

Neutral triple gauge boson vertices and fermionic UV models

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Based on: R. Cepedello, F. Esser, M. Hirsch and V. Sanz

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Contents

I.

Introduction

II.

NTGCs and Effective Field Theory

III.

NTGCs in (fermionic) UV models

IV.

Experimental limits on NTGCs

V.

Conclusions



I.

Introduction

Anomalous NTGCs

Considering Bose symmetry and gauge invariance, the following CP-conserving vertices can be written down (after EWSB):

Gounaris et al. 1999 & 2000

$$ie\Gamma_{ZZV}^{\alpha\beta\mu}(q_1, q_2, q_3) = e \frac{(q_3^2 - m_V^2)}{m_Z^2} \left[f_5^V \epsilon^{\mu\alpha\beta\rho} (q_1 - q_2)_\rho \right],$$

$$ie\Gamma_{Z\gamma V}^{\alpha\beta\mu}(q_1, q_2, q_3) = e \frac{(q_3^2 - m_V^2)}{m_Z^2} \left[h_3^V \epsilon^{\mu\alpha\beta\rho} q_{2,\rho} + \frac{h_4^V}{m_Z^2} q_3^\alpha \epsilon^{\mu\beta\rho\sigma} q_{3,\rho} q_{2,\sigma} \right].$$

Anomalous NTGCs

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A few comments:

⇒ **aNTGC vanish** for all 3 bosons **on-shell**

⇒ “form factors” f_5^V , h_3^V and h_4^V in principle independent parameters

(→ h_4^V can not be generated at 1-loop and $d = 8$)

⇒ There are also CP-violating vertices (ignored in this talk!)

⇒ More vertices possible for more than one boson off-shell

→ **experimentally irrelevant** (therefore ignored here)

⇒ experimentally **at LHC** cleanest final state is $ZZ \rightarrow 4l$



II.

NTGCs and EFT

SMEFT and NTGCs

In **Greens basis** for SMEFT list all operators at $d = 6$ containing only bosons
MatchMakerEFT (1908.05295):

X^3		$X^2 H^2$		$H^2 D^4$	
\mathcal{O}_{3G}	$f^{ABC} G_\mu^{A\nu} G_\nu^{B\rho} G_\rho^{C\mu}$	\mathcal{O}_{HG}	$G_{\mu\nu}^A G^{A\mu\nu} (H^\dagger H)$	\mathcal{R}_{DH}	$(D_\mu D^\mu H)^\dagger (D_\nu D^\nu H)$
$\mathcal{O}_{\widetilde{3G}}$	$f^{ABC} \widetilde{G}_\mu^{A\nu} G_\nu^{B\rho} G_\rho^{C\mu}$	$\mathcal{O}_{H\widetilde{G}}$	$\widetilde{G}_{\mu\nu}^A G^{A\mu\nu} (H^\dagger H)$	$H^4 D^2$	
\mathcal{O}_{3W}	$\epsilon^{IJK} W_\mu^{I\nu} W_\nu^{J\rho} W_\rho^{K\mu}$	\mathcal{O}_{HW}	$W_{\mu\nu}^I W^{I\mu\nu} (H^\dagger H)$	$\mathcal{O}_{H\Box}$	$(H^\dagger H) \Box (H^\dagger H)$
$\mathcal{O}_{\widetilde{3W}}$	$\epsilon^{IJK} \widetilde{W}_\mu^{I\nu} W_\nu^{J\rho} W_\rho^{K\mu}$	$\mathcal{O}_{H\widetilde{W}}$	$\widetilde{W}_{\mu\nu}^I W^{I\mu\nu} (H^\dagger H)$	\mathcal{O}_{HD}	$(H^\dagger D^\mu H)^\dagger (H^\dagger D_\mu H)$
$X^2 D^2$		\mathcal{O}_{HB}	$B_{\mu\nu} B^{\mu\nu} (H^\dagger H)$	\mathcal{R}'_{HD}	$(H^\dagger H) (D_\mu H)^\dagger (D^\mu H)$
\mathcal{R}_{2G}	$-\frac{1}{2} (D_\mu G^{A\mu\nu}) (D^\rho G_{\rho\nu}^A)$	$\mathcal{O}_{H\widetilde{B}}$	$\widetilde{B}_{\mu\nu} B^{\mu\nu} (H^\dagger H)$	\mathcal{R}''_{HD}	$(H^\dagger H) D_\mu (H^\dagger \overleftrightarrow{D}^\mu H)$
\mathcal{R}_{2W}	$-\frac{1}{2} (D_\mu W^{I\mu\nu}) (D^\rho W_{\rho\nu}^I)$	\mathcal{O}_{HWB}	$W_{\mu\nu}^I B^{\mu\nu} (H^\dagger \sigma^I H)$	H^6	
\mathcal{R}_{2B}	$-\frac{1}{2} (\partial_\mu B^{\mu\nu}) (\partial^\rho B_{\rho\nu})$	$\mathcal{O}_{H\widetilde{W}B}$	$\widetilde{W}_{\mu\nu}^I B^{\mu\nu} (H^\dagger \sigma^I H)$	\mathcal{O}_H	$(H^\dagger H)^3$
		$H^2 X D^2$			
		\mathcal{R}_{WDH}	$D_\nu W^{I\mu\nu} (H^\dagger i \overleftrightarrow{D}_\mu^I H)$		
		\mathcal{R}_{BDH}	$\partial_\nu B^{\mu\nu} (H^\dagger i \overleftrightarrow{D}_\mu H)$		

⇒ Only X^3 , $X^2 D^2$ and $H^2 X^2$ contains three (or more) gauge bosons

⇒ No NTGC in any of these operators

⇒ Need to go to $d = 8$ operators!

$d = 8$ basis for NTGCs

The following four $d = 8$ operators generate the NTGCs shown previously:

$$\mathcal{O}_{DB\tilde{B}} = i \frac{c_{DB\tilde{B}}}{\Lambda^4} H^\dagger \tilde{B}_{\mu\nu} (D^\rho B_{\nu\rho}) D_\mu H + \text{h.c.},$$

$$\mathcal{O}_{DW\tilde{W}} = i \frac{c_{DW\tilde{W}}}{\Lambda^4} H^\dagger \tilde{W}_{\mu\nu} (D^\rho W_{\nu\rho}) D_\mu H + \text{h.c.},$$

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All are type:
 $X\tilde{Y}H^2D^2$

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All are type:
 $X\tilde{Y}H^2D^2$

\Rightarrow All 4 operators are necessary to describe f_5^Z , f_5^γ , h_3^Z and h_3^γ as independent parameters

\Rightarrow We disagree with previous literature (Degrande, 2013) that defines:

$$\mathcal{O}_{\tilde{B}W} = i \frac{c_{\tilde{B}W}}{\Lambda^4} H^\dagger \tilde{B}_{\mu\nu} (W_{\mu\rho}) \{D^\rho, D_\mu\} H + \text{h.c.},$$

\Rightarrow This operator can not describe all of f_5^Z , f_5^γ , h_3^Z and h_3^γ

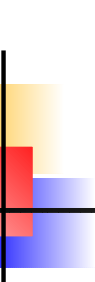
Form factors

The relation to the form factors are:

$$\begin{aligned}
 f_5^Z &= \frac{v^2 m_Z^2}{\Lambda^4} \frac{1}{c_W s_W} \left[s_W^2 c_{DB\tilde{B}} + c_W^2 c_{DW\tilde{W}} + \frac{1}{2} c_W s_W (c_{DW\tilde{B}} + c_{DB\tilde{W}}) \right], \\
 f_5^\gamma &= \frac{v^2 m_Z^2}{\Lambda^4} \frac{1}{c_W s_W} \left[c_W s_W (-c_{DB\tilde{B}} + c_{DW\tilde{W}}) - \frac{1}{2} (s_W^2 c_{DW\tilde{B}} - c_W^2 c_{DB\tilde{W}}) \right], \\
 h_3^Z &= \frac{v^2 m_Z^2}{\Lambda^4} \frac{1}{c_W s_W} \left[c_W s_W (-c_{DB\tilde{B}} + c_{DW\tilde{W}}) + \frac{1}{2} (c_W^2 c_{DW\tilde{B}} - s_W^2 c_{DB\tilde{W}}) \right], \\
 h_3^\gamma &= \frac{v^2 m_Z^2}{\Lambda^4} \frac{1}{c_W s_W} \left[c_W^2 c_{DB\tilde{B}} + s_W^2 c_{DW\tilde{W}} - \frac{1}{2} c_W s_W (c_{DW\tilde{B}} + c_{DB\tilde{W}}) \right],
 \end{aligned}$$

Note that in the special case of $c_{DW\tilde{B}} = c_{DB\tilde{W}}$ (true in many models):

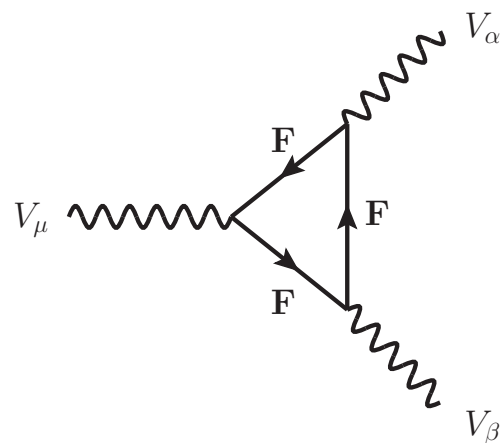
$$\begin{aligned}
 f_5^\gamma &= \frac{v^2 m_Z^2}{\Lambda^4} \frac{1}{c_W s_W} \left[c_W s_W (-c_{DB\tilde{B}} + c_{DW\tilde{W}}) + \frac{1}{2} c_{DW\tilde{B}} (c_W^2 - s_W^2) \right], \\
 h_3^Z &= \frac{v^2 m_Z^2}{\Lambda^4} \frac{1}{c_W s_W} \left[c_W s_W (-c_{DB\tilde{B}} + c_{DW\tilde{W}}) + \frac{1}{2} c_{DW\tilde{B}} (c_W^2 - s_W^2) \right],
 \end{aligned}$$



III.

NTGCs in (fermionic) models

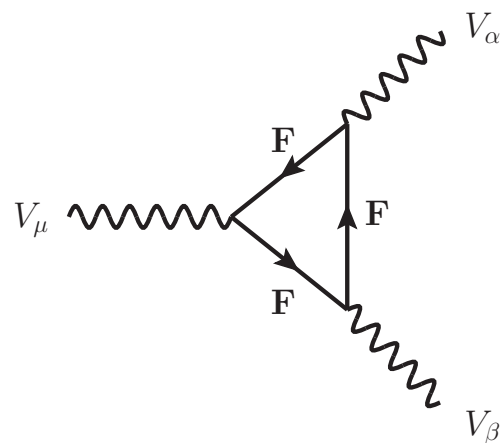
Fermionic loops



Fermionic triangle diagram
in the mass eigenstate basis

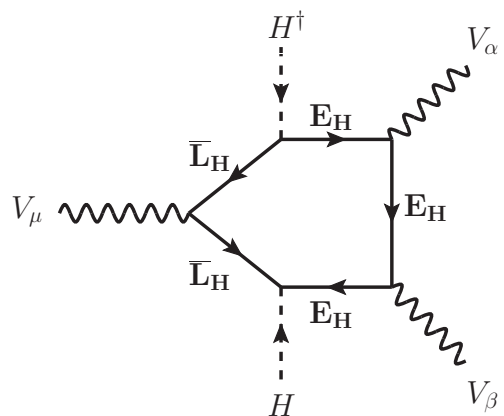
Adequate for calculation
in broken phase

Fermionic loops



Fermionic triangle diagram
in the mass eigenstate basis

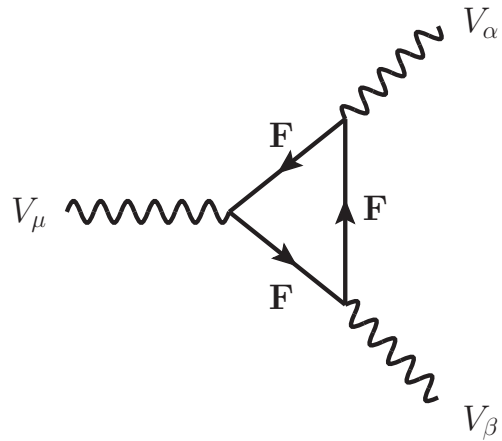
Adequate for calculation
in broken phase



Fermionic pentagon diagram
in the weak eigenstate basis

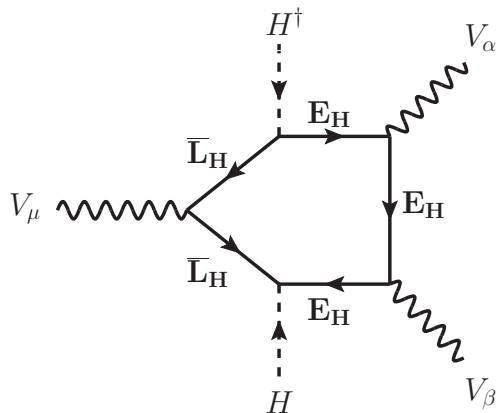
Adequate for calculation
in SMEFT

Fermionic loops



Fermionic triangle diagram
in the mass eigenstate basis

Adequate for calculation
in broken phase



Fermionic pentagon diagram
in the weak eigenstate basis

Adequate for calculation
in SMEFT

$L_H = F_{1,2,-1/2}$ and $E_H = F_{1,1,-1}$
“vector-like” leptons

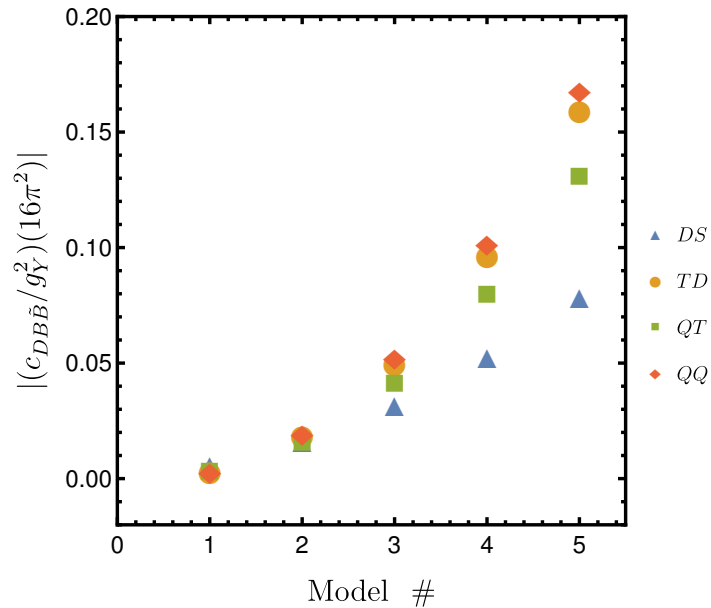
Diracology:

$$\text{tr}[\gamma_\mu \gamma_\nu \gamma_\rho \gamma_\sigma P_{L/R}] = 2(g_{\mu\nu} g_{\rho\sigma} - g_{\mu\rho} g_{\nu\sigma} + g_{\mu\sigma} g_{\nu\rho} \pm i\epsilon^{\mu\nu\rho\sigma}).$$

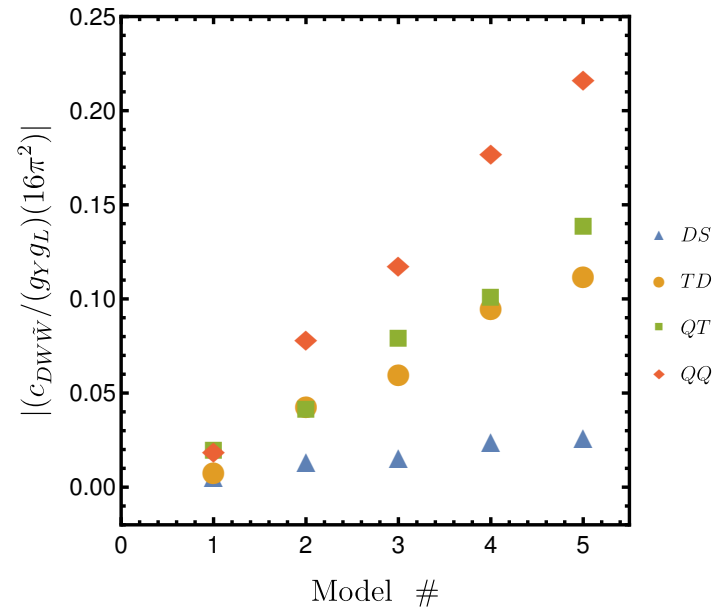
⇒ left-handed and right-handed couplings must differ!

Models and matching

Matching for: $c_{DB\tilde{B}}$



$c_{DB\tilde{W}}$



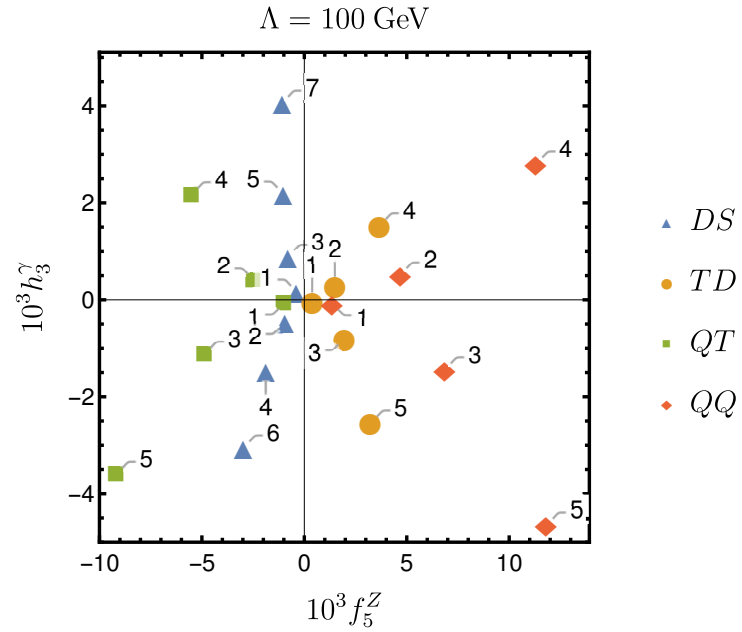
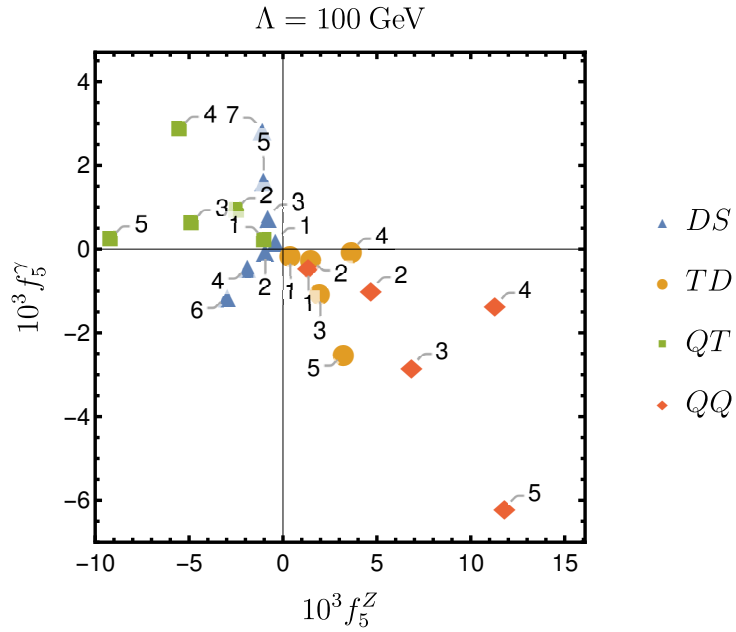
⇒ All models need **two fermions** with $\Delta(Y) = (1/2)$

⇒ Product of $SU_L(2)$ needs to contain doublet:

→ 2×1 (DS), 3×2 (TD), 4×3 (QT), 5×4 (QQ), ...

→ Model # increasing hypercharge from $1 = (0, -1/2)$ to $5 = (-2, -5/2)$
(except DS, which starts at $1 = (-1/2, -1)$)

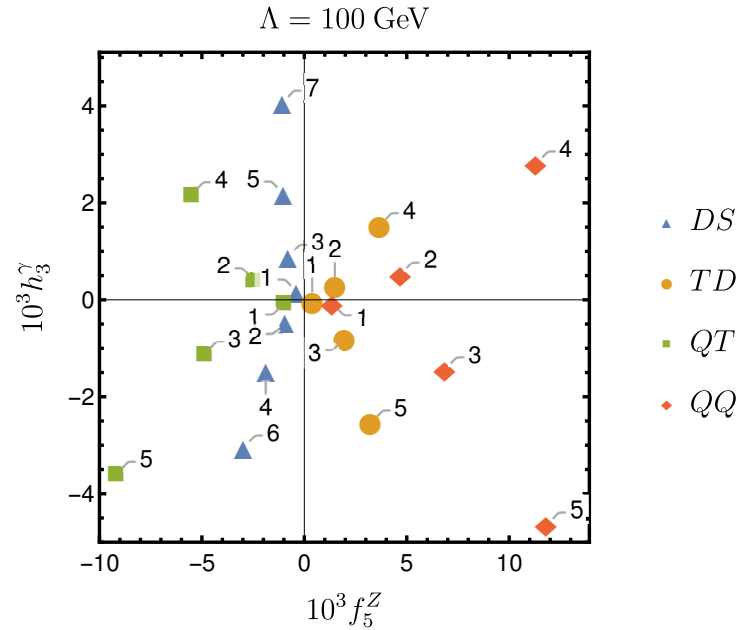
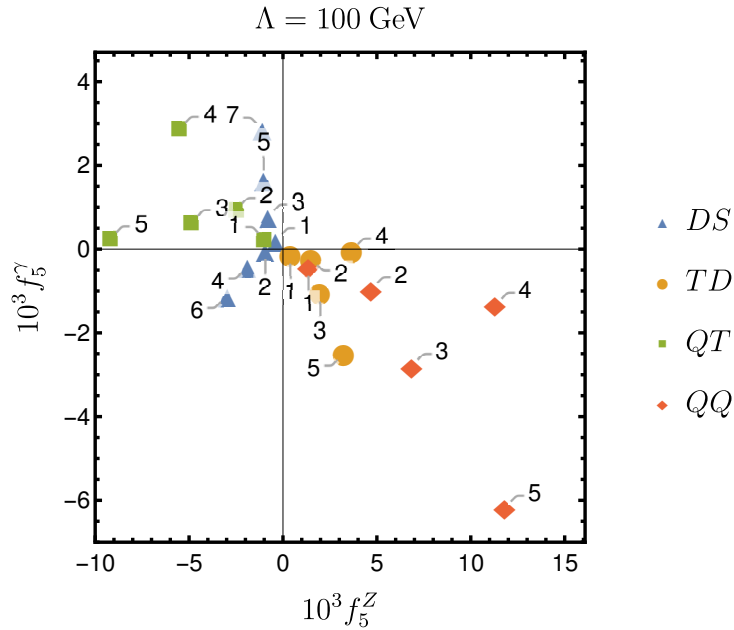
Form factors



\Rightarrow For $\Lambda = 0.1 \text{ TeV}$ typically (few) 10^{-3}

\Rightarrow All models make different predictions!

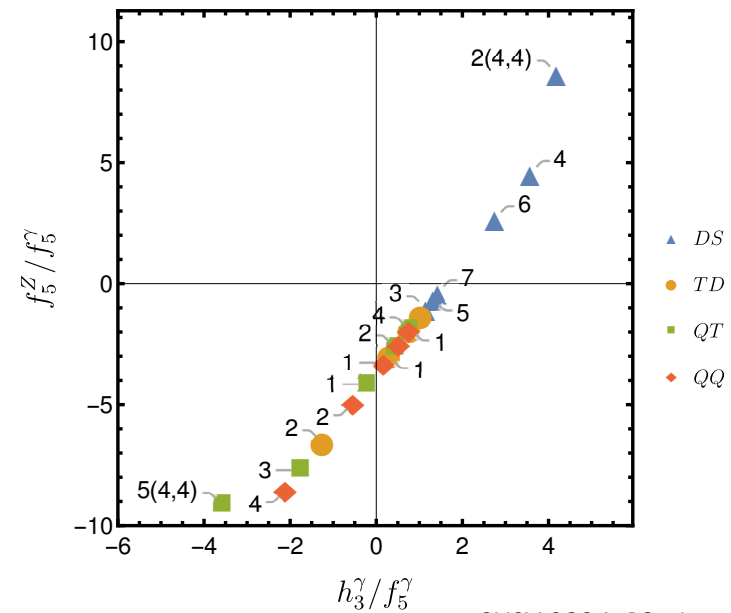
Form factors



\Rightarrow For $\Lambda = 0.1 \text{ TeV}$ typically (few) 10^{-3}

\Rightarrow All models make different predictions!

Plot double-ratio:
 Independent of Λ
 ALL models lie along a line!



$d = 6$ versus $d = 8$

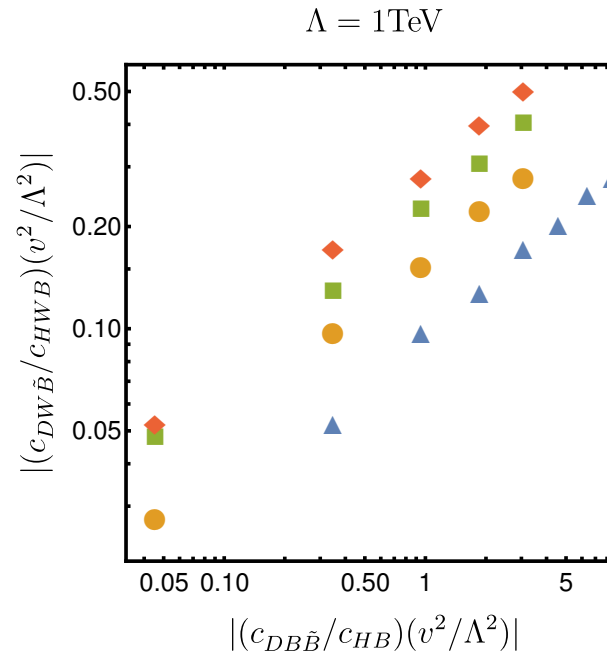
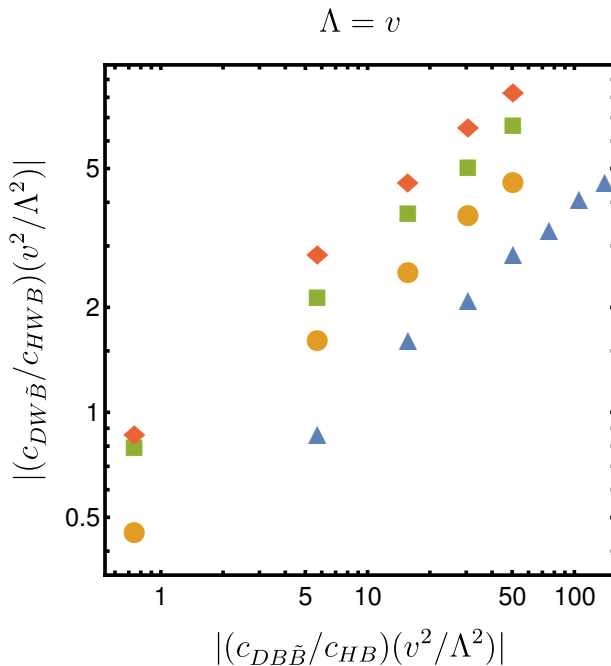
All models that generate NTGCs also will generate the following $d = 6$ operators:

$$\begin{aligned}\mathcal{O}_{HB} &= \frac{c_{HB}}{\Lambda^2} H^\dagger H B_{\mu\nu} B^{\mu\nu}, \\ \mathcal{O}_{HWB} &= \frac{c_{HWB}}{\Lambda^2} H^\dagger H W_{\mu\nu} B^{\mu\nu}, \\ \mathcal{O}_{HW} &= \frac{c_{HW}}{\Lambda^2} H^\dagger H W_{\mu\nu} W^{\mu\nu},\end{aligned}$$

$d = 6$ versus $d = 8$

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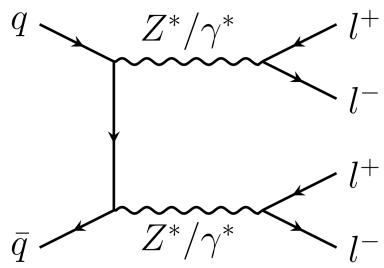
Plot:
double-ratio
 $(d = 6)/(d = 8)$



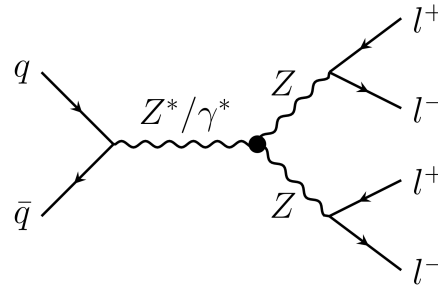
IV.

Experimental limits on NTGCs

Measuring $ZZ \rightarrow 4l$

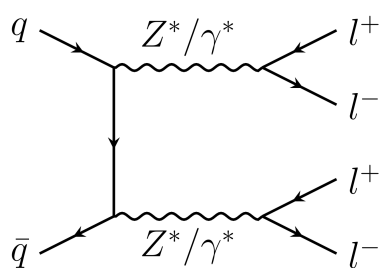


SM background

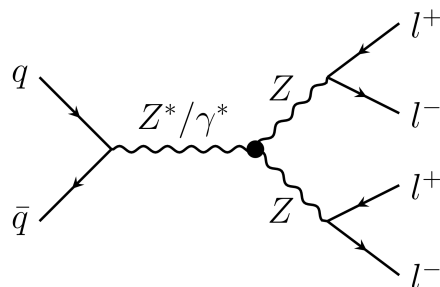


NTGC signal

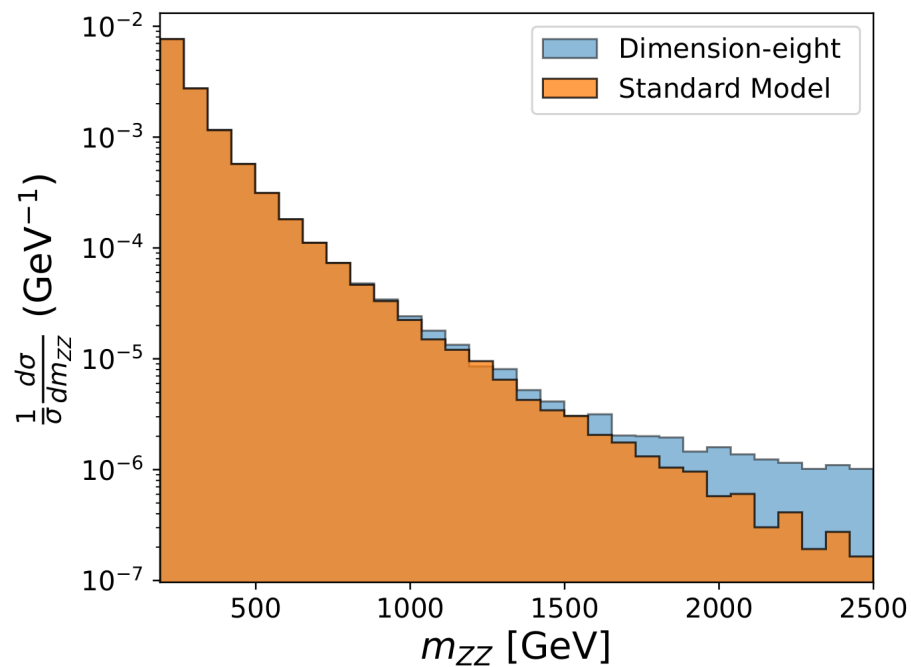
Measuring $ZZ \rightarrow 4l$



SM background

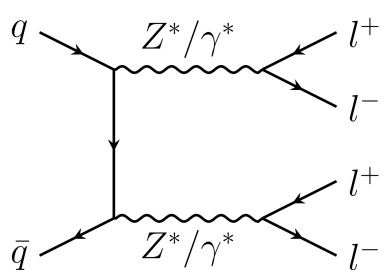


NTGC signal

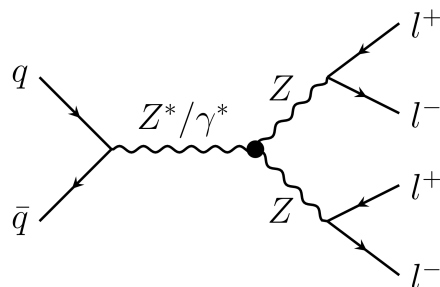


X-section versus invariant mass of 4 lepton system

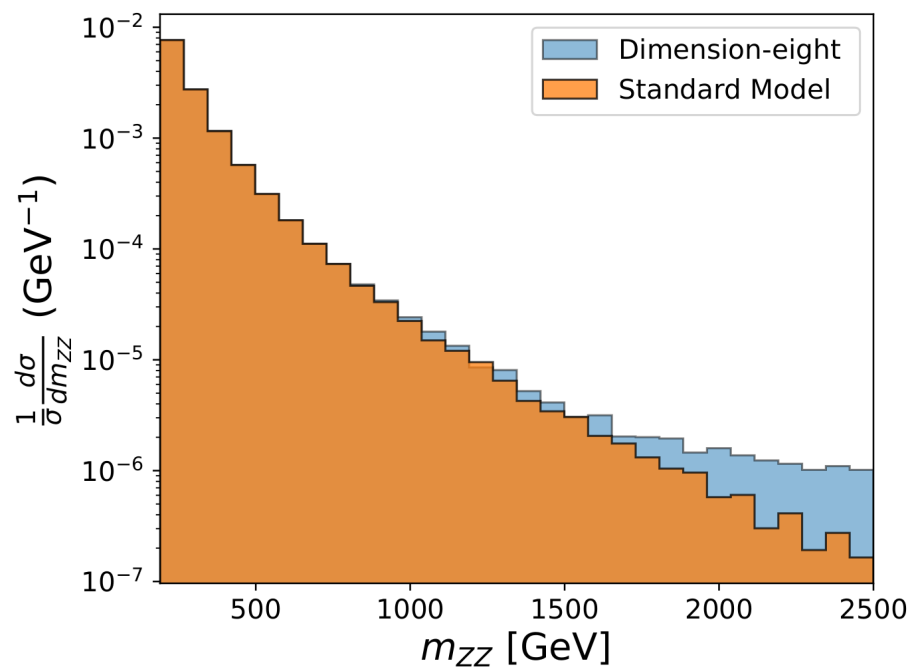
Measuring $ZZ \rightarrow 4l$



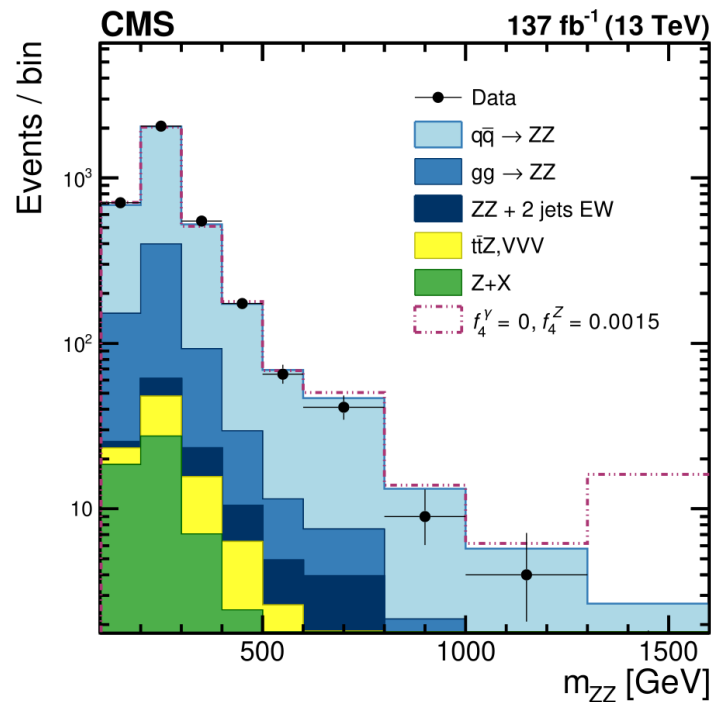
SM background



NTGC signal

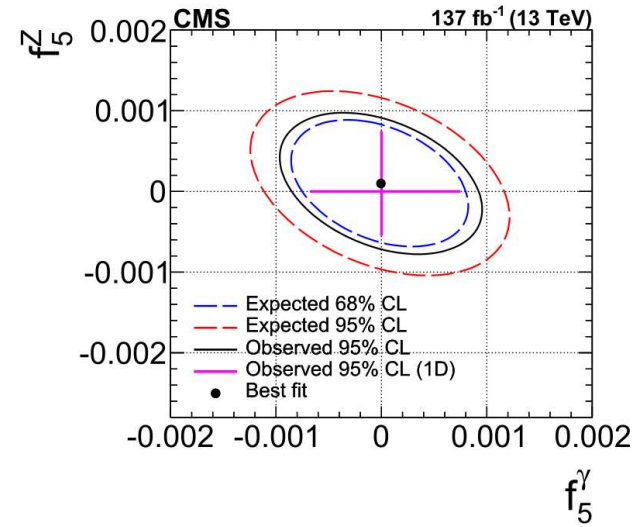


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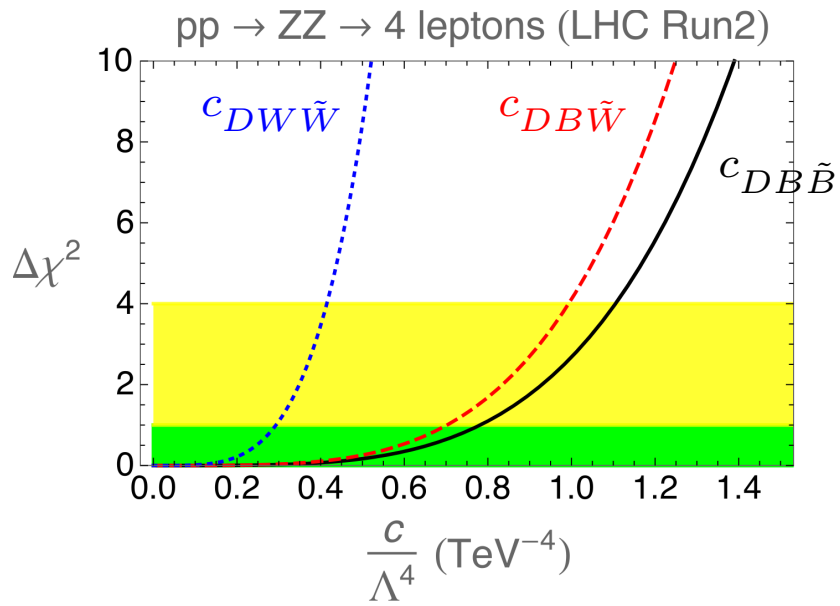
CMS, 2021

Limits

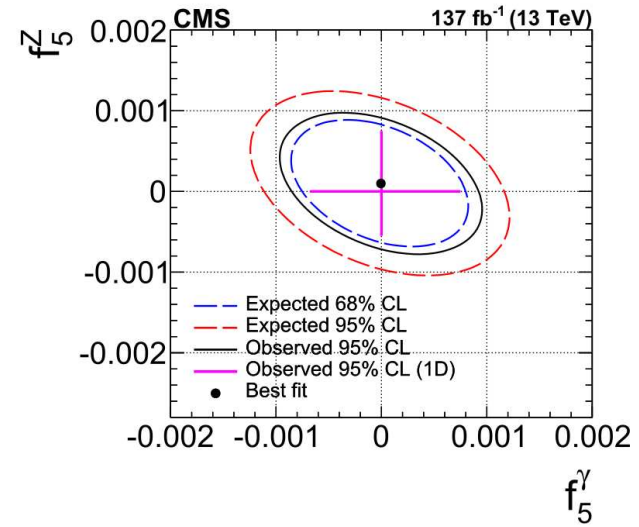


Limits on f_5^Z and f_5^γ from CMS

Limits

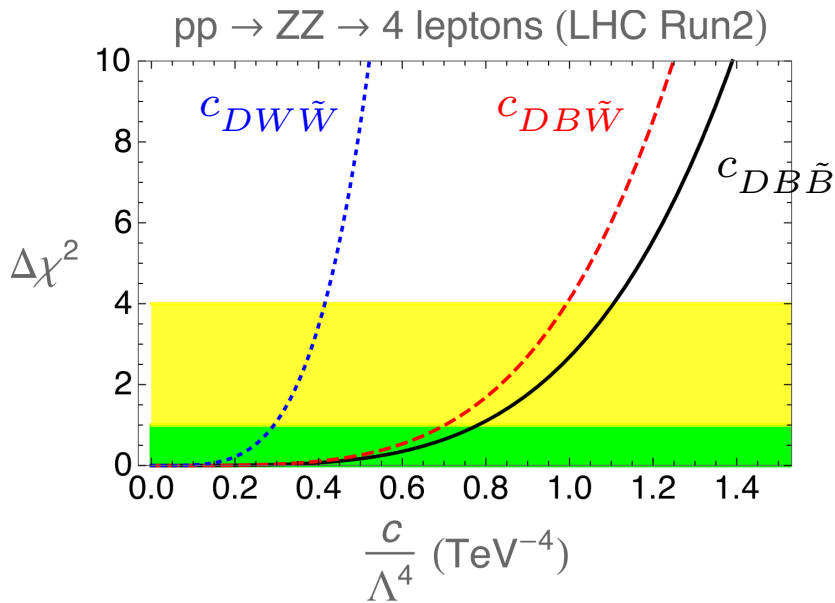


Limits on $d = 8$ operators

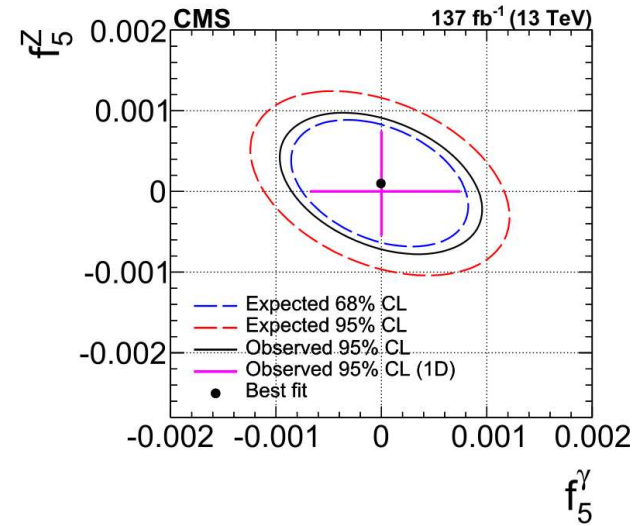


Limits on f_5^Z and f_5^γ from CMS

Limits



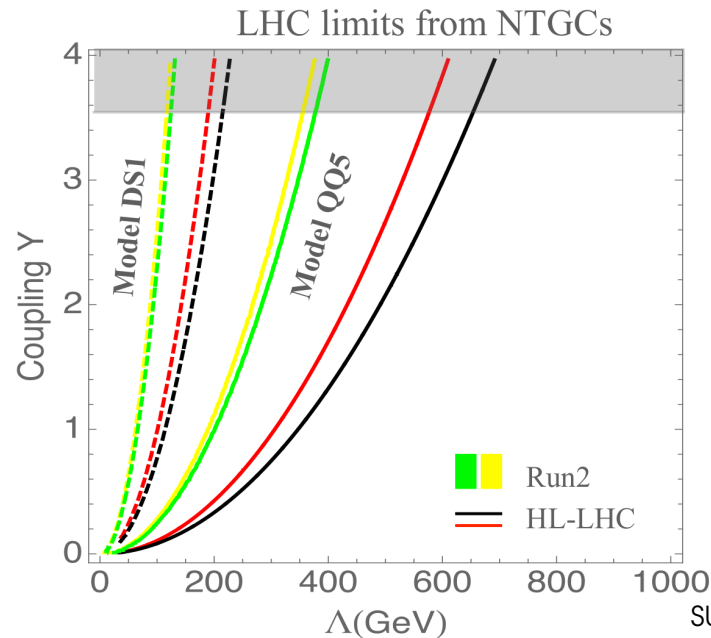
Limits on $d = 8$ operators



Limits on f_5^Z and f_5^γ from CMS

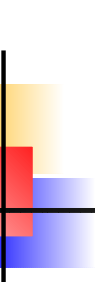
Limits on models
very weak!

$\Lambda \gtrsim 100$ GeV (now)
 \gtrsim (few) 100 GeV for 3/ab



Conclusions

- ⇒ Complete and correct $d = 8$ operator basis for NTGCs found
- ⇒ Correcting existing literature!
- ⇒ Limits derived on operators from experimental searches by ATLAS and CMS: $\Lambda \sim \mathcal{O}(1)$ TeV, depending on operator
- ⇒ Set of fermionic UV models for NTGCs constructed
- ⇒ Limits derived on models, but ...
 - Limits on models very weak! Not even $\Lambda = 100$ GeV in some cases (EFT assumption not valid!)



Backup

General matching formulas

It is actually possible to derive an analytic formula for \tilde{c}_{DAB} for all models:

$$\begin{aligned}\tilde{c}_{DB\tilde{B}} &= \frac{1}{160} (-1)^{(\mathbf{r}_1 \bmod 2)} \operatorname{sgn}(y_2^2 - y_1^2) \sqrt{2\mathbf{r}_1\mathbf{r}_2} \left(y_1^2 + y_2^2 + \frac{4}{3} y_2 y_1 \right), \\ \tilde{c}_{DW\tilde{W}} &= \frac{1}{160} (-1)^{(\mathbf{r}_1 \bmod 2)} \operatorname{sgn}(y_2^2 - y_1^2) \sqrt{2\mathbf{r}_1\mathbf{r}_2} \frac{1}{12} \left[(\mathbf{r}_1^2 - 1) + (\mathbf{r}_2^2 - 1) + \frac{4}{3} (\mathbf{r}_1\mathbf{r}_2 - 2) \right], \\ \tilde{c}_{DW\tilde{B}} &= \frac{1}{48} (-1)^{(\mathbf{r}_1 \bmod 2)} \sqrt{2\mathbf{r}_1\mathbf{r}_2} \frac{1}{12} (y_1 + y_2) \left[(\mathbf{r}_1 + \mathbf{r}_2) + \frac{3}{5} (y_1 - y_2) \right], \\ \tilde{c}_{DB\tilde{W}} &= \tilde{c}_{DW\tilde{B}}.\end{aligned}$$

NTGCs - Lagrangian

The following **effective Lagrangian** generates all NTGC CPC vertices:

Gounaris et al., 1999

$$\mathcal{L}_{\text{NP}}^{\text{CPC}} = \frac{e}{2m_Z^2} \left[f_5^\gamma (\partial^\sigma F_{\sigma\mu}) \tilde{X}^{\mu\beta} Z_\beta + f_5^Z (\partial^\sigma Z_{\sigma\mu}) \tilde{X}^{\mu\beta} Z_\beta \right. \\ \left. - h_3^\gamma (\partial^\sigma F_{\sigma\mu}) \tilde{F}^{\mu\beta} Z_\beta - h_3^Z (\partial^\sigma Z_{\sigma\mu}) \tilde{F}^{\mu\beta} Z_\beta \right. \\ \left. + \frac{h_4^\gamma}{2m_Z^2} [\square (\partial^\sigma F^{\rho\alpha})] \tilde{F}_{\rho\alpha} Z_\sigma + \frac{h_4^Z}{2m_Z^2} [(\square + m_Z^2) (\partial^\sigma Z^{\rho\alpha})] \tilde{F}_{\rho\alpha} Z_\sigma \right],$$

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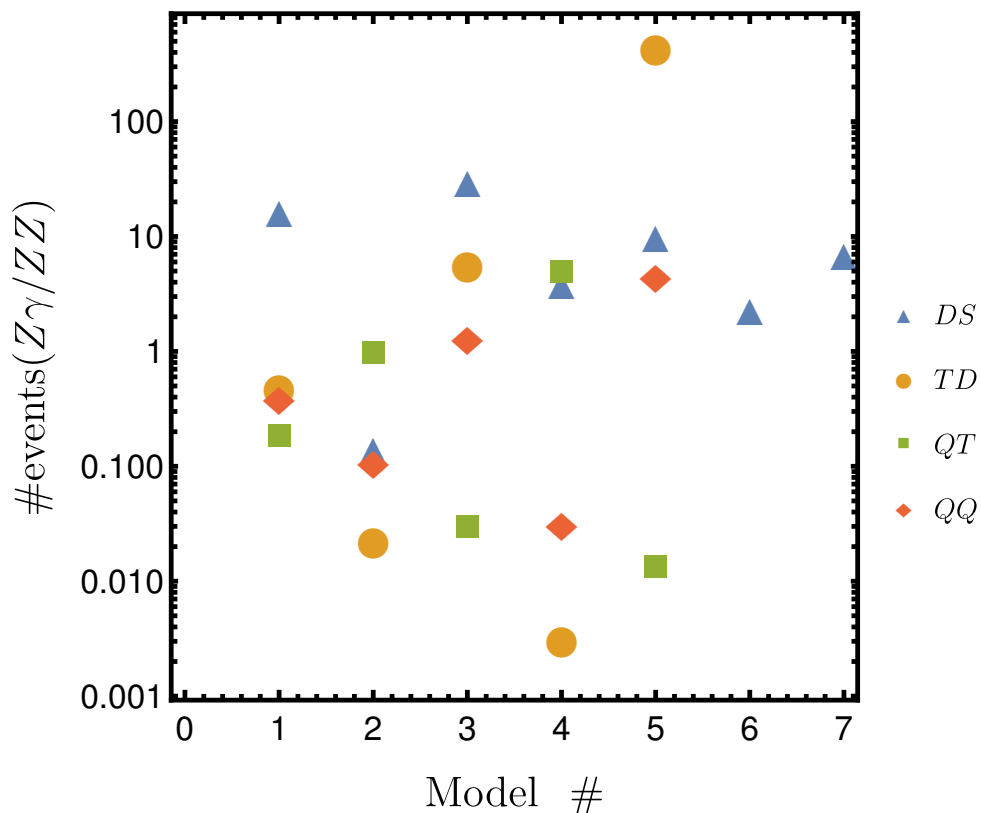
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⇒ What type of **SMEFT operator** can produce this **Lagrangian**?

γZ versus ZZ

Ratio of number of events in γZ and ZZ final states versus model:



Large variation ...

but $> 1/2$ of models have
 $(\gamma Z)/(ZZ) > 1$