;

Bordone et al. 2401.18007; ...)

$$\Delta C_9^{\text{univ.}}(\bar{s}\gamma_\alpha P_L b)(\bar{\ell}\gamma^\alpha \ell)$$

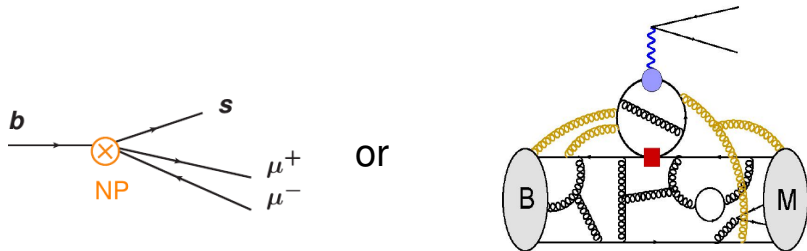
$$\Delta C_{10}^{\text{univ.}}(\bar{s}\gamma_\alpha P_L b)(\bar{\ell}\gamma^\alpha \gamma_5 \ell)$$

- ▶ LFU ratios don't give constraints (by construction)
- ▶ $B_s \rightarrow \mu^+ \mu^-$ branching ratio in agreement with SM
- ▶ $b \rightarrow s\mu\mu$ observables (P'_5 and semileptonic BRs) prefer non-standard C_9
- ▶ our fit finds a $\sim 3\sigma$ preference for new physics in C_9

$$\Delta C_9^{\text{univ.}} \simeq -0.80 \pm 0.22$$

$$\Delta C_{10}^{\text{univ.}} \simeq +0.12 \pm 0.20$$

New Physics or Underestimated Hadronic Effects?



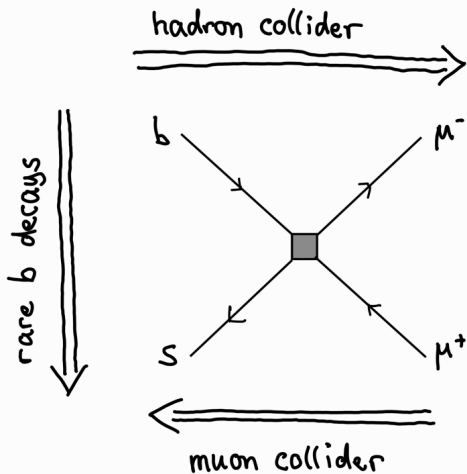
It is very difficult to distinguish lepton flavor universal new physics in C_9 from a long distance hadronic effect (“charm loops”)

$$\Delta C_9^{\text{univ.}}(\bar{s}\gamma_\alpha P_L b)(\bar{\ell}\gamma^\alpha \ell)$$

Lot's of activity to better understand the “charm loops”:
lattice QCD, QCD factorization, dispersion relations, unitarity bounds, data driven methods, generic parameterizations, models, ...

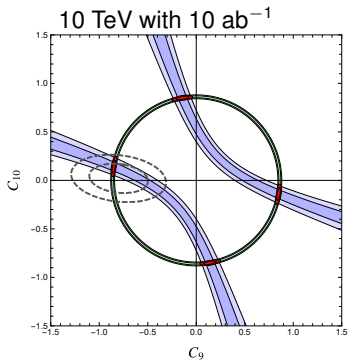
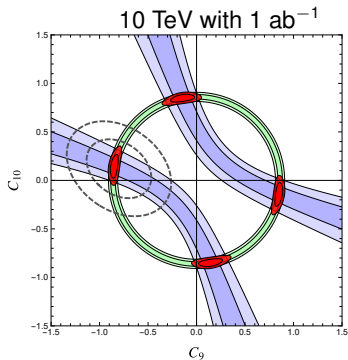
Ciuchini et al. 2212.10516; Gubernari, Reboud, van Dyk, Virto 2206.03797, 2305.06301;
LHCb 2312.09102, 2405.17347; Isidori, Polonski, Tinari 2405.17551 ... many others

Probing Hints of New Physics From All Directions



Sensitivity of a Muon Collider

WA, Gadam, Profumo 2203.07495 and 2306.15017



- ▶ If there is new physics in $b \rightarrow s\mu\mu$, a **10 TeV muon collider** would clearly see it, and one does not need to worry about long distance QCD.

...but hopefully things get sorted out earlier...

(see also Huang et al. 2103.01617; Asadi et al. 2104.05720; Azatov et al. 2205.13552)

Beyond Anomalies

- ▶ “Bread and Butter”: continue to **improve well established probes**:
e.g. mass differences in meson mixing; $B \rightarrow X_s \gamma$; $B_s \rightarrow \mu^+ \mu^-$; ...
- requires high precision hadronic matrix elements from the lattice
- requires high precision CKM input (\rightarrow need to sort out V_{cb} and V_{ub} !)

Beyond Anomalies

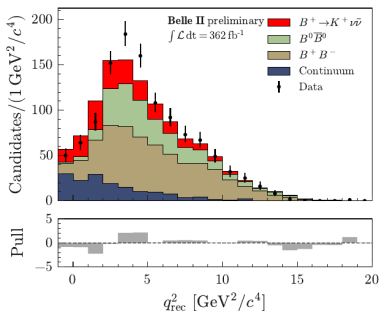
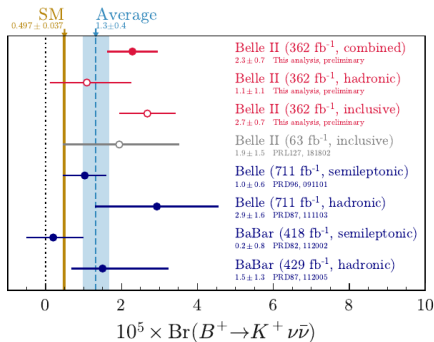
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- ▶ Explore new processes where **$O(1)$ NP effects** are still possible.
 \rightarrow obtain qualitatively new information on a new types of processes

Examples for the near future:

- $B \rightarrow K^{(*)} \nu \bar{\nu}$ (new intriguing results from Belle II)
- CP violation in $D^0 - \bar{D}^0$ oscillations
- rare kaon decays $K \rightarrow \pi \nu \bar{\nu}$, ...
- $b \rightarrow d \ell \ell$ decays

Evidence for $B \rightarrow K\nu\bar{\nu}$

Belle II 2311.14647

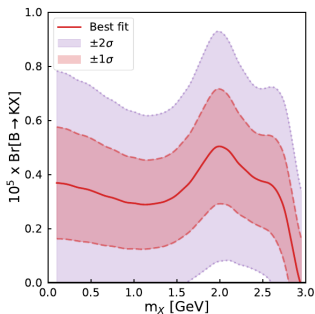


- ▶ Evidence for $B \rightarrow K\nu\bar{\nu}$ at 3.5σ above background and 2.7σ above the SM prediction.
- ▶ Excess of events is particularly pronounced around $q^2 \simeq 4 \text{ GeV}^2$.

A Hint for Light New Physics?

- Instead of fitting the excess with a continuous 3-body spectrum from $B \rightarrow K\nu\bar{\nu}$ one gets a better fit with a new resonance $B \rightarrow KX$

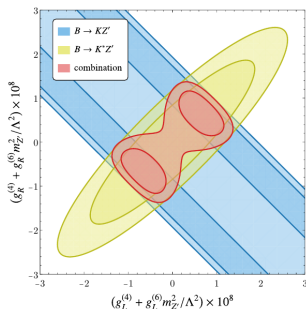
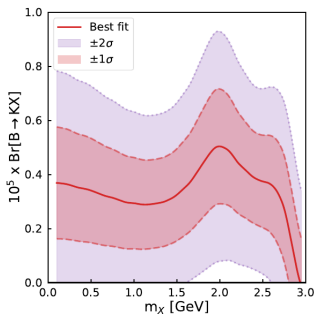
WA, Crivellin, Haigh, Inguglia, Martin Camalich 2311.14629



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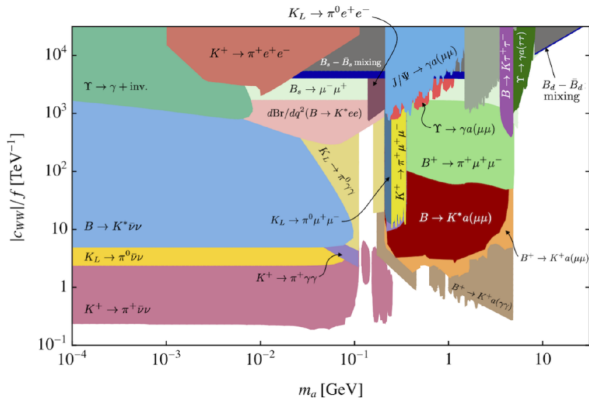


- ▶ Could be for example a Z' or ALP with mass around 2 GeV
- ▶ Constraints from $B \rightarrow K^* \nu\bar{\nu}$ narrow down couplings

see also Bause et al. 2309.00075; Allwicher et al. 2309.02246; Felkl et al. 2309.02940;
McKeen et al. 2312.00982; Fridell et al. 2312.12507; Ho et al. 2401.10112; Gabrielli et al. 2402.05901;
Hou et al 2402.19208; Bolton et al. 2403.13887; He et al 2403.12485; Marzocca et al 2404.06533;
Eguren et al 2405.00108; Buras et al. 2405.06742; ...

Flavor and Axion Like Particles

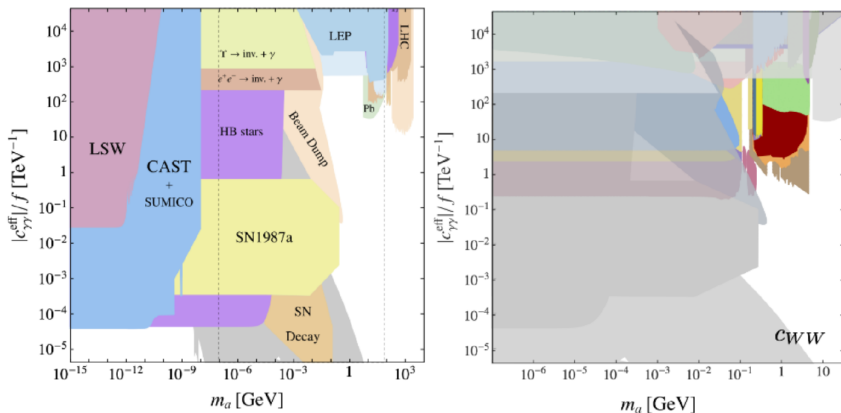
- ▶ Flavor transitions are excellent probes of ALPs, even if the ALPs don't have flavor violating couplings to begin with
- ▶ SM loops necessarily induce flavor violating couplings!



(Bauer, Neubert, Renner, Schnubel, Thamm 2110.10698)

Flavor and Axion Like Particles II

(Bauer, Neubert, Renner, Schnubel, Thamm 2110.10698)

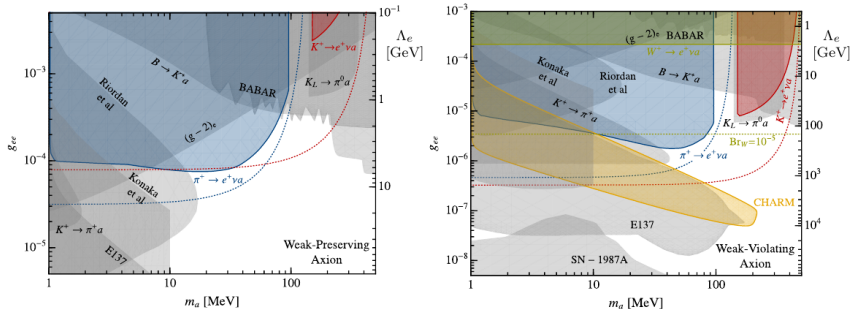


- Flavor violating processes have unique sensitivity to ALPs in the GeV mass range, which is weakly constrained otherwise

Probing Leptophilic ALPs with Pion Decays

- ▶ Can use pion decays to probe ALPs in the 10 – 100 MeV mass range
- ▶ SM decay $\pi^+ \rightarrow e^+ \nu$ is helicity suppressed; adding an ALP can lift the helicity suppression: $\pi^+ \rightarrow e^+ \nu a$

WA, Dror, Gori 2209.00665



- ▶ Planned next generation rare pion decay experiment **PIONEER** (2203.01981) will have unprecedented sensitivity to this decay mode

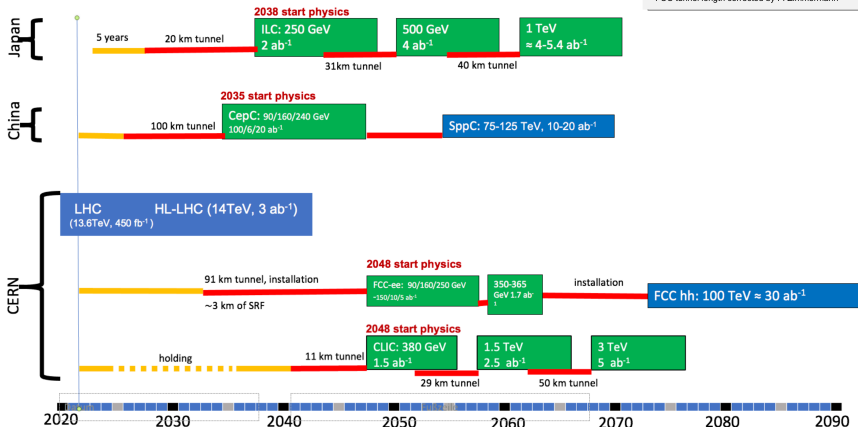
The Future of Flavor

Indicative scenarios of future colliders [considered by ESG]

- Proton collider
- Electron collider
- Muon collider

- Construction/Transformation
- Preparation / R&D

Original from ESPP by Ursula Bassler
 Updated July 25, 2022 by Meenkshi Narain
 FCC tunnel length corrected by F. Zimmermann



[Karl Jacobs @ 2nd ECFA meeting on e^+e^- Higgs, electroweak, and top factories
 Oct 11-13, 2023, Paestum, Italy]

Circular e^+e^- Colliders are Flavor Factories

Running on the Z pole allows one to probe the flavor structure of Z couplings with extreme precision.

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⇒ unique sensitivity to a large number of flavor processes that are not accessible at LHC(b) or Belle II

b Hadrons from 10^{13} Z bosons

FCC-ee Snowmass Whitepaper 2203.06520

Particle production (10^9)	B^0/\bar{B}^0	B^+/B^-	B_s^0/\bar{B}_s^0	B_c^+/\bar{B}_c^-	$\Lambda_b/\bar{\Lambda}_b$	$c\bar{c}$	$\tau^+\tau^-$
Belle II	27.5	27.5	n/a	n/a	n/a	65	45
FCC-ee	620	620	150	4	130	600	170

► FCC-ee/CEPC vs. Belle II:

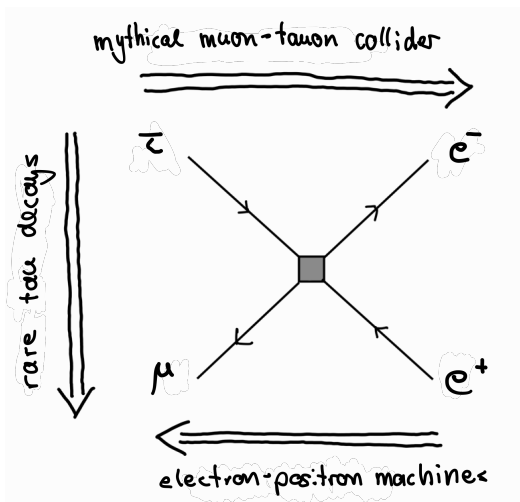
- order of magnitude more B^+ and B^0 , unique opportunities for B_s , B_c , and Λ_b .
- $b\bar{b}$ from Z decays are **highly boosted**.

► FCC-ee/CEPC vs. LHCb:

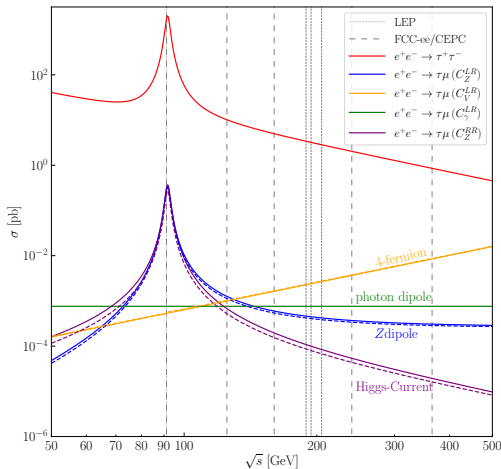
- lower yields at e^+e^- colliders, but **cleaner environment**.
- much easier access to final states with neutrals (π^0 , γ , neutrinos).

$$B_s \rightarrow \tau\tau, \quad B \rightarrow K^*\tau\tau, \quad B_s \rightarrow \phi\nu\bar{\nu}, \quad B_c \rightarrow \tau\nu, \dots$$

Probing Lepton Flavor Violation at FCC-ee/CEPC



Characteristic Dependence on \sqrt{s}

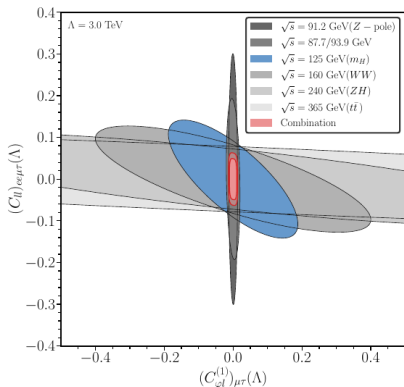


WA, Munbodh, Oh 2305.03869
 (in the plot $\Lambda_{\text{NP}} = 3 \text{ TeV}$, $C_i = 1$)

- ▶ $\tau^+\tau^-$ background falls like $1/s$
- ▶ $\tau\mu$ production increases linearly with s for 4-fermion operators
- ▶ $\tau\mu$ production is flat in s for dipole operators
- ▶ $\tau\mu$ production falls like $1/s$ for Higgs current operators
- ▶ resonance at $s = m_Z^2$ if Z-mediated

Sensitivity Projections for FCC-ee/CEPC

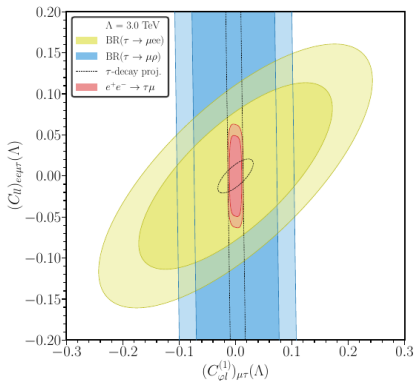
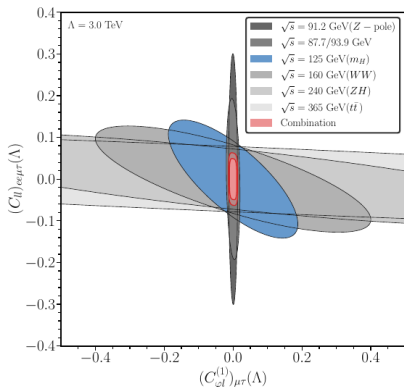
WA, Munbodh, Oh 2305.03869



- The Z -pole searches and the high- \sqrt{s} searches are **complementary**.

Sensitivity Projections for FCC-ee/CEPC

WA, Munbodh, Oh 2305.03869



- ▶ The Z-pole searches and the high- \sqrt{s} searches are **complementary**.
- ▶ Expected **FCC-ee/CEPC sensitivity** rivals the one from current and future searches for **LFV τ decays**.

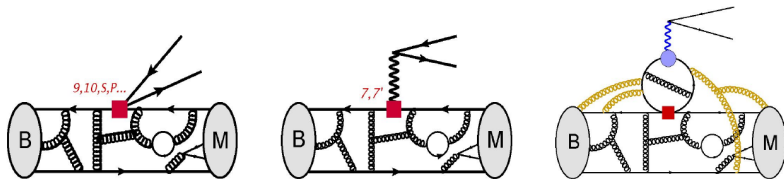
- ▶ Flavor physics will continue to play a crucial role in the search for new physics (→ “Leave no stone unturned”)
- ▶ Anomalies need to be followed up in every way possible.
- ▶ Beyond anomalies, one can expect qualitatively new insights into flavored new physics from a number of processes in the near future.
- ▶ In the far future, circular e^+e^- colliders could enable a very impactful flavor program.

¡Buena Pesca!



Back Up

$b \rightarrow sll$ Amplitudes

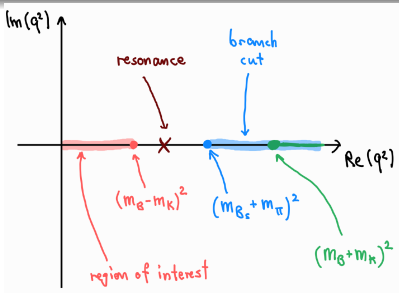


$$\mathcal{A}_\lambda^{L,R} = \mathcal{N}_\lambda \left\{ (C_9 \mp C_{10}) \mathcal{F}_\lambda(q^2) + \frac{2m_b M_B}{q^2} \left[C_7 \mathcal{F}_\lambda^T(q^2) - 16\pi^2 \frac{M_B}{m_b} \mathcal{H}_\lambda(q^2) \right] \right\} + \mathcal{O}(\alpha^2)$$

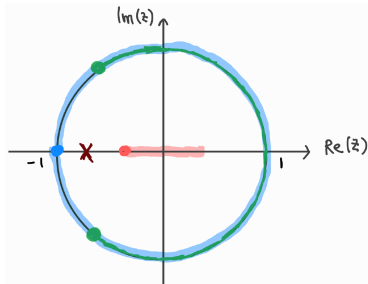
- ▶ Local (Form Factors): $\mathcal{F}_\lambda^{(\tau)}(q^2) = \langle \bar{M}_\lambda(k) | \bar{s} \Gamma_\lambda^{(\tau)} b | \bar{B}(k+q) \rangle$
- ▶ Non-Local: $\mathcal{H}_\lambda(q^2) = i \mathcal{P}_\mu^\lambda \int d^4x e^{iq \cdot x} \langle \bar{M}_\lambda(k) | T \{ j_{\text{em}}^\mu(x), C_i \mathcal{O}_i(0) \} | \bar{B}(q+k) \rangle$

(talk by Javier Virto at Flavour@TH workshop, CERN May 11, 2023)

Parameterization of the Local Form Factors



\Rightarrow



- The form factors can be parameterized by a power series in z with bounded coefficients.

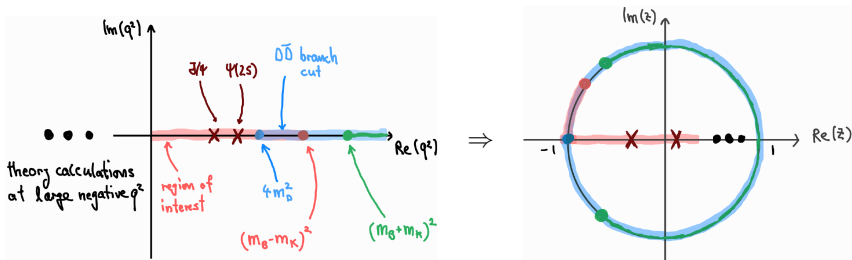
Boyd, Grinstein, Lebed hep-ph/9412324; Caprini, Lellouch, Neubert hep-ph/9712417;

Bourrely, Caprini, Lellouch 0807.2722; ...

Flynn, Juttner, Tsang 2303.11285; Gubernari, Reboud, van Dyk, Virto 2305.06301

$$\mathcal{F}(q^2) = \frac{1}{\mathcal{B}_{\mathcal{F}}(z)\phi_{\mathcal{F}}(z)} \sum_k \alpha_k^{\mathcal{F}} p_k^{\mathcal{F}}(z) , \quad \sum_{\mathcal{F},k} |\alpha_k^{\mathcal{F}}|^2 < 1$$

Parameterization of the Charm Loop



- ▶ Proposed parameterization analogous to the local form factors.
- ▶ Works for q^2 below the $D\bar{D}$ branch cut.

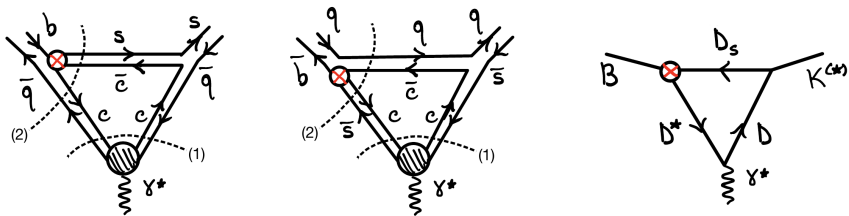
Bobeth, Chruszcz, van Dyk, Virto 1707.07305; Gubernari, van Dyk, Virto 2011.09813;
Gubernari, Reboud, van Dyk, Virto 2206.03797

$$\mathcal{H}(q^2) = \frac{1}{\mathcal{B}_{\mathcal{H}}(z)\phi_{\mathcal{H}}(z)} \sum_k \beta_k^{\mathcal{H}} p_k^{\mathcal{H}}(z) , \quad \sum_{\mathcal{H},k} |\beta_k^{\mathcal{H}}|^2 < 1$$

Additional Charm Loop Effects?

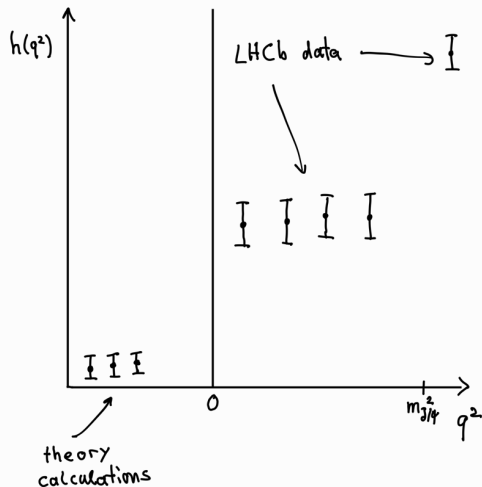
- ▶ The charm loop also gives “triangle diagrams” involving e.g. intermediate $D_s \bar{D}$ states

Ciuchini, Fedele, Franco, Paul, Silvestrini, Valli 2212.10516



- ▶ E.g. decay $B \rightarrow D_s D^*$ followed by rescattering $D_s D^* \rightarrow K^{(*)} \gamma^*$
- ▶ How disruptive are they to the proposed parameterization?

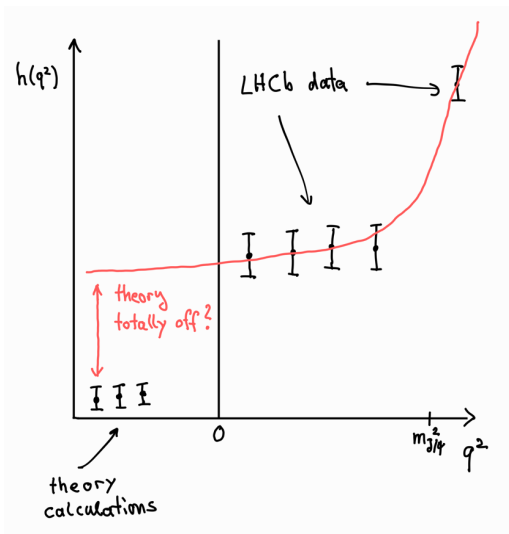
Cartoon Picture of the “Charm Loop”



[Note: This is highly oversimplified]

Fit the charm loop parameterization to data and/or theory calculations

Cartoon Picture of the “Charm Loop”

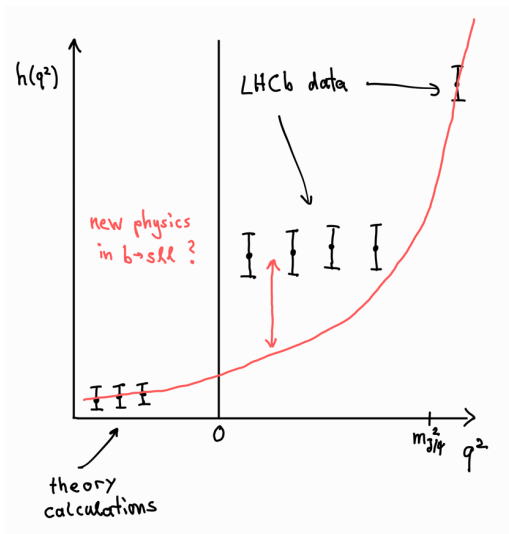


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How **reliable** are the theory calculations?

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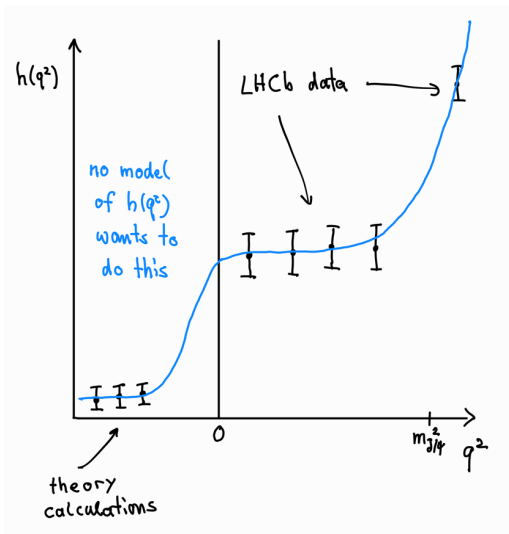


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Fit the charm loop parameterization to data and/or theory calculations

How **reliable** are the theory calculations?

Is the parameterization **robust** / **sufficiently generic**?