

New Constraints on Gauged $U(1)_{L_\mu-L_\tau}$ Models via Z-Z' Mixing

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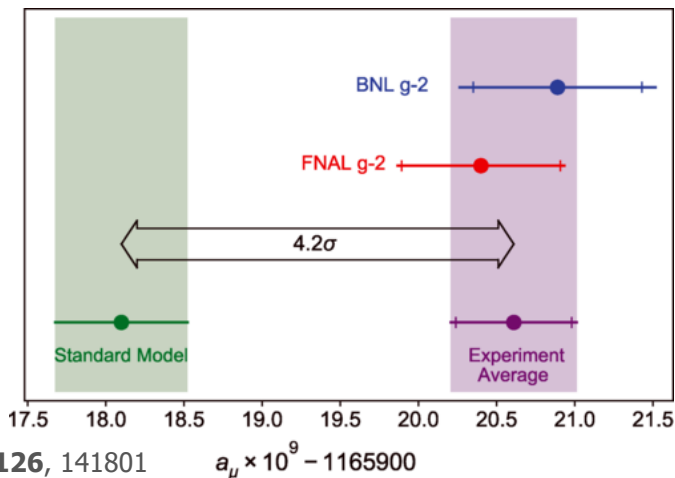
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Based on [arXiv : 2401.17613]

working with K. Asai(ICRR), S. Okawa(KEK), and K. Tsumura(Kyushu U.).

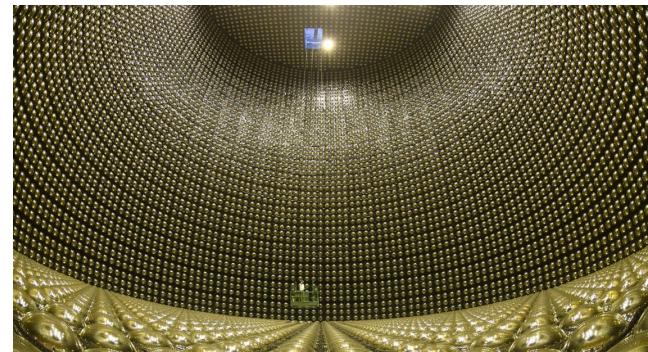
Background

- The discrepancy of muon $g-2$ between the SM and experimental results.
 - $U(1)_{L_\mu-L_\tau}$ gauge models can explain.
- The recent experiments of the neutrino oscillation become more precise.
 - Simple $U(1)_{L_\mu-L_\tau}$ gauge models seem hard to describe the neutrino physics.



Cited from
Phys. Rev. Lett. **126**, 141801

$a_\mu \times 10^9 - 1165900$



Cited from <https://www-sk.icrr.u-tokyo.ac.jp/sk/>

Purpose

- To find the $U(1)_{L_\mu-L_\tau}$ gauge models which are consistent to the latest neutrino experiments.
- To get new (model dependent) constraints on the $U(1)_{L_\mu-L_\tau}$ gauge models.

Cited from NuFIT v5.2

| | Normal Ordering (best fit) | | Inverted Ordering ($\Delta\chi^2 = 6.4$) | | |
|--------------------------|---|---------------------------------|--|---------------------------------|-------------------------------|
| | bfp $\pm 1\sigma$ | 3σ range | bfp $\pm 1\sigma$ | 3σ range | |
| with SK atmospheric data | $\sin^2 \theta_{12}$ | $0.303^{+0.012}_{-0.012}$ | 0.270 \rightarrow 0.341 | $0.303^{+0.012}_{-0.011}$ | 0.270 \rightarrow 0.341 |
| | $\theta_{12}/^\circ$ | $33.41^{+0.75}_{-0.72}$ | 31.31 \rightarrow 35.74 | $33.41^{+0.75}_{-0.72}$ | 31.31 \rightarrow 35.74 |
| | $\sin^2 \theta_{23}$ | $0.451^{+0.019}_{-0.016}$ | 0.408 \rightarrow 0.603 | $0.569^{+0.016}_{-0.021}$ | 0.412 \rightarrow 0.613 |
| | $\theta_{23}/^\circ$ | $42.2^{+1.1}_{-0.9}$ | 39.7 \rightarrow 51.0 | $49.0^{+1.0}_{-1.2}$ | 39.9 \rightarrow 51.5 |
| | $\sin^2 \theta_{13}$ | $0.02225^{+0.00056}_{-0.00059}$ | 0.02052 \rightarrow 0.02398 | $0.02223^{+0.00058}_{-0.00058}$ | 0.02048 \rightarrow 0.02416 |
| | $\theta_{13}/^\circ$ | $8.58^{+0.11}_{-0.11}$ | 8.23 \rightarrow 8.91 | $8.57^{+0.11}_{-0.11}$ | 8.23 \rightarrow 8.94 |
| | $\delta_{CP}/^\circ$ | 232^{+36}_{-26} | 144 \rightarrow 350 | 276^{+22}_{-29} | 194 \rightarrow 344 |
| | $\frac{\Delta m_{21}^2}{10^{-5} \text{ eV}^2}$ | $7.41^{+0.21}_{-0.20}$ | 6.82 \rightarrow 8.03 | $7.41^{+0.21}_{-0.20}$ | 6.82 \rightarrow 8.03 |
| | $\frac{\Delta m_{3\ell}^2}{10^{-3} \text{ eV}^2}$ | $+2.507^{+0.026}_{-0.027}$ | +2.427 \rightarrow +2.590 | $-2.486^{+0.025}_{-0.028}$ | -2.570 \rightarrow -2.406 |

Minimal $U(1)_{L_\mu-L_\tau}$ gauge model

- Fields : SM + three right-handed neutrino N_i + one scalar field.
- Symmetry : SM gauge $\times U(1)_{L_\mu-L_\tau}$ gauge.

| | | | |
|------------------------------|-------------------------------|----------------------|----------------------|
| Lepton | $(\ell_e \ell_\mu \ell_\tau)$ | $(e_R \mu_R \tau_R)$ | $(N_e N_\mu N_\tau)$ |
| $U(1)_{L_\mu-L_\tau}$ charge | (0 +1 -1) | (0 +1 -1) | (0 +1 -1) |

| | | | |
|--------|------------------------------|------------------------------|---------------------------|
| Scalar | Φ_{+1} SU(2) doublet | Φ_{-1} SU(2) doublet | σ SU(2) singlet |
| charge | +1 | -1 | +1 |

Results for Analysis of Neutrino Mass Matrix Structure

- Model independent result set by neutrino mass matrix.
- Each models have their own mass matrix structure.

Our work
(previous work[Phys. Rev. D 99 (2019) 05502])

| Model | Normal ordering | Inverted ordering |
|---------------------------------|---|-----------------------------------|
| $\text{SM} + N_i + \sigma_{+1}$ | Viable in 2σ (Viable at 3σ) | Excluded (Excluded) |
| $\text{SM} + N_i + \Phi_{+1}$ | Excluded (Excluded) | Viable at 3σ (Excluded) |
| $\text{SM} + N_i + \Phi_{-1}$ | Excluded (Excluded) | Excluded (Excluded) |

→ Are there any other constraints on the viable model?

Z-Z' Mixing

- The additional $U(1)_{L_\mu-L_\tau}$ gauge symmetry induces Z-Z' mixing.

$$\mathcal{L}_{\text{gauge}} = -\frac{1}{4}B_{\mu\nu}B^{\mu\nu} - \frac{1}{4}Z'_{\mu\nu}Z'^{\mu\nu} + \frac{1}{2}\frac{\varepsilon}{\cos\theta_W}B_{\mu\nu}Z'^{\mu\nu}$$

$$\mathcal{L}_{\varepsilon_Z} = \frac{1}{2} \begin{pmatrix} Z_\mu & Z'_\mu \end{pmatrix} \begin{pmatrix} 1 & -\varepsilon_Z \\ -\varepsilon_Z & m_{Z'}^2/m_Z^2 \end{pmatrix} \begin{pmatrix} Z^\mu \\ Z'^\mu \end{pmatrix}$$



$$\varepsilon_Z \equiv \frac{m_{Z'}}{m_Z}\delta$$

$$\mathcal{L} \supset Z'_\mu \left(g_{Z'} J_{L_\mu-L_\tau}^\mu + \varepsilon e J_{\text{em}}^\mu + \varepsilon_Z g_Z J_{\text{NC}}^\mu \right)$$

→ G_F and $\sin^2\theta_W$ are changed.

(Now we ignore the kinetic mixing $\varepsilon \sim g_{Z'}/70$ which is much smaller than ε_Z in our interest parameter space.)

Atomic Parity Violation (APV)

- The weak charge of Cs is given by the measurements of APV;

$$Q_W^{\text{exp}}(^{133}_{55}\text{Cs}) = -72.94(43)$$

- The weak charge of Cs based on SM is changed by Z-Z' mixing;

$$Q_W(^{133}_{55}\text{Cs}) \simeq Q_W^{\text{SM}}(^{133}_{55}\text{Cs}) (1 + \delta^2)$$

$$\longrightarrow |\delta|^2 \lesssim 5.7 \times 10^{-3} \quad (90\% \text{ CL})$$

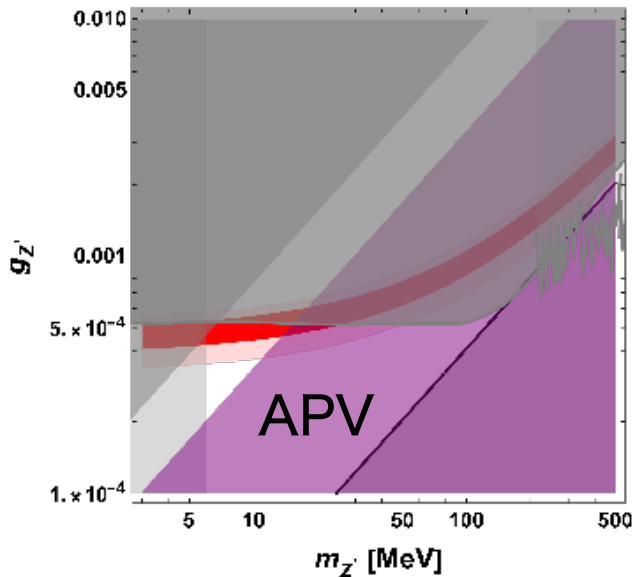
Flavor Changing Meson Decay

- Flavor changing meson decays provide a good probe of a light Z' boson.
- Branching ratio of $K^+ \rightarrow \pi^+ Z'$ is written by;

$$\text{Br}(K^+ \rightarrow \pi^+ Z') \simeq 1.6 \times 10^{-4} |\delta|^2$$

$$\rightarrow |\delta| \lesssim 2.5 \times 10^{-4} \sqrt{\frac{\text{Br}(K^+ \rightarrow \pi^+ Z')_{\text{exp}}}{1 \times 10^{-11}}}$$

Constraint on Model with Φ_{+1}



In this model $\delta = \frac{1}{v} \frac{m_{Z'}}{g_{Z'}}$.

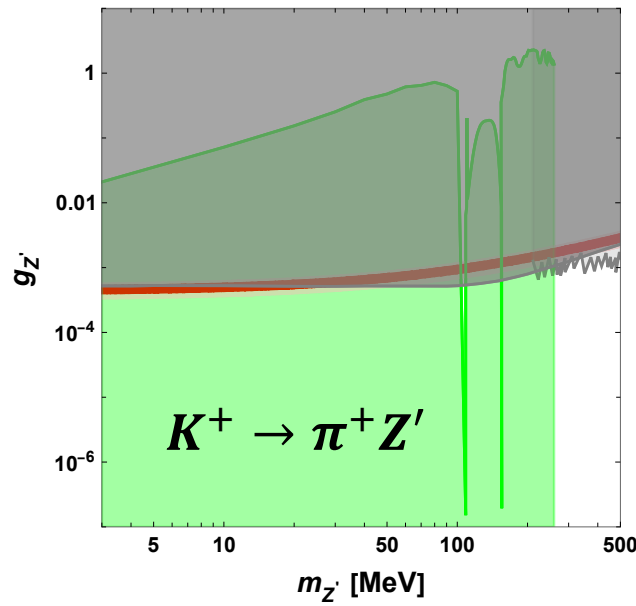


Cs APV :

$$g_{Z'} \gtrsim 5.4 \times 10^{-4} \left(\frac{m_{Z'}}{10 \text{ MeV}} \right)$$

$K^+ \rightarrow \pi^+ Z'$:

$$g_{Z'} \gtrsim 1.6 \times 10^{-1} \sqrt{\frac{1 \times 10^{-11}}{\text{Br}(K^+ \rightarrow \pi^+ Z')_{\text{exp}}}} \left(\frac{m_{Z'}}{10 \text{ MeV}} \right)$$



- The gray shaded region are excluded by the well-known constraints (from BABAR, NA64 μ , white dwarf cooling, and effective number of neutrinos).
- The red region gives the proper correction to muon g-2.
- There is **no region** which gives proper correction to muon g-2.

Constraint on Model with Φ_{+1} and a SU(2) singlet scalar σ_{+1}

In this model, $\delta = \frac{\text{sign}(Q_\Phi)}{1 + \tan^2 \theta} \frac{1}{v} \frac{m_{Z'}}{g_{Z'}}$.

↓ $\tan \theta \equiv \frac{v_\sigma}{v_\Phi}$ ($v_{\Phi(\sigma)}$ means VEV of $\Phi(\sigma)$).

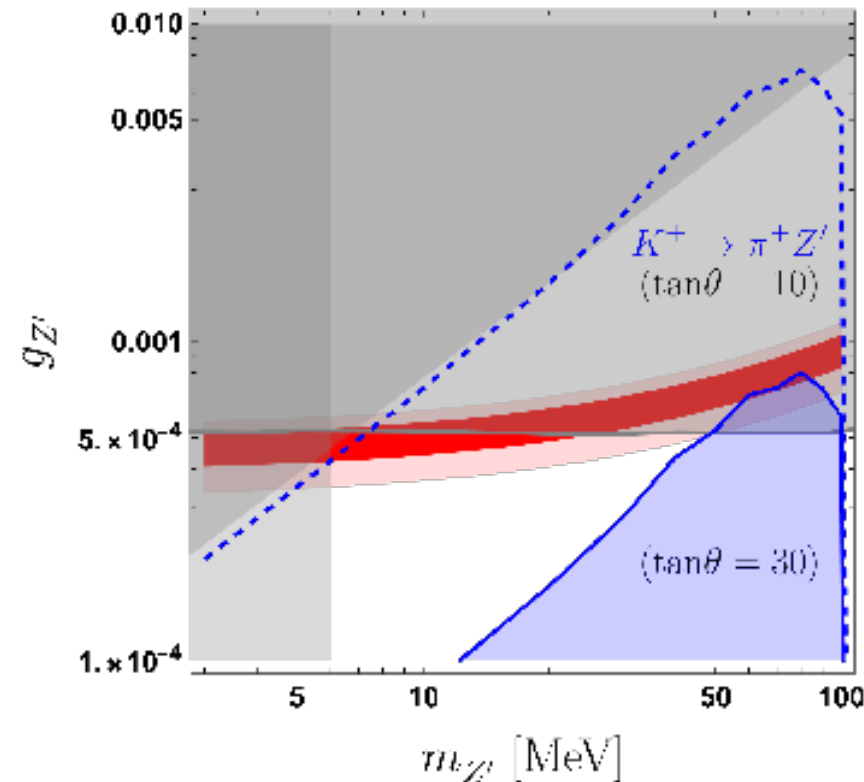
Cs APV : much smaller than the flavor changing meson decay.

$$g_{Z'} \gtrsim \frac{5.4 \times 10^{-4}}{1 + \tan^2 \theta} \left(\frac{m_{Z'}}{10 \text{ MeV}} \right)$$

$K^+ \rightarrow \pi^+ Z'$:

$$g_{Z'} \gtrsim \frac{1.6 \times 10^{-1}}{1 + \tan^2 \theta} \sqrt{\frac{1 \times 10^{-11}}{\text{Br}(K^+ \rightarrow \pi^+ Z')_{\text{exp}}}} \left(\frac{m_{Z'}}{10 \text{ MeV}} \right)$$

- Model gives proper correction to the muon g-2 discrepancy when $\tan \theta \equiv \frac{v_\sigma}{v_2} \gtrsim 10$.



Conclusion

- We revisited the minimal $U(1)_{L_\mu-L_\tau}$ gauge model based on the latest NuFITv5.2 data. As the results, the model with $SU(2)$ doublet scalar Φ_{+1} was viable at 3σ in case of Inverted ordering while the model was excluded in the previous work.
- Considering the constraints from Z - Z' mixing (APV and flavor changing meson decay process), the model with Φ_{+1} is completely excluded in the region which give the explanation to muon $g-2$.
- The model with Φ and σ is viable when $\tan \theta \equiv \frac{v_\sigma}{v_2} \gtrsim 10$.

BACKUP

NuFITv4.0

| | Normal Ordering (best fit) | | Inverted Ordering ($\Delta\chi^2 = 9.3$) | | |
|--------------------------|---|---------------------------------|--|---------------------------------|-------------------------------|
| | bfp $\pm 1\sigma$ | 3σ range | bfp $\pm 1\sigma$ | 3σ range | |
| with SK atmospheric data | $\sin^2 \theta_{12}$ | $0.310^{+0.013}_{-0.012}$ | 0.275 \rightarrow 0.350 | $0.310^{+0.013}_{-0.012}$ | 0.275 \rightarrow 0.350 |
| | $\theta_{12}/^\circ$ | $33.82^{+0.78}_{-0.76}$ | 31.61 \rightarrow 36.27 | $33.82^{+0.78}_{-0.75}$ | 31.62 \rightarrow 36.27 |
| | $\sin^2 \theta_{23}$ | $0.582^{+0.015}_{-0.019}$ | 0.428 \rightarrow 0.624 | $0.582^{+0.015}_{-0.018}$ | 0.433 \rightarrow 0.623 |
| | $\theta_{23}/^\circ$ | $49.7^{+0.9}_{-1.1}$ | 40.9 \rightarrow 52.2 | $49.7^{+0.9}_{-1.0}$ | 41.2 \rightarrow 52.1 |
| | $\sin^2 \theta_{13}$ | $0.02240^{+0.00065}_{-0.00066}$ | 0.02044 \rightarrow 0.02437 | $0.02263^{+0.00065}_{-0.00066}$ | 0.02067 \rightarrow 0.02461 |
| | $\theta_{13}/^\circ$ | $8.61^{+0.12}_{-0.13}$ | 8.22 \rightarrow 8.98 | $8.65^{+0.12}_{-0.13}$ | 8.27 \rightarrow 9.03 |
| | $\delta_{CP}/^\circ$ | 217^{+40}_{-28} | 135 \rightarrow 366 | 280^{+25}_{-28} | 196 \rightarrow 351 |
| | $\frac{\Delta m_{21}^2}{10^{-5} \text{ eV}^2}$ | $7.39^{+0.21}_{-0.20}$ | 6.79 \rightarrow 8.01 | $7.39^{+0.21}_{-0.20}$ | 6.79 \rightarrow 8.01 |
| | $\frac{\Delta m_{3\ell}^2}{10^{-3} \text{ eV}^2}$ | $+2.525^{+0.033}_{-0.031}$ | +2.431 \rightarrow +2.622 | $-2.512^{+0.034}_{-0.031}$ | -2.606 \rightarrow -2.413 |

From <http://www.nu-fit.org/?q=node/177>

Neutrino Mass Matrix

- In general,

$$\mathcal{M}_{\nu_L} = U_{\text{PMNS}} \text{diag}(m_1 \ m_2 \ m_3) U_{\text{PMNS}}^T \equiv \mathcal{M}_{\nu_L}^{\text{gen}}.$$

$$U_{\text{PMNS}} \equiv \begin{pmatrix} V_{11} & V_{12} & V_{13} \\ V_{21} & V_{22} & V_{23} \\ V_{31} & V_{32} & V_{33} \end{pmatrix} \begin{pmatrix} 1 & & \\ & e^{\frac{i\alpha_2}{2}} & \\ & & e^{\frac{i\alpha_3}{2}} \end{pmatrix}.$$

m_i :light neutrino mass α_i :Majorana phase

V_{ij} :matrix component including mixing angles and CP phase

- Through the seesaw mechanism

$$\mathcal{M}_{\nu_L} \simeq -\mathcal{M}_D \mathcal{M}_R^{-1} \mathcal{M}_D^T.$$

→ Some equations arise by comparing these.

Two Zero Texture (Minor) Structure Mass Matrix

- Classification of structures;

$$\mathbf{B}_3: \begin{pmatrix} * & 0 & * \\ 0 & 0 & * \\ * & * & * \end{pmatrix}, \mathbf{B}_4: \begin{pmatrix} * & * & 0 \\ * & * & * \\ 0 & * & 0 \end{pmatrix}, \mathbf{C}: \begin{pmatrix} * & * & * \\ * & 0 & * \\ * & * & 0 \end{pmatrix}$$

- Thorough the seesaw mechanism, the neutrino mass matrix (or its inverted one) often has such structure.
 - Two components of \mathcal{M}_{ν_L} are zero → Two zero texture
 - Two components of $\mathcal{M}_{\nu_L}^{-1}$ are zero → Two zero minor

The mass matrix with such structures
give us two equations. → Predictions

Light Neutrino Mass

$$m_3 = \sqrt{\frac{\Delta m_{31}^2}{1 - \frac{1}{|R_3(\theta_{12}, \theta_{13}, \theta_{23}, \delta_{CP})|^2}}}$$

Neutrino mass
in case of NO.

$$m_1 = \sqrt{m_3^2 - \Delta m_{31}^2}$$

$$m_2 = \sqrt{m_1^2 + \Delta m_{21}^2} = \sqrt{m_3^2 + \Delta m_{21}^2 - \Delta m_{31}^2}$$

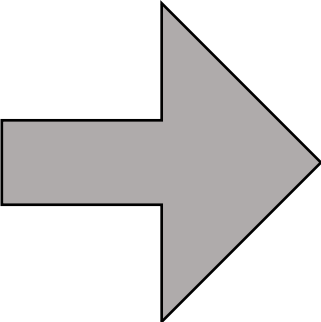
- These masses can be described in terms of θ_{23} .

($\theta_{12}, \theta_{13}, \Delta m_{21}^2, \Delta m_{31}^2$ are fixed as best fit value of NuFITv5.2)

Analysis of SM+N_i+Φ₊₁ Model

- **B₃ texture** : $(\mathcal{M}_{\nu L})_{[1,2],[2,2]} = 0$

$$\left\{ \begin{array}{l} (\mathcal{M}_\nu^{\text{gen}})_{12} = m_1 V_{11} V_{21} + m_2 e^{i\alpha_2} V_{12} V_{22} + m_3 e^{i\alpha_3} V_{13} V_{23} = 0 \quad (= (\mathcal{M}_\nu)_{12}). \\ (\mathcal{M}_\nu^{\text{gen}})_{22} = m_1 V_{21}^2 + m_2 e^{i\alpha_2} V_{22}^2 + m_3 e^{i\alpha_3} V_{23}^2 = 0 \quad (= (\mathcal{M}_\nu)_{22}). \end{array} \right.$$

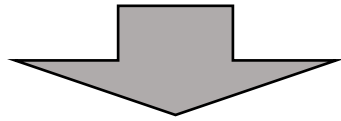


$$\left\{ \begin{array}{l} e^{i\alpha_2} \equiv \frac{m_1}{m_2} R_2(\theta_{12}, \theta_{13}, \theta_{23}, \delta) \equiv \frac{R_2}{|R_2|} \\ e^{i\alpha_3} \equiv \frac{m_1}{m_3} R_3(\theta_{12}, \theta_{13}, \theta_{23}, \delta) \equiv \frac{R_3}{|R_3|} \end{array} \right.$$

V_{ij} : components of PMNS matrix , θ_{ij} : mixing angle, δ : CP phase,
 α : Majorana phase

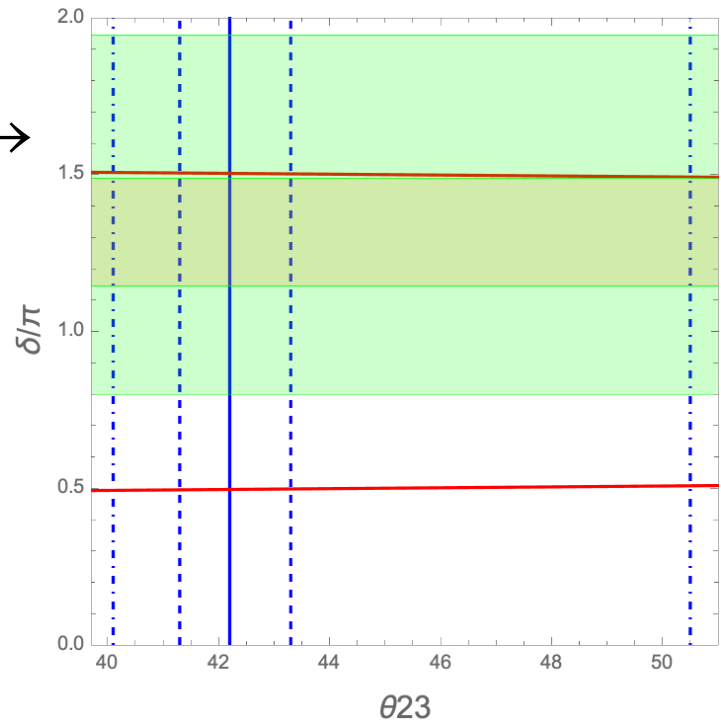
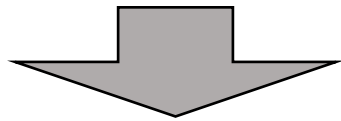
- To rewrite mass-squared difference in case of Normal ordering(NO)

$$\left\{ \begin{array}{l} \Delta m_{31}^2 = m_3^2 - m_1^2 = m_1^2 (|R_3|^2 - 1) \\ \Delta m_{21}^2 = m_2^2 - m_1^2 = m_1^2 (|R_2|^2 - 1) \end{array} \right.$$



$$(|R_2|^2 - 1) = \frac{\Delta m_{21}^2}{\Delta m_{31}^2} (|R_3|^2 - 1) \rightarrow$$

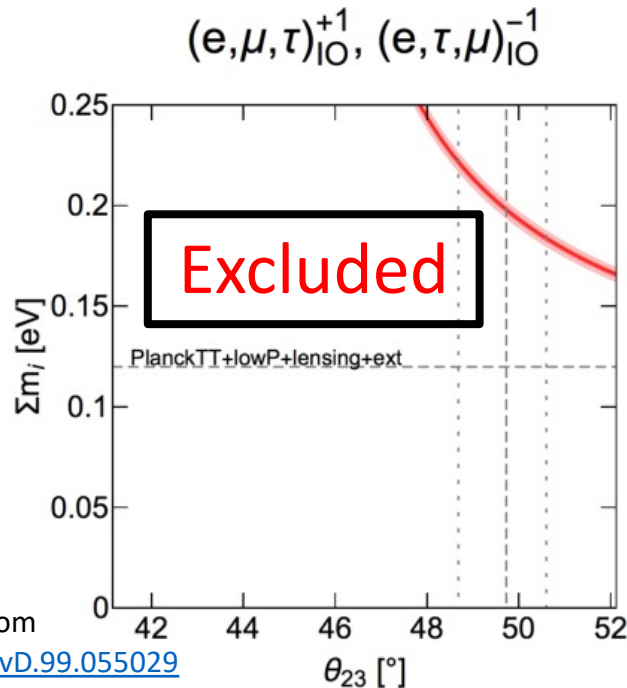
By fixing $\theta_{12}, \theta_{13}, \Delta m_{21}^2, \Delta m_{31}^2$ as the best-fit value of NuFITv5.2, θ_{23} -dependence of δ are found.



Neutrino mass and Majorana phase can be written by θ_{23} !

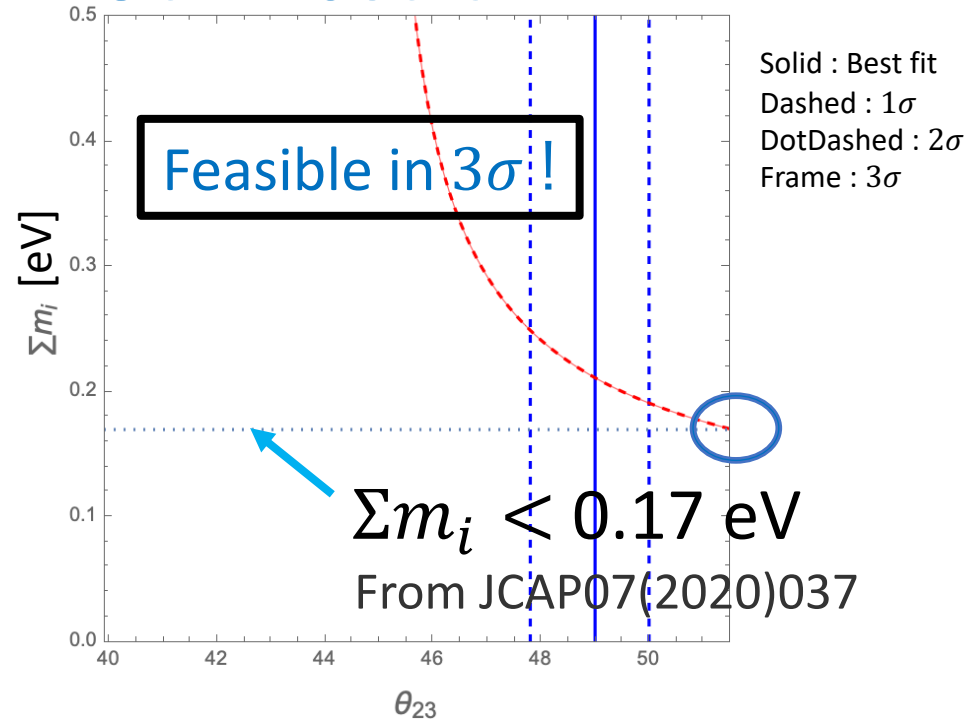
Result of Analysis (B_3 Texture)

- Previous Work



Cited from
[PhysRevD.99.055029](https://arxiv.org/abs/1905.05502)
 Based on NuFIT4.0

- Our Result

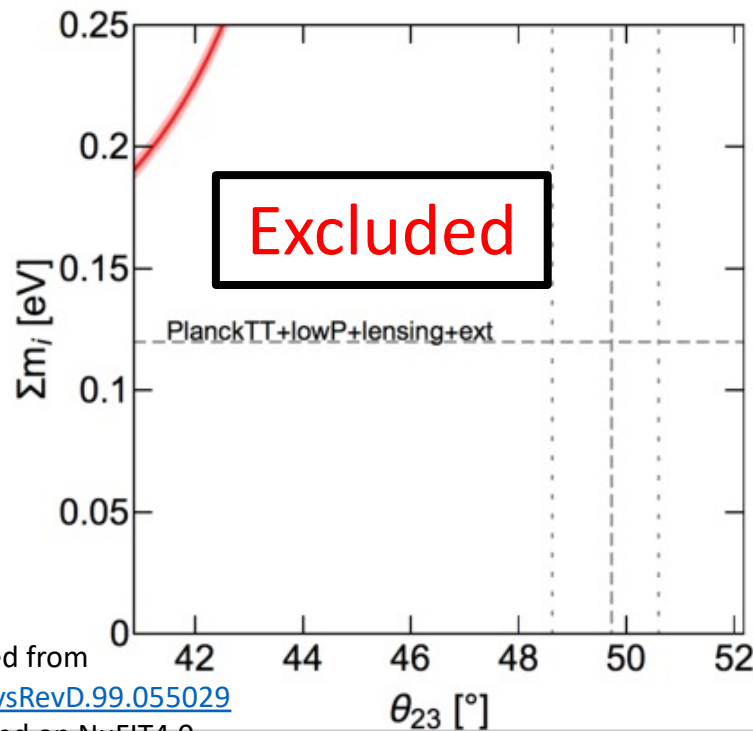


- B_3 -type mass matrix in Inverted ordering is revived.
- The range of θ_{23} shift to left in the latest NuFITv5.2.
- The mass sum constraint is relaxed because of being had considered mass ordering in the analysis.

Analysis Result (NO)

- Previous Work

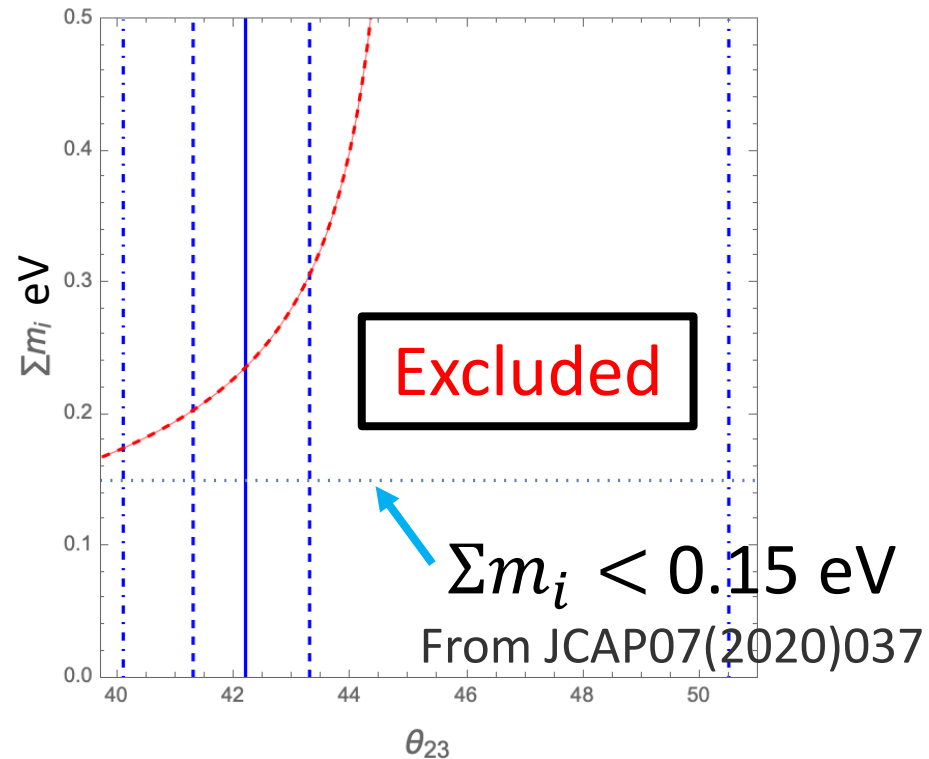
$$(e, \mu, \tau)_{NO}^{+1}, (e, \tau, \mu)_{NO}^{-1}$$



Cited from
[PhysRevD.99.055029](https://arxiv.org/abs/1907.08787)
 Based on NuFIT4.0

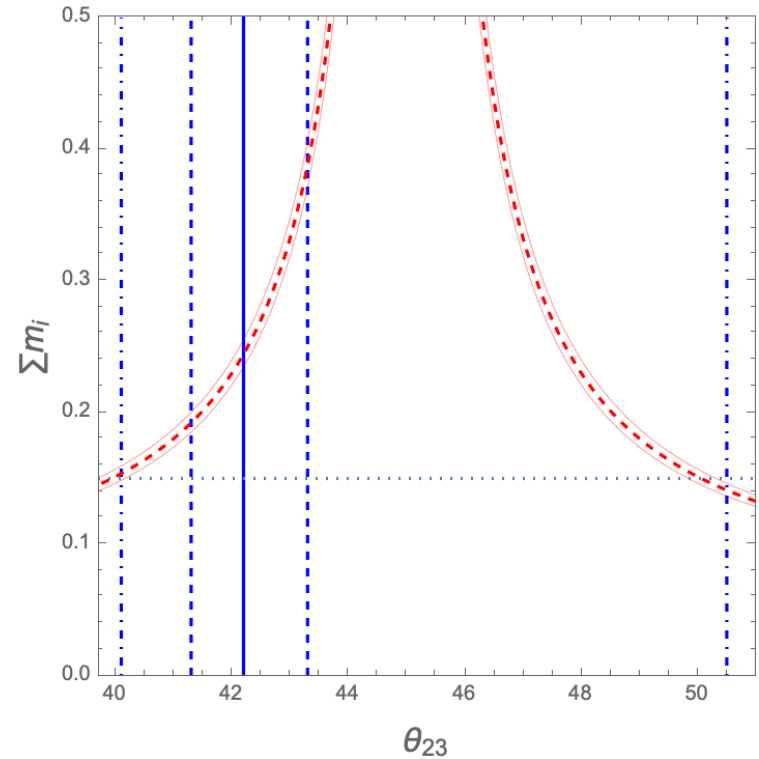
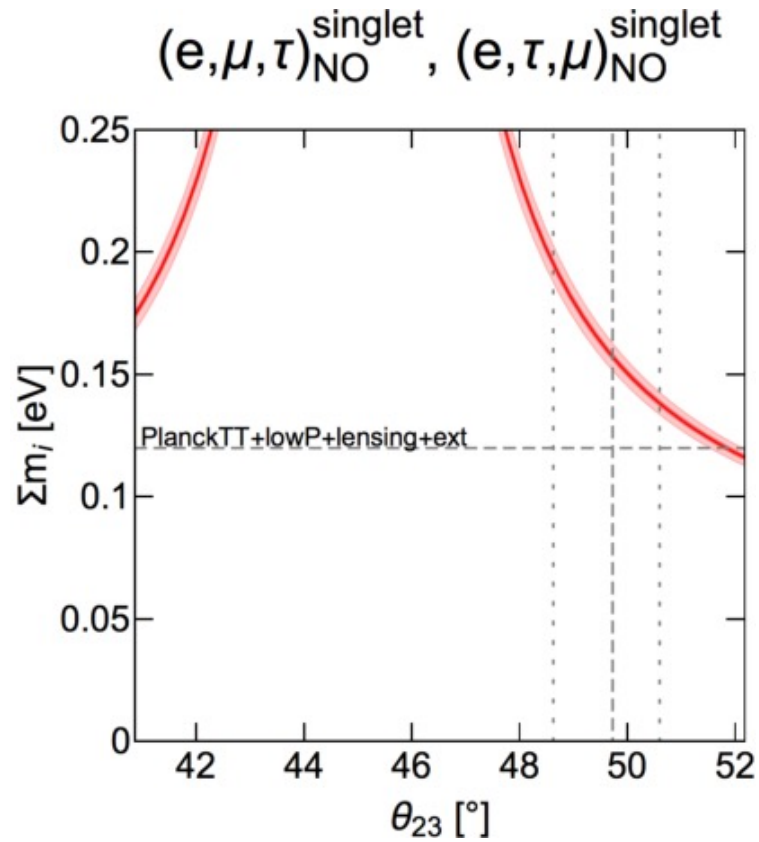
- Our Result

Solid : Best fit
 Dashed : 1σ
 DotDashed : 2σ
 Frame : 3σ

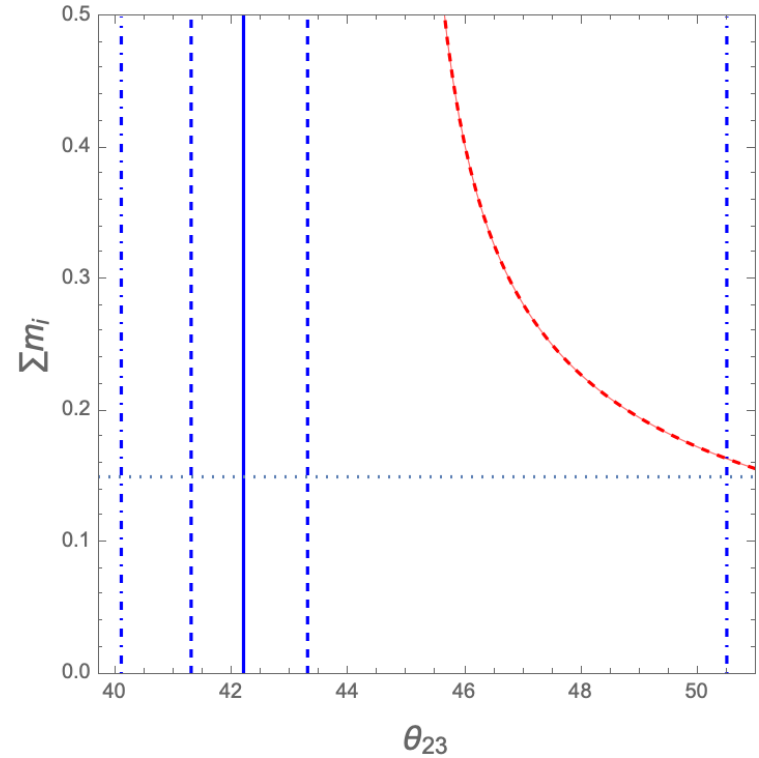
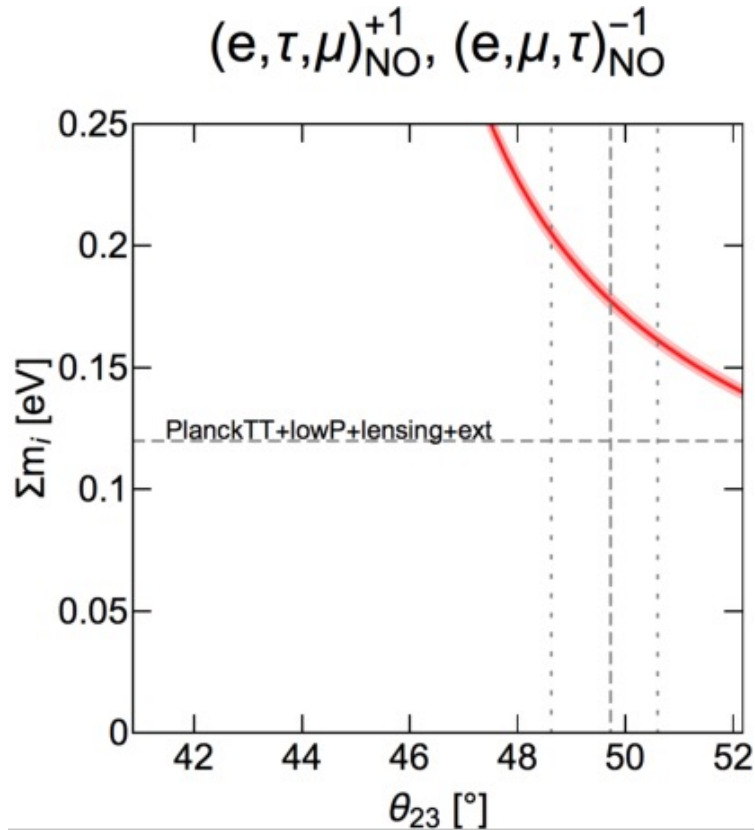


- The range of θ_{23} shift to left in the latest NuFITv5.2.
- The mass sum constraint are relaxed by considering mass ordering.

C Minor (NO)

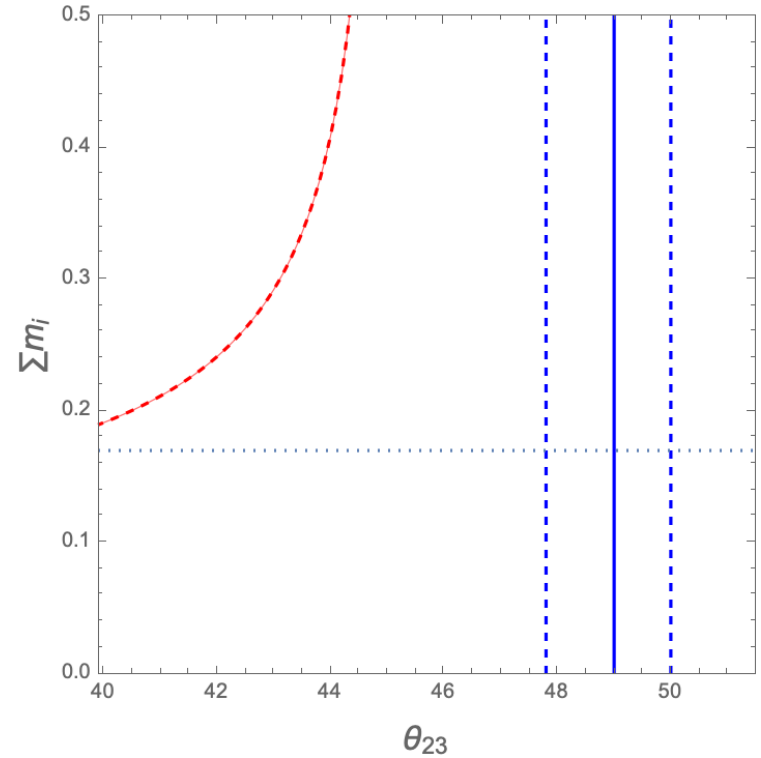
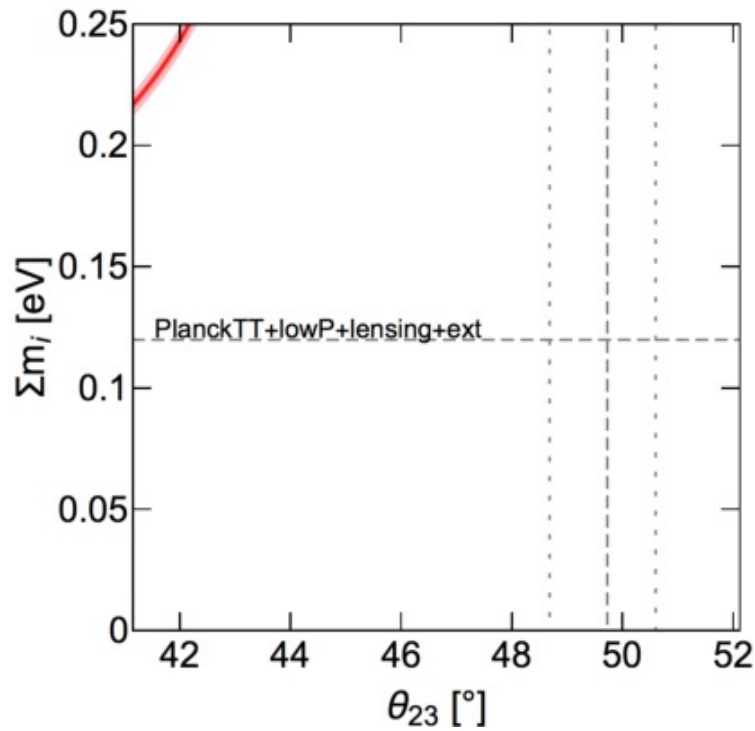


B4 Texture (NO)

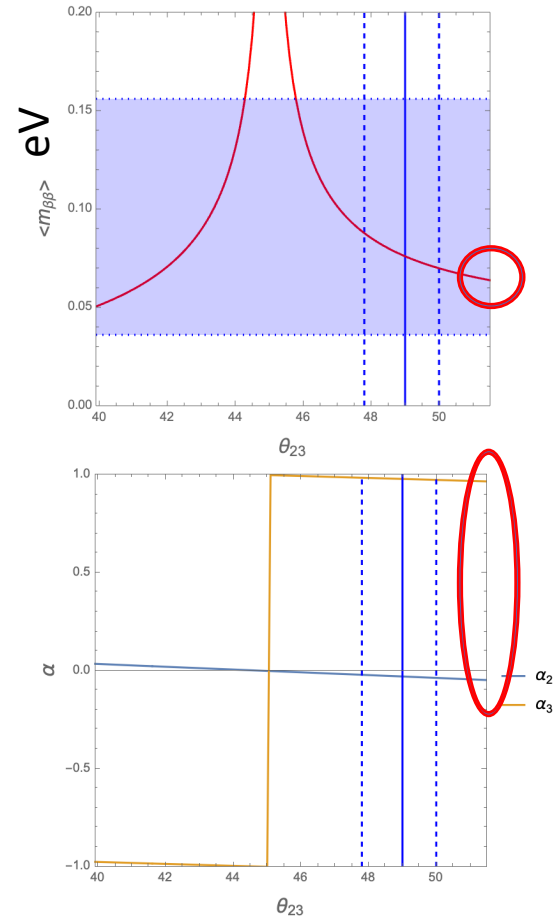
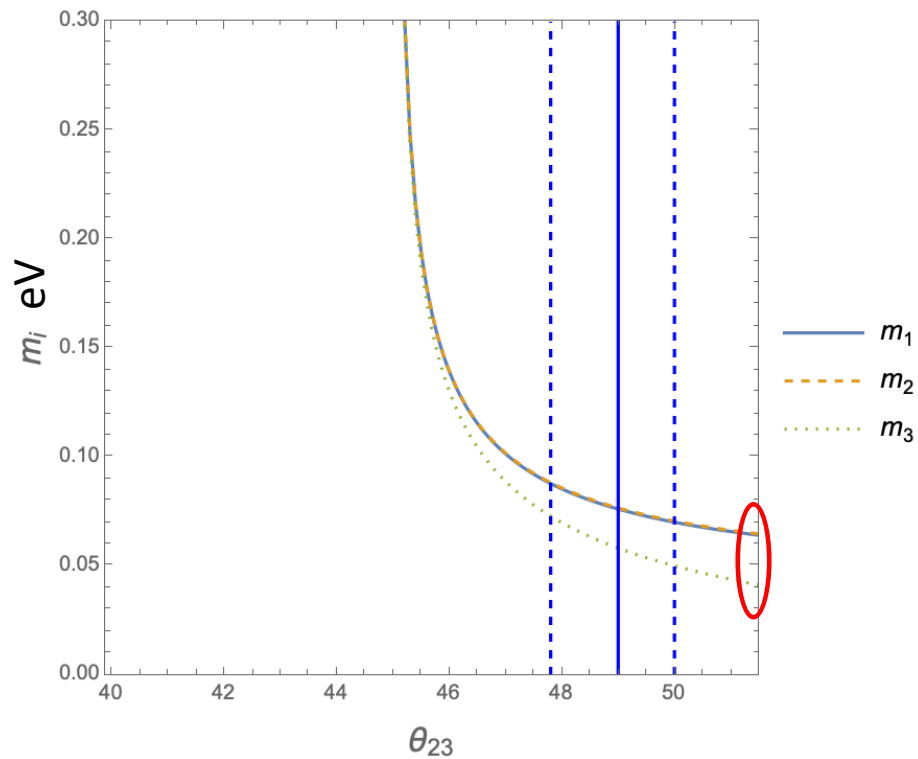


B4 Texture (IO)

$(e, \tau, \mu)_{IO}^{+1}, (e, \mu, \tau)_{IO}^{-1}$



Result of analysis 2



| | m_1 [eV] | m_2 [eV] | m_3 [eV] | α_2/π | α_3/π | $\langle m_{\beta\beta} \rangle$ [eV] |
|---|------------|------------|------------|----------------|----------------|---------------------------------------|
| \mathbf{B}_3 texture (IO) | 0.064 | 0.065 | 0.041 | -0.05 | 0.96 | 0.064 |