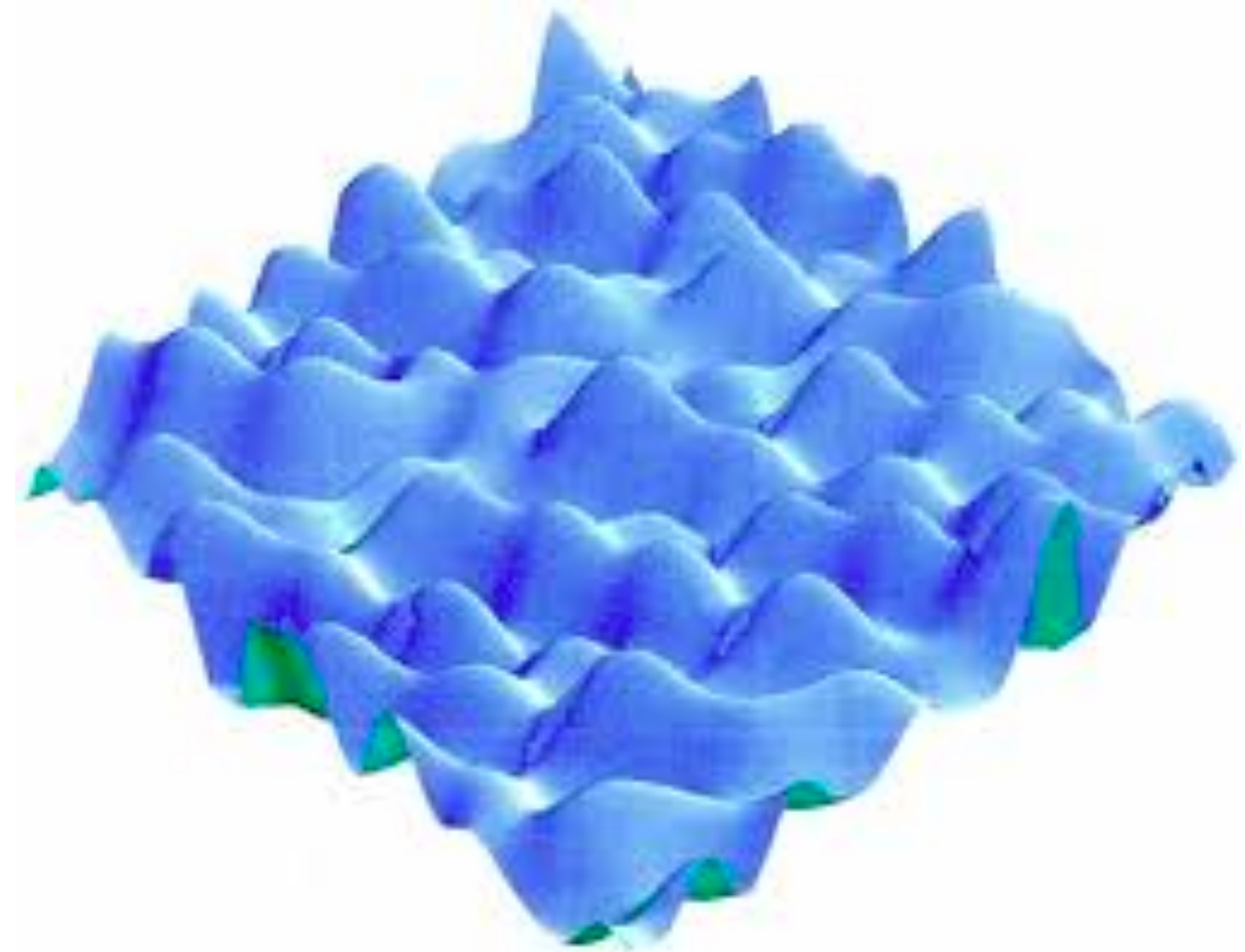


Weak scale SUSY emergent from the string landscape

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University of Oklahoma

Based on arXiv:2202.07046, 2206.14839

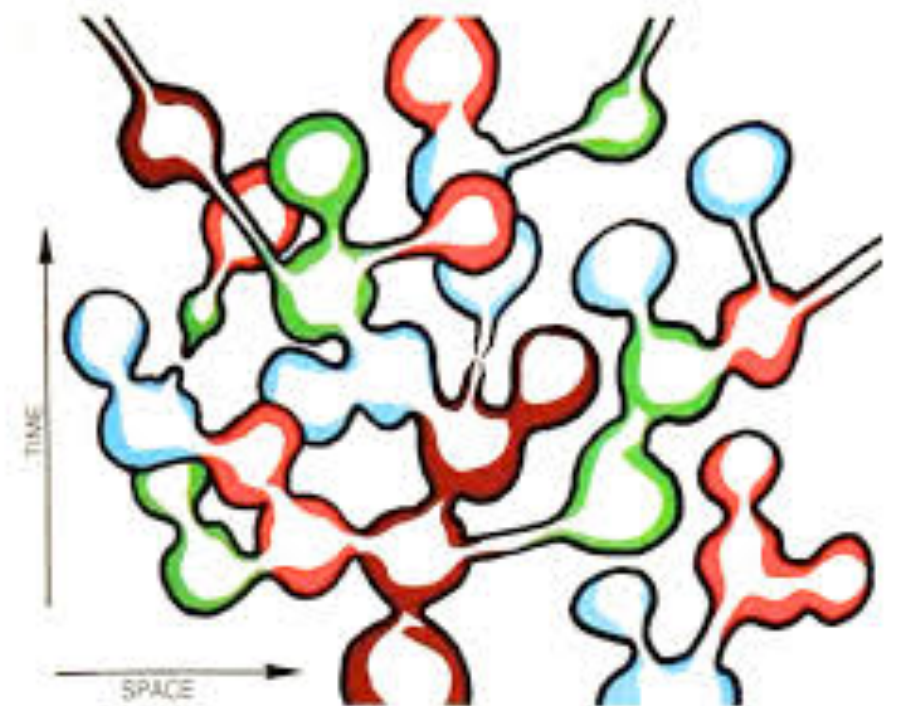
with V. Barger, D. Martinez and S. Salam



Popular review: Bousso&Polchinski, Sci. Am., 291 (2004) 60.

String theory provides most promising setting for unifying quantum gravity with SM

- 1974- string theory proposed as theory of quantum gravity
- 1984- 1st string revolution: anomalies, heterotic string compactified on CY manifold
- 1995-1996: 2nd string revolution: D-branes and string dualities
- 2001: **3rd string revolution**: emergence of string landscape (Bousso&Polchinski, Douglas et al., Susskind); $\sim 10^{500}$ flux compactifications \Rightarrow eternal inflation model with each pocket universe with different 4-d laws of physics:
- Setting for Weinberg's anthropic solution to CC problem

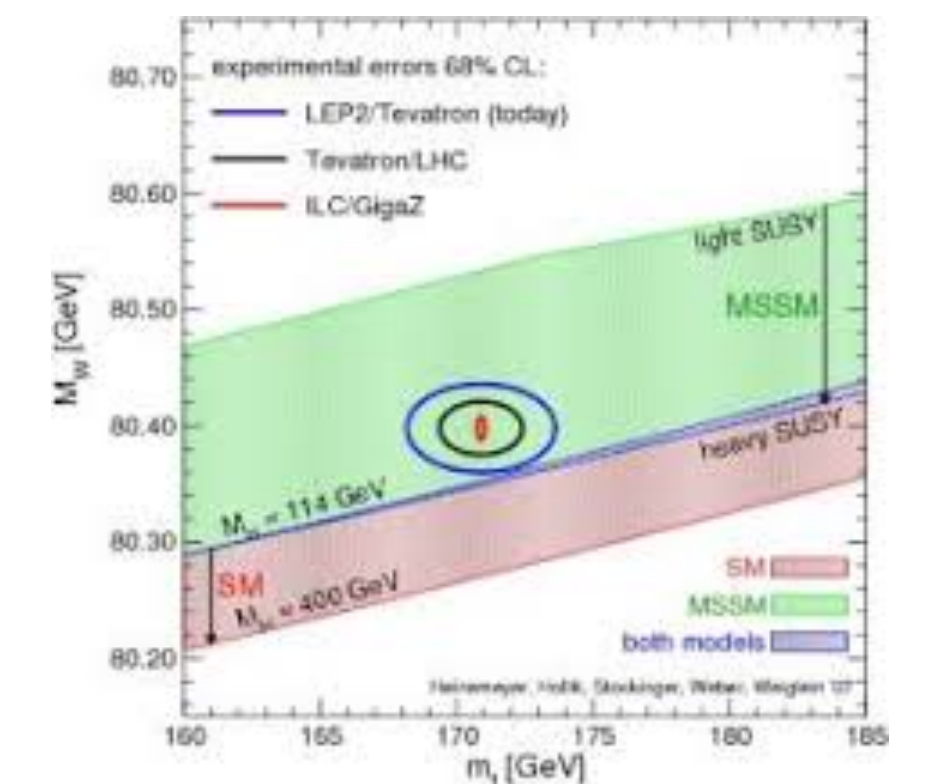
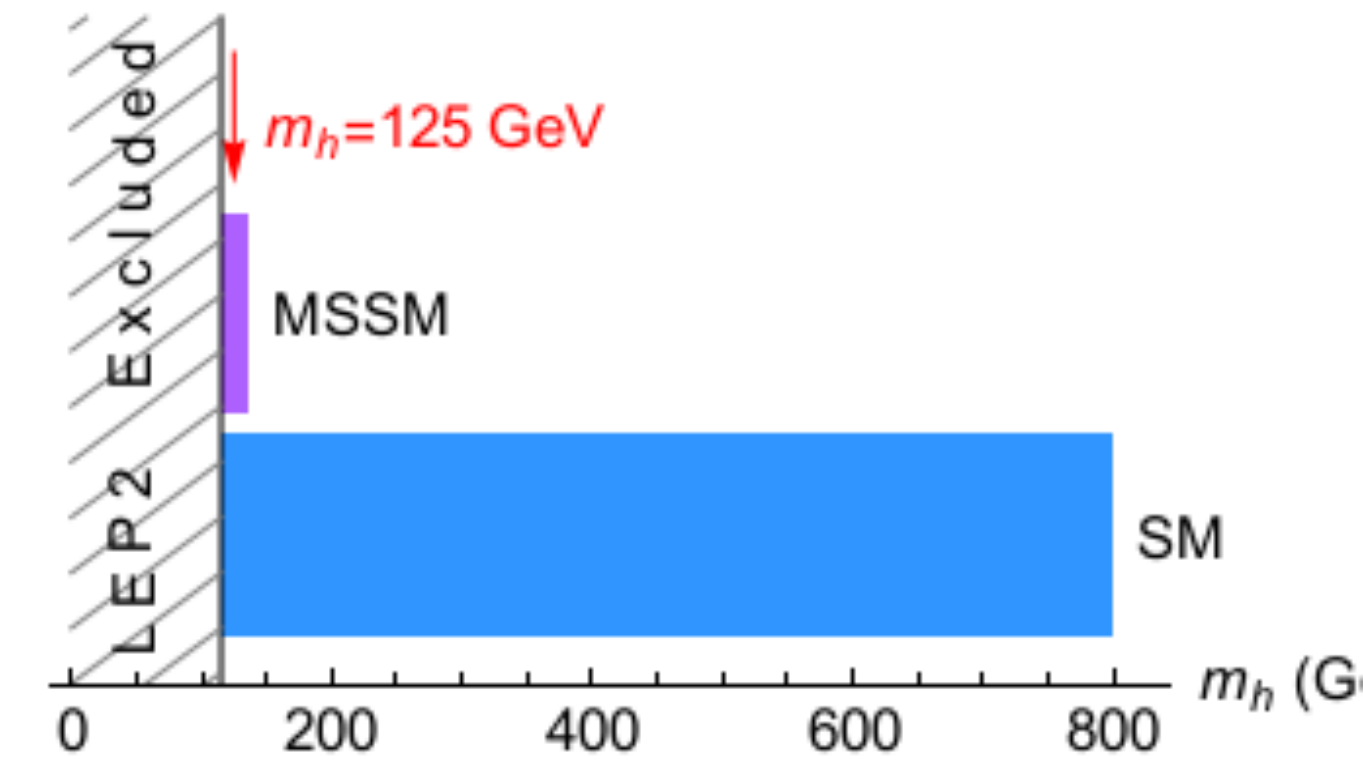
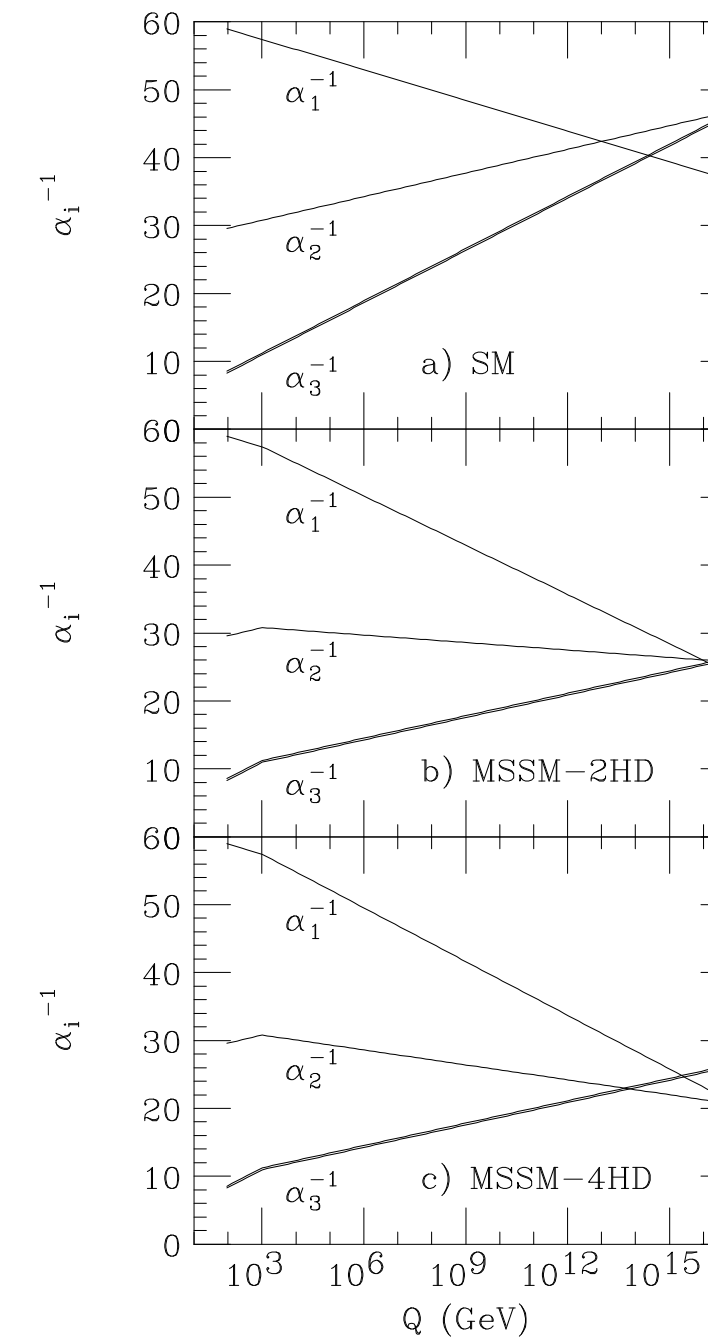


landscape phenomenology

- Candelas et al. (1984) compactification on Ricci-flat manifold with special holonomy (CY 3-fold) => conserved Killing spinor=> leaves some remnant spacetime SUSY in LE-EFT
- Acharya conjecture: only such compactifications lead to stable universes- otherwise Witten's bubble of nothing
- Dine et al.: metastable vacua actually stable under SUSY; broken SUSY leaves lifetime longer than age of universe $\Gamma \sim \exp(-m_P^2/m_{3/2}^2)$
- hard (but perhaps not impossible) to construct viable compactifications without remnant spacetime SUSY
- But is SUSY broken at some high scale $\gg m(\text{weak})$ or $\sim m(\text{weak})$ in which case solution to gauge hierarchy problem

WSS supported by virtual quantum effects

- gauge coupling unification
- $m(\text{top})$ in right range for radiative EWSB
- $m(\text{Higgs}) \sim 125 \text{ GeV}$ in accord with MSSM
- precision EW (Sven): prefers heavy SUSY over SM



WSS from the landscape: the setup

- flux compactification leaving remnant N=1 spacetime SUSY
- moduli stabilized ala KKLT via flux and non-perturbative effects
- parsimony: assume SM+SUSY (=MSSM) plus e.g. possible PQ sector/moduli in LE-EFT
- friendly/predictive patch of landscape: only CC and SUSY breaking scale scan (ADK)
- different dependence of soft terms on hidden sector: gaugino masses, scalar masses, A and B terms effectively scan independently
- Douglas: CC selection doesn't impact SUSY breaking selection

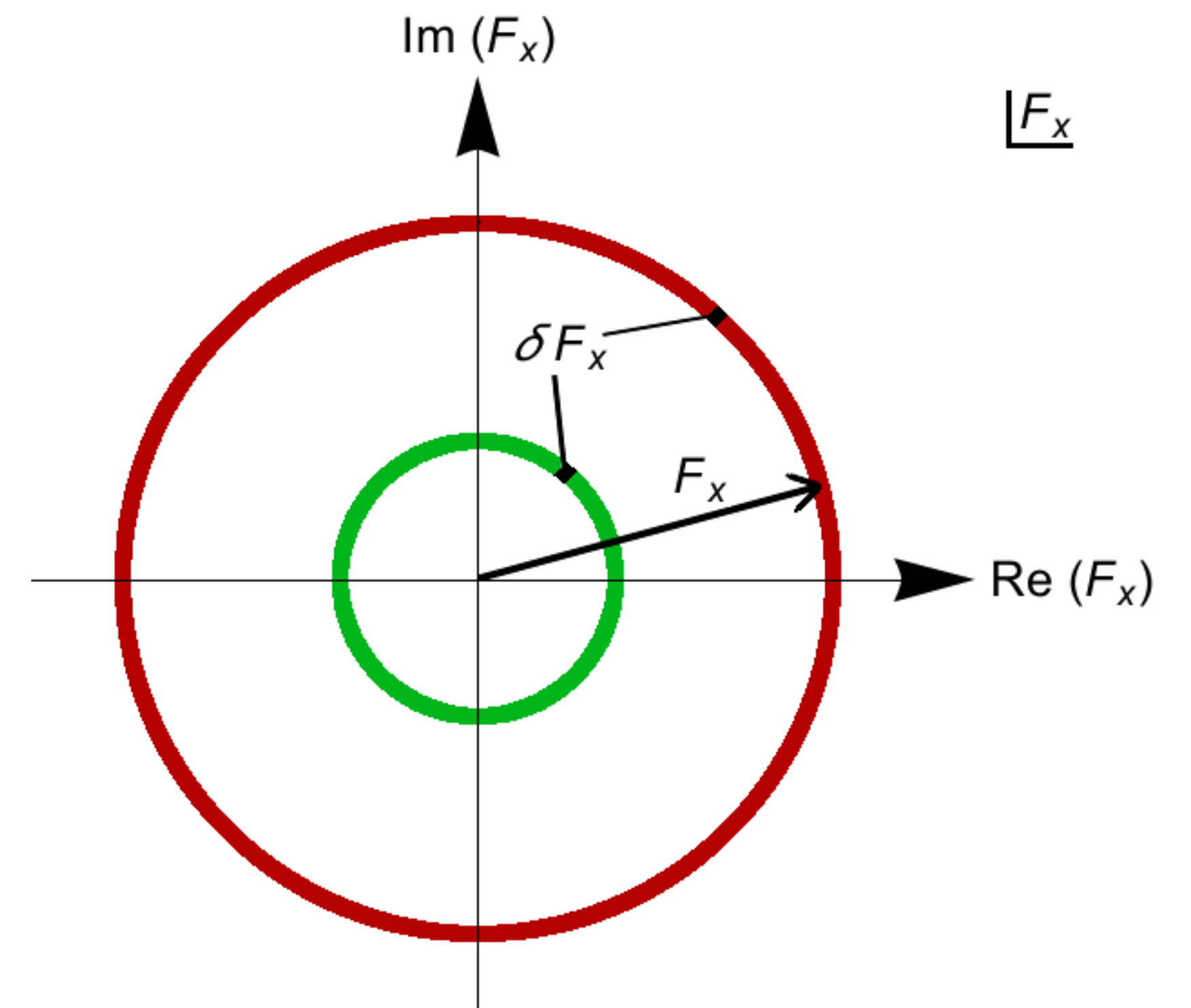
$$dN_{vac} \sim f_{SUSY} \cdot f_{EW SB} \cdot dm_{SUSY}^2$$

Gravity/moduli mediated SUSY breaking

$$m_{hidden}^4 = \sum_i F_i^\dagger F_i + \sum_\alpha D_\alpha D_\alpha \sim (10^{11} \text{ GeV})^4$$

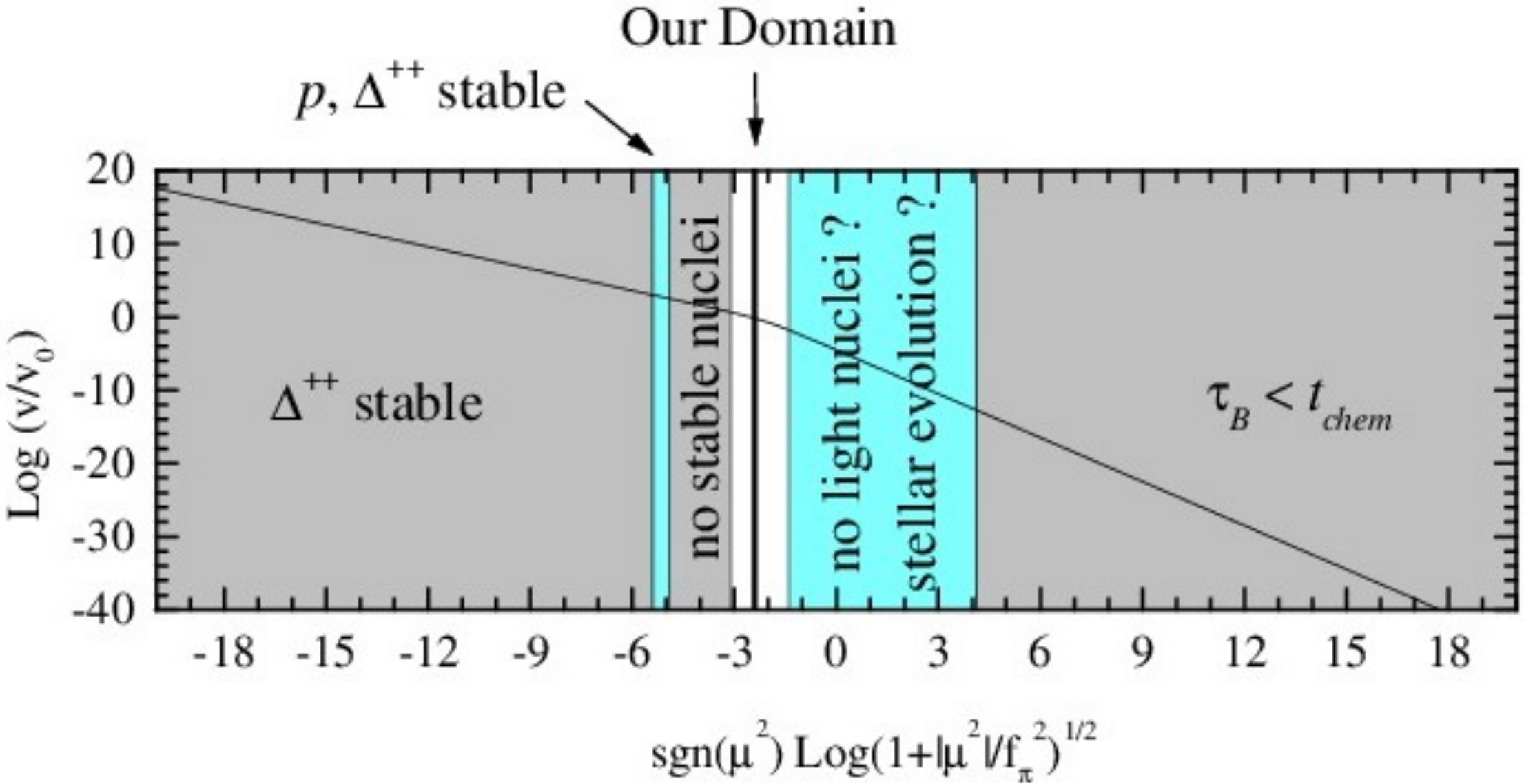
$$m_{soft} \sim m_{hidden}^2 / m_P \sim 1 \text{ TeV}$$

- textbook case SUSY breaking via single F-term **distributed randomly as complex number**: linear draw to large soft terms
- more SUSY breaking fields: stronger power-law draw to large soft terms (Douglas)



$$f_{SUSY} \sim m_{soft}^{2n_F + n_D - 1}$$

Agrawal, Barr, Donoghue, Seckel (ABDS) window (1997)



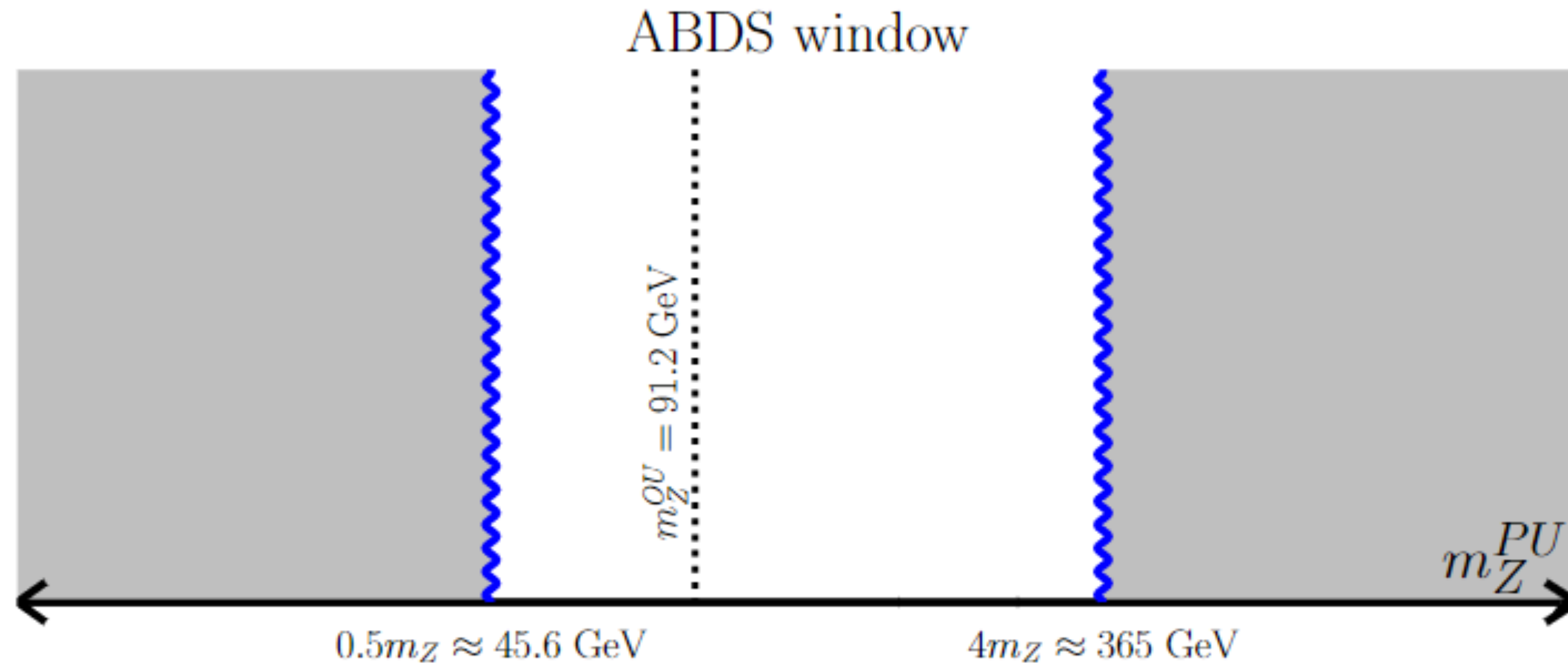
for complex nuclei (hence atoms) to exist (atomic principle/anthropic requirement), then

$$m_{weak}^{OU}/2 < \sim m_{weak}^{PU} < \sim (2 - 5)m_{weak}^{OU}$$

value of mweak in pocket universe close to the value measured in our universe; otherwise, no atoms as we know them!

an underrated prediction of MSSM: magnitude of weak scale

$$m_Z^2/2 = \frac{m_{H_d}^2 + \Sigma_d^d - (m_{H_u}^2 + \Sigma_u^u) \tan^2 \beta}{\tan^2 \beta - 1} - \mu^2 \simeq -m_{H_u}^2 - \mu^2 - \Sigma_u^u(\tilde{t}_{1,2})$$



(usually this relation is finetuned in μ or m_{H_u} to make $m(Z) \sim 91.2 \text{ GeV}$)

Recent work: place on more quantitative footing:
scan soft SUSY breaking parameters in NUHM3 model
as $m(\text{soft})^n$ along with $f(\text{EWFT})$ penalty^{*,**}

We scan according to m_{soft}^n over:

- $m_0(1, 2) : 0.1 - 40 \text{ TeV},$
- $m_0(3) : 0.1 - 20 \text{ TeV},$
- $m_{1/2} : 0.5 - 10 \text{ TeV},$
- $A_0 : 0 - -60 \text{ TeV},$
- $m_A : 0.3 - 10 \text{ TeV},$
- $\tan \beta : 3 - 60 \quad (\text{flat})$

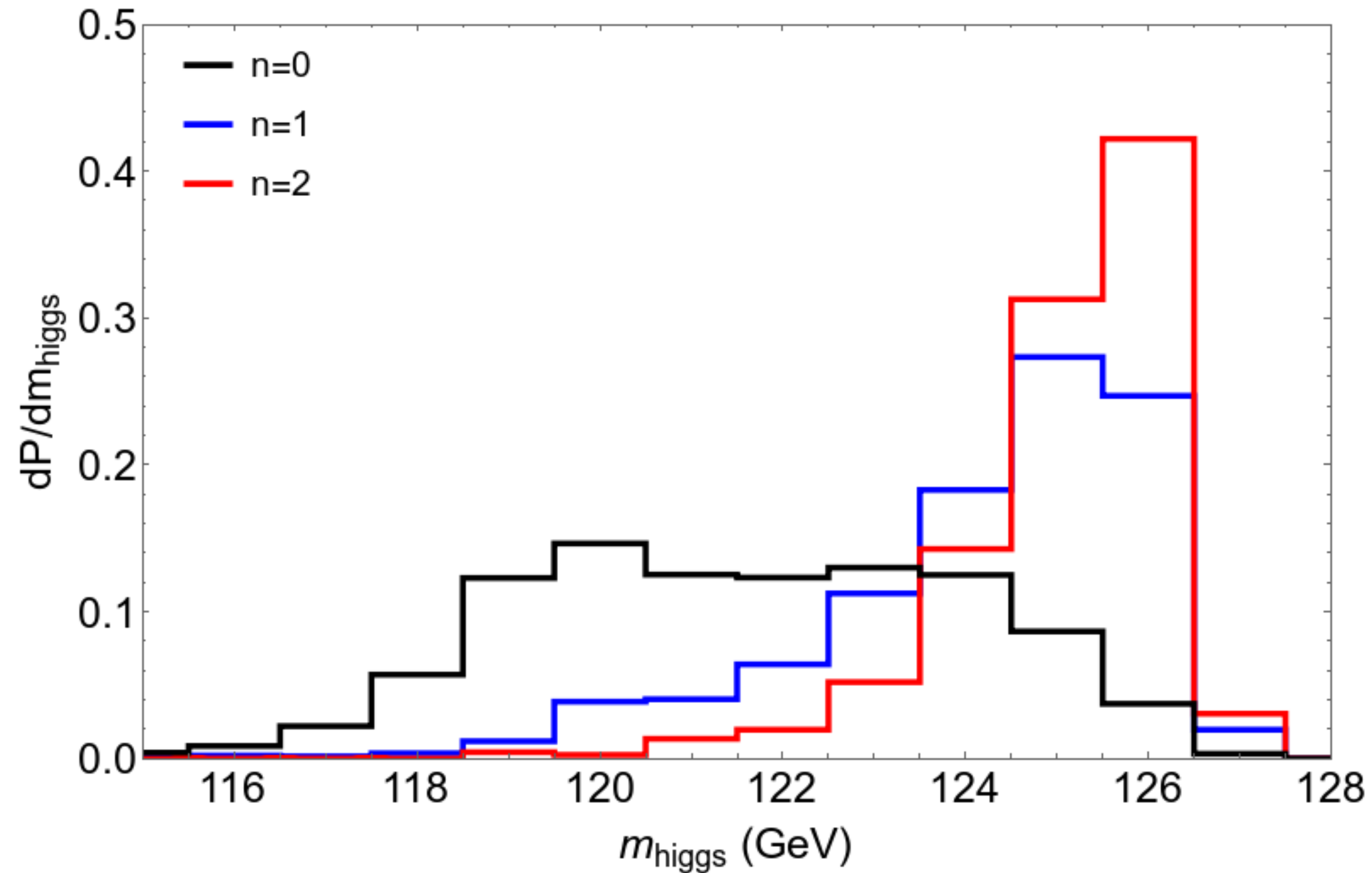
$\mu=150 \text{ GeV}$ (fixed)

(*upper limits chosen beyond those set by ABDS window)

**then rescale to $m_Z=91.2 \text{ GeV}$ to gain predictions for Our Universe (OU))

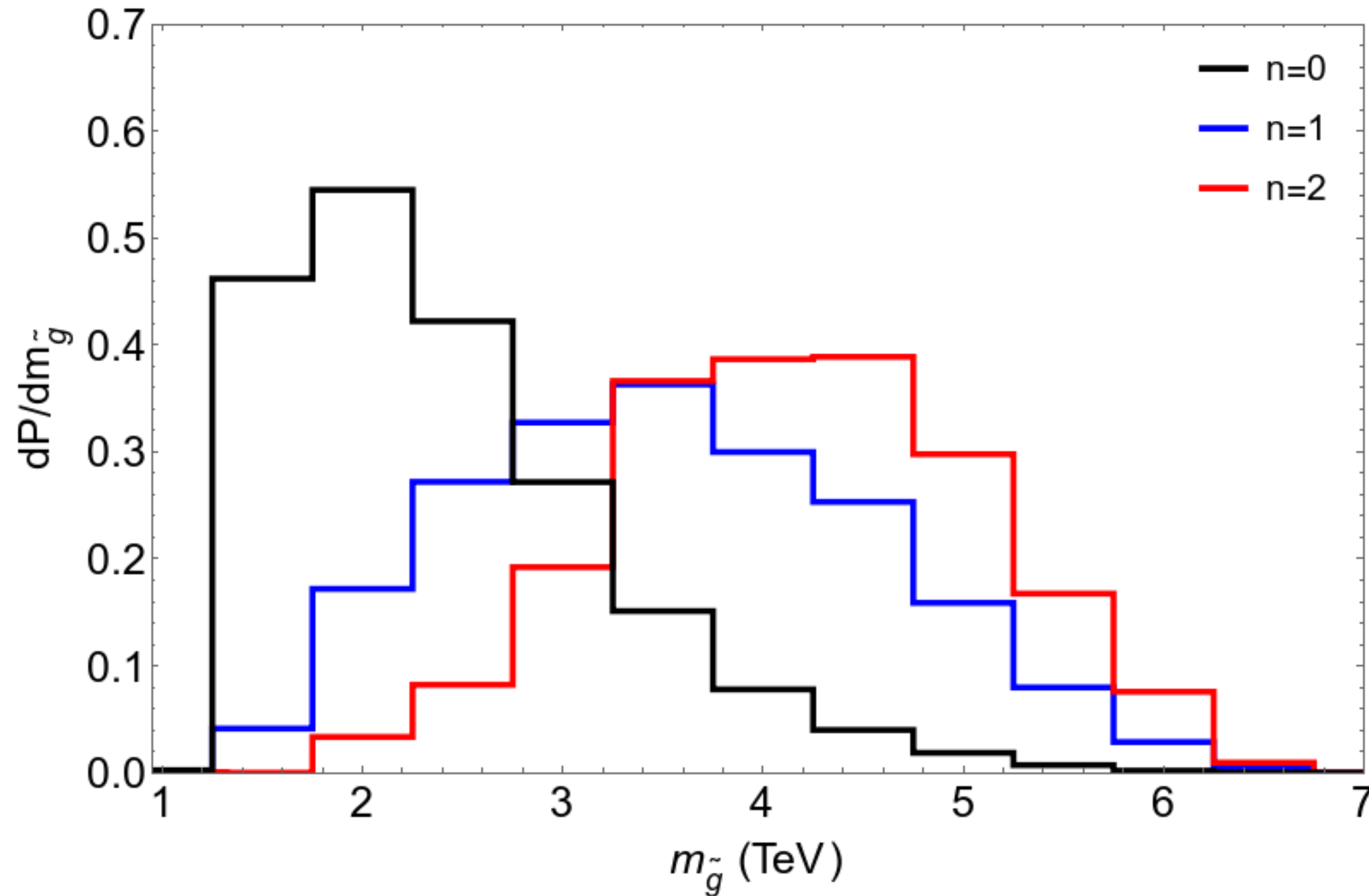
Making the picture more quantitative:

$$dN_{vac}[m_{hidden}^2, m_{weak}, \Lambda] = f_{SUSY}(m_{hidden}^2) \cdot f_{EWFT} \cdot f_{cc} dm_{hidden}^2$$



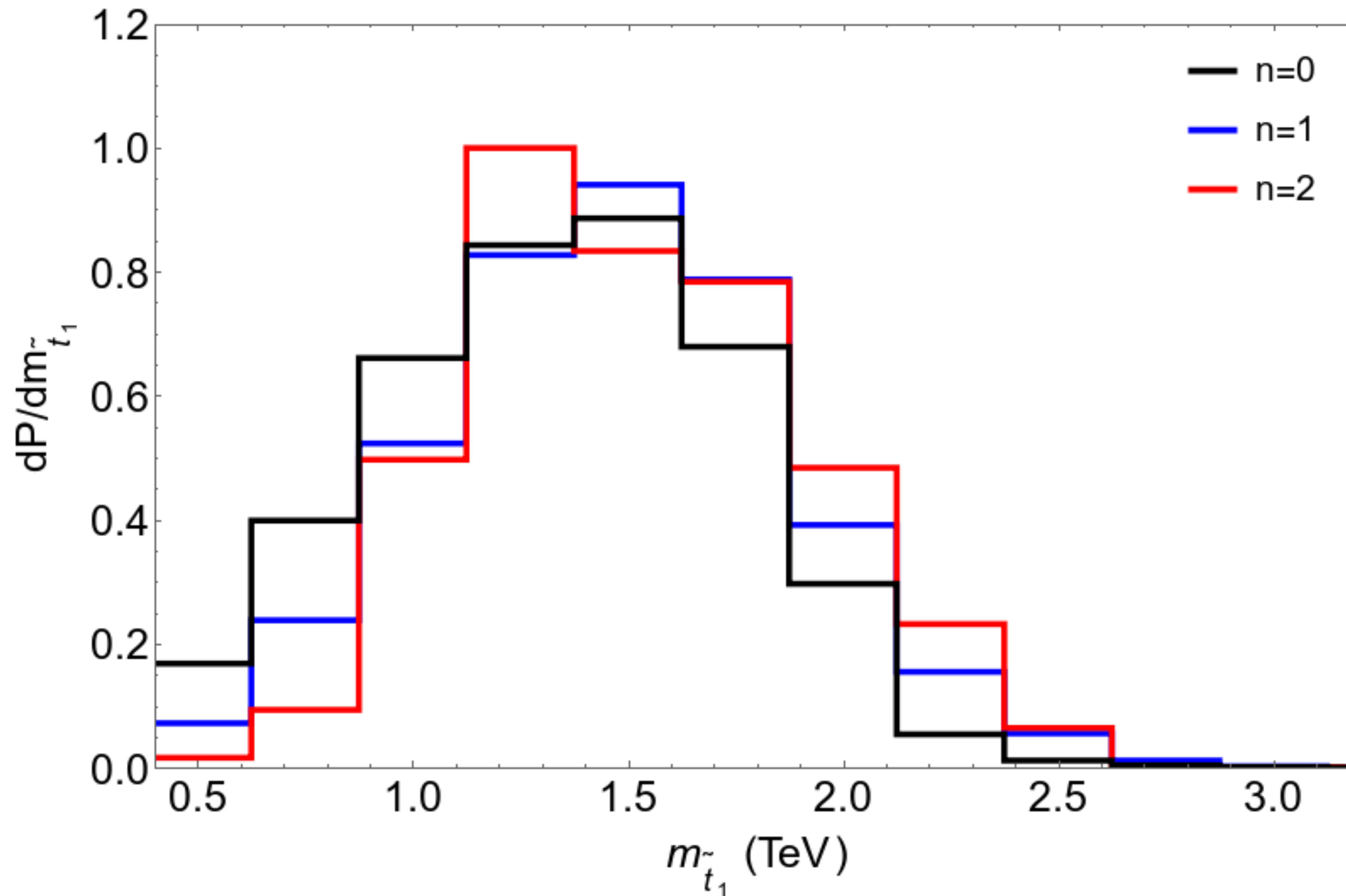
$m(h) \sim 125$ most favored for $n=1,2$

What is corresponding distribution for gluino mass?



gluino typically beyond LHC 14 reach
(need higher energy hadron collider)

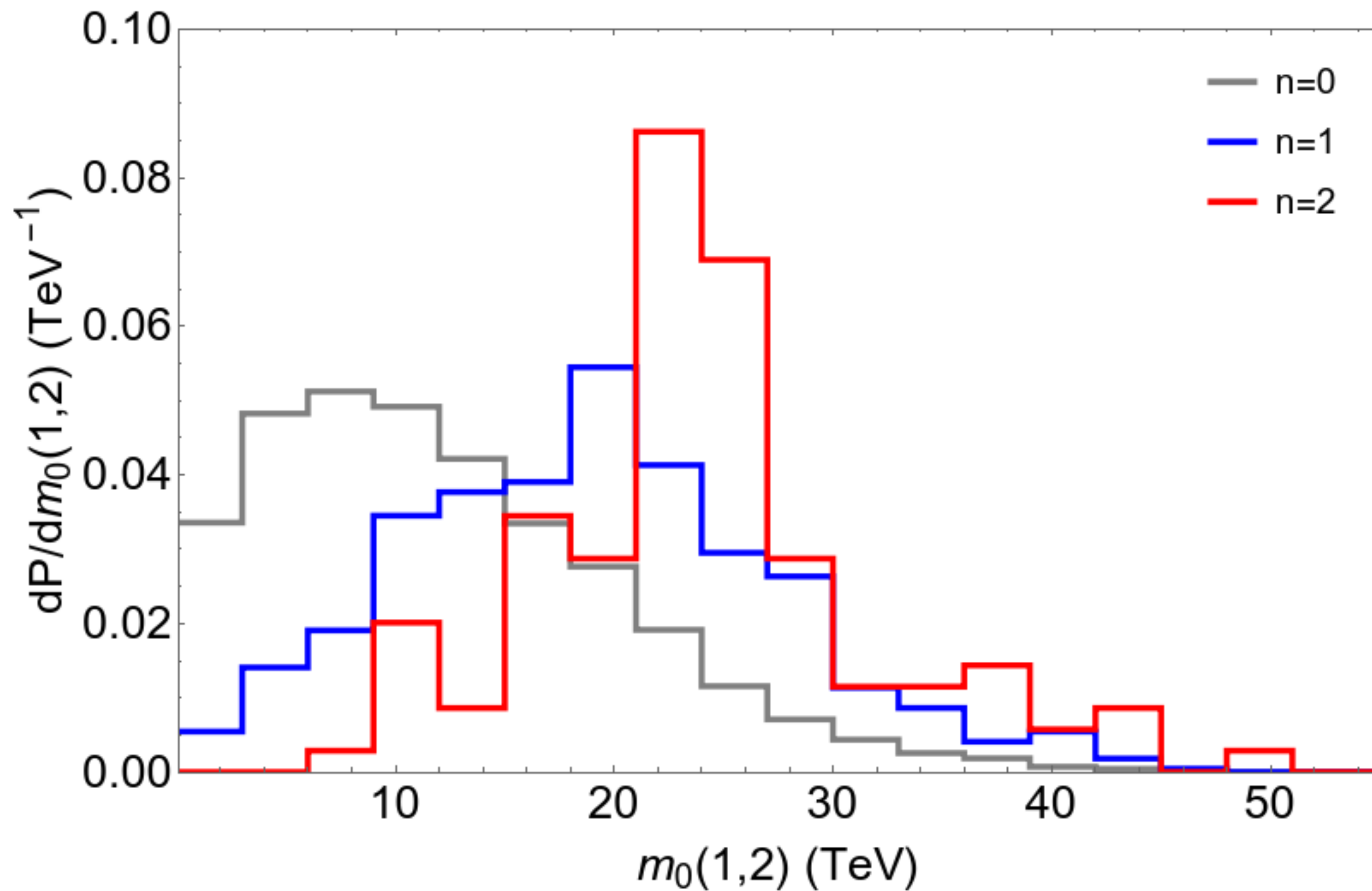
and top-squark mass $m(t_1)$?



$m(t_1)$ typically beyond present LHC reach

first/second generation sfermions pulled to 10–40 TeV:
landscape mixed decoupling/quasi-degeneracy sol'n
to SUSY flavor/CP problems

HB, Barger, Sengupta, arXiv:1910.00090



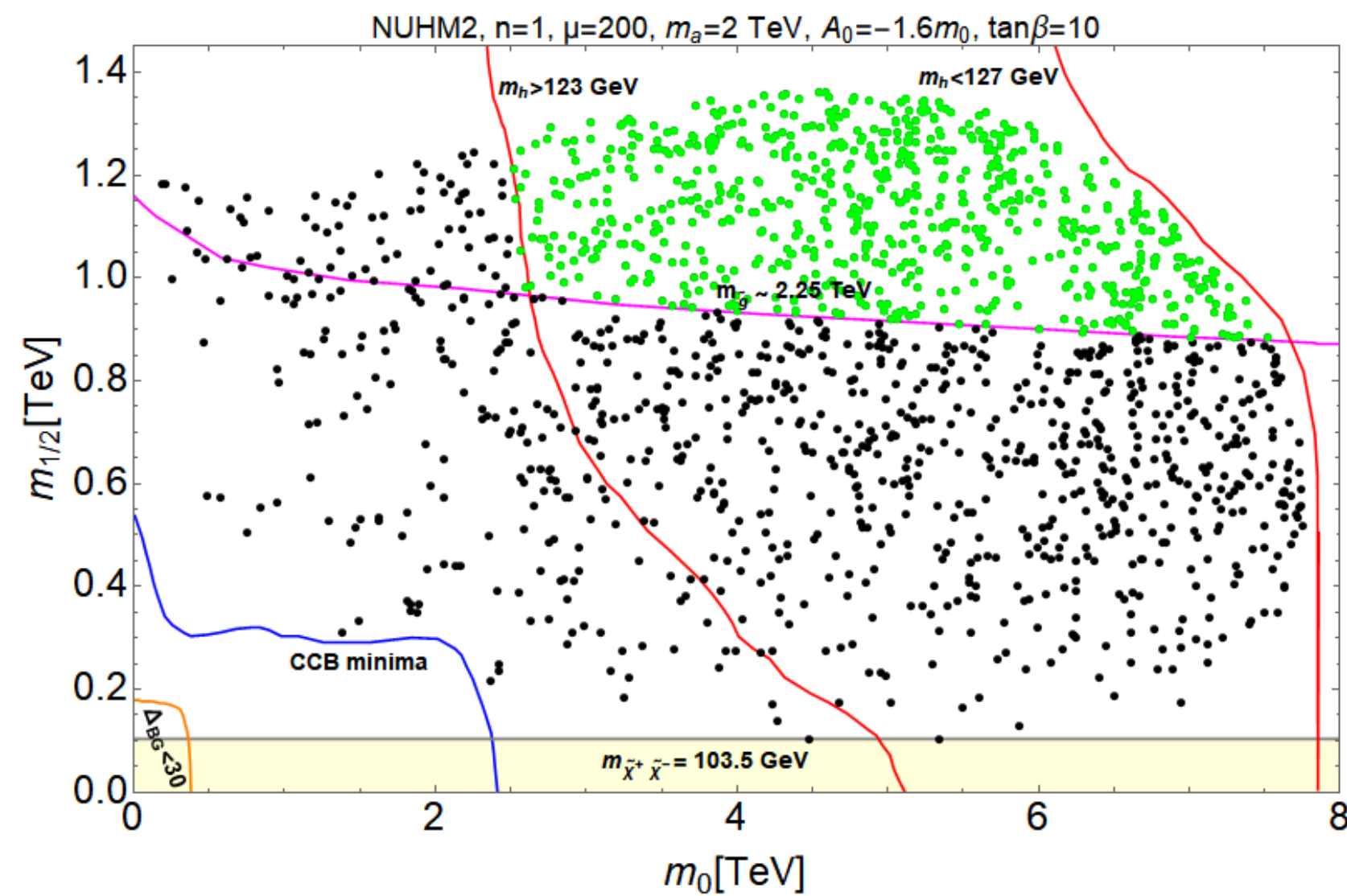
Stringy naturalness (Douglas): higher density of points are more stringy natural!

conventional natural: favor low m_0, m_h

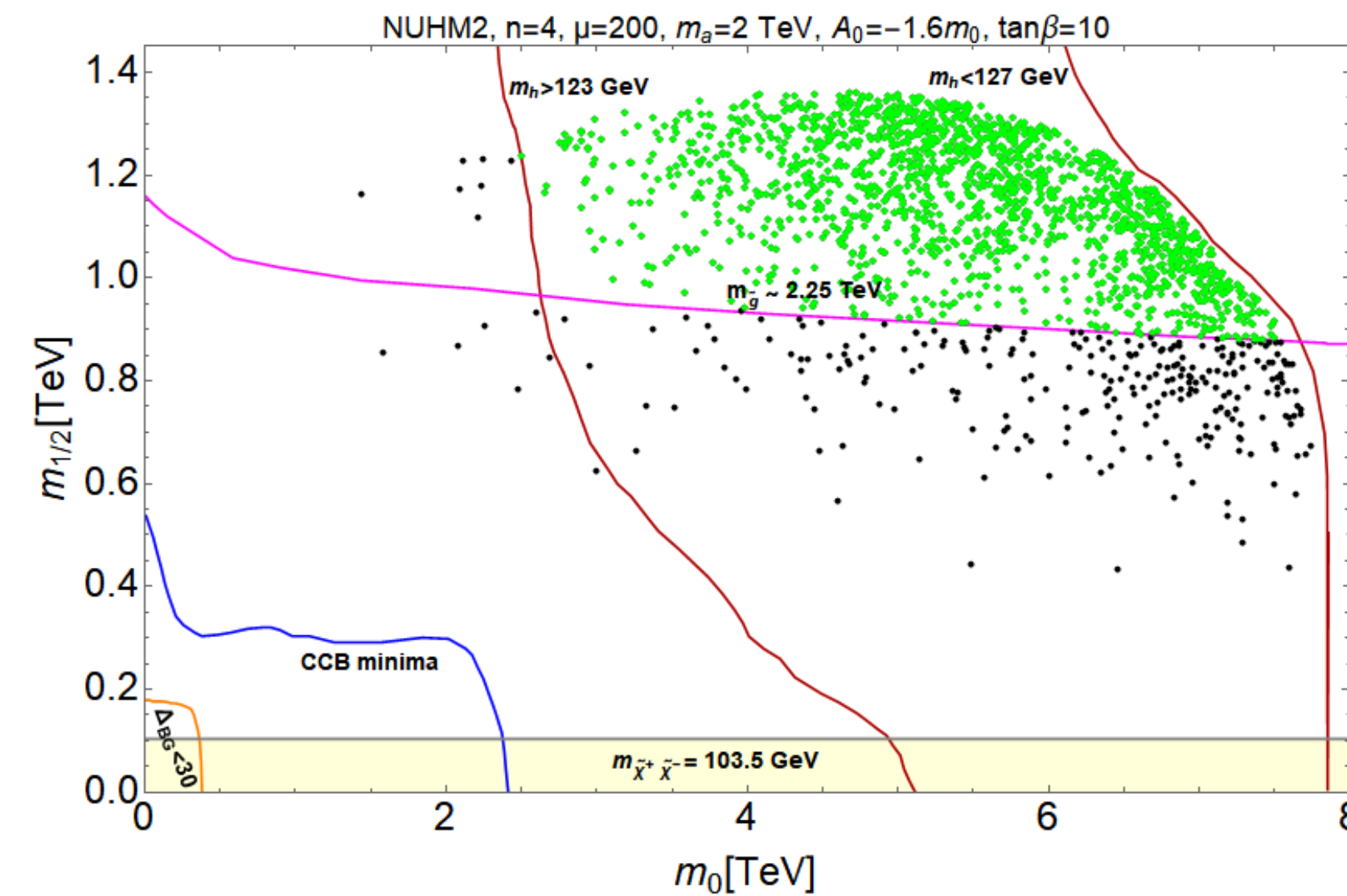
stringy naturalness: favor high m_0, m_h so long as $m(\text{weak}) \sim 100 \text{ GeV}$

HB, Barger, Salam, arXiv:1906.07741

Living dangerously: Arkani-Hamed, Dimopoulos, Kachru, hep-ph/0501082



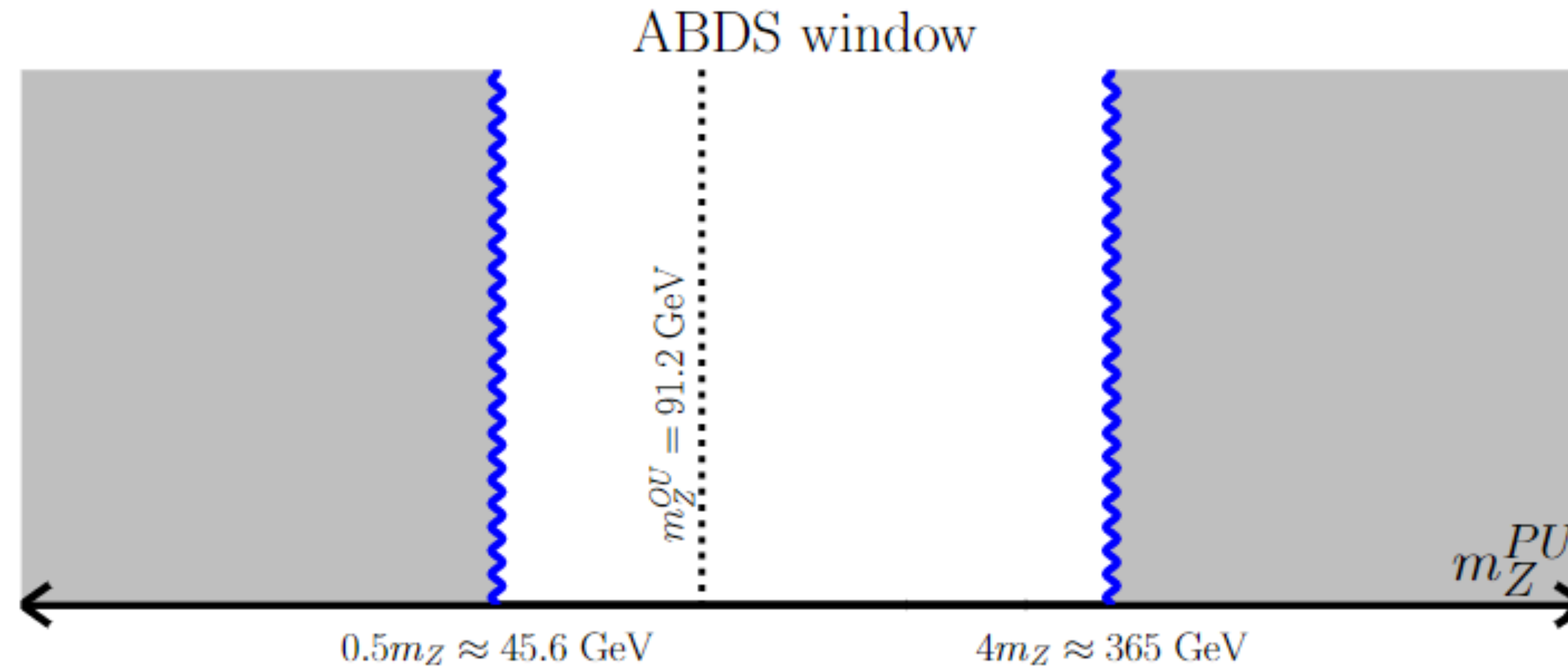
$$m(\text{soft})^1$$



$$m(\text{soft})^4$$

Under stringy naturalness, a 3 TeV gluino is more natural than a 300 GeV gluino!

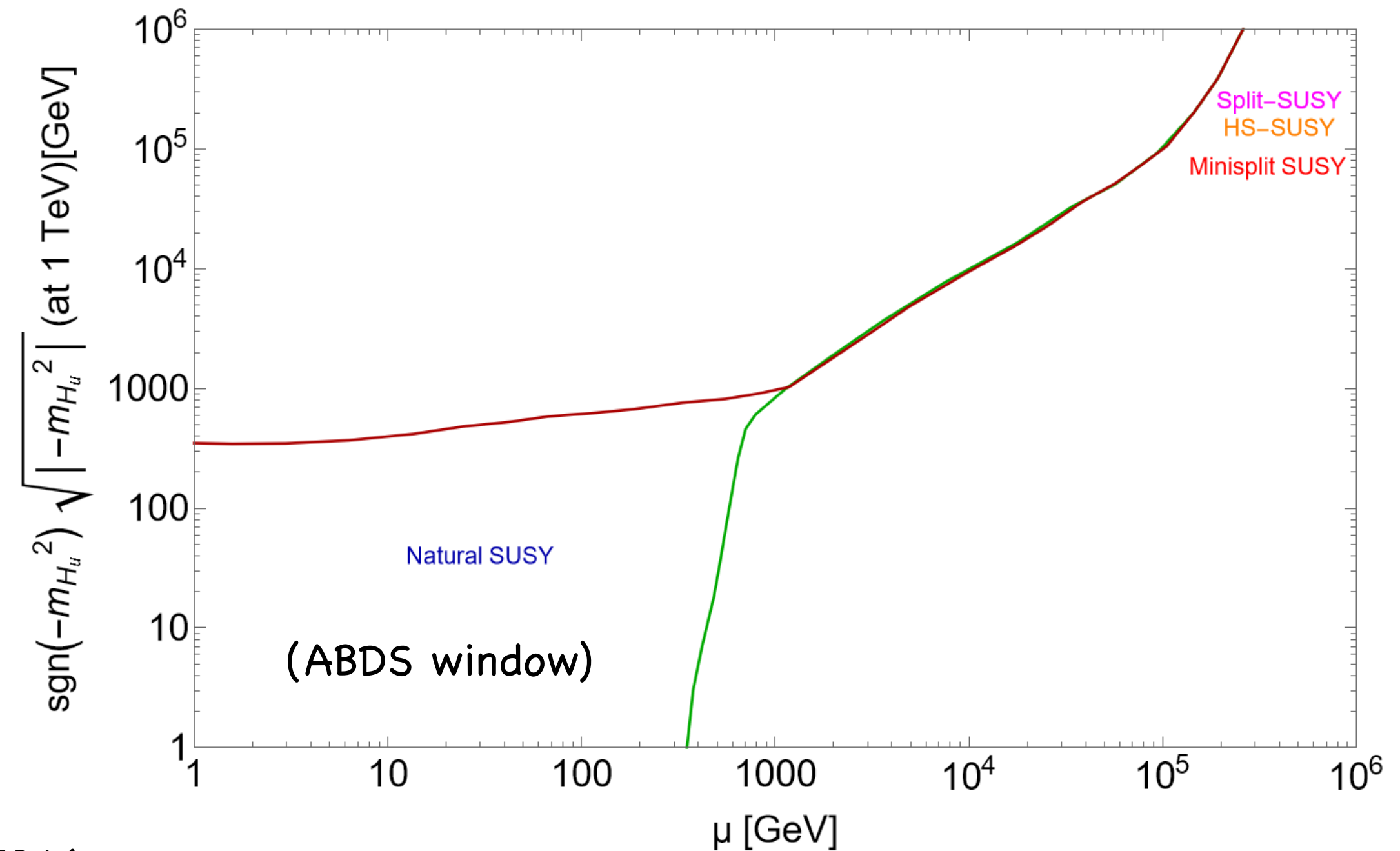
why natural SUSY more likely than finetuned SUSY from landscape



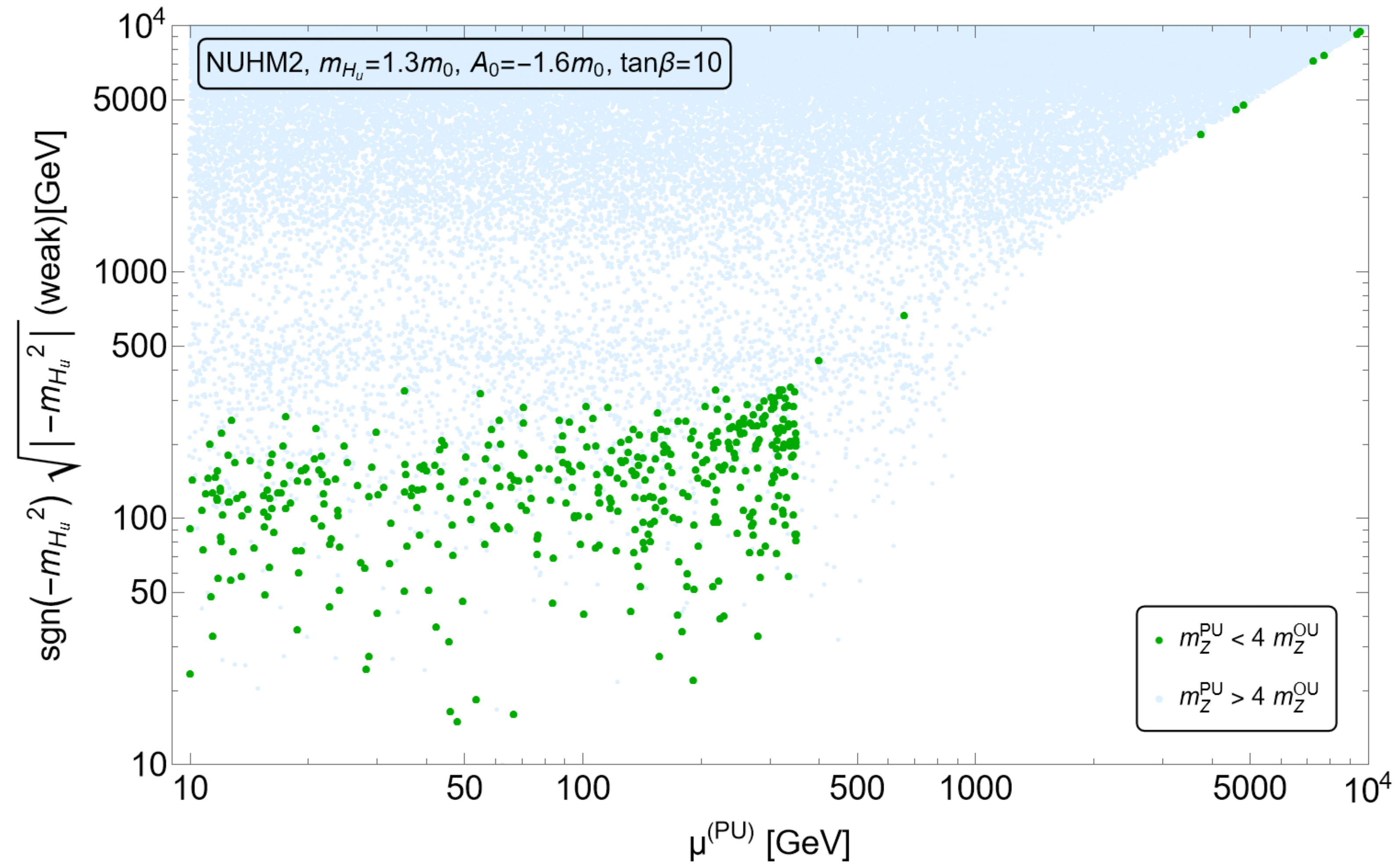
ABDS window in μ vs. m_{H_u} plane

volume of scan space shrinks to teensy values for finetuned SUSY

radiatively-driven natural SUSY (RNS) most likely to emerge from landscape!

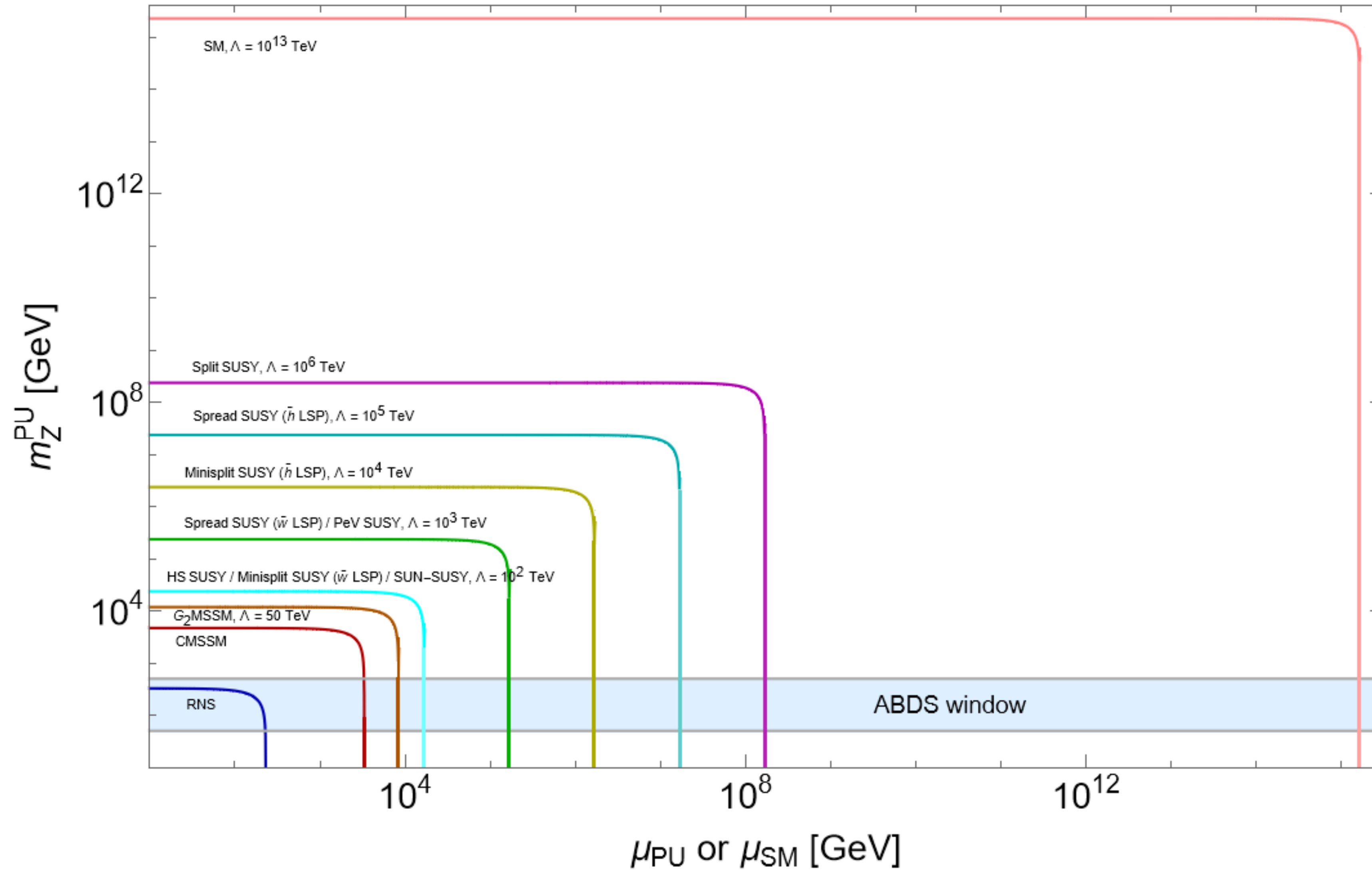


toy multiverse selection shows it is rare to generate finetuned solutions compared to natural spectra



no mu scale preferred: mu distributed uniformly across the decades

'length' of range of mu s.t. $m_Z(\text{PU})$ within ABDS window can be simple measure of probability

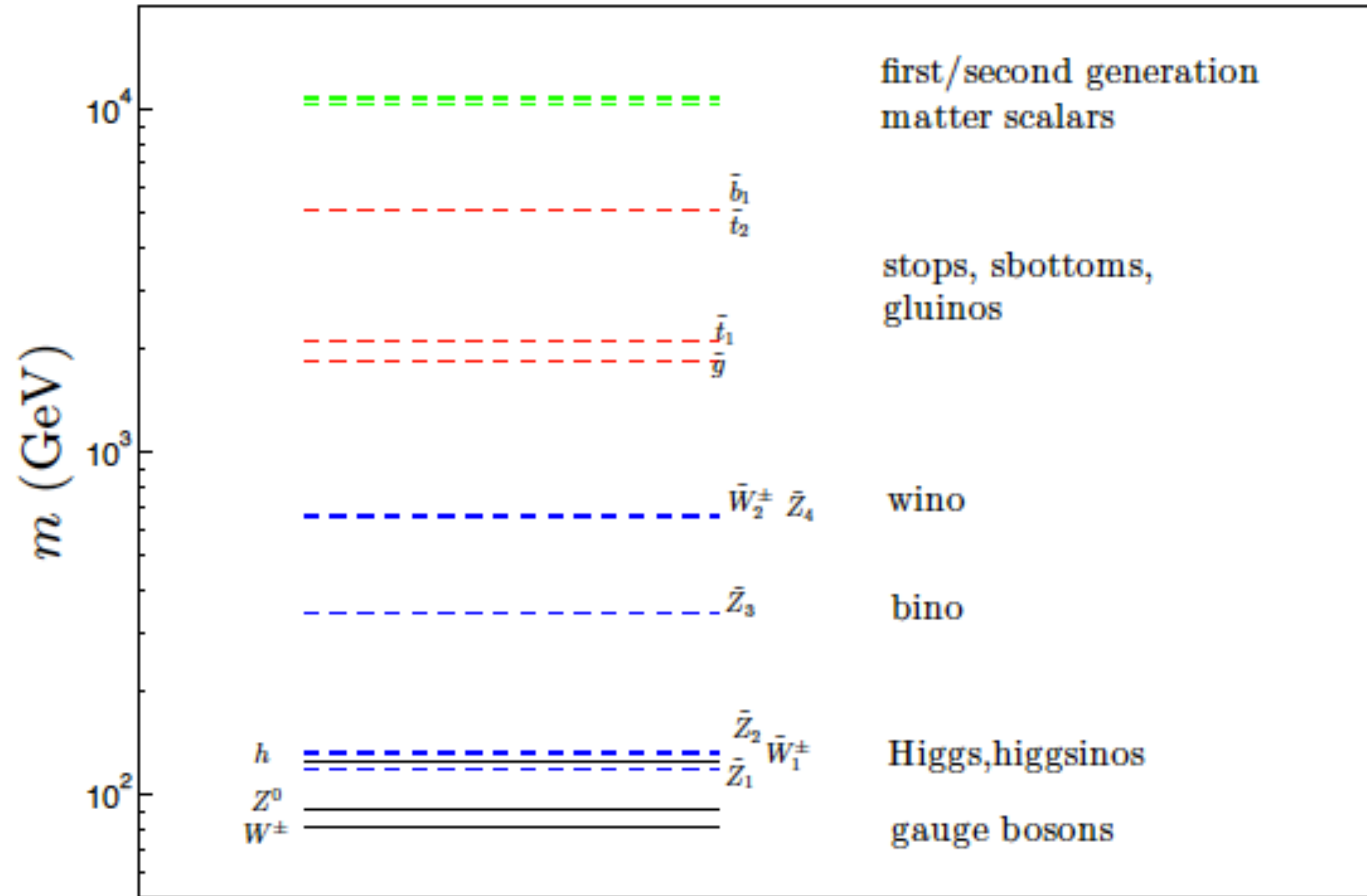


computed relative probabilities for
various natural and unnatural SUSY models

model	$\bar{m}(1,2)$	$\bar{m}(3)$	gauginos	higgsinos	m_h	P_μ
SM	-	-	-	-	-	$7 \cdot 10^{-27}$
CMSSM ($\Delta_{EW} = 2641$)	~ 1	~ 1	~ 1	~ 1	0.1 – 0.13	$5 \cdot 10^{-3}$
PeV SUSY	$\sim 10^3$	$\sim 10^3$	~ 1	$1 - 10^3$	0.125 – 0.155	$5 \cdot 10^{-6}$
Split SUSY	$\sim 10^6$	$\sim 10^6$	~ 1	~ 1	0.13 – 0.155	$7 \cdot 10^{-12}$
HS-SUSY	$\gtrsim 10^2$	$\gtrsim 10^2$	$\gtrsim 10^2$	$\gtrsim 10^2$	0.125 – 0.16	$6 \cdot 10^{-4}$
Spread (\bar{h} LSP)	10^5	10^5	10^2	~ 1	0.125 – 0.15	$9 \cdot 10^{-10}$
Spread (\bar{w} LSP)	10^3	10^3	~ 1	$\sim 10^2$	0.125 – 0.14	$5 \cdot 10^{-6}$
Mini-Split (\bar{h} LSP)	$\sim 10^4$	$\sim 10^4$	$\sim 10^2$	~ 1	0.125 – 0.14	$8 \cdot 10^{-8}$
Mini-Split (\bar{w} LSP)	$\sim 10^2$	$\sim 10^2$	~ 1	$\sim 10^2$	0.11 – 0.13	$4 \cdot 10^{-4}$
SUN-SUSY	$\sim 10^2$	$\sim 10^2$	~ 1	$\sim 10^2$	0.125	$4 \cdot 10^{-4}$
G ₂ MSSM	30 – 100	30 – 100	~ 1	~ 1	0.11 – 0.13	$2 \cdot 10^{-3}$
RNS/landscape	5 – 40	0.5 – 3	~ 1	0.1 – 0.35	0.123 – 0.126	1.4

Table 1: A survey of some unnatural and natural SUSY models along with general expectations for sparticle and Higgs mass spectra in TeV units. We also show relative probability measure P_μ for the model to emerge from the landscape. For RNS, we take $\mu_{\text{min}} = 10$ GeV.

Typical spectrum for low Δ_{EW} models



There is a Little Hierarchy, but it is **no problem**

$$\mu \ll m_{3/2}$$

higgsinos likely the lightest superparticles!

Conclusions

- 3rd string revolution: the landscape => CC problem, then SUSY?
- rather general arguments for power-law draw to large soft terms for SUSY in landscape
- derived value of weak scale within ABDS window
- predict $m(h) \sim 125$ GeV with sparticles beyond present LHC reach
- landscape favors EW natural SUSY over finetuned
- stringy naturalness: $m_{\text{gluino}} \sim 3$ TeV more natural than $m_{\text{gl}} \sim 300$ GeV!
- LHC signatures in talk #2

A final takeaway

- ATLAS/CMS have done a good job ruling out much of the 20th century SUSY parameter space
- This has influenced the European strategy and US Snowmass/P5 reports with regard to construction of future facilities
- However, we see that ATLAS/CMS are only beginning to explore the 21st century SUSY parameter space as expected from the string landscape
- Probably need hadron collider with $\sqrt{s} \sim 30\text{--}50$ TeV and e^+e^- collider with $\sqrt{s} > 2m(\text{higgsino})$