

Prospects of nuclear-coupled dark-matter detection via correlation spectroscopy of I_2^+ and Ca^+

E.M., Gilad Perez and Ziv Meir – arXiv:2404.00616 [phys.atom-ph]

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Atomic Spectroscopy

measurements with (very) high precision

e.g.
$${}^{40}\text{Ca}^+ (4\text{s}\,{}^2\text{S}_{1/2} - 3d\,{}^2\text{D}_{5/2})$$
: [BIPM (2020)]

$$f(^{40}\text{Ca}^+) = 411\,042\,129\,776\,400.4(7)\,\text{Hz}$$
 $\left(\delta f/f = 1.8 \times 10^{-15}\right)$

• optical clock frequency comparison at 18 digits

[BACON, Nature 591 (2021)]

 $f_{\rm Yb}/f_{\rm Sr} = 1.207\,507\,039\,343\,337\,848\,2(82)$

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 $f_{\rm Yb}/f_{\rm Sr} = 1.207\,507\,039\,343\,337\,848\,2(82)$

 \implies Can be used to probe ultra-light dark matter!

Ultra-Light Dark Matter

o behaves like classical field

$$\langle N \rangle \sim n \lambda_{\rm dB}^3 \sim \frac{\rho}{m} (mv)^{-3} \gg 1 \qquad \Longrightarrow \qquad m_{\rm DM} \lesssim 1 \, {\rm eV}$$

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 \circ oscillates coherently $\left(au_{
m coh}\sim rac{2\pi}{mv^2}\sim 10^6\,$ oscillationsight)

$$\ddot{\phi} + 3H\dot{\phi} + m^2\phi^2 = 0 \implies \phi(t) \sim a^{-\frac{3}{2}}\cos(mt+\delta)$$

Nuclear-Coupled Dark Matter

$$\circ \text{ scalar } \phi: \qquad \mathcal{L} \supset -\frac{\beta_s}{2 g_s} \underbrace{\sqrt{4\pi G_N}}_{\kappa} d_g \phi G^a_{\mu\nu} G^{a \,\mu\nu} \\ \implies \qquad \Lambda_{\text{QCD}} = \Lambda^{(0)}_{\text{QCD}} \left(1 + \kappa \, d_g \, \phi\right)$$

Nuclear-Coupled Dark Matter

axion *a*:

$$\begin{array}{ll} \circ \text{ scalar } \phi & : & \mathcal{L} \supset -\frac{\beta_s}{2 \, g_s} \underbrace{\sqrt{4\pi G_N}}_{\kappa} \, d_g \, \phi \, G^a_{\mu\nu} G^{a \, \mu\nu} \\ & \Longrightarrow & \Lambda_{\rm QCD} = \Lambda^{(0)}_{\rm QCD} \left(1 + \kappa \, d_g \, \phi\right) \end{array}$$

$$\mathcal{L} \supset \frac{g_s^2}{32\pi^2} \frac{a}{f_a} \tilde{G}^a_{\mu\nu} G^{a\,\mu\nu}$$

$$\implies \qquad m_\pi^2 \sim m_\pi^{2\,(0)} \left(1 - \frac{m_u m_d}{4 \,(m_u + m_d)^2} \,\frac{a^2}{f_a^2} \right)$$

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Nuclear-Coupled Dark Matter

○ scalar
$$\phi$$
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$$\implies \quad \text{oscillating nucleon mass:} \\ \frac{\Delta m_N}{m_N} \propto d_g \, \phi \propto \cos(m_\phi t) \qquad \text{or} \qquad \frac{\Delta m}{m_N}$$

$$\frac{\Delta m_N}{m_N} \propto \frac{a^2}{f_a^2} \propto \cos(2\,m_a t)$$

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Clock Comparison Experiments



[credit: Arvanitaki, Huang, Van Tilburg, PRD 91 (2015)]

Clock Comparison Experiments



correlation spectroscopy: directly manipulate $A \otimes B$ bipartite system

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Why I_2^+ ? \implies correlation spectroscopy



dense rovibrational spectrum \implies



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Ramsey interrogation :

[Ramsey, PR 76 (1949), PR 78 (1950)]

 $\frac{\pi}{2} \text{ pulse} + \text{free evolution} + \frac{\pi}{2} \text{ pulse} + \text{measurement}$

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T

 $\cdot y$

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 $\frac{\pi}{2}$ pulse + free evolution + $\frac{\pi}{2}$ pulse + measurement

$$\varphi_i = 2\pi (f_i - f_L) T_R + \varphi_N + \tilde{\varphi}_i$$



$$|g
angle + e^{iarphi_1} |e
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 $\langle \sigma_z \rangle \propto \cos \varphi_1$

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• Correlation spectroscopy:

[Chwalla et al., APB 89 (2007)]

operate on product state of the two clocks

$$\left[\left(1 + e^{i\varphi_1} \right) |g\rangle + \left(1 - e^{i\varphi_1} \right) |e\rangle \right] \otimes \left[\left(1 + e^{i\varphi_2} \right) |g\rangle + \left(1 - e^{i\varphi_2} \right) |e\rangle \right]$$

$$\left\langle \sigma_z \otimes \sigma_z \right\rangle \propto \cos \left(\varphi_1 - \varphi_2 \right) = \cos \left[2\pi (f_1 - f_2) T_R + \Delta \tilde{\varphi} \right]$$

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Scalar Dark Matter



Axion Dark Matter



Conclusion

- $\odot~I_2^+/Ca^+$ molecular-ion vs. atomic-ion clock comparison
- comparison via correlation spectroscopy
- o for scalar ULDM:
 - ca. 2 to 3 orders of magnitude improvement compared to current clock bounds
 - strongest bounds for $m_\phi \lesssim 10^{-19}\,{
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- o for axion ULDM:
 - ca. 1 order of magnitude improvement
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Thank you for your attention!

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