

# HIGGS MASS PREDICTIONS IN THE CP-VIOLATING HIGH-SCALE NMSSM

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**SUSY 2024**

Theory meets Experiment

THE 31TH INTERNATIONAL CONFERENCE ON SUPERSYMMETRY  
AND UNIFICATION OF FUNDAMENTAL INTERACTIONS

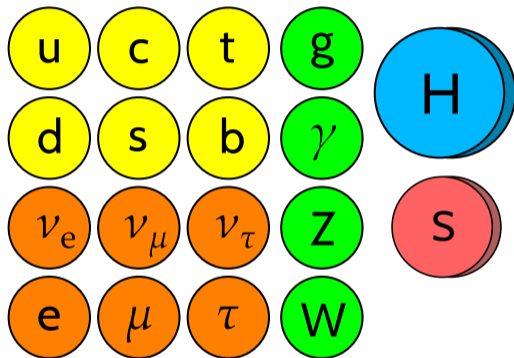
Madrid, 11 June 2024

# Outline

- 1 Next-to-Minimal Supersymmetric Standard Model
- 2 Higgs mass calculations in the EFT approach
- 3 Numerical results
- 4 Summary

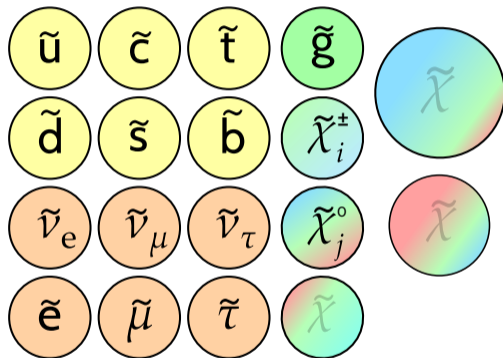
# The Next-to-Minimal Supersymmetric Standard Model

## Standard Model particles\*\*



● Quarks   
 ● Leptons   
 ● Gauge bosons   
 ● Higgs   
 ● Singlet Higgs

## Supersymmetric partners



● Squarks   
 ● Sleptons   
 ● Gluino   
 ● Neutralinos & charginos

# The Next-to-Minimal Supersymmetric Standard Model

## Complex Next-to-Minimal Supersymmetric Standard Model

### Superpotential of the $\mathbb{Z}_3$ -symmetric NMSSM

$$\mathcal{W}_{\text{NMSSM}} = [y_e \hat{H}_d \cdot \hat{L} \hat{E}^c + y_d \hat{H}_d \cdot \hat{Q} \hat{D}^c - y_u \hat{H}_u \cdot \hat{Q} \hat{U}^c] - \lambda \hat{S} \hat{H}_d \cdot \hat{H}_u + \frac{1}{3} \kappa \hat{S}^3$$

- Complex scalar singlet extension of the MSSM ( $\lambda, \kappa$  complex, e.g.  $\lambda = |\lambda| e^{i\varphi_\lambda}$ )

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- ▶  $\mu$  parameter is generated dynamically:

$$\mu_{\text{eff}} = \frac{e^{i\varphi_s} v_s \lambda}{\sqrt{2}}$$

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### Higgs sector

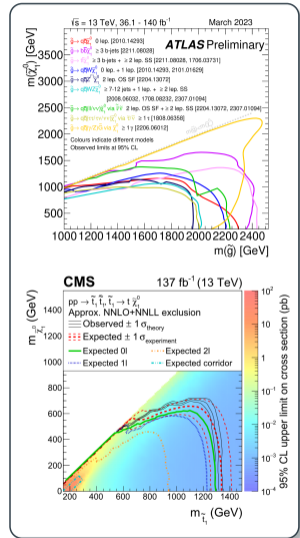
$$H_d = \begin{pmatrix} \frac{v_d + h_d + i a_d}{\sqrt{2}} \\ h_d^- \end{pmatrix}, \quad H_u = e^{i\varphi_u} \begin{pmatrix} h_u^+ \\ \frac{v_u + h_u + i a_u}{\sqrt{2}} \end{pmatrix}, \quad S = \frac{e^{i\varphi_S}}{\sqrt{2}} (v_S + h_S + i a_S)$$

$$\tan \beta = \frac{v_u}{v_d}$$

$$v = \sqrt{v_u^2 + v_d^2} = 246 \text{ GeV}$$

$h_d, h_u, h_S, a_d, a_u, a_S$  and  $h_d^\pm, h_u^\pm$  mixing to  $h_1, h_2, h_3, h_4, h_5, G^0$  and  $h^\pm, G^\pm$

# Supersymmetry – out of reach?





# Higgs mass calculations in the NMSSM

Constraining the **NMSSM parameter space** with the  $m_h^{\text{SM}}$  measurement

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### Fixed-order status:

- Full 1L, 2L, in different renormalisation schemes ( $\overline{\text{DR}}$ , mixed OS- $\overline{\text{DR}}$ )

*[Ellwanger et al. '93, '05][Elliot et al. '93][Pandita '93][King, White '95][Degrassi, Slavich '10][Staub et al. '10][Drechsel et al. '17][Ham et al. '01-'07][Funakubo, Tao '04][Cheung et al. '10][Goodsell, Staub '17][Domingo et al. '17][Goodsell et al. '15][Ender et al. '12][Graf et al. '12][Mühlleitner et al. '14][Dao et al. '19-'21]*

- Tools: FlexibleSUSY *[Athron et al.]*, NMSSMCALC *[Baglio et al.]*, NMSSMTools *[Ellwanger et al.]*, SOFTSUSY *[Allanach et al.]*, SARAH/Spheno *[Porod, Staub]*

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### Status of effective field theory (EFT) approach:

- ▶ Pole-mass matching in FlexibleEFTHiggs *[Athron et al. '17]*, SARAH/Spheno *[Staub, Porod '17]*
- ▶ Automated full 1L EFT matching in SARAH *[Gabelmann et al. '18-'19]*
- ▶ Full 1L + (NMSSM-specific) 2L EFT matching in the real NMSSM *[Bagnaschi, Goodsell, Slavich '22]*

# Higgs mass calculations at higher orders

## Fixed-order calculations for the Higgs mass:

- ▶ Full perturbative series truncated at fixed order
- ▶ Reliable for not too high SUSY masses
- ▶ Dominant corrections from top/stop sector, e.g. at 1-loop:  $\Delta M_h^2 \sim Y_t \ln \frac{m_{\tilde{t}}^2}{m_t^2}$

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If SUSY masses (e.g. stops) are heavy: large separation of scales

$$\text{EW scale: } m_t \sim v \ll \text{SUSY scale: } m_{\tilde{t}} \sim M_{\text{SUSY}} \Rightarrow \ln \frac{M_{\text{SUSY}}^2}{v^2} \gg 1$$

Large logs  $\ln \frac{M_{\text{SUSY}}^2}{v^2}$  from higher orders are relevant and **need to be resummed!**

# Effective field theory approach to calculating $M_h$

Assuming **all** SUSY particles are heavy:

Consider the SM as a (renormalisable) **effective field theory (EFT)** valid at the EW scale  $\sim m_t \sim v$ , and the NMSSM as its **UV completion** at the high scale  $\sim M_{\text{SUSY}}$

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→ Non-log terms  $\mathcal{O}(v/M_{\text{SUSY}})$  only included partially: **EFT valid for  $v/M_{\text{SUSY}} \ll 1!$**

# Matching the NMSSM parameters to the SM

**Matching conditions** relate the SM and NMSSM couplings such that both theories describe the **same physics at the high scale  $Q = Q_{\text{match}}$**

$V^{\text{SM}} \supset \lambda^{\text{SM}} |H|^4$

$$\lambda^{\text{SM}}(Q) \stackrel{!}{=} \lambda^{\text{NMSSM}}(Q), \quad Y_i^{\text{SM}}(Q) \stackrel{!}{=} Y_i^{\text{NMSSM}}(Q), \quad g_j^{\text{SM}}(Q) \stackrel{!}{=} g_j^{\text{NMSSM}}(Q), \quad \dots$$

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In general (at the scale  $Q_{\text{match}}$ ):



$n$ -loop  $m$ -point amplitudes with the same external (light) states should yield the same results

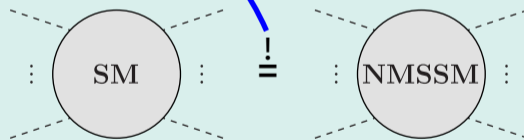
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$$“M_h^{\text{SM}} = M_h^{\text{NMSSM}}”$$

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⇒ Compare both approaches, estimate size of  $\mathcal{O}(v/M_{\text{SUSY}})$  terms

# Quartic-coupling matching

Evaluated for  $v_u, v_d \rightarrow 0$  ( $\tan \beta = \text{const.}, v_S \neq 0$ ) and vanishing ext. momentum

$$\lambda^{\text{SM}}(Q_{\text{match}}) = \lambda^{\text{NMSSM}}(Q_{\text{match}})$$

with

$$\lambda^{\text{NMSSM}}(Q_{\text{match}}) = \lambda_h^{\text{NMSSM,tree}} + \Delta\lambda_h^{\text{NMSSM,1L}} + \Delta\lambda_h^{\text{MSSM,2L}}$$



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[SusyHD: Pardo Vega, Villadoro '15]

$\frac{m}{M}$  light  
scale  
= heavy  
scale

$$= \left[ \text{tree} + \left( \text{1-loop } m \rightarrow 0 \right) \right] + \left( \text{2-loop } M \right)$$

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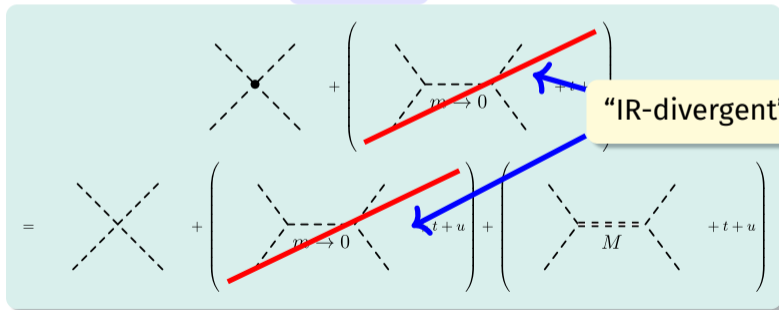
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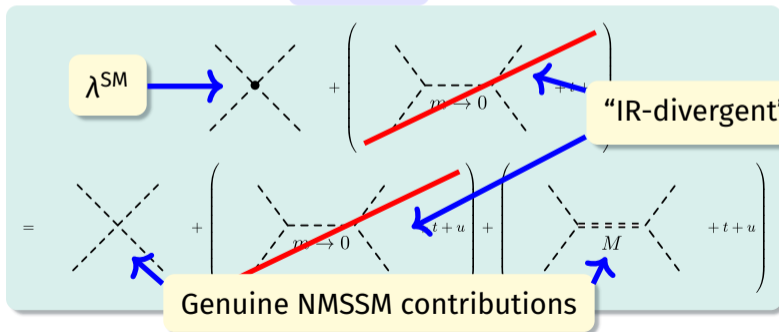
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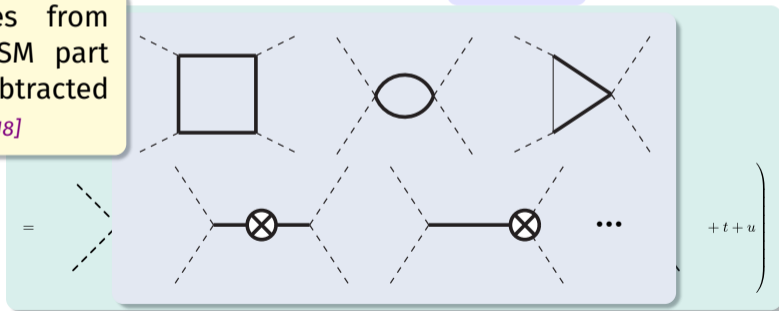
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1L amplitudes from SARAH with SM part already subtracted

[Gabelmann et al. '18]

$\stackrel{M}{=} =$  heavy scale



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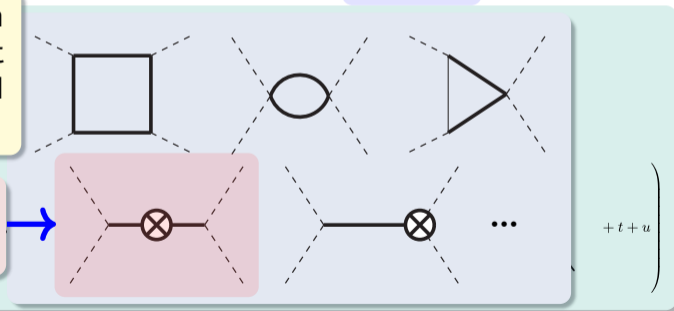
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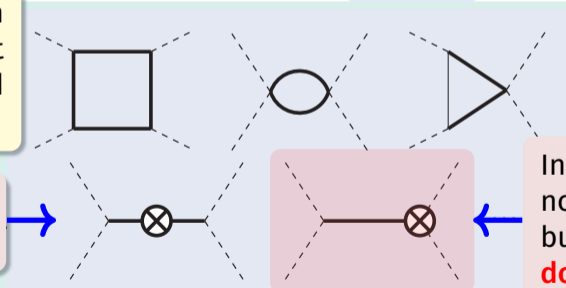
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In CPV case:  
non-vanishing contribution from **CP-odd doublet tadpole!**

[SusyHD: Pardo Vega, Villadoro '15]

# Pole-mass matching

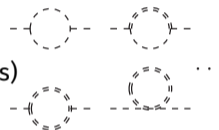
Demand that the pole masses of the SM-like Higgs states are the same:

$$(M_h^{\text{SM}})^2 = (M_h^{\text{NMSSM}})^2$$

e.g. [Athron et al. '16][Braathen et al. '18]

with  $(M_h^X)^2 = (m_h^X)^2 - \hat{\Sigma}_h^{\text{SM}}((M_h^X)^2)$

$(m_h^X$ : SM(-like)  $\overline{\text{MS}}$  ( $\overline{\text{DR}}$ ) Higgs mass in the SM (NMSSM);  $\hat{\Sigma}_h^X$ : renormalised self energies)



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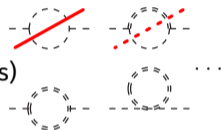
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Use  $\overline{\text{MS}}$  relation  $(m_h^{\text{SM}})^2 = 2(v^{\text{SM}})^2 \lambda_h^{\text{SM}}$  and solve for  $\lambda_h^{\text{SM}}$ :

$$\lambda_h^{\text{SM}} = \frac{1}{2(v^{\text{NMSSM}})^2} \left[ (m_h^{\text{NMSSM}})^2 (1 - 2\Delta\hat{\Sigma}'_h) - \Delta\hat{\Sigma}_h \right]$$

with

$$\Delta\hat{\Sigma}_h^{(r)} \equiv \hat{\Sigma}_h^{\text{NMSSM}(r)}(0) - \hat{\Sigma}_h^{\text{SM}(r)}(0)$$

⇒ Consistent expansion at 1L, captures leading terms when expanding in  $v/M_{\text{SUSY}}$



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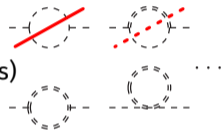
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$(m_h^X$ : SM(-like)  $\overline{\text{MS}}$  ( $\overline{\text{DR}}$ ) Higgs mass in the SM (NMSSM);  $\hat{\Sigma}_h^X$ : renormalised self energies)



Use  $\overline{\text{MS}}$  relation  $(m_h^{\text{SM}})^2 = 2(v^{\text{SM}})^2 \lambda_h^{\text{SM}}$  and solve for  $\lambda_h^{\text{SM}}$ :

$$\lambda_h^{\text{SM}} = \frac{1}{2(v^{\text{NMSSM}})^2} \left[ (m_h^{\text{NMSSM}})^2 (1 - 2\Delta\hat{\Sigma}'_h) - \Delta\hat{\Sigma}_h \right]$$

with

$$\Delta\hat{\Sigma}_h^{(r)} \equiv \hat{\Sigma}_h^{\text{NMSSM}(r)}(0) - \hat{\Sigma}_h^{\text{SM}(r)}(0)$$

⇒ Consistent expansion at 1L, captures leading terms when expanding in  $v/M_{\text{SUSY}}$

$$(v^{\text{SM}})^2 = (v^{\text{NMSSM}})^2 + \delta v^2$$

# Pole-mass matching

Demand that the pole masses of the SM-like Higgs states are the same:

$$(M_h^{\text{SM}})^2 = (M_h^{\text{NMSSM}})^2$$

e.g. [Athron et al. '16][Braathen et al. '18]

with  $(M_h^X)^2 = (m_h^X)^2 - \hat{\Sigma}_h^{\text{SM}}((M_h^X)^2)$

**Numerical limit of  $v \rightarrow 0$ :** excellent agreement with  $\lambda_h^{\text{SM}}$  from **quartic-coupling matching!**

(: renormalised self energies)

for  $\lambda_h^{\text{SM}}$ :

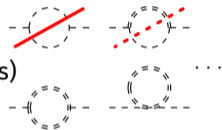
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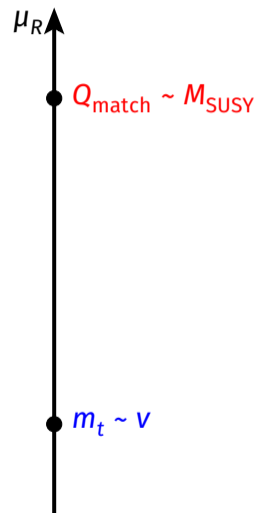
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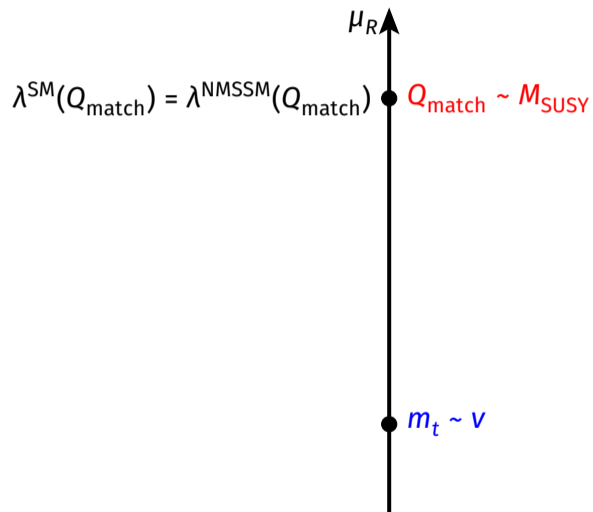
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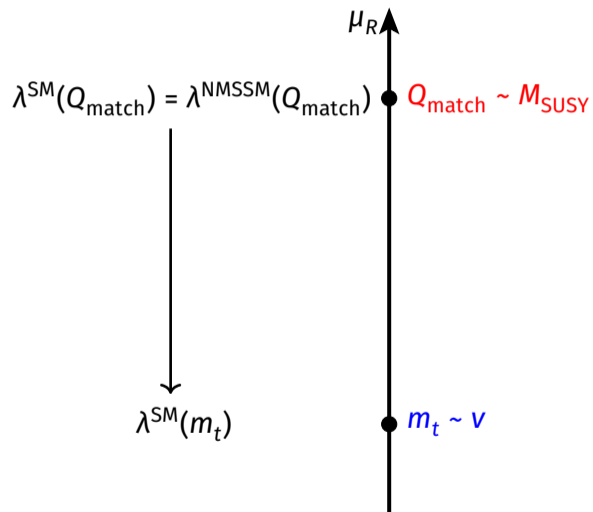
# Renormalisation-group running of $\lambda^{\text{SM}}$



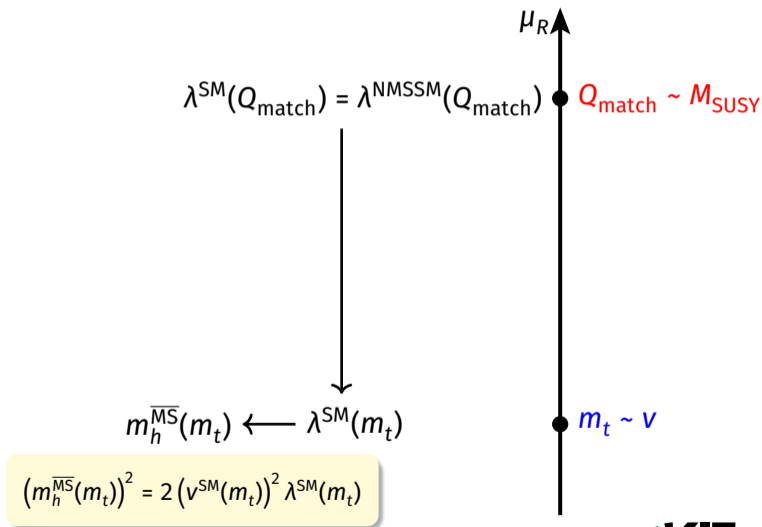
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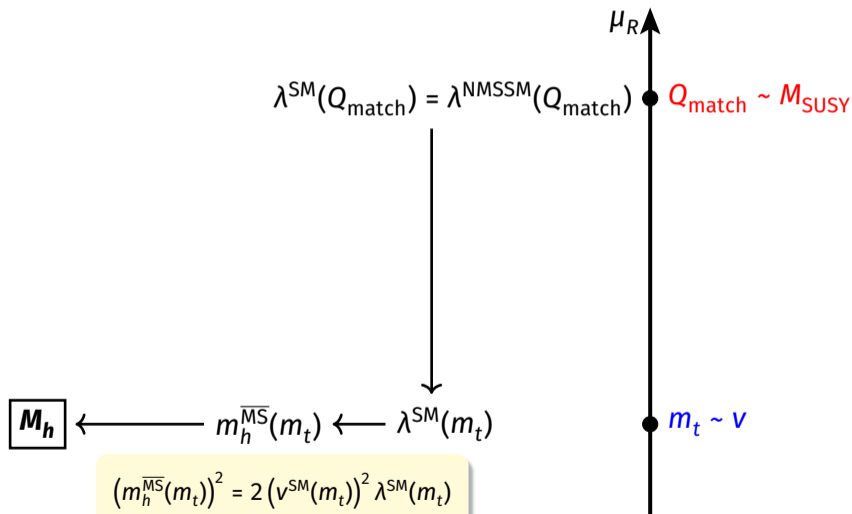
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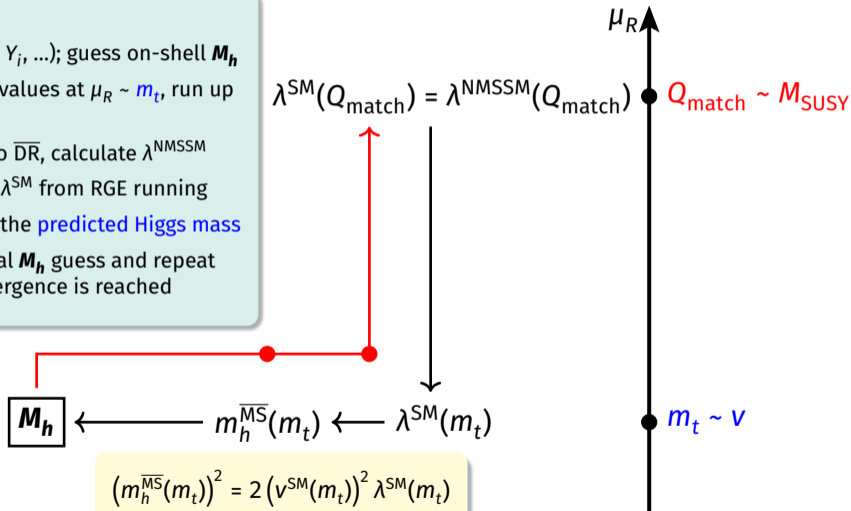
# Renormalisation-group running of $\lambda^{\text{SM}}$



# Renormalisation-group running of $\lambda^{\text{SM}}$

## Iterative procedure

1. Choose SM inputs ( $g_i, Y_i, \dots$ ); guess on-shell  $M_h$
2. Convert inputs to  $\overline{\text{MS}}$  values at  $\mu_R \sim m_t$ , run up to  $\mu_R = Q_{\text{match}}$
3. Convert parameters to  $\overline{\text{DR}}$ , calculate  $\lambda^{\text{NMSSM}}$
4. Compare  $\lambda^{\text{NMSSM}}$  with  $\lambda^{\text{SM}}$  from RGE running
- 5a. If same, chosen  $M_h$  is the **predicted Higgs mass**
- 5b. If not same, vary initial  $M_h$  guess and repeat procedure until convergence is reached

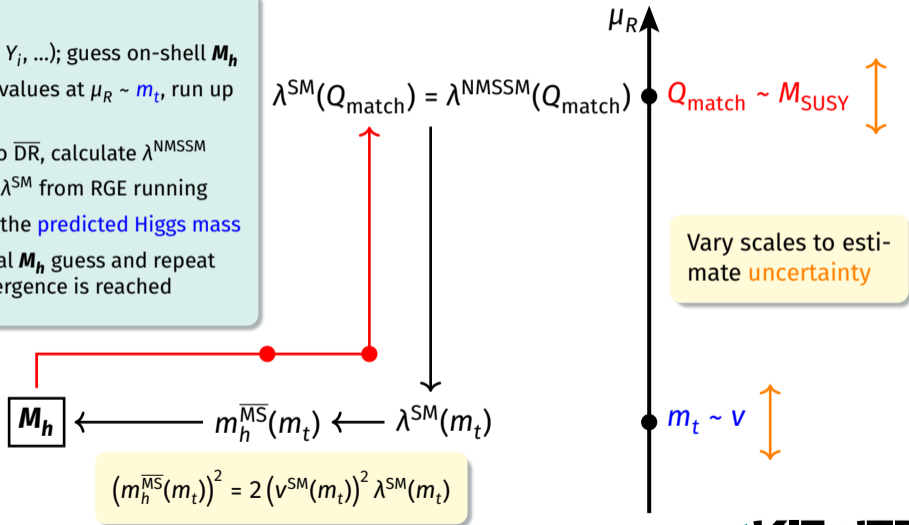




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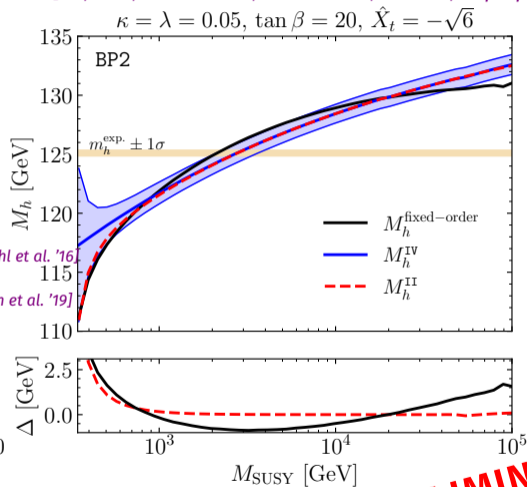
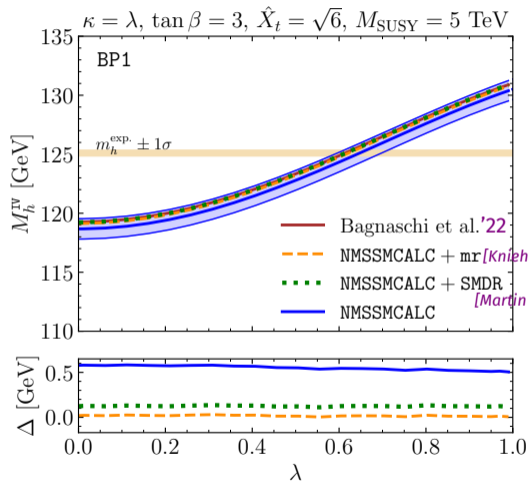
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# Comparison with previous works

[CB, Dao, Gabelmann, Mühlleitner, Rzehak; in preparation]

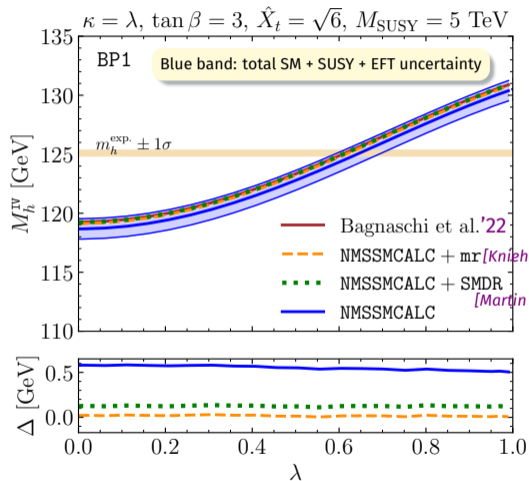


**PRELIMINARY**

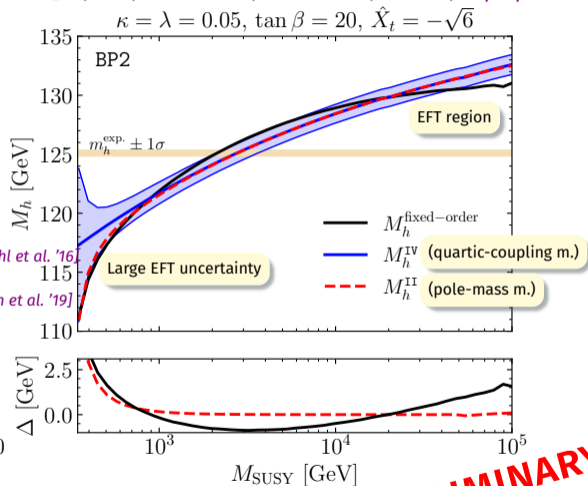
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$\Delta$ : difference to red line (left), difference to blue line (right)

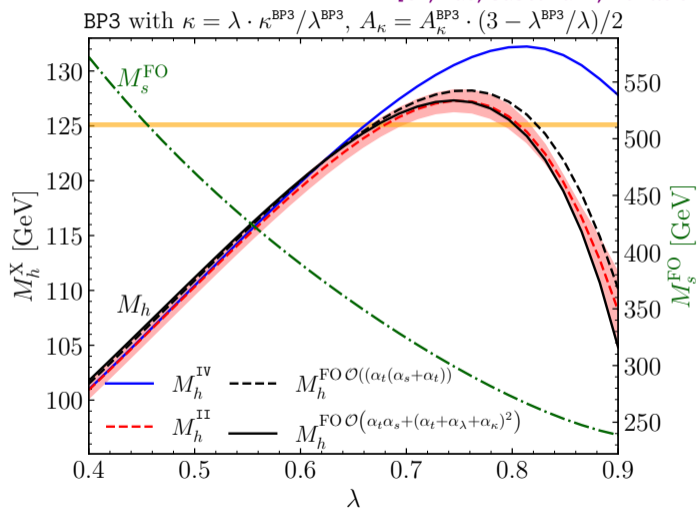


**PRELIMINARY**

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# The case of a light singlet

[CB, Dao, Gabelmann, Mühlleitner, Rzehak; in preparation]

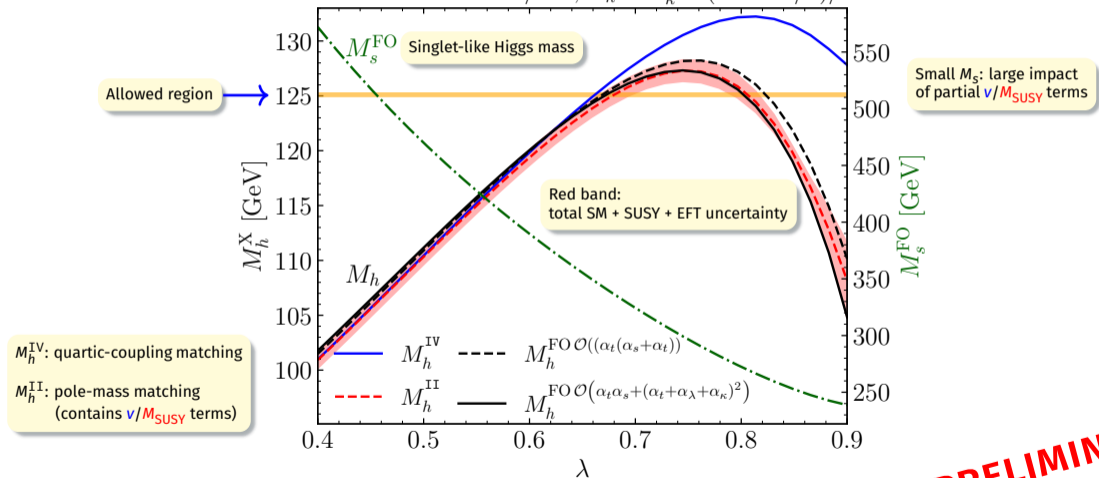


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[CB, Dao, Gabelmann, Mühlleitner, Rzehak; in preparation]

BP3 with  $\kappa = \lambda \cdot \kappa^{\text{BP3}} / \lambda^{\text{BP3}}$ ,  $A_\kappa = A_\kappa^{\text{BP3}} \cdot (3 - \lambda^{\text{BP3}} / \lambda) / 2$

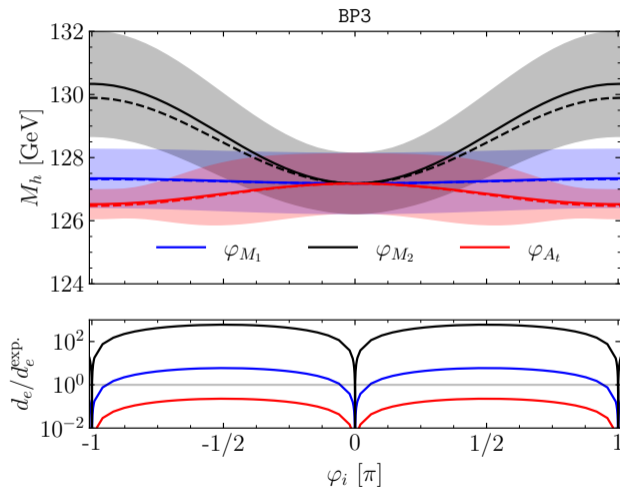


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# Effects of CP-violating phases

[CB, Dao, Gabelmann, Mühlleitner, Rzehak; in preparation]

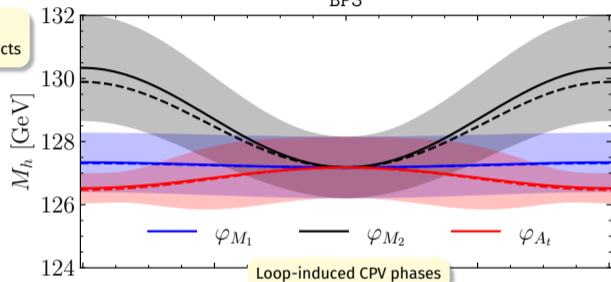


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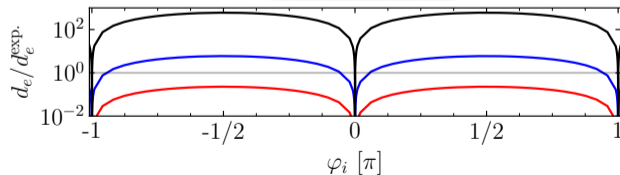
[CB, Dao, Gabelmann, Mühlleitner, Rzehak; in preparation]

BP3

Bands: total uncertainty  
Solid vs. dashed:  $v/M_{\text{SUSY}}$  effects



$d_e/d_e^{\text{exp.}} > 1$ :  
excluded by electron EDMs



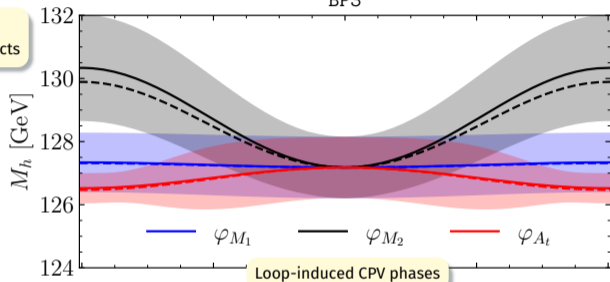
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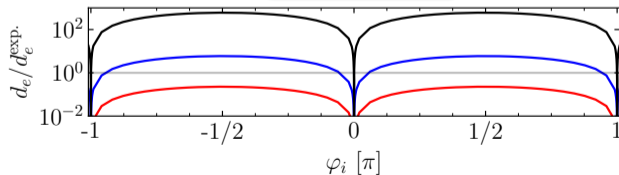
BP3

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Tree-level CPV phases, e.g.  $\varphi_\lambda$ ,  
more strongly constrained

$d_e/d_e^{\text{exp.}} > 1$ :  
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# Summary

## Calculation of the SM-like Higgs mass in the EFT approach for the CPV NMSSM

- ▶ Applicable for heavy SUSY masses ( $v/M_{\text{SUSY}} \ll 1$ ), resums large logarithms
- ▶ Implementation at full 1L (+2L MSSM) via **quartic-coupling** & **pole-mass matching**
  - Excellent agreement found for CPC and CPV case in  $v \rightarrow 0$  limit ✓
  - Estimate of partial  $v/M_{\text{SUSY}}$  contributions for EFT uncertainty

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- ▶ **Spectrum calculator** of 1L & 2L Higgs masses, self couplings, decay widths
- ▶ For the **CP-conserving** and **CP-violating** NMSSM
- ▶ ...and more: electron eDMs, muon  $g - 2$ ,  $\rho$  parameter,  $W$  mass

EFT implementation to appear soon!

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THANK YOU FOR YOUR ATTENTION! 😊

# Backup

# Quartic-coupling matching: tree-level contribution

$$\begin{aligned}
 \lambda_h^{\text{NMSSM,tree}} = & \underbrace{\frac{1}{8}(g_1^2 + g_2^2) \cos^2 2\beta}_{\text{MSSM } D\text{-terms}} + \underbrace{\frac{1}{4}|\lambda|^2 \sin^2 2\beta}_{\text{NMSSM } F\text{-terms}} \\
 & - \frac{1}{48|\kappa|^2 M_S^2 (3M_S^2 + M_{A_S}^2)} \left( 3|\kappa|^2 M_{H^\pm}^2 (1 - \cos 4\beta) \right. \\
 & \quad \left. + (3M_S^2 + M_{A_S}^2) (|\kappa||\lambda| \cos \varphi_y \sin 2\beta - 2|\lambda|^2) \right)^2 \\
 & \quad \quad \quad \underbrace{\hspace{15em}}_{s/t/u\text{-channel } S} \\
 & - \underbrace{\frac{3}{16M_{A_S}^2} |\lambda|^2 (3M_S^2 + M_{A_S}^2) \sin^2 2\beta \sin^2 \varphi_y}_{s/t/u\text{-channel } A_S}
 \end{aligned}$$

# Benchmark points

BP1: [Bagnaschi et al. '22]

BP2: [Slavich et al. '20]

	$\tan \beta$	$\lambda$	$\kappa$	$M_1$	$M_2$	$M_3$	$A_t$	$A_\lambda$	$A_\kappa$	$\mu_{eff.}$	$m_{\tilde{Q}_{L3}}$	$m_{\tilde{t}_{R3}}$
BP1	3.0	0.6	0.6	1.0	2.0	2.5	12.75	0.3	-2.0	1.5	5.0	5.0
BP2	20.0	0.05	0.05	3.0	3.0	3.0	-7.20	-2.85	-1.0	3.0	3.0	3.0
BP3	1.27	0.73	0.62	0.24	1.18	2.3	-0.39	0.06	-1.44	0.49	1.79	1.51

	$M_h^{II}$	$M_h^{IV}$	$m_{h_2}$	$m_{h_3}$	$m_{A_1}$	$m_{A_2}$	$m_{H^+}$
BP1	124.29 ( $h_u$ )	124.31 ( $h_u$ )	2407.6 ( $h_s$ )	2971.8 ( $h_d$ )	2905.7 ( $a$ )	3000.2 ( $a_s$ )	2967.1
BP2	125.26 ( $h_u$ )	125.28 ( $h_u$ )	2996.4 ( $h_d$ )	5744.4 ( $h_s$ )	2985.3 ( $a_s$ )	3010.5 ( $a$ )	2997.8
BP3	127.18 ( $h_u$ )	129.47 ( $h_u$ )	305.5 ( $h_s$ )	659.5 ( $h_d$ )	663.8 ( $a$ )	1308.7 ( $a_s$ )	658.4