



Complete One-loop Renormalization-group Equations in the Seesaw Effective Field Theories

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Based on JHEP 05 (2023) 044 in collaboration with Y.-L. Wang and S. Zhou

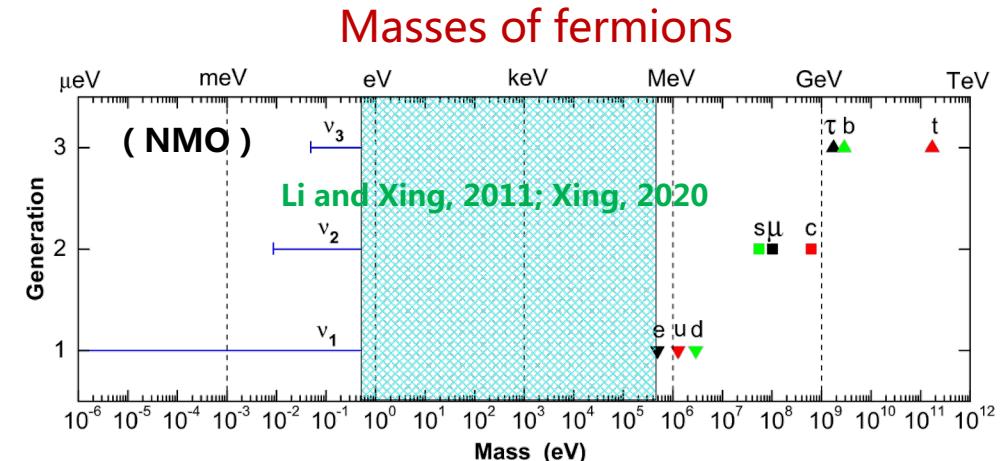
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Seesaw Mechanisms

Global-fit results for neutrino oscillation parameters:

with SK atmospheric data	Esteban et al., 2020		NuFIT 5.3 (2024)	
	Normal Ordering (best fit)		Inverted Ordering ($\Delta\chi^2 = 9.1$)	
	bfp $\pm 1\sigma$	3σ range	bfp $\pm 1\sigma$	3σ range
$\sin^2 \theta_{12}$	$0.307^{+0.012}_{-0.011}$	$0.275 \rightarrow 0.344$	$0.307^{+0.012}_{-0.011}$	$0.275 \rightarrow 0.344$
$\theta_{12}/^\circ$	$33.67^{+0.73}_{-0.71}$	$31.61 \rightarrow 35.94$	$33.67^{+0.73}_{-0.71}$	$31.61 \rightarrow 35.94$
$\sin^2 \theta_{23}$	$0.454^{+0.019}_{-0.016}$	$0.411 \rightarrow 0.606$	$0.568^{+0.016}_{-0.021}$	$0.412 \rightarrow 0.611$
$\theta_{23}/^\circ$	$42.3^{+1.1}_{-0.9}$	$39.9 \rightarrow 51.1$	$48.9^{+0.9}_{-1.2}$	$39.9 \rightarrow 51.4$
$\sin^2 \theta_{13}$	$0.02224^{+0.00056}_{-0.00057}$	$0.02047 \rightarrow 0.02397$	$0.02222^{+0.00069}_{-0.00057}$	$0.02049 \rightarrow 0.02420$
$\theta_{13}/^\circ$	$8.58^{+0.11}_{-0.11}$	$8.23 \rightarrow 8.91$	$8.57^{+0.13}_{-0.11}$	$8.23 \rightarrow 8.95$
$\delta_{\text{CP}}/^\circ$	232^{+39}_{-25}	$139 \rightarrow 350$	273^{+24}_{-26}	$195 \rightarrow 342$
$\frac{\Delta m_{21}^2}{10^{-5} \text{ eV}^2}$	$7.41^{+0.21}_{-0.20}$	$6.81 \rightarrow 8.03$	$7.41^{+0.21}_{-0.20}$	$6.81 \rightarrow 8.03$
$\frac{\Delta m_{3\ell}^2}{10^{-3} \text{ eV}^2}$	$+2.505^{+0.024}_{-0.026}$	$+2.426 \rightarrow +2.586$	$-2.487^{+0.027}_{-0.024}$	$-2.566 \rightarrow -2.407$



Flavor mixing of fermions

$$\begin{pmatrix} \textcolor{red}{\square} & \textcolor{red}{\square} & \textcolor{red}{\square} \\ \textcolor{red}{\square} & \textcolor{red}{\square} & \textcolor{red}{\square} \\ \textcolor{red}{\square} & \textcolor{red}{\square} & \textcolor{red}{\square} \end{pmatrix} \quad \begin{pmatrix} \textcolor{blue}{\square} & \textcolor{blue}{\square} & \textcolor{blue}{\square} \\ \textcolor{blue}{\square} & \textcolor{blue}{\square} & \textcolor{blue}{\square} \\ \textcolor{blue}{\square} & \textcolor{blue}{\square} & \textcolor{blue}{\square} \end{pmatrix}$$

U_{PMNS} V_{CKM}

Xing, 2020

The origin of neutrino masses is quite different from that of charged fermions

Seesaw mechanisms extending the Standard Model with **fermion singlets**, **scalar triplet**, or **fermion triplets**

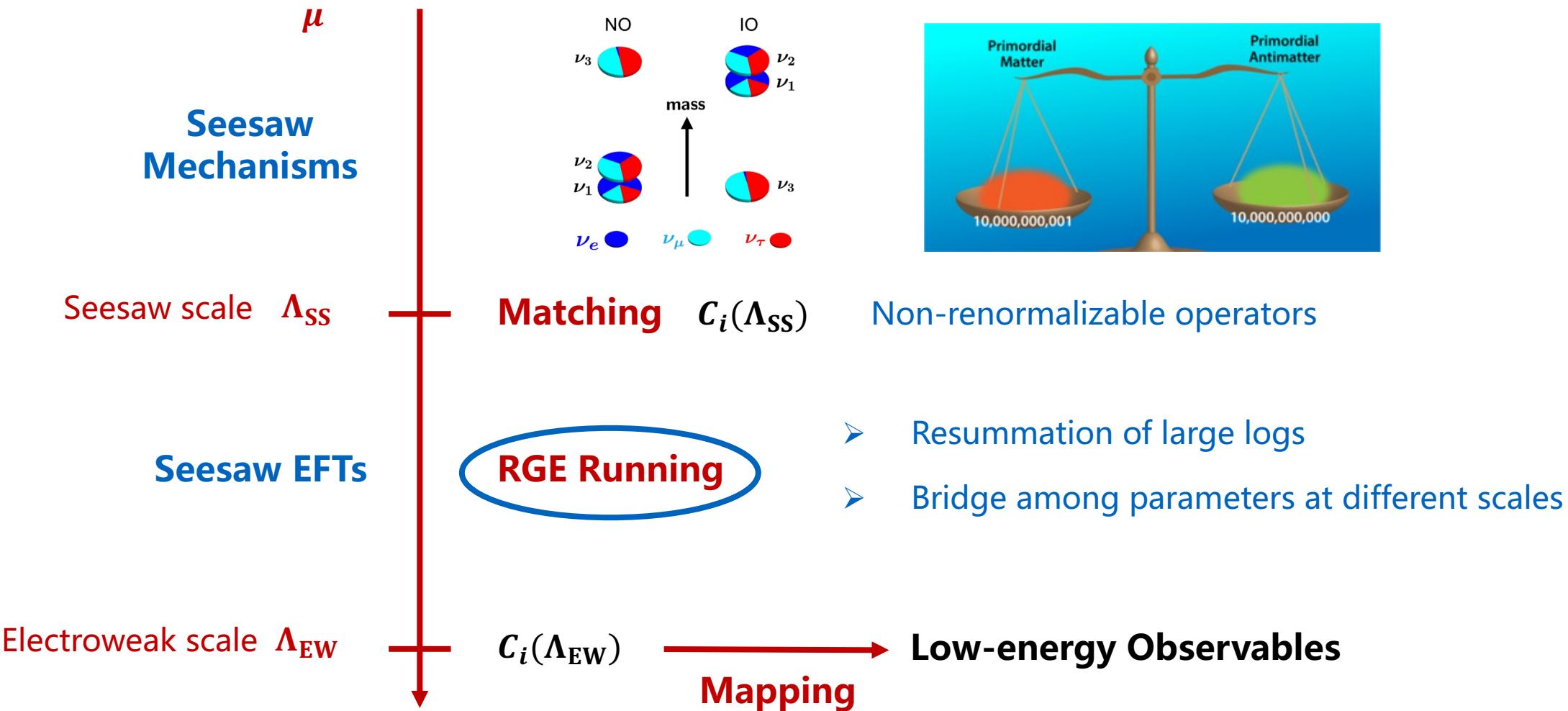
Minkowski, 1977; Yanagida, 1979; ... Konetschny and Kummer, 1977; ... Foot et al., 1989; Ma, 1998

■ The **simplest** and the **most natural** ways to explain **tiny neutrino masses**

Fukugita and Yanagida, 1986

■ A bonus is to account for the matter-antimatter asymmetry of the Universe via **leptogenesis**

Seesaw Effective Field Theories



To achieve the **complete** one-loop RGEs up to $O(1/\Lambda_{\text{SS}}^2)$ in the **seesaw EFTs** induced by **seesaw mechanisms**

Seesaw Effective Field Theories

The type-I seesaw mechanism extending the SM with **three singlet right-handed neutrinos**

$$\mathcal{L}_{\text{SS}} = \mathcal{L}_{\text{SM}} + \overline{N_R} i\cancel{\partial} N_R - \left(\frac{1}{2} \overline{N_R^c} M_R N_R + \overline{\ell_L} Y_\nu \tilde{H} N_R + \text{h.c.} \right)$$

Minkowski, 1977; Yanagida, 1979;
Gell-Mann et al., 1979; Glashow, 1980;
Mohapatra, Senjanovic, 1980

Integrating out **heavy** right-handed neutrinos at the **tree level** (i.e., the tree-level matching)



The tree-level seesaw EFT up to $\mathcal{O}(1/\Lambda_{\text{SS}}^2)$: Broncano, Gavela and Jenkins, 2003a; 2003b; Abada et al, 2007

$$\mathcal{L}_{\text{SEFT}} = \mathcal{L}_{\text{SM}} + \frac{1}{2} \left(C_5^{\alpha\beta} \mathcal{O}_{\alpha\beta}^{(5)} + \text{h.c.} \right) + C_{H\ell}^{(1)\alpha\beta} \mathcal{O}_{H\ell}^{(1)\alpha\beta} + C_{H\ell}^{(3)\alpha\beta} \mathcal{O}_{H\ell}^{(3)\alpha\beta}$$

- Dim-5 operator

$$\mathcal{O}_{\alpha\beta}^{(5)} = \overline{\ell_{\alpha L}} \tilde{H} \tilde{H}^T \ell_{\beta L}^c$$

The Weinberg operator

S. Weinberg, 1979



Neutrino masses

- Dim-6 operators

The Warsaw basis
Grzadkowski, 2010

$$\mathcal{O}_{\alpha\beta}^{(1)} = (\overline{\ell_{\alpha L}} \gamma^\mu \ell_{\beta L}) (H^\dagger i \overset{\leftrightarrow}{D}_\mu H)$$

$$\mathcal{O}_{\alpha\beta}^{(3)} = (\overline{\ell_{\alpha L}} \gamma^\mu \sigma^I \ell_{\beta L}) (H^\dagger i \overset{\leftrightarrow}{D}_\mu^I H)$$



Unitarity violation of the lepton flavor mixing

Broncano, Gavela and Jenkins, 2003a; 2003b; 2005;
Antusch et al., 2006; Abada et al., 2007

The corresponding **Wilson coefficients** at the matching scale $\mu_M \sim \Lambda_{\text{SS}} = \mathcal{O}(M_R)$

$$C_5(\mu_M) = Y_\nu M_R^{-1} Y_\nu^T$$

$$C_{H\ell}^{(1)}(\mu_M) = -C_{H\ell}^{(3)}(\mu_M) = \frac{1}{4} Y_\nu M_R^{-2} Y_\nu^\dagger$$

One-loop RGEs in the Seesaw EFT

The general structure of the RGEs up to $\mathcal{O}(1/\Lambda_{\text{SS}}^2)$

$$16\pi^2 \mu \frac{dC_i^{(5)}}{d\mu} = \gamma'_{ij} C_j^{(5)}$$

Chankowski, Pluciennik, 1993;
Babu, Leung, Pantaleone, 1993;
Antusch et al., 2001

$$16\pi^2 \mu \frac{dC_i^{(6)}}{d\mu} = \gamma_{ij} C_j^{(6)} + \widehat{\gamma}_{jk}^i C_j^{(5)} C_k^{(5)}$$

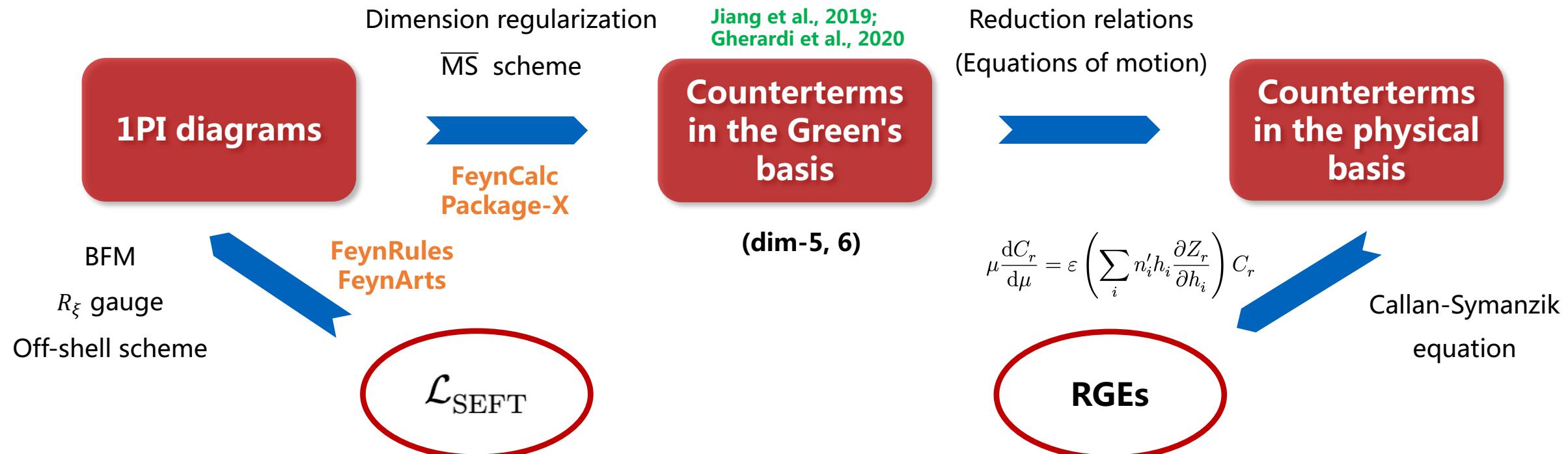
Jenkins et al., 2013; 2014;
Alonso et al., 2014a; 2014b

Wang, DZ, and Zhou, 2023

Incomplete and
not fully correct!!!

Broncano, Gavela, Jenkins, 2005;
Davidson, Gorbahn, Leak, 2018

The procedure for calculations:



Crosscheck has been done by taking advantage of the package **Matchmakereft**
Carmona et al., 2022

One-loop RGEs in the Seesaw EFT

The general structure of the RGEs up to $\mathcal{O}(1/\Lambda_{\text{SS}}^2)$

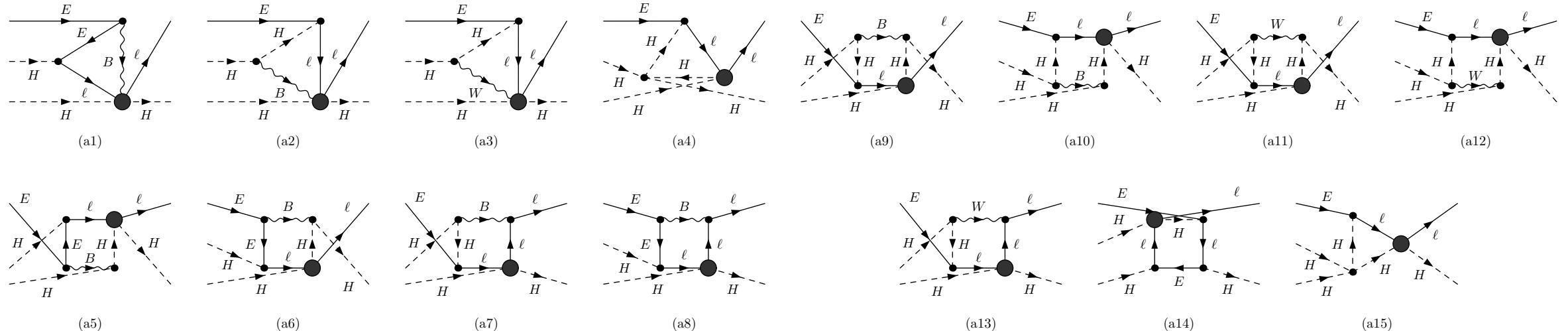
Wang, DZ, and Zhou, 2023

$$16\pi^2 \mu \frac{dC_i^{(5)}}{d\mu} = \gamma'_{ij} C_j^{(5)}$$

$$16\pi^2 \mu \frac{dC_i^{(6)}}{d\mu} = \gamma_{ij} C_j^{(6)} + \widehat{\gamma}_{jk}^i C_j^{(5)} C_k^{(5)}$$

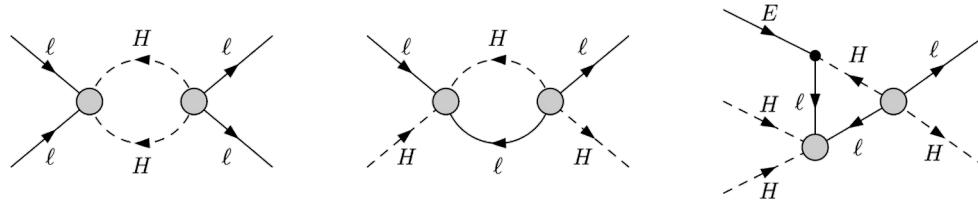
Two types of contributions:

- Single insertion of the dim-5 or dim-6 operators

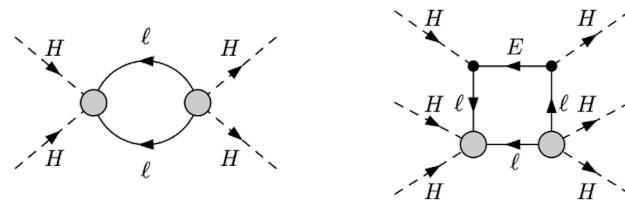


An example for diagrams renormalizing $\mathcal{O}_{eH} = \overline{\ell_L} H E_R (H^\dagger H)$ vertex

- Double insertions of the dim-5 operator



Flavor dependent



Flavor independent

All possible diagrams

One-loop RGEs in the Seesaw EFT

■ The SM couplings:

Wang, DZ, and Zhou, 2023

$$16\pi^2 \mu \frac{dg_1}{d\mu} = \frac{41}{6} g_1^3 ,$$

$$16\pi^2 \mu \frac{dg_2}{d\mu} = -\frac{19}{6} g_2^3 ,$$

$$16\pi^2 \mu \frac{dg_s}{d\mu} = -7g_s^3 .$$

$$16\pi^2 \mu \frac{dY_l}{d\mu} = \left[-\frac{15}{4}g_1^2 - \frac{9}{4}g_2^2 + T + \frac{3}{2}Y_l Y_l^\dagger - \boxed{2m^2 (C_{H\ell}^{(1)} + 3C_{H\ell}^{(3)})} \right] Y_l ,$$

$$16\pi^2 \mu \frac{dY_u}{d\mu} = \left[-\frac{17}{12}g_1^2 - \frac{9}{4}g_2^2 - 8g_s^2 + T + \frac{3}{2} (Y_u Y_u^\dagger - Y_d Y_d^\dagger) \right] Y_u ,$$

$$16\pi^2 \mu \frac{dY_d}{d\mu} = \left[-\frac{5}{12}g_1^2 - \frac{9}{4}g_2^2 - 8g_s^2 + T - \frac{3}{2} (Y_u Y_u^\dagger - Y_d Y_d^\dagger) \right] Y_d ,$$

$$16\pi^2 \mu \frac{dm^2}{d\mu} = \left(-\frac{3}{2}g_1^2 - \frac{9}{2}g_2^2 + 12\lambda + 2T \right) m^2$$

$$\begin{aligned} 16\pi^2 \mu \frac{d\lambda}{d\mu} &= 24\lambda^2 - 3\lambda (g_1^2 + 3g_2^2) + \frac{3}{8} (g_1^2 + g_2^2)^2 + \frac{3}{4}g_2^4 + 4\lambda T \\ &\quad - 2\text{tr} \left[(Y_l Y_l^\dagger)^2 + 3(Y_u Y_u^\dagger)^2 + 3(Y_d Y_d^\dagger)^2 \right] \\ &\quad + m^2 \text{tr} \left(\boxed{2C_5 C_5^\dagger} - \boxed{\frac{8}{3}g_2^2 C_{H\ell}^{(3)} + 8C_{H\ell}^{(3)} Y_l Y_l^\dagger} \right) . \end{aligned}$$

■ The Weinberg operator:

$$16\pi^2 \mu \frac{dC_5}{d\mu} = (-3g_2^2 + 4\lambda + 2T) C_5 - \frac{3}{2} Y_l Y_l^\dagger C_5 - \frac{3}{2} C_5 \left(Y_l Y_l^\dagger \right)^T$$

■ Dim-6 operators:

$$\begin{aligned} 16\pi^2 \mu \frac{dC_{H\ell}^{(1)}}{d\mu} &= \boxed{-\frac{3}{2} C_5 C_5^\dagger} + \frac{2}{3} g_1^2 \text{tr} \left(C_{H\ell}^{(1)} \right) \mathbb{1} \\ &\quad + \left[\frac{1}{3} g_1^2 + 2\text{tr} \left(Y_l Y_l^\dagger + 3Y_u Y_u^\dagger + 3Y_d Y_d^\dagger \right) \right] C_{H\ell}^{(1)} \\ &\quad + \frac{1}{2} \left[(4C_{H\ell}^{(1)} + 9C_{H\ell}^{(3)}) Y_l Y_l^\dagger + Y_l Y_l^\dagger (4C_{H\ell}^{(1)} + 9C_{H\ell}^{(3)}) \right] , \end{aligned}$$

$$\begin{aligned} 16\pi^2 \mu \frac{dC_{H\ell}^{(3)}}{d\mu} &= \boxed{C_5 C_5^\dagger} + \frac{2}{3} g_2^2 \text{tr} \left(C_{H\ell}^{(3)} \right) \mathbb{1} \\ &\quad + \left[-\frac{17}{3} g_2^2 + 2\text{tr} \left(Y_l Y_l^\dagger + 3Y_u Y_u^\dagger + 3Y_d Y_d^\dagger \right) \right] C_{H\ell}^{(3)} \\ &\quad + \frac{1}{2} \left[(3C_{H\ell}^{(1)} + 2C_{H\ell}^{(3)}) Y_l Y_l^\dagger + Y_l Y_l^\dagger (3C_{H\ell}^{(1)} + 2C_{H\ell}^{(3)}) \right] , \end{aligned}$$

⋮

One-loop RGEs in the Seesaw EFT

■ Dim-6 operators:

H^6 and H^4D^2		ψ^2H^3		$(\bar{L}L)(\bar{L}L)$	
\mathcal{O}_H	$(H^\dagger H)^3$	$\mathcal{O}_{eH}^{\alpha\beta}$	$(\overline{\ell_{\alpha L}} H E_{\beta R}) (H^\dagger H)$	$\mathcal{O}_{\ell\ell}^{\alpha\beta\gamma\lambda}$	$(\overline{\ell_{\alpha L}} \gamma^\mu \ell_{\beta L}) (\overline{\ell_{\gamma L}} \gamma_\mu \ell_{\lambda L})$
$\mathcal{O}_{H\square}$	$(H^\dagger H) \square (H^\dagger H)$	$\mathcal{O}_{uH}^{\alpha\beta}$	$(\overline{Q_{\alpha L}} \tilde{H} U_{\beta R}) (H^\dagger H)$	$\mathcal{O}_{\ell q}^{(1)\alpha\beta\gamma\lambda}$	$(\overline{\ell_{\alpha L}} \gamma^\mu \ell_{\beta L}) (\overline{Q_{\gamma L}} \gamma_\mu Q_{\lambda L})$
\mathcal{O}_{HD}	$(H^\dagger D^\mu H)^* (H^\dagger D_\mu H)$	$\mathcal{O}_{dH}^{\alpha\beta}$	$(\overline{Q_{\alpha L}} H D_{\beta R}) (H^\dagger H)$	$\mathcal{O}_{\ell q}^{(3)\alpha\beta\gamma\lambda}$	$(\overline{\ell_{\alpha L}} \gamma^\mu \sigma^I \ell_{\beta L}) (\overline{Q_{\gamma L}} \gamma_\mu \sigma^I Q_{\lambda L})$
ψ^2H^2D		$(\bar{L}L)(\bar{R}R)$			
$\mathcal{O}_{H\ell}^{(1)\alpha\beta}$	$(\overline{\ell_{\alpha L}} \gamma^\mu \ell_{\beta L}) (H^\dagger i\overleftrightarrow{D}_\mu H)$	$\mathcal{O}_{Hq}^{(3)\alpha\beta}$	$(\overline{Q_{\alpha L}} \gamma^\mu \sigma^I Q_{\beta L}) (H^\dagger i\overleftrightarrow{D}_\mu^I H)$	$\mathcal{O}_{\ell e}^{\alpha\beta\gamma\lambda}$	$(\overline{\ell_{\alpha L}} \gamma^\mu \ell_{\beta L}) (\overline{E_{\gamma R}} \gamma_\mu E_{\lambda R})$
$\mathcal{O}_{He}^{(3)\alpha\beta}$	$(\overline{\ell_{\alpha L}} \gamma^\mu \sigma^I \ell_{\beta L}) (H^\dagger i\overleftrightarrow{D}_\mu^I H)$	$\mathcal{O}_{Hu}^{\alpha\beta}$	$(\overline{U_{\alpha R}} \gamma^\mu U_{\beta R}) (H^\dagger i\overleftrightarrow{D}_\mu H)$	$\mathcal{O}_{\ell u}^{\alpha\beta\gamma\lambda}$	$(\overline{\ell_{\alpha L}} \gamma^\mu \ell_{\beta L}) (\overline{U_{\gamma R}} \gamma_\mu U_{\lambda R})$
$\mathcal{O}_{He}^{\alpha\beta}$	$(\overline{E_{\alpha R}} \gamma^\mu E_{\beta R}) (H^\dagger i\overleftrightarrow{D}_\mu H)$	$\mathcal{O}_{Hd}^{\alpha\beta}$	$(\overline{D_{\alpha R}} \gamma^\mu D_{\beta R}) (H^\dagger i\overleftrightarrow{D}_\mu H)$	$\mathcal{O}_{\ell d}^{\alpha\beta\gamma\lambda}$	$(\overline{\ell_{\alpha L}} \gamma^\mu \ell_{\beta L}) (\overline{D_{\gamma R}} \gamma_\mu D_{\lambda R})$
$\mathcal{O}_{Hq}^{(1)\alpha\beta}$	$(\overline{Q_{\alpha L}} \gamma^\mu Q_{\beta L}) (H^\dagger i\overleftrightarrow{D}_\mu H)$				

- The contribution from **double insertions** of the dim-5 operator to the RGE of quartic Higgs coupling λ is **new**, while those to RGEs of $\mathcal{O}_{H\square}, \mathcal{O}_{HD}, \mathcal{O}_H, \mathcal{O}_{eH}, \mathcal{O}_{uH}, \mathcal{O}_{dH}$ missed a factor of **1/2** in **Davidson et al., 2018**, now are **corrected** and **completed**
- Contributions from **double insertions** of the dim-5 operator are **generic** and thus valid not only for the seesaw EFTs but also for the SMEFT in general
- Results for the **type-II** and **type-III** seesaw EFTs are also achieved and included in **Wang, DZ, and Zhou, 2023**

- In the type-I seesaw EFT, **17** dim-6 operators can be generated via their one-loop RGEs apart from the 2 tree-level ones
- Contributions from **single insertions** of the dim-5 and dim-6 operators are consistent with results in the literature
Antusch et al., 2001; Jenkins et al., 2013

One-loop RGEs of Flavor Mixing Parameters

After spontaneous symmetry breaking

$$\begin{aligned} \mathcal{L}_{\text{SEFT}}^l &= - \left[\overline{l_{\alpha L}} (M_l)_{\alpha\beta} l_{\beta R} + \frac{1}{2} \overline{\nu_{\alpha L}} (M_\nu)_{\alpha\beta} \nu_{\beta L}^c + \text{h.c.} \right] \\ &\quad + \left[\frac{g_2}{\sqrt{2}} \overline{l_{\alpha L}} \gamma^\mu (\delta_{\alpha\beta} - \tilde{\eta}_{\alpha\beta}) \nu_{\beta L} W_\mu^- + \text{h.c.} \right] + \frac{g_2}{2 c_W} \overline{\nu_{\alpha L}} \gamma^\mu (\delta_{\alpha\beta} - \tilde{\eta}'_{\alpha\beta}) \nu_{\beta L} Z_\mu \\ &\quad - \frac{g_2}{2 c_W} \overline{l_{\alpha L}} \gamma^\mu \left[(1 - 2 s_W^2) \delta_{\alpha\beta} + (\tilde{\eta}' - 2 \tilde{\eta})_{\alpha\beta} \right] l_{\beta L} Z_\mu + \frac{g_2}{c_W} s_W^2 \overline{l_{\alpha R}} \gamma^\mu l_{\alpha R} Z_\mu , \end{aligned}$$

$$\begin{aligned} &= - \left(\overline{l_L} \widehat{M}_l l_R + \frac{1}{2} \overline{\nu_L} \widehat{M}_\nu \nu_L^c + \text{h.c.} \right) \xrightarrow{\text{CC: Non-unitarity}} \\ &\quad + \left(\frac{g_2}{\sqrt{2}} \overline{l_L} \gamma^\mu V \nu_L W_\mu^- + \text{h.c.} \right) + \frac{g_2}{2 c_W} \overline{\nu_L} \gamma^\mu N^\dagger N \nu_L Z_\mu \xrightarrow{\text{NC}} \\ &\quad - \frac{g_2}{2 c_W} \overline{l_L} \gamma^\mu \left[(1 - 2 s_W^2) + (\eta' - 2 \eta) \right] l_L Z_\mu + \frac{g_2}{c_W} s_W^2 \overline{l_R} \gamma^\mu l_R Z_\mu \end{aligned}$$

$V \equiv (\mathbb{1} - \eta) \cdot U \cdot Q$ Non-unitary Pontecorvo-Maki-Nakagawa-Sakata matrix

$N \equiv (\mathbb{1} - \eta'/2) \cdot U \cdot Q$ Flavor-changing NC

$$\eta \equiv P^\dagger U_l^\dagger \tilde{\eta} U_l P = \begin{pmatrix} \eta_{ee} & |\eta_{e\mu}| e^{+i\phi_{e\mu}} & |\eta_{e\tau}| e^{+i\phi_{e\tau}} \\ |\eta_{e\mu}| e^{-i\phi_{e\mu}} & \eta_{\mu\mu} & |\eta_{\mu\tau}| e^{+i\phi_{\mu\tau}} \\ |\eta_{e\tau}| e^{-i\phi_{e\tau}} & |\eta_{\mu\tau}| e^{-i\phi_{\mu\tau}} & \eta_{\tau\tau} \end{pmatrix} \quad \eta' \equiv P^\dagger U_l^\dagger \tilde{\eta}' U_l P = \begin{pmatrix} \eta'_{ee} & |\eta'_{e\mu}| e^{+i\phi'_{e\mu}} & |\eta'_{e\tau}| e^{+i\phi'_{e\tau}} \\ |\eta'_{e\mu}| e^{-i\phi'_{e\mu}} & \eta'_{\mu\mu} & |\eta'_{\mu\tau}| e^{+i\phi'_{\mu\tau}} \\ |\eta'_{e\tau}| e^{-i\phi'_{e\tau}} & |\eta'_{\mu\tau}| e^{-i\phi'_{\mu\tau}} & \eta'_{\tau\tau} \end{pmatrix}$$

$$M_l = Y_l v / \sqrt{2} \quad M_\nu = -C_5 v^2 / 2$$

$$\tilde{\eta} \equiv -C_{Hl}^{(3)} v^2 \quad \tilde{\eta}' \equiv (C_{Hl}^{(1)} - C_{Hl}^{(3)}) v^2$$

$$s_W \equiv \sin \theta_W$$

$$U_\nu^\dagger M_\nu U_\nu^* = \widehat{M}_\nu \equiv \text{diag}\{m_1, m_2, m_3\}$$

$$U_l^\dagger M_l U_l' = \widehat{M}_l \equiv \text{diag}\{m_e, m_\mu, m_\tau\}$$

$$V' \equiv U_l^\dagger U_\nu \quad \text{Definition}$$

$$V' \equiv P \cdot U \cdot Q \quad \text{Parametrization}$$

$$P \equiv \text{diag}\{e^{i\phi_e}, e^{i\phi_\mu}, e^{i\phi_\tau}\}$$

$$Q \equiv \text{diag}\{e^{i\rho}, e^{i\sigma}, 1\}$$

$$\eta'(\mu_M) = 2\eta(\mu_M)$$

One-loop RGEs of Flavor Mixing Parameters

RGEs for **eigenvalues**:

$$\dot{y}_\alpha = \left[\left(-\frac{15}{4}g_1^2 - \frac{9}{4}g_2^2 + T \right) + \frac{3}{2}y_\alpha^2 - 2\frac{m^2}{v^2}(\eta' - 4\eta)_{\alpha\alpha} \right] y_\alpha$$

$$\dot{\kappa}_i = [(-3g_2^2 + 4\lambda + 2T) - 3\text{Re}\mathcal{S}_{ii}] \kappa_i$$

RGEs for lepton flavor **mixing parameters** and those in the **NC** interactions:

$$\begin{aligned} V'_{\alpha i} &= \sum_{\beta} \left(\dot{U}_l^\dagger U_l \right)_{\alpha\beta} V'_{\beta i} + \sum_j V'_{\alpha j} \left(U_\nu^\dagger \dot{U}_\nu \right)_{ji} \\ &= \sum_{\beta \neq \alpha} 2 \frac{m^2}{v^2} y_{\alpha\beta} (\eta' - 4\eta)_{\alpha\beta} e^{i(\phi_\alpha - \phi_\beta)} V'_{\beta i} - \sum_{j \neq i} \frac{3}{2} V'_{\alpha j} \frac{1}{\kappa_i^2 - \kappa_j^2} [(\kappa_i^2 + \kappa_j^2) \mathcal{S}_{ji} + 2\kappa_i \kappa_j \mathcal{S}_{ji}^*] \end{aligned}$$

$$\begin{aligned} \dot{\eta}_{\alpha\beta} &= i(\dot{\phi}_\beta - \dot{\phi}_\alpha) \eta_{\alpha\beta} + \sum_{\gamma \neq \alpha} 2 \frac{m^2}{v^2} y_{\alpha\gamma} (\eta' - 4\eta)_{\alpha\gamma} \eta_{\gamma\beta} + \sum_{\varrho \neq \beta} 2 \frac{m^2}{v^2} y_{\beta\varrho} \eta_{\alpha\varrho} (\eta' - 4\eta)_{\varrho\beta} \\ &\quad - \sum_i \kappa_i^2 v^2 U_{\alpha i} U_{\beta i}^* + \frac{2}{3} g_2^2 \text{tr}(\eta) \delta_{\alpha\beta} + \left(-\frac{17}{3} g_2^2 + 2T \right) \eta_{\alpha\beta} + \frac{1}{2} (y_\alpha^2 + y_\beta^2) (5\eta_{\alpha\beta} - 3\eta'_{\alpha\beta}) \end{aligned}$$

$$\begin{aligned} \dot{\eta}'_{\alpha\beta} &= i(\dot{\phi}_\beta - \dot{\phi}_\alpha) \eta'_{\alpha\beta} + \sum_{\gamma \neq \alpha} 2 \frac{m^2}{v^2} y_{\alpha\gamma} (\eta' - 4\eta)_{\alpha\gamma} \eta'_{\gamma\beta} + \sum_{\varrho \neq \beta} 2 \frac{m^2}{v^2} y_{\beta\varrho} \eta'_{\alpha\varrho} (\eta' - 4\eta)_{\varrho\beta} \\ &\quad + \frac{2}{3} (g_2^2 - g_1^2) \text{tr}(\eta) \delta_{\alpha\beta} + \frac{2}{3} g_1^2 \text{tr}(\eta') \delta_{\alpha\beta} - \frac{1}{3} (17g_2^2 + g_1^2) \eta_{\alpha\beta} + \left(\frac{1}{3} g_1^2 + 2T \right) \eta'_{\alpha\beta} \\ &\quad - \frac{5}{2} \sum_i \kappa_i^2 v^2 U_{\alpha i} U_{\beta i}^* + \frac{1}{2} (y_\alpha^2 + y_\beta^2) (\eta'_{\alpha\beta} - 8\eta_{\alpha\beta}), \end{aligned}$$

$$\kappa \equiv C_5$$

$$U_l^\dagger Y_l U_l' = \text{diag}\{y_e, y_\mu, y_\tau\}$$

$$U_\nu^\dagger \kappa U_\nu^* = \text{diag}\{\kappa_1, \kappa_2, \kappa_3\}$$

$$\mathcal{S} \equiv V'^\dagger \widehat{Y}_l^2 V'$$

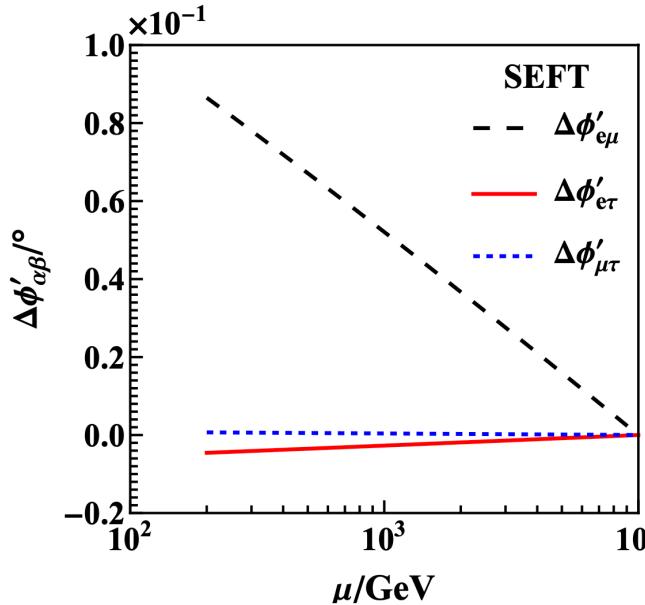
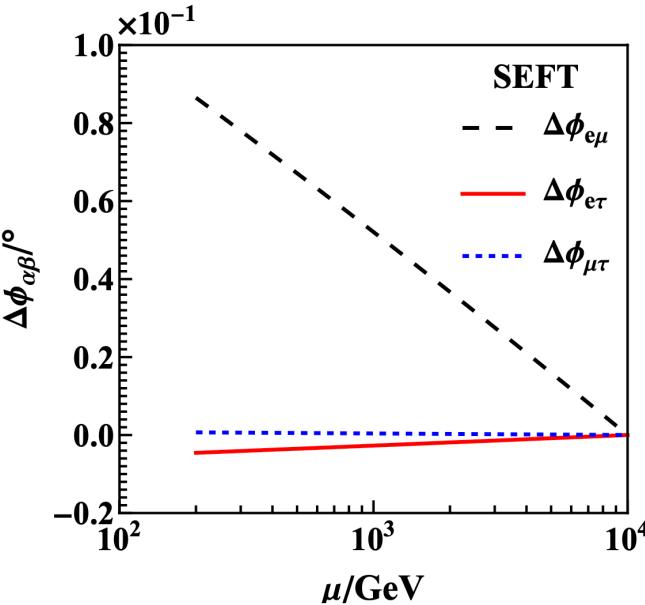
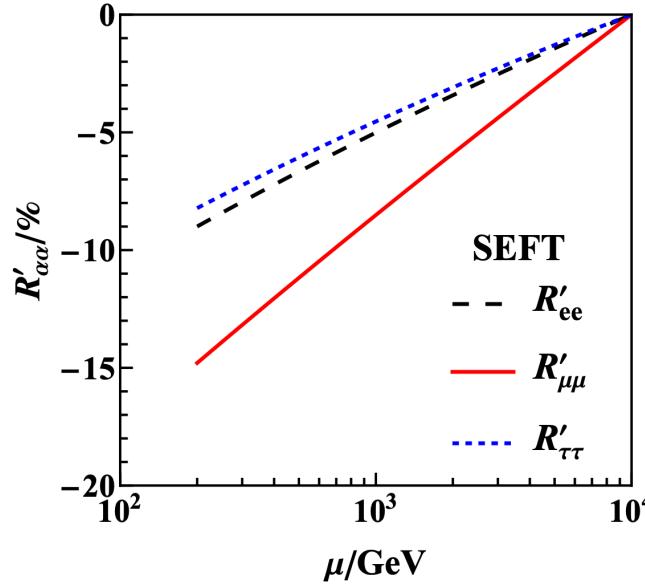
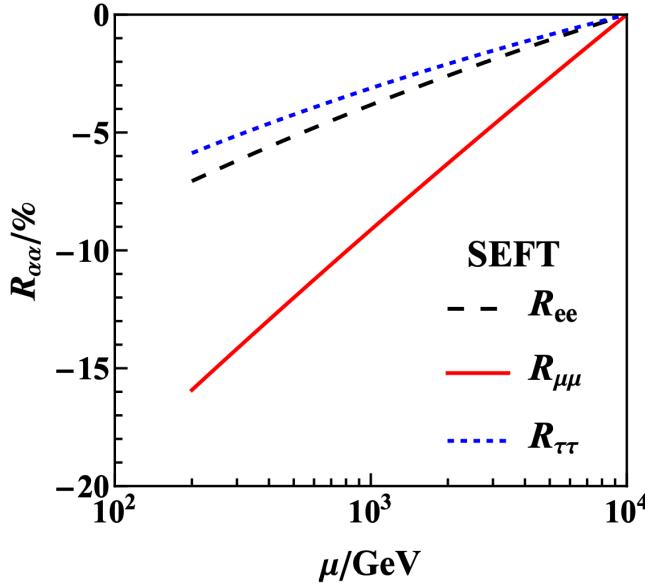
The non-unitary PMNS matrix

$$V \equiv (\mathbb{1} - \eta) \cdot U \cdot Q$$

- Lepton masses
- Mixing angles and phases in V'
- Unitarity-violating and FCNC parameters, i.e., magnitude and arguments of η and η'

See, **Wang, DZ, and Zhou, 2023**

Numerical Analysis



$$R_{\alpha\beta}^{(I)} \equiv [|\eta_{\alpha\beta}^{(I)}(\mu)| - |\eta_{\alpha\beta}^{(I)}(\mu_M)|]/|\eta_{\alpha\beta}^{(I)}(\mu_M)| \times 100\%$$

$$R_{\alpha\alpha} \sim \Delta t \left[\frac{2}{3} g_2^2 \frac{\text{tr}(\eta)}{\eta_{\alpha\alpha}} - \frac{17}{3} g_2^2 + 6y_t^2 \right],$$

$$R'_{\alpha\alpha} \sim \Delta t \left[\frac{1}{3} (g_1^2 + g_2^2) \frac{\text{tr}(\eta)}{\eta_{\alpha\alpha}} + \frac{1}{6} (g_1^2 - 17g_2^2) + 6y_t^2 \right],$$

$$\Delta t = \ln(\mu_B/\mu_M)/(16\pi^2) < 0$$

$$0 < \eta_{\mu\mu} < \eta_{ee} < \eta_{\tau\tau} \Rightarrow R_{\mu\mu}^{(I)} < R_{ee}^{(I)} < R_{\tau\tau}^{(I)} < 0$$

$$\Delta\phi_{\alpha\beta}^{(I)} \equiv \phi_{\alpha\beta}^{(I)}(\mu) - \phi_{\alpha\beta}^{(I)}(\mu_M)$$

$$\Delta\phi_{e\mu}^{(I)} \sim \left[8 \frac{m^2}{v^2} \left| \frac{\eta_{e\tau}\eta_{\mu\tau}}{\eta_{e\mu}} \right| s_{e\mu} + \frac{3}{4} \zeta_{12}^{-1} y_\tau^2 s_{23}^2 s_{2(\rho-\sigma)} \right] \Delta t,$$

$$\Delta\phi_{e\tau}^{(I)} \sim \frac{3}{4} \zeta_{12}^{-1} y_\tau^2 s_{23}^2 s_{2(\rho-\sigma)} \Delta t,$$

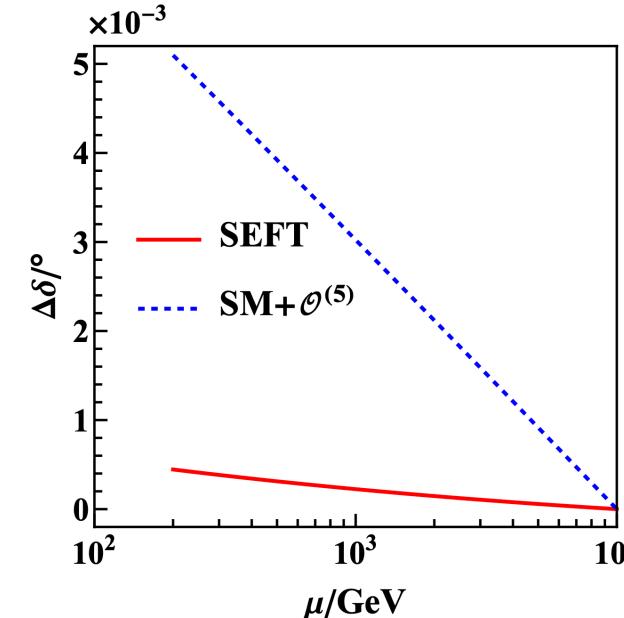
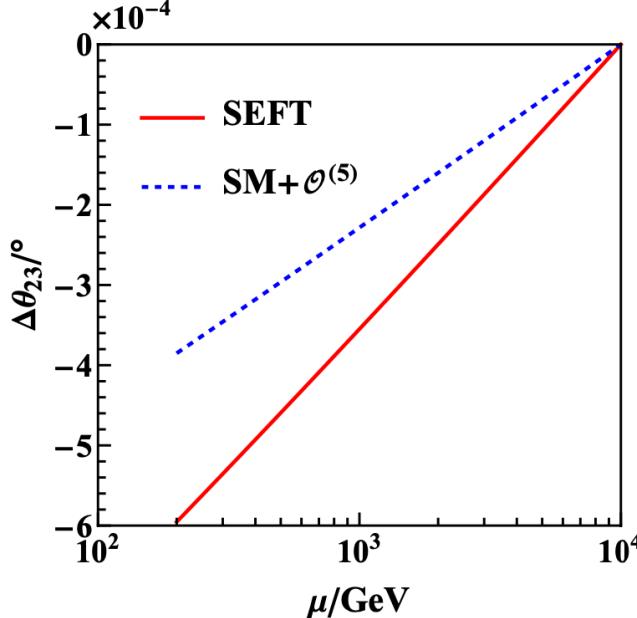
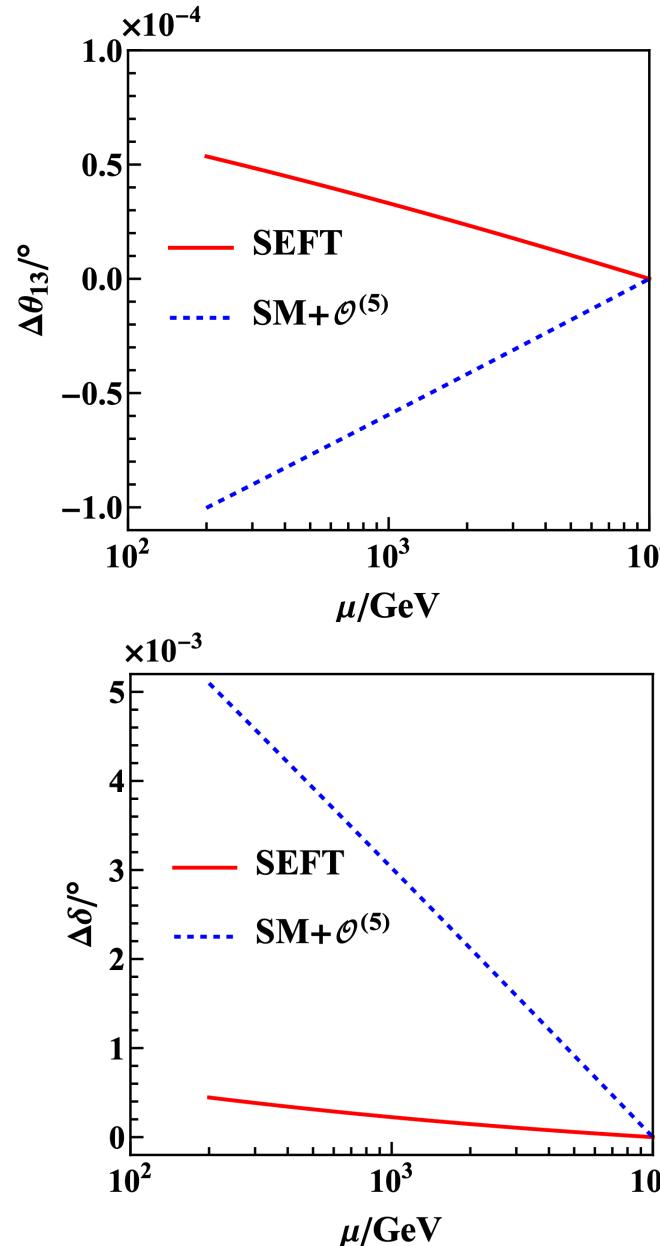
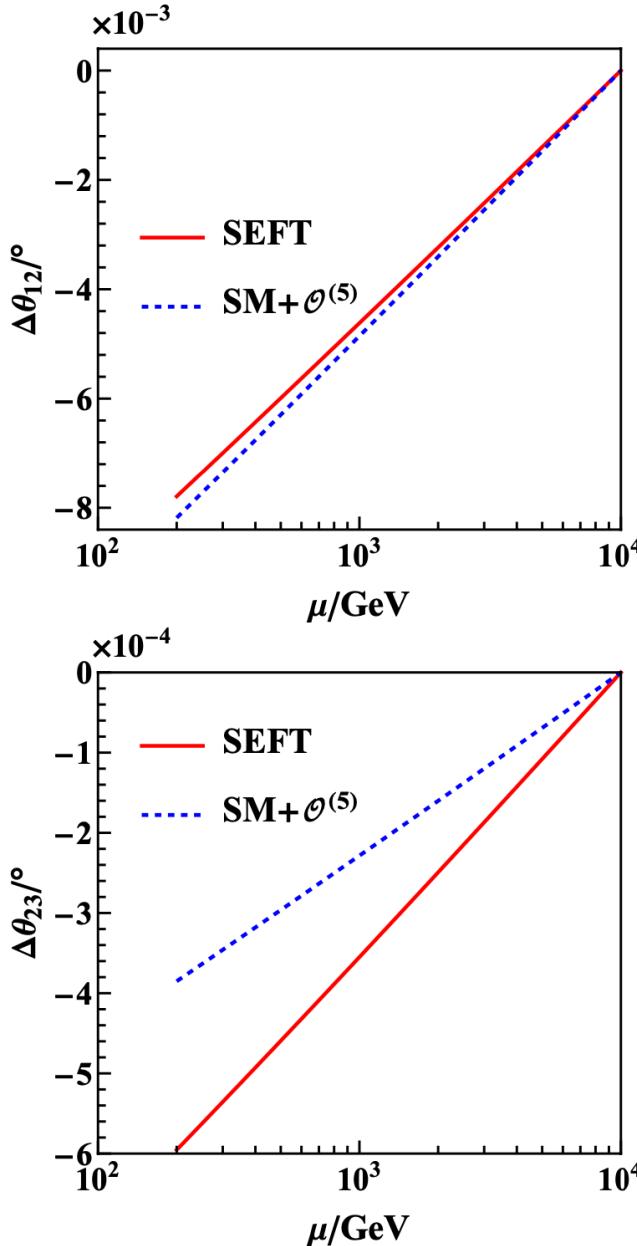
$$|\Delta\phi_{\mu\tau}^{(I)}| \sim \mathcal{O}(|\eta|, \zeta_{23}^{-1} y_\tau^2) |\Delta t| \sim 10^{-4} |\Delta t|,$$



$$\zeta_{ij} \equiv (\kappa_i - \kappa_j)/(\kappa_i + \kappa_j)$$

$$\Delta\phi_{\alpha\beta} \approx \Delta\phi'_{\alpha\beta}$$

Numerical Analysis



$$\Delta\theta_{13} \sim -\Delta t \left[4 \frac{m^2}{v^2} |\eta_{e\tau}| c_{23} c_{e\tau+\delta} + \frac{3}{8} \zeta_{23}^{-1} y_\tau^2 \sin 2\theta_{12} \sin 2\theta_{23} s_{\rho-\sigma} s_{\delta+\rho+\sigma} \right]$$

- Two terms have opposite signs and the absolute value of the first one is slightly larger than that of the second one, thus opposite running directions
- But depends on the initial inputs
- Approximately and analytically, the running behaviors of all these parameters can be well-understood
- The non-unitary parameters may significantly affect the running of leptonic flavor mixing parameters

Summary

- We derived the **complete** set of **one-loop RGEs** for the SM couplings and Wilson coefficients of operators up to dim-6 and $\mathcal{O}(1/\Lambda_{\text{SS}}^2)$ in seesaw EFTs
- Besides two tree-level-generated dim-6 operators, **17** dim-6 operators can be generated by the **one-loop** RGEs in the **type-I** seesaw EFT
- We gave the **explicit expressions** of the RGEs of all the **parameters** involved in the charged- and neutral-current interactions of leptons
- Together with the one-loop matching results at Λ_{SS} , these one-loop RGEs establish a **self-consistent** EFT framework to investigate **low-energy phenomena** of seesaw models up to $\mathcal{O}(1/\Lambda_{\text{SS}}^2)$ at the **one-loop level**

THANKS FOR YOUR ATTENTION
GRACIAS / DANKE / 谢谢

Backup

The Greens basis in the SMEFT

Jiang et al., 2019; Gherardi et al., 2020; Carmona et al., 2022

$\psi^2 D^3$		$\psi^2 XD$		$\psi^2 DH^2$	
$\mathcal{R}_{qD}^{\alpha\beta}$	$\frac{i}{2}\overline{Q_{\alpha L}}\{D_\mu D^\mu, \not{D}\} Q_{\beta L}$	$\mathcal{R}_{Gq}^{\alpha\beta}$	$(\overline{Q_{\alpha L}}T^A\gamma^\mu Q_{\beta L}) D^\nu G_{\mu\nu}^A$	$\mathcal{O}_{Hq}^{(1)\alpha\beta}$	$(\overline{Q_{\alpha L}}\gamma^\mu Q_{\beta L}) (H^\dagger i\overleftrightarrow{D}_\mu H)$
$\mathcal{R}_{uD}^{\alpha\beta}$	$\frac{i}{2}\overline{U_{\alpha R}}\{D_\mu D^\mu, \not{D}\} U_{\beta R}$	$\mathcal{R}'_{Gq}^{\alpha\beta}$	$\frac{1}{2}(\overline{Q_{\alpha L}}T^A\gamma^\mu i\overleftrightarrow{D}^\nu Q_{\beta L}) G_{\mu\nu}^A$	$\mathcal{R}'_{Hq}^{(1)\alpha\beta}$	$(\overline{Q_{\alpha L}}i\overleftrightarrow{D}^\nu Q_{\beta L}) (H^\dagger H)$
$\mathcal{R}_{dD}^{\alpha\beta}$	$\frac{i}{2}\overline{D_{\alpha R}}\{D_\mu D^\mu, \not{D}\} D_{\beta R}$	$\mathcal{R}_{\tilde{G}q}^{\alpha\beta}$	$\frac{1}{2}(\overline{Q_{\alpha L}}T^A\gamma^\mu i\overleftrightarrow{D}^\nu Q_{\beta L}) \tilde{G}_{\mu\nu}^A$	$\mathcal{R}''_{Hq}^{(1)\alpha\beta}$	$(\overline{Q_{\alpha L}}\gamma^\mu Q_{\beta L}) \partial_\mu (H^\dagger H)$
$\mathcal{R}_{eD}^{\alpha\beta}$	$\frac{i}{2}\overline{\ell_{\alpha L}}\{D_\mu D^\mu, \not{D}\} \ell_{\beta L}$	$\mathcal{R}_{Wq}^{\alpha\beta}$	$(\overline{Q_{\alpha L}}\sigma^I\gamma^\mu Q_{\beta L}) D^\nu W_{\mu\nu}^I$	$\mathcal{O}_{Hq}^{(3)\alpha\beta}$	$(\overline{Q_{\alpha L}}\gamma^\mu \sigma^I Q_{\beta L}) (H^\dagger i\overleftrightarrow{D}_\mu^I H)$
\mathcal{R}_{eD}	$\frac{i}{2}\overline{E_{\alpha R}}\{D_\mu D^\mu, \not{D}\} E_{\beta R}$	$\mathcal{R}'_{Wq}^{\alpha\beta}$	$\frac{1}{2}(\overline{Q_{\alpha L}}\sigma^I\gamma^\mu i\overleftrightarrow{D}^\nu Q_{\beta L}) W_{\mu\nu}^I$	$\mathcal{R}'_{Hq}^{(3)\alpha\beta}$	$(\overline{Q_{\alpha L}}i\overleftrightarrow{D}^\nu Q_{\beta L}) (H^\dagger \sigma^I H)$
$\psi^2 HD^2$		$\mathcal{R}_{\tilde{W}q}^{\alpha\beta}$	$\frac{1}{2}(\overline{Q_{\alpha L}}\sigma^I\gamma^\mu i\overleftrightarrow{D}^\nu Q_{\beta L}) \widetilde{W}_{\mu\nu}^I$	$\mathcal{R}''_{Hq}^{(3)\alpha\beta}$	$(\overline{Q_{\alpha L}}\sigma^I\gamma^\mu Q_{\beta L}) D_\mu (H^\dagger \sigma^I H)$
$\mathcal{R}_{uHD1}^{\alpha\beta}$	$(\overline{Q_{\alpha L}}U_{\beta R}) D_\mu D^\mu \tilde{H}$	$\mathcal{R}_{Bq}^{\alpha\beta}$	$(\overline{Q_{\alpha L}}\gamma^\mu Q_{\beta L}) \partial^\nu B_{\mu\nu}$	$\mathcal{O}_{Hu}^{\alpha\beta}$	$(\overline{U_{\alpha R}}\gamma^\mu U_{\beta R}) (H^\dagger i\overleftrightarrow{D}_\mu H)$
$\mathcal{R}_{uHD2}^{\alpha\beta}$	$(\overline{Q_{\alpha L}}i\sigma_{\mu\nu} D^\mu U_{\beta R}) D^\nu \tilde{H}$	$\mathcal{R}'_{Bq}^{\alpha\beta}$	$\frac{1}{2}(\overline{Q_{\alpha L}}\gamma^\mu i\overleftrightarrow{D}^\nu Q_{\beta L}) B_{\mu\nu}$	$\mathcal{R}'_{Hu}^{\alpha\beta}$	$(\overline{U_{\alpha R}}i\overleftrightarrow{D}^\nu U_{\beta R}) (H^\dagger H)$
$\mathcal{R}_{uHD3}^{\alpha\beta}$	$(\overline{Q_{\alpha L}}D_\mu D^\mu U_{\beta R}) \tilde{H}$	$\mathcal{R}_{\tilde{B}q}^{\alpha\beta}$	$\frac{1}{2}(\overline{Q_{\alpha L}}\gamma^\mu i\overleftrightarrow{D}^\nu Q_{\beta L}) \tilde{B}_{\mu\nu}$	$\mathcal{R}''_{Hu}^{\alpha\beta}$	$(\overline{U_{\alpha R}}\gamma^\mu U_{\beta R}) \partial_\mu (H^\dagger H)$
$\mathcal{R}_{uHD4}^{\alpha\beta}$	$(\overline{Q_{\alpha L}}D_\mu U_{\beta R}) D^\mu \tilde{H}$	$\mathcal{R}_{Gu}^{\alpha\beta}$	$(\overline{U_{\alpha R}}T^A\gamma^\mu U_{\beta R}) D^\nu G_{\mu\nu}^A$	$\mathcal{O}_{Hd}^{\alpha\beta}$	$(\overline{D_{\alpha R}}\gamma^\mu D_{\beta R}) (H^\dagger i\overleftrightarrow{D}_\mu H)$
$\mathcal{R}_{dHD1}^{\alpha\beta}$	$(\overline{Q_{\alpha L}}D_{\beta R}) D_\mu D^\mu H$	$\mathcal{R}'_{Gu}^{\alpha\beta}$	$\frac{1}{2}(\overline{U_{\alpha R}}T^A\gamma^\mu i\overleftrightarrow{D}^\nu U_{\beta R}) G_{\mu\nu}^A$	$\mathcal{R}'_{Hd}^{\alpha\beta}$	$(\overline{D_{\alpha R}}i\overleftrightarrow{D}^\nu D_{\beta R}) (H^\dagger H)$
$\mathcal{R}_{dHD2}^{\alpha\beta}$	$(\overline{Q_{\alpha L}}i\sigma_{\mu\nu} D^\mu D_{\beta R}) D^\nu H$	$\mathcal{R}''_{Gu}^{\alpha\beta}$	$\frac{1}{2}(\overline{U_{\alpha R}}T^A\gamma^\mu i\overleftrightarrow{D}^\nu U_{\beta R}) \tilde{G}_{\mu\nu}^A$	$\mathcal{R}''_{Hd}^{\alpha\beta}$	$(\overline{D_{\alpha R}}\gamma^\mu D_{\beta R}) \partial_\mu (H^\dagger H)$
$\mathcal{R}_{dHD3}^{\alpha\beta}$	$(\overline{Q_{\alpha L}}D_\mu D^\mu D_{\beta R}) H$	$\mathcal{R}_{Bu}^{\alpha\beta}$	$(\overline{U_{\alpha R}}\gamma^\mu U_{\beta R}) \partial^\nu B_{\mu\nu}$	$\mathcal{O}_{Hud}^{\alpha\beta}$	$i(\overline{U_{\alpha R}}\gamma^\mu D_{\beta R}) (\tilde{H}^\dagger D_\mu H)$
$\mathcal{R}_{dHD4}^{\alpha\beta}$	$(\overline{Q_{\alpha L}}D_\mu D_{\beta R}) D^\mu H$	$\mathcal{R}'_{Bu}^{\alpha\beta}$	$\frac{1}{2}(\overline{U_{\alpha R}}\gamma^\mu i\overleftrightarrow{D}^\nu U_{\beta R}) B_{\mu\nu}$	$\mathcal{O}_{H\ell}^{(1)\alpha\beta}$	$(\overline{\ell_{\alpha L}}\gamma^\mu \ell_{\beta L}) (H^\dagger i\overleftrightarrow{D}_\mu H)$
$\mathcal{R}_{eHD1}^{\alpha\beta}$	$(\overline{\ell_{\alpha L}}E_{\beta R}) D_\mu D^\mu H$	$\mathcal{R}''_{Bu}^{\alpha\beta}$	$\frac{1}{2}(\overline{U_{\alpha R}}\gamma^\mu i\overleftrightarrow{D}^\nu U_{\beta R}) \tilde{B}_{\mu\nu}$	$\mathcal{R}'_{H\ell}^{(1)\alpha\beta}$	$(\overline{\ell_{\alpha L}}i\overleftrightarrow{D}^\nu \ell_{\beta L}) (H^\dagger H)$
$\mathcal{R}_{eHD2}^{\alpha\beta}$	$(\overline{\ell_{\alpha L}}i\sigma_{\mu\nu} D^\mu E_{\beta R}) D^\nu H$	$\mathcal{R}_{Gd}^{\alpha\beta}$	$(\overline{D_{\alpha R}}T^A\gamma^\mu D_{\beta R}) D^\nu G_{\mu\nu}^A$	$\mathcal{R}''_{H\ell}^{(1)\alpha\beta}$	$(\overline{\ell_{\alpha L}}\gamma^\mu \ell_{\beta L}) \partial_\mu (H^\dagger H)$
$\mathcal{R}_{eHD3}^{\alpha\beta}$	$(\overline{\ell_{\alpha L}}D_\mu D^\mu E_{\beta R}) H$	$\mathcal{R}'_{Gd}^{\alpha\beta}$	$\frac{1}{2}(\overline{D_{\alpha R}}T^A\gamma^\mu i\overleftrightarrow{D}^\nu D_{\beta R}) G_{\mu\nu}^A$	$\mathcal{O}_{H\ell}^{(3)\alpha\beta}$	$(\overline{\ell_{\alpha L}}\gamma^\mu \ell_{\beta L}) (H^\dagger i\overleftrightarrow{D}_\mu^I H)$
$\mathcal{R}_{eHD4}^{\alpha\beta}$	$(\overline{\ell_{\alpha L}}D_\mu E_{\beta R}) D^\mu H$	$\mathcal{R}''_{Gd}^{\alpha\beta}$	$\frac{1}{2}(\overline{D_{\alpha R}}T^A\gamma^\mu i\overleftrightarrow{D}^\nu D_{\beta R}) \tilde{G}_{\mu\nu}^A$	$\mathcal{R}'_{H\ell}^{(3)\alpha\beta}$	$(\overline{\ell_{\alpha L}}i\overleftrightarrow{D}^\nu \ell_{\beta L}) (H^\dagger \sigma^I H)$
$\psi^2 XH$		$\mathcal{R}_{Bd}^{\alpha\beta}$	$(\overline{D_{\alpha R}}\gamma^\mu D_{\beta R}) \partial^\nu B_{\mu\nu}$	$\mathcal{R}''_{H\ell}^{(3)\alpha\beta}$	$(\overline{\ell_{\alpha L}}\sigma^I\gamma^\mu \ell_{\beta L}) D_\mu (H^\dagger \sigma^I H)$
$\mathcal{O}_{uG}^{\alpha\beta}$	$(\overline{Q_{\alpha L}}\sigma^{\mu\nu} T^A U_{\beta R}) \tilde{H} G_{\mu\nu}^A$	$\mathcal{R}'_{Bd}^{\alpha\beta}$	$\frac{1}{2}(\overline{D_{\alpha R}}\gamma^\mu i\overleftrightarrow{D}^\nu D_{\beta R}) B_{\mu\nu}$	$\mathcal{O}_{He}^{\alpha\beta}$	$(\overline{E_{\alpha R}}\gamma^\mu E_{\beta R}) (H^\dagger i\overleftrightarrow{D}_\mu H)$
$\mathcal{O}_{uW}^{\alpha\beta}$	$(\overline{Q_{\alpha L}}\sigma^{\mu\nu} U_{\beta R}) \sigma^I \tilde{H} W_{\mu\nu}^I$	$\mathcal{R}''_{Bd}^{\alpha\beta}$	$\frac{1}{2}(\overline{D_{\alpha R}}\gamma^\mu i\overleftrightarrow{D}^\nu D_{\beta R}) \tilde{B}_{\mu\nu}$	$\mathcal{R}'_{He}^{\alpha\beta}$	$(\overline{E_{\alpha R}}i\overleftrightarrow{D}^\nu E_{\beta R}) (H^\dagger H)$
$\mathcal{O}_{uB}^{\alpha\beta}$	$(\overline{Q_{\alpha L}}\sigma^{\mu\nu} U_{\beta R}) \tilde{H} B_{\mu\nu}$	$\mathcal{R}_{W\ell}^{\alpha\beta}$	$(\overline{\ell_{\alpha L}}\sigma^I\gamma^\mu \ell_{\beta L}) D^\nu W_{\mu\nu}^I$	$\mathcal{R}''_{He}^{\alpha\beta}$	$(\overline{E_{\alpha R}}\gamma^\mu E_{\beta R}) \partial_\mu (H^\dagger H)$

$\mathcal{O}_{dG}^{\alpha\beta}$	$(\overline{Q_{\alpha L}}\sigma^{\mu\nu} T^A D_{\beta R}) H G_{\mu\nu}^A$	$\mathcal{R}'_{W\ell}^{\alpha\beta}$	$\frac{1}{2}(\overline{\ell_{\alpha L}}\sigma^I\gamma^\mu i\overleftrightarrow{D}^\nu \ell_{\beta L}) W_{\mu\nu}^I$	$\psi^2 H^3$
$\mathcal{O}_{dW}^{\alpha\beta}$	$(\overline{Q_{\alpha L}}\sigma^{\mu\nu} D_{\beta R}) \sigma^I H W_{\mu\nu}^I$	$\mathcal{R}''_{W\ell}^{\alpha\beta}$	$\frac{1}{2}(\overline{\ell_{\alpha L}}\sigma^I\gamma^\mu i\overleftrightarrow{D}^\nu \ell_{\beta L}) \widetilde{W}_{\mu\nu}^I$	$\mathcal{O}_{uH}^{\alpha\beta}$
$\mathcal{O}_{dB}^{\alpha\beta}$	$(\overline{Q_{\alpha L}}\sigma^{\mu\nu} D_{\beta R}) H B_{\mu\nu}$	$\mathcal{R}_{Be}^{\alpha\beta}$	$(\overline{\ell_{\alpha L}}\gamma^\mu \ell_{\beta L}) \partial^\nu B_{\mu\nu}$	$\mathcal{O}_{dH}^{\alpha\beta}$
$\mathcal{O}_{eW}^{\alpha\beta}$	$(\overline{\ell_{\alpha L}}\sigma^{\mu\nu} E_{\beta R}) \sigma^I H W_{\mu\nu}^I$	$\mathcal{R}'_{Be}^{\alpha\beta}$	$\frac{1}{2}(\overline{\ell_{\alpha L}}\gamma^\mu i\overleftrightarrow{D}^\nu \ell_{\beta L}) B_{\mu\nu}$	$\mathcal{O}_{eH}^{\alpha\beta}$
$\mathcal{O}_{eB}^{\alpha\beta}$	$(\overline{\ell_{\alpha L}}\sigma^{\mu\nu} E_{\beta R}) H B_{\mu\nu}$	$\mathcal{R}''_{Be}^{\alpha\beta}$	$\frac{1}{2}(\overline{\ell_{\alpha L}}\gamma^\mu i\overleftrightarrow{D}^\nu \ell_{\beta L}) \widetilde{B}_{\mu\nu}$	

Table A.1: Two-fermion operators in the Green's basis in the SMEFT.

X^3		$X^2 H^2$		$H^2 D^4$	
\mathcal{O}_{3G}	$f^{ABC} G_\mu^{A\nu} G_\nu^{B\rho} G_\rho^{C\mu}$	\mathcal{O}_{HG}	$G_{\mu\nu}^A G^{A\mu\nu} H^\dagger H$	\mathcal{R}_{DH}	$(D_\mu D^\mu H)^\dagger (D_\nu D^\nu H)$
$\mathcal{O}_{3\widetilde{G}}$	$f^{ABC} \widetilde{G}_\mu^{A\nu} G_\nu^{B\rho} G_\rho^{C\mu}$	$\mathcal{O}_{H\widetilde{G}}$	$\widetilde{G}_{\mu\nu}^A G^{A\mu\nu} H^\dagger H$		$H^4 D^2$
\mathcal{O}_{3W}	$\epsilon^{IJK} W_\mu^{I\nu} W_\nu^{J\rho} W_\rho^{K\mu}$	\mathcal{O}_{HW}	$W_{\mu\nu}^I W^{I\mu\nu} H^\dagger H$	$\mathcal{O}_{H\square}$	$(H^\dagger H) \square (H^\dagger H)$
$\mathcal{O}_{3\widetilde{W}}$	$\epsilon^{IJK} \widetilde{W}_\mu^{I\nu} W_\nu^{J\rho} W_\rho^{K\mu}$	$\mathcal{O}_{H\widetilde{W}}$	$\widetilde{W}_{\mu\nu}^I W^{I\mu\nu} H^\dagger H$	\mathcal{O}_{HD}	$(H^\dagger D^\mu H)^* (H^\dagger D_\mu H)$
$X^2 D^2$		H^6		\mathcal{O}_H	$(H^\dagger H)^3$
\mathcal{R}_{2G}	$-\frac{1}{2}(D_\mu G^{A\mu\nu}) (D^\rho G_{\rho\nu}^A)$	$\mathcal{O}_{H\widetilde{B}}$	$B_{\mu\nu} B^{\mu\nu} H^\dagger H$		
\mathcal{R}_{2W}	$-\frac{1}{2}(D_\mu W^{I\mu\nu}) (D^\rho W_{\rho\nu}^I)$	\mathcal{O}_{HWB}	$W_{\mu\nu}^I B^{\mu\nu} (H^\dagger \sigma^I H)$		
\mathcal{R}_{2B}	$-\frac{1}{2}(\partial_\mu B^{\mu\nu}) (\partial^\rho B_{\rho\nu})$	$\mathcal{O}_{H\widetilde{W}B}$	$\widetilde{W}_{\mu\nu}^I B^{\mu\nu} (H^\dagger \sigma^I H)$		
$H^2 X D^2$		$H^2 X D^2$		\mathcal{R}_{WDH}	$D_\nu W^{I\mu\nu} (H^\dagger i\overleftrightarrow{D}_\mu^I H)$
				\mathcal{R}_{BDH}	$\partial_\nu B^{\mu\nu} (H^\dagger i\overleftrightarrow{D}_\mu H)$

Table A.2: Bosonic operators in the Green's basis in the SMEFT.

Backup

The Greens basis in the SMEFT Jiang et al., 2019; Gherardi et al., 2020; Carmona et al., 2022

Four-quark	Four-lepton	Semileptonic			
$\mathcal{O}_{qq}^{(1)\alpha\beta\gamma\lambda}$	$(\overline{Q_{\alpha L}}\gamma^\mu Q_{\beta L})(\overline{Q_{\gamma L}}\gamma_\mu Q_{\lambda L})$	$\mathcal{O}_{\ell\ell}^{\alpha\beta\gamma\lambda}$	$(\overline{\ell_{\alpha L}}\gamma^\mu \ell_{\beta L})(\overline{\ell_{\gamma L}}\gamma_\mu \ell_{\lambda L})$	$\mathcal{O}_{\ell q}^{(1)\alpha\beta\gamma\lambda}$	$(\overline{\ell_{\alpha L}}\gamma^\mu \ell_{\beta L})(\overline{Q_{\gamma L}}\gamma_\mu Q_{\lambda L})$
$\mathcal{O}_{qq}^{(3)\alpha\beta\gamma\lambda}$	$(\overline{Q_{\alpha L}}\gamma^\mu \sigma^I Q_{\beta L})(\overline{Q_{\gamma L}}\gamma_\mu \sigma^I Q_{\lambda L})$	$\mathcal{O}_{ee}^{\alpha\beta\gamma\lambda}$	$(\overline{E_{\alpha R}}\gamma^\mu E_{\beta R})(\overline{E_{\gamma R}}\gamma_\mu E_{\lambda R})$	$\mathcal{O}_{\ell q}^{(3)\alpha\beta\gamma\lambda}$	$(\overline{\ell_{\alpha L}}\gamma^\mu \sigma^I \ell_{\beta L})(\overline{Q_{\gamma L}}\gamma_\mu \sigma^I Q_{\lambda L})$
$\mathcal{O}_{uu}^{\alpha\beta\gamma\lambda}$	$(\overline{U_{\alpha R}}\gamma^\mu U_{\beta R})(\overline{U_{\gamma R}}\gamma_\mu U_{\lambda R})$	$\mathcal{O}_{\ell e}^{\alpha\beta\gamma\lambda}$	$(\overline{\ell_{\alpha L}}\gamma^\mu \ell_{\beta L})(\overline{E_{\gamma R}}\gamma_\mu E_{\lambda R})$	$\mathcal{O}_{eu}^{\alpha\beta\gamma\lambda}$	$(\overline{E_{\alpha R}}\gamma^\mu E_{\beta R})(\overline{U_{\gamma R}}\gamma_\mu U_{\lambda R})$
$\mathcal{O}_{dd}^{\alpha\beta\gamma\lambda}$	$(\overline{D_{\alpha R}}\gamma^\mu D_{\beta R})(\overline{D_{\gamma R}}\gamma_\mu D_{\lambda R})$			$\mathcal{O}_{ed}^{\alpha\beta\gamma\lambda}$	$(\overline{E_{\alpha R}}\gamma^\mu E_{\beta R})(\overline{D_{\gamma R}}\gamma_\mu D_{\lambda R})$
$\mathcal{O}_{ud}^{(1)\alpha\beta\gamma\lambda}$	$(\overline{U_{\alpha R}}\gamma^\mu U_{\beta R})(\overline{D_{\gamma R}}\gamma_\mu D_{\lambda R})$			$\mathcal{O}_{qe}^{\alpha\beta\gamma\lambda}$	$(\overline{Q_{\alpha L}}\gamma^\mu Q_{\beta L})(\overline{E_{\gamma R}}\gamma_\mu E_{\lambda R})$
$\mathcal{O}_{ud}^{(8)\alpha\beta\gamma\lambda}$	$(\overline{U_{\alpha R}}\gamma^\mu T^A U_{\beta R})(\overline{D_{\gamma R}}\gamma_\mu T^A D_{\lambda R})$			$\mathcal{O}_{\ell u}^{\alpha\beta\gamma\lambda}$	$(\overline{\ell_{\alpha L}}\gamma^\mu \ell_{\beta L})(\overline{U_{\gamma R}}\gamma_\mu U_{\lambda R})$
$\mathcal{O}_{qu}^{(1)\alpha\beta\gamma\lambda}$	$(\overline{Q_{\alpha L}}\gamma^\mu Q_{\beta L})(\overline{U_{\gamma R}}\gamma_\mu U_{\lambda R})$			$\mathcal{O}_{\ell d}^{\alpha\beta\gamma\lambda}$	$(\overline{\ell_{\alpha L}}\gamma^\mu \ell_{\beta L})(\overline{D_{\gamma R}}\gamma_\mu D_{\lambda R})$
$\mathcal{O}_{qu}^{(8)\alpha\beta\gamma\lambda}$	$(\overline{Q_{\alpha L}}\gamma^\mu T^A Q_{\beta L})(\overline{U_{\gamma R}}\gamma_\mu T^A U_{\lambda R})$			$\mathcal{O}_{\ell edq}^{\alpha\beta\gamma\lambda}$	$(\overline{\ell_{\alpha L}}E_{\beta R})(\overline{D_{\gamma R}}Q_{\lambda L})$
$\mathcal{O}_{qd}^{(1)\alpha\beta\gamma\lambda}$	$(\overline{Q_{\alpha L}}\gamma^\mu Q_{\beta L})(\overline{D_{\gamma R}}\gamma_\mu D_{\lambda R})$			$\mathcal{O}_{\ell equ}^{(1)\alpha\beta\gamma\lambda}$	$(\overline{\ell_{\alpha L}}E_{\beta R})\epsilon^{ab}(\overline{Q_{\gamma L}^b}U_{\lambda R})$
$\mathcal{O}_{qd}^{(8)\alpha\beta\gamma\lambda}$	$(\overline{Q_{\alpha L}}\gamma^\mu T^A Q_{\beta L})(\overline{D_{\gamma R}}\gamma_\mu T^A D_{\lambda R})$			$\mathcal{O}_{\ell equ}^{(3)\alpha\beta\gamma\lambda}$	$(\overline{\ell_{\alpha L}}\sigma^{\mu\nu} E_{\beta R})\epsilon^{ab}(\overline{Q_{\gamma L}^b}\sigma_{\mu\nu} U_{\lambda R})$
$\mathcal{O}_{quqd}^{(1)\alpha\beta\gamma\lambda}$	$(\overline{Q_{\alpha L}^a}U_{\beta R})\epsilon^{ab}(\overline{Q_{\gamma L}^b}D_{\lambda R})$				
$\mathcal{O}_{quqd}^{(8)\alpha\beta\gamma\lambda}$	$(\overline{Q_{\alpha L}^a}T^A U_{\beta R})\epsilon^{ab}(\overline{Q_{\gamma L}^b}T^A D_{\lambda R})$				

B- and L-number violating	
$\mathcal{O}_{duq}^{\alpha\beta\gamma\lambda}$	$\epsilon^{ABC}\epsilon^{ab}\left[(D_{\alpha R}^A)^T C U_{\beta R}^B\right]\left[(Q_{\gamma L}^{Ca})^T C \ell_{\lambda L}^b\right]$
$\mathcal{O}_{qqu}^{\alpha\beta\gamma\lambda}$	$\epsilon^{ABC}\epsilon^{ab}\left[(Q_{\alpha L}^{Aa})^T C Q_{\beta L}^{Bb}\right]\left[(U_{\gamma R}^C)^T C E_{\lambda R}\right]$
$\mathcal{O}_{qqq}^{\alpha\beta\gamma\lambda}$	$\epsilon^{ABC}\epsilon^{ad}\epsilon^{be}\left[(Q_{\alpha L}^{Aa})^T C Q_{\beta L}^{Bb}\right]\left[(Q_{\gamma L}^{Ce})^T C \ell_{\lambda L}^d\right]$
$\mathcal{O}_{duu}^{\alpha\beta\gamma\lambda}$	$\epsilon^{ABC}\left[(D_{\alpha R}^A)^T C U_{\beta R}^B\right]\left[(U_{\gamma R}^C)^T C E_{\lambda R}\right]$

Table A.3: Baryon and lepton number conserving four-fermion operators in the Green's basis (and also in the Warsaw basis) in the SMEFT.

Backup

The tree-level Lagrange up to dimension-six for three seesaw mechanisms

$$\begin{aligned} \mathcal{L}_{\text{I,EFT}}^{\text{tree}} = & \mathcal{L}_{\text{SM}} + \left[\frac{1}{2} (Y_\nu M^{-1} Y_\nu^\text{T})_{\alpha\beta} \overline{\ell}_{\alpha L} \tilde{H} \tilde{H}^\text{T} \ell_{\beta L}^c + \text{h.c.} \right] + \frac{1}{4} (Y_\nu M^{-2} Y_\nu^\dagger)_{\alpha\beta} \\ & \times \left[(\overline{\ell}_{\alpha L} \gamma^\mu \ell_{\beta L}) (H^\dagger i \overleftrightarrow{D}_\mu H) - (\overline{\ell}_{\alpha L} \gamma^\mu \tau^I \ell_{\beta L}) (H^\dagger i \overleftrightarrow{D}_\mu^I H) \right], \quad (3.58) \end{aligned}$$

$$\begin{aligned} \mathcal{L}_{\text{II,EFT}}^{\text{tree}} = & \mathcal{L}_{\text{SM}} + 2\lambda_\Delta^2 \left(1 + \frac{2m^2}{M_\Delta^2} \right) (H^\dagger H)^2 - \left[\frac{\lambda_\Delta (Y_\Delta)_{\alpha\beta}}{M_\Delta} \overline{\ell}_{\alpha L} \tilde{H} \tilde{H}^\text{T} \ell_{\beta L}^c + \text{h.c.} \right] \\ & + \frac{(Y_\Delta)_{\alpha\gamma} (Y_\Delta^\dagger)_{\beta\delta}}{4M_\Delta^2} (\overline{\ell}_{\alpha L} \gamma^\mu \ell_{\beta L}) (\overline{\ell}_{\gamma L} \gamma_\mu \ell_{\delta L}) + \frac{2(4\lambda - \lambda_3 + \lambda_4 - 8\lambda_\Delta^2) \lambda_\Delta^2}{M_\Delta^2} \\ & \times (H^\dagger H)^3 + \frac{2\lambda_\Delta^2}{M_\Delta^2} \left[(H^\dagger H) \square (H^\dagger H) + 2(H^\dagger D^\mu H)^\dagger (H^\dagger D_\mu H) \right] \\ & + \frac{2\lambda_\Delta^2}{M_\Delta^2} \left[(\overline{\ell}_L Y_l H E_R + \overline{Q}_L Y_u \tilde{H} U_R + \overline{Q}_L Y_d H D_R) (H^\dagger H) + \text{h.c.} \right], \quad (3.59) \end{aligned}$$

$$\begin{aligned} \mathcal{L}_{\text{III,EFT}}^{\text{tree}} = & \mathcal{L}_{\text{SM}} + \left[\frac{1}{2} (Y_\Sigma M_\Sigma^{-1} Y_\Sigma^\text{T})_{\alpha\beta} \overline{\ell}_{\alpha L} \tilde{H} \tilde{H}^\text{T} \ell_{\beta L}^c + \text{h.c.} \right] + \frac{1}{4} (Y_\Sigma M_\Sigma^{-2} Y_\Sigma^\dagger)_{\alpha\beta} \\ & \times \left[3(\overline{\ell}_{\alpha L} \gamma^\mu \ell_{\beta L}) (H^\dagger i \overleftrightarrow{D}_\mu H) + (\overline{\ell}_{\alpha L} \gamma^\mu \tau^I \ell_{\beta L}) (H^\dagger i \overleftrightarrow{D}_\mu^I H) \right] \\ & + \left[\overline{\ell}_L Y_\Sigma M_\Sigma^{-2} Y_\Sigma^\dagger Y_l H E_R (H^\dagger H) + \text{h.c.} \right]. \quad (3.60) \end{aligned}$$

Backup

Examples for RGEs of parameters

$$\begin{aligned}\dot{\eta}_{\alpha\alpha} = & \sum_{\gamma \neq \alpha} 4 \frac{m^2}{v^2} y_{\alpha\gamma} [|\eta'_{\alpha\gamma}| \cos(\phi'_{\alpha\gamma} - \phi_{\alpha\gamma}) - 4 |\eta_{\alpha\gamma}|] |\eta_{\alpha\gamma}| + \frac{2}{3} g_2^2 \text{tr}(\eta) \\ & - \sum_i \kappa_i^2 v^2 |U_{\alpha i}|^2 + y_\alpha^2 (5\eta_{\alpha\alpha} - 3\eta'_{\alpha\alpha}) + \left(-\frac{17}{3}g_2^2 + 2T\right) \eta_{\alpha\alpha},\end{aligned}$$

$$\begin{aligned}\dot{\eta}'_{\alpha\alpha} = & \sum_{\gamma \neq \alpha} 4 \frac{m^2}{v^2} y_{\alpha\gamma} [|\eta'_{\alpha\gamma}| - 4 |\eta_{\alpha\gamma}| \cos(\phi_{\alpha\gamma} - \phi'_{\alpha\gamma})] |\eta'_{\alpha\gamma}| + \frac{2}{3} g_1^2 \text{tr}(\eta') + \frac{2}{3} (g_2^2 - g_1^2) \text{tr}(\eta) \\ & - \frac{5}{2} \sum_i \kappa_i^2 v^2 |U_{\alpha i}|^2 + y_\alpha^2 (\eta'_{\alpha\alpha} - 8\eta_{\alpha\alpha}) + \left(\frac{1}{3}g_1^2 + 2T\right) \eta'_{\alpha\alpha} - \frac{1}{3} (17g_2^2 + g_1^2) \eta_{\alpha\alpha},\end{aligned}$$

$$\begin{aligned}\dot{\phi}_{\mu\tau} = & +\kappa_1^2 v^2 |\eta_{\mu\tau}|^{-1} \{ [s_{23}c_{23}(c_{12}^2 s_{13}^2 - s_{12}^2) + s_{12}c_{12}s_{13}(c_{23}^2 - s_{23}^2)c_\delta] s_{\mu\tau} + s_{12}c_{12}s_{13}s_\delta c_{\mu\tau} \} \\ & +\kappa_2^2 v^2 |\eta_{\mu\tau}|^{-1} \{ [s_{23}c_{23}(s_{12}^2 s_{13}^2 - c_{12}^2) - s_{12}c_{12}s_{13}(c_{23}^2 - s_{23}^2)c_\delta] s_{\mu\tau} - s_{12}c_{12}s_{13}s_\delta c_{\mu\tau} \} \\ & +\kappa_3^2 v^2 |\eta_{\mu\tau}|^{-1} c_{13}^2 s_{23}c_{23}s_{\mu\tau} + \frac{3}{2} (y_\mu^2 + y_\tau^2) |\eta'_{\mu\tau}| |\eta_{\mu\tau}|^{-1} \sin(\phi_{\mu\tau} - \phi'_{\mu\tau}) \\ & +2 \frac{m^2}{v^2} |\eta_{\mu\tau}|^{-1} \{ y_{\mu\tau} (\eta_{\tau\tau} - \eta_{\mu\mu}) |\eta'_{\mu\tau}| \sin(\phi'_{\mu\tau} - \phi_{\mu\tau}) \\ & - y_{e\mu} |\eta_{e\tau}| [|\eta'_{e\mu}| \sin(\phi_{e\tau} - \phi'_{e\mu} - \phi_{\mu\tau}) - 4 |\eta_{e\mu}| \sin(\phi_{e\tau} - \phi_{e\mu} - \phi_{\mu\tau})] \\ & - y_{e\tau} |\eta_{e\mu}| [|\eta'_{e\tau}| \sin(\phi'_{e\tau} - \phi_{e\mu} - \phi_{\mu\tau}) - 4 |\eta_{e\tau}| \sin(\phi_{e\tau} - \phi_{e\mu} - \phi_{\mu\tau})] \} \\ & -2 \frac{m^2}{v^2} s_{13}c_{13}^{-1} [y_{e\mu} s_{23}^{-1} (|\eta'_{e\mu}| s'_{e\mu+\delta} - 4 |\eta_{e\mu}| s_{e\mu+\delta}) - y_{e\tau} c_{23}^{-1} (|\eta'_{e\tau}| s'_{e\tau+\delta} - 4 |\eta_{e\tau}| s_{e\tau+\delta})] \\ & -2 \frac{m^2}{v^2} y_{\mu\tau} s_{23}^{-1} c_{23}^{-1} (c_{23}^2 - s_{23}^2) (|\eta'_{\mu\tau}| s'_{\mu\tau} - 4 |\eta_{\mu\tau}| s_{\mu\tau}) \\ & +C_\kappa \zeta_{13}^{-1} s_{12} [s_{12} (y_\mu^2 - y_\tau^2) c_\rho - c_{12} s_{13} s_{23}^{-1} c_{23}^{-1} (y_e^2 - s_{23}^2 y_\mu^2 - c_{23}^2 y_\tau^2) c_{\rho+\delta}] s_\rho \\ & -C_\kappa \zeta_{13} s_{12} [s_{12} (y_\mu^2 - y_\tau^2) s_\rho - c_{12} s_{13} s_{23}^{-1} c_{23}^{-1} (y_e^2 - s_{23}^2 y_\mu^2 - c_{23}^2 y_\tau^2) s_{\rho+\delta}] c_\rho \\ & +C_\kappa \zeta_{23}^{-1} c_{12} [c_{12} (y_\mu^2 - y_\tau^2) c_\sigma + s_{12} s_{13} s_{23}^{-1} c_{23}^{-1} (y_e^2 - s_{23}^2 y_\mu^2 - c_{23}^2 y_\tau^2) c_{\sigma+\delta}] s_\sigma \\ & -C_\kappa \zeta_{23} c_{12} [c_{12} (y_\mu^2 - y_\tau^2) s_\sigma + s_{12} s_{13} s_{23}^{-1} c_{23}^{-1} (y_e^2 - s_{23}^2 y_\mu^2 - c_{23}^2 y_\tau^2) s_{\sigma+\delta}] c_\sigma,\end{aligned}$$

Backup

The standard parametrization of the mixing matrix [Workman et al. \[Particle Data Group\], 2022](#)

$$U = \begin{pmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -s_{12}c_{23} - c_{12}s_{13}s_{23}e^{i\delta} & +c_{12}c_{23} - s_{12}s_{13}s_{23}e^{i\delta} & c_{13}s_{23} \\ +s_{12}s_{23} - c_{12}s_{13}c_{23}e^{i\delta} & -c_{12}s_{23} - s_{12}s_{13}c_{23}e^{i\delta} & c_{13}c_{23} \end{pmatrix}$$

Inputs for parameters

m_u/MeV	1.2504	m_e/MeV	0.5239	g_1	0.3589	η_{ee}	1.25×10^{-3}
m_d/MeV	2.7176	m_μ/GeV	0.1104	g_2	0.6468	$\eta_{\mu\mu}$	2.21×10^{-4}
m_s/MeV	54.120	m_τ/GeV	1.8748	g_s	1.1525	$\eta_{\tau\tau}$	2.81×10^{-3}
m_c/GeV	0.6299	m_1/eV	0.05	λ	0.1235	$ \eta_{e\mu} $	1.20×10^{-5}
m_b/GeV	2.8731	m_2/eV	0.05074	m^2/GeV^2	-8672.61	$ \eta_{e\tau} $	1.35×10^{-3}
m_t/GeV	173.075	m_3/eV	0.07079	δ^q	1.144	$ \eta_{\mu\tau} $	6.13×10^{-4}
$\sin \theta_{12}^q$	0.2250	$\sin \theta_{12}$	0.5505	δ	3.438	$\phi_{e\mu}$	$\pi/3$
$\sin \theta_{23}^q$	0.04182	$\sin \theta_{23}$	0.7563	ρ	$\pi/6$	$\phi_{e\tau}$	$\pi/3$
$\sin \theta_{13}^q$	0.00369	$\sin \theta_{13}$	0.1484	σ	$\pi/4$	$\phi_{\mu\tau}$	$\pi/3$

[Alam and Martin, 2023](#)

[Workman et al, 2022](#)

[Esteban et al, 2020](#)

[Fernandez-Martinez,
Hernandez-Garcia and
Lopez-Pavon, 2016](#)

$$\eta'(\mu_M) = 2\eta(\mu_M)$$

$$v = \sqrt{-m^2/\lambda}$$

Table 2: Summary of the input values of all the relevant parameters at the benchmark energy scale $\mu_B = 200 \text{ GeV}$. See the main text for further explanations.

$\mu_M = O(M_R) = 10^4 \text{ GeV}$ but $O(Y_\nu) \sim 1$ Sizable η and η' Underlying symmetries

See, e.g., Kersten and Smirnov, 2007; Abada et al., 2007