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Dark linear seesaw mechanism

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Motivation

The Standard Model cannot explain:

- Neutrino flavour oscillations which imply non-zero neutrino masses
- Observed Dark Matter abundance

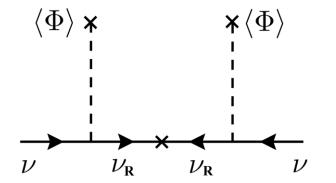
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Straightforward and elegant solution for neutrino masses:

Type-I Seesaw Model



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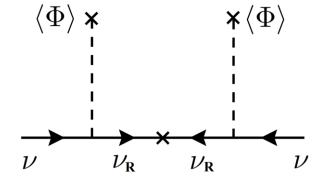
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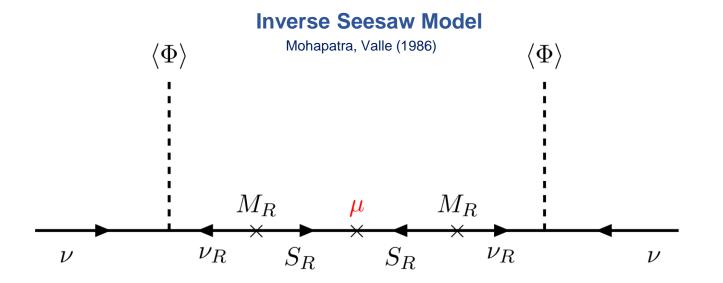
$$\mathbf{M}_{
u}\simeq -\mathbf{M}_{D}\mathbf{M}_{R}^{-1}\mathbf{M}_{D}^{T}\sim rac{v_{\Phi}^{2}}{\mathbf{M}_{R}}$$

A major drawback of this model is the **large mass scale** of the right-handed neutrinos, far away from the reach of current experiments ...

Low-scale seesaw mechanisms

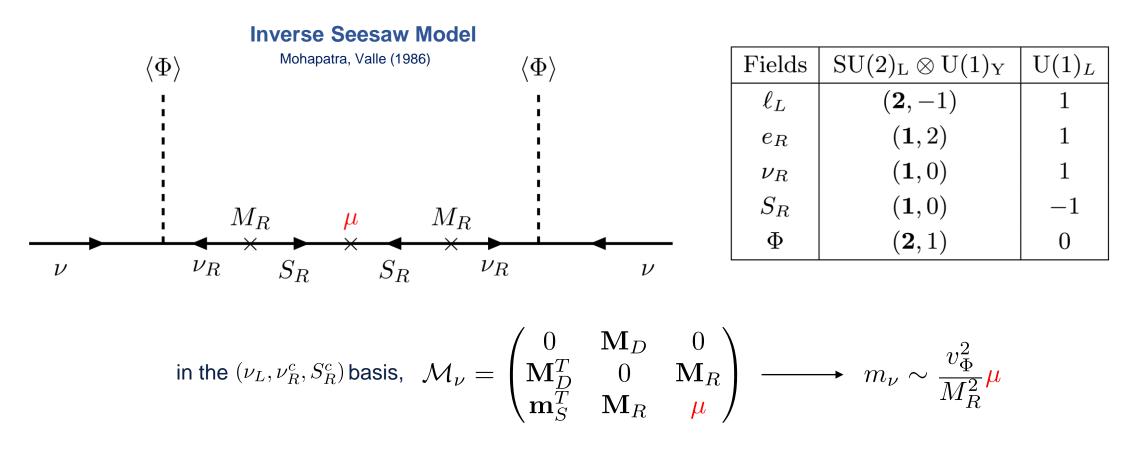
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ℓ_L	(2 ,-1)	1
e_R	(1, 2)	1
$ u_R$	(1 ,0)	1
S_R	(1 ,0)	-1
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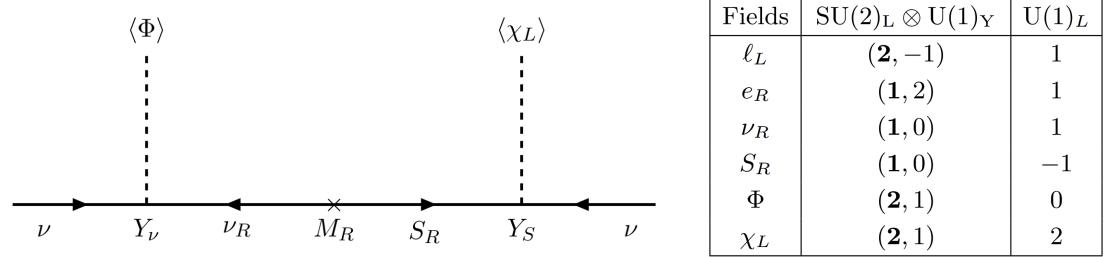
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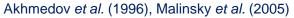


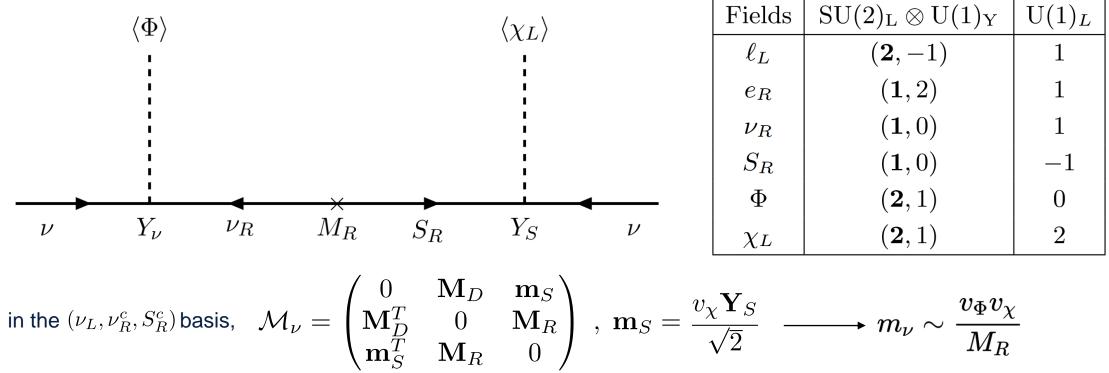
μ violates lepton number and can be naturally small in the t'Hooft sense.

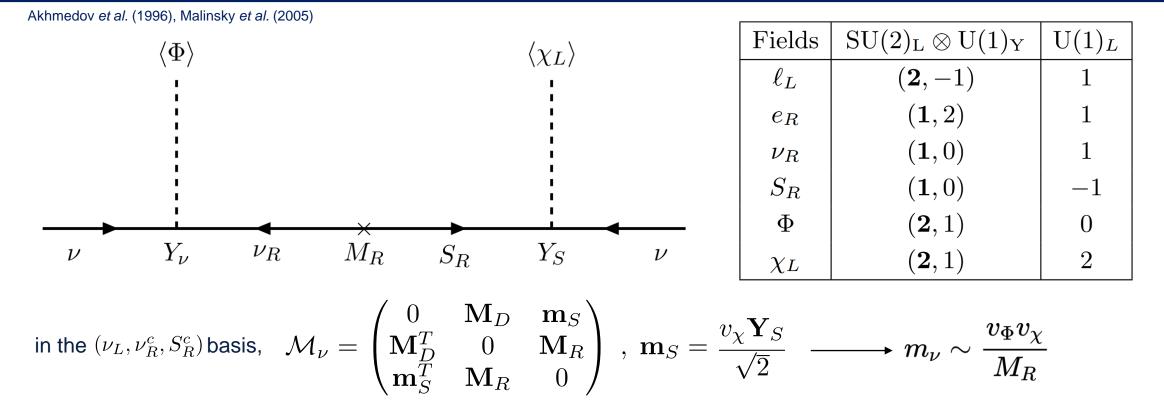
Hence, neutrino masses are suppressed, and extremely heavy mediators like in the Type-I seesaw are not required.

Akhmedov et al. (1996), Malinsky et al. (2005)

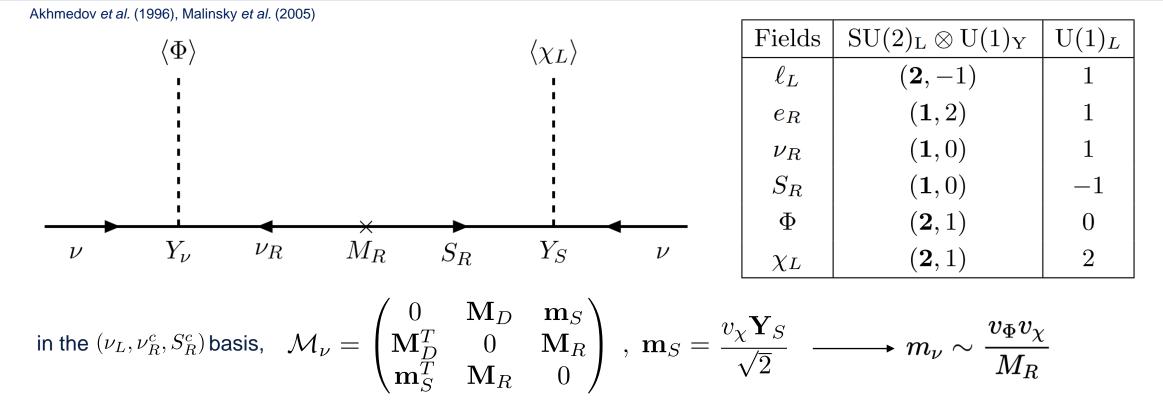








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Is there a way to incorporate **dark matter** into these models?

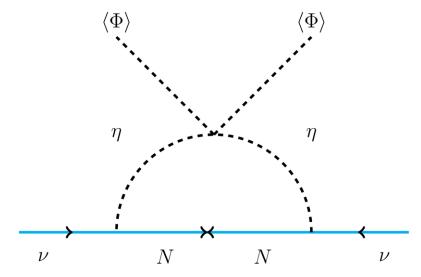
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L	$(1,2,-\frac{1}{2})$	+1
ℓ	$(1, 1, -\tilde{1})$	+1
Φ	$(1, 2, \frac{1}{2})$	+1
N	(1, 1, 0)	-1
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Scotogenic Model Tao (1996), Ma (2006)

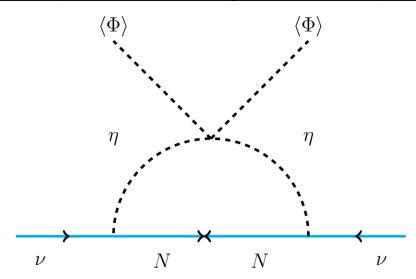
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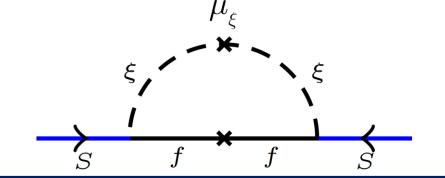
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Scotogenic Model



Dark Inverse Seesaw Model Mandal et al. (2021)

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S	(1 ,0)	-1	+
f	(1 ,0)	0	—
Φ	(2 ,1)	0	+
ξ	(1, 1)	1	



We propose a novel model where the **low-scale linear seesaw** mechanism is seeded by a **dark sector** accounting for both **neutrino flavour oscillations** and the observed **dark matter** abundance.

Henrique Brito Câmara – SUSY24 – June 13, 2024

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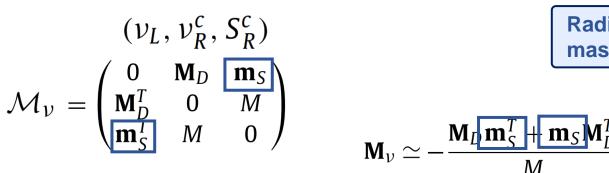
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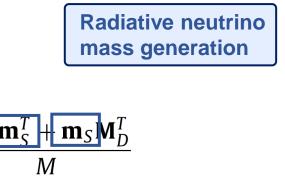
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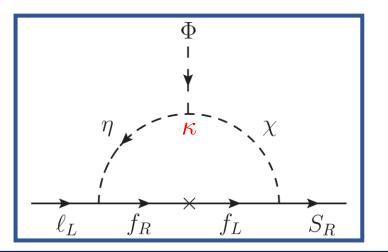
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Numerical scan

Sector	Parameters	Scan range
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We also applied the following experimental constraints:

- LEP: $m_{\zeta_i} > m_Z/2, \ m_{\zeta_i} + m_{\zeta_j} > m_Z \text{ and } m_{\zeta^+} > 70 \text{ GeV}$
- LHC Higgs data: $BR(h \to inv) \le 0.19$ and $R_{\gamma\gamma} = \frac{BR(h \to \gamma\gamma)}{BR(h \to \gamma\gamma)_{SM}} = 0.99^{+0.15}_{-0.14}$
- cLFV: ${\rm BR}(\mu \to e \gamma) < 4.2 \times 10^{-13}$ [MEG]

Dark matter

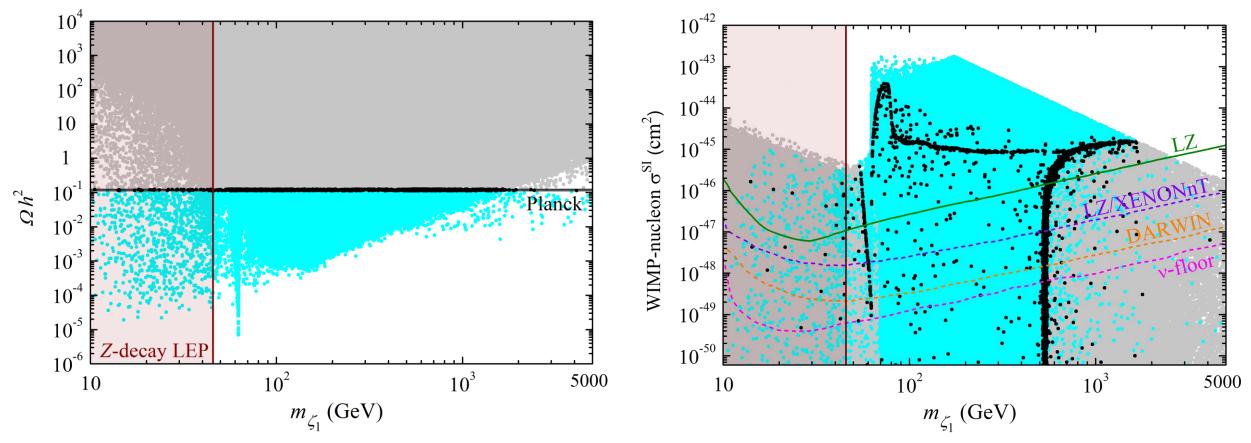
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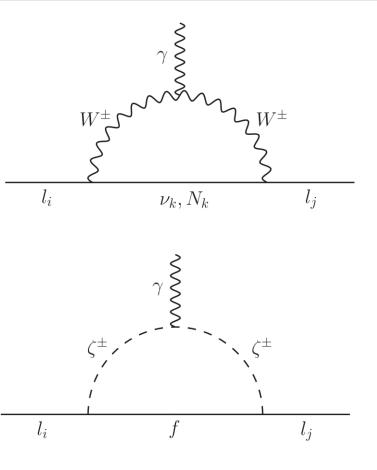
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Relic abundance

Direct detection

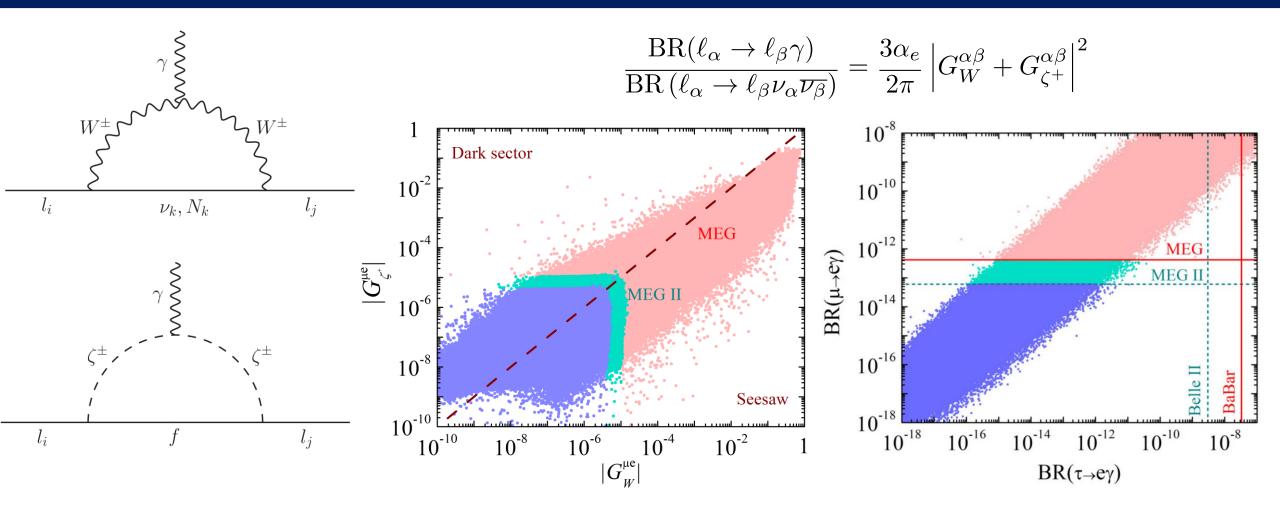


Charged lepton flavour violation



$$\frac{\mathrm{BR}(\ell_{\alpha} \to \ell_{\beta} \gamma)}{\mathrm{BR}\left(\ell_{\alpha} \to \ell_{\beta} \nu_{\alpha} \overline{\nu_{\beta}}\right)} = \frac{3\alpha_{e}}{2\pi} \left| G_{W}^{\alpha\beta} + G_{\zeta^{+}}^{\alpha\beta} \right|^{2}$$

Charged lepton flavour violation



The new particles can mediate **charged lepton flavour violating decays** with sizable branching ratios. The model can be **tested/probed** through these processes at various current and upcoming experiments.

Concluding remarks

- The **Type-I Seesaw** is by far the **simplest solution** to the neutrino mass problem. However, the large mass scale of the righthanded neutrinos makes it far away from the reach of current experiments.
- Low-scale solutions, such as the inverse and linear seesaws, despite being more complicated offer more testability prospects at ongoing experiments.
- Small neutrino masses can also be induced at the quantum level in radiative neutrino mass models. Particles entering the loops may also be viable DM candidates stabilized by some symmetry as a simple Z_2 .
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