

Seeking a coherent explanation of LHC excesses for compressed spectra



Taylor Murphy

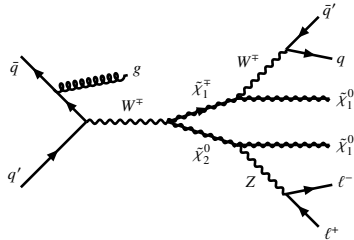
**Laboratoire de Physique
Théorique et Hautes Énergies**
Sorbonne Université & CNRS

11 June 2024

Based on arXiv:2404.12423
in collaboration with **D. Agin**
B. Fuks
and **M. D. Goodsell**

SMALL EXCESSES IN $2\ell + E_T^{\text{MISS}} \dots$

- ATLAS and CMS target compressed EWinos in soft leptons $(2/3\ell) + E_T^{\text{miss}}$ channels [1, 2, 3, 4, 5]



- Results interpreted for pure higgsino or wino LSP
- 2ℓ analyses show 1–2 σ excesses corresponding to $\Delta m(\tilde{\chi}_2^0, \tilde{\chi}_1^0) = 10\text{--}30 \text{ GeV}$
 - Overlapping excesses in the monojet channel [6, 7, 8]?

RECASTING THE ATLAS 2/3 ℓ ANALYSES



- ATLAS-SUSY-2018-16: 2 ℓ and 1 ℓ 1 T electroweakino channels are recast in HACKANALYSIS (HA) v2
- Multiple technical challenges:
 - RJR variables require RESTFRAMES [9], thus `root`
 - Results are sensitive to $m_{\ell\ell}$ distributions so that PYTHIA 8 cannot be used to handle the three-body electroweakino decays
 - Signal efficiencies are of $\mathcal{O}(10^{-5})$, requiring large samples
- ATLAS-SUSY-2019-09: off-shell WZ 3 ℓ selection is recast in HA
 - Lepton reconstruction efficiencies are scraped
 - Object-based E_T^{miss} significance computed in HA using momentum resolution uncertainties taken from ATLAS [10, 11]
- Validations use ATLAS cutflows and reproduced exclusion plots; agreement is excellent

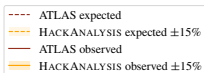
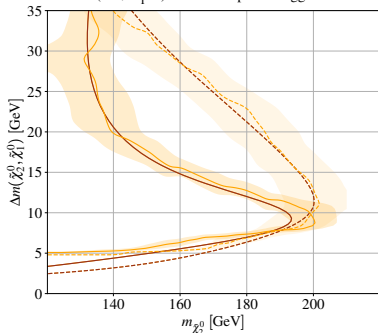
SIMULATION AND ANALYSIS TOOLCHAIN

- Relevant processes simulated in each model as follows:
 - Hard events with ≤ 2 hard jets generated using MG5_AMC, including full decay chains in matrix elements
 - Matching, showering, and hadronization performed with PYTHIA 8
 - Normalizations computed at best available accuracy; for MSSM we use RESUMMINO [12] at NLO + NLL
- Samples range in size from 3.2 M to 20 M events per parameter point, with $\mathcal{O}(10)$ – $\mathcal{O}(100)$ points considered for each model
- ATLAS and CMS monojet searches implemented in MADANALYSIS 5 and available on the Public Analysis Database (PAD), but ported to HA v2 to streamline the workflow
- Efficiencies computed by HA; statistical analysis performed by SPEY [13]

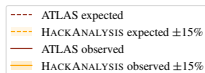
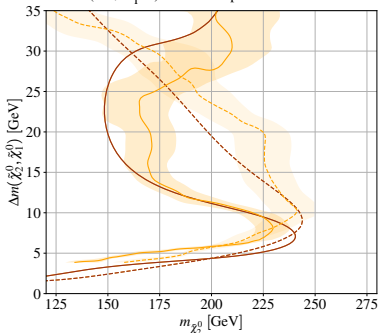
ATLAS-SUSY-2018-16 VALIDATION



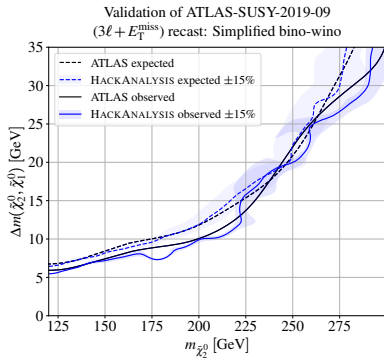
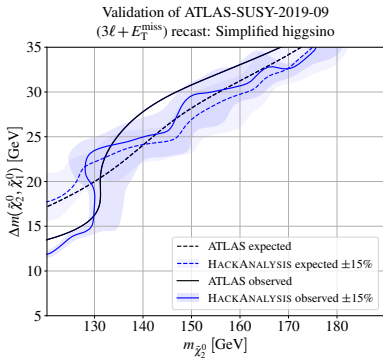
Validation of ATLAS-SUSY-2018-16
($2\ell + E_T^{\text{miss}}$) recast: Simplified higgsino



Validation of ATLAS-SUSY-2018-16
($2\ell + E_T^{\text{miss}}$) recast: Simplified bino-wino



ATLAS-SUSY-2019-09 VALIDATION



WHAT'S THE PLAN?

- With validated ATLAS $2/3\ell$ recasts (and monojets, with caveats), we can try to find suitable models
 - Why not also try to produce a dark matter candidate?
 - Other minimal(ish) SUSY scenarios seem like good candidates
 - What about non-SUSY models with different topologies? $m_{\ell\ell}$ distributions should differ
- Four candidate models: scan parameter spaces to find suitable spectra, compute LHC cross sections and (where applicable) Ωh^2 , apply recasts
- For each model and each analysis compute expected/observed limits at 95% CL
- Identify a best-fit point ($\hat{\mu} \approx 1$) and a significant point (lowest p -value); if we get lucky, we can find points with good fit and significance for all analyses

CANDIDATE 1: MSSM WITH DECOUPLED HIGGSINOS

- Already looked at simplified (pure) higgsinos; ATLAS interpreted soft leptons for simplified bino-wino LSP
 - Bino-winos not symmetric like higgsinos; instead $m_{\tilde{\chi}_1^0} < m_{\tilde{\chi}_1^\pm} = m_{\tilde{\chi}_2^0}$
 - Mixing and decays fixed, curiously with 2–10% higgsino admixture; physical masses changed by hand
- (Mostly) bino LSP is a viable DM candidate, with coannihilations important in the parameter space relevant to soft-lepton analyses
- Straightforward to perform a “realistic” scan over (M_1, M_2)
- Higgsinos decoupled with $\mu = 2$ TeV; first- and second-generation sfermions taken to $\mathcal{O}(10)$ TeV for simplicity

CANDIDATE 2: NMSSM WITH SINGLINO LSP



- MSSM + singlet superfield \rightarrow five neutralinos:

$$\mathcal{W} \supset \mu \mathcal{H}_u \mathcal{H}_d \rightarrow \mathcal{W}_{\text{NMSSM}} \supset \lambda \mathcal{S} \mathcal{H}_u \mathcal{H}_d + \frac{1}{3} \kappa \mathcal{S}^3$$

- Light S a bit lighter than light higgsinos (with small mixing) solves higgsino underabundance problem
- Singlino-higgsino compression and mixing tuned with λ, κ
- Light $\tilde{\chi}_3^0$ offers additional signal processes
- More parameters impose greater computing needs for scanning—in the end we want

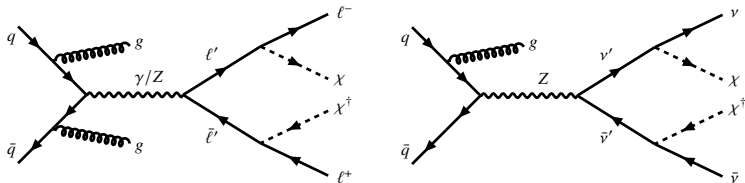
$$\lambda \langle S \rangle \in [100, 250] \text{ GeV}, \quad \kappa \in [0.001, 0.02], \quad \lambda \approx 2\kappa + \epsilon$$

CANDIDATE 3: VLL DOUBLET

- SM + scalar DM χ + weak doublet $\Psi^\top = (\nu', \ell')$ of vector-like leptons:

$$\mathcal{L} \supset \lambda_I \chi (\bar{\nu}' \nu_{LI} + \bar{\ell}' \ell_{LI}) + \text{H.c.}$$

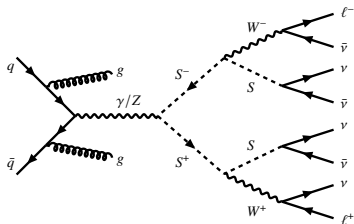
- VLL pair production produces 2ℓ /monojet signals



- VLL and DM masses independently tuned, look for $\Delta m \sim 0\text{--}30$ GeV
- χ annihilates to ν, ℓ and coannihilates with ν', ℓ'

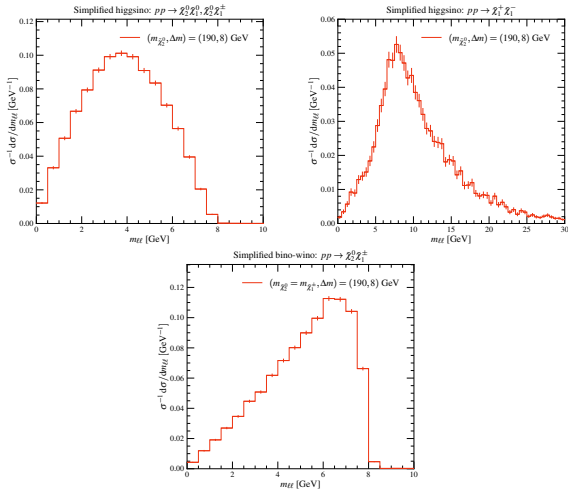
CANDIDATE 4: TYPE-II SEESAW

- SM + $Y = 1$ weak triplet $\Delta \rightarrow h, S, S^\pm, S^{\pm\pm}$
- Tree-level generation of Majorana neutrino masses
- If $m_S < m_{S^\pm} < m_{S^{\pm\pm}}$, $v_\Delta \ll 1$ GeV, and h is SM like, $S \rightarrow \nu\bar{\nu}$
- Charged scalar decays via off-shell W bosons produce 2ℓ /monojet signals

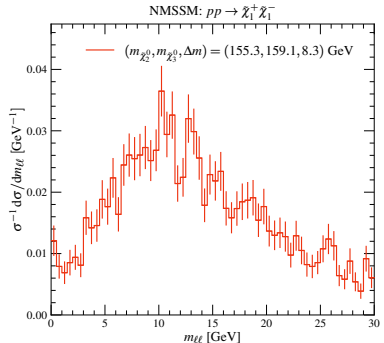
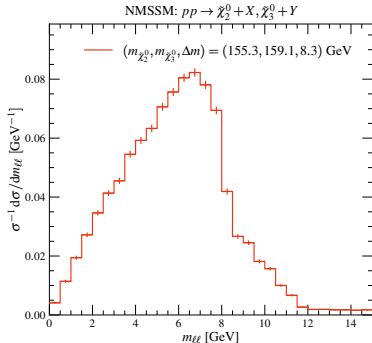


- Triplet scalar splitting controlled by $\mathcal{L} \supset -\lambda_4 \Phi^\dagger \Delta \Delta^\dagger \Phi$
- Scalars known to be weakly constrained for small $\lambda_4 < 0$

$m_{\ell\ell}$ DISTRIBUTIONS: SIMPLIFIED MSSM

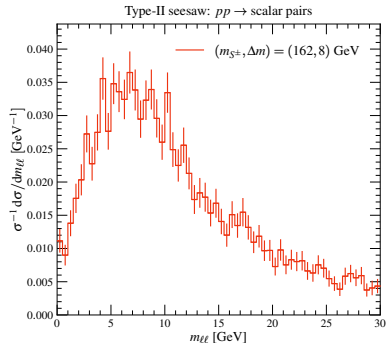
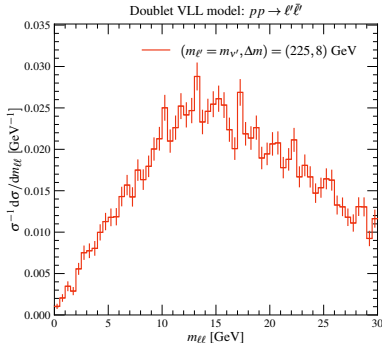


$m_{\ell\ell}$ DISTRIBUTIONS: NMSSM

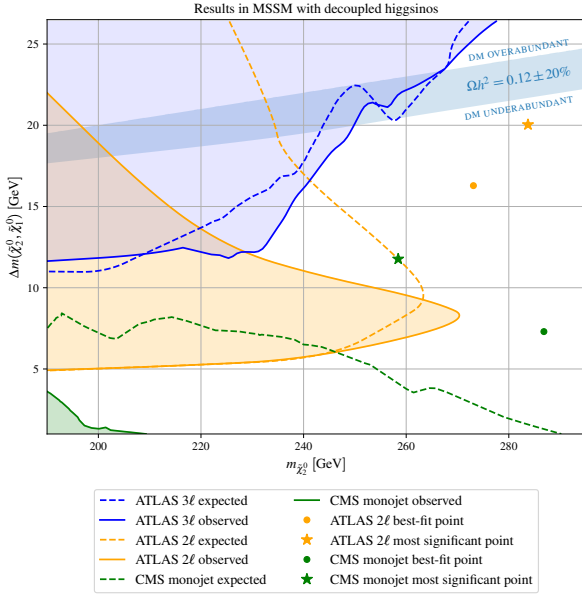


- $X \in \{\tilde{\chi}_1^\pm, \tilde{\chi}_3^0\}, Y \in \{\tilde{\chi}_1^\pm, \tilde{\chi}_2^0\}$
- $\tilde{\chi}^+ \tilde{\chi}^-$ similar to higgsinos but $\tilde{\chi}^0$ processes have higgsino and bino-wino features

$m_{\ell\ell}$ DISTRIBUTIONS: NON-SUSY



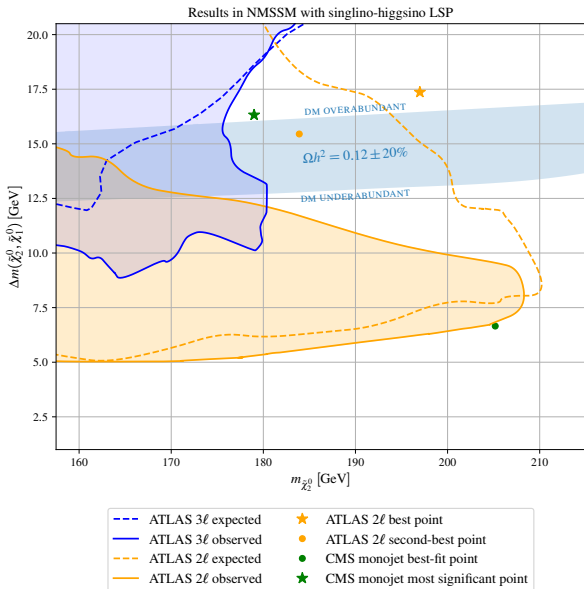
- Much harder leptons from direct $\ell' \rightarrow \chi\ell$ decays
- Type-II seesaw somewhat similar to $\tilde{\chi}^+\tilde{\chi}^-$ since $\ell^+\ell^-$ come from off-shell W^\pm



MSSM SIGNIFICANT POINTS

Point ($m_{\tilde{\chi}_2^0}, \Delta m$) [GeV]	● ATLAS 2ℓ best fit (273, 16.3)	★ ATLAS 2ℓ most significant (284, 20.0)	● CMS monojet best fit (287, 7.30)	★ CMS monojet most significant (258, 11.8)
($p, \hat{\mu}$) [ATLAS 2ℓ]	(0.047, 1.12)	(0.041, 1.26)	(> 0.5, < 0.1)	(0.290, 0.30)
($p, \hat{\mu}$) [ATLAS 3ℓ]	(> 0.5, < 0.1)	(0.426, < 0.1)	(> 0.5, < 0.1)	(> 0.5, < 0.1)
($p, \hat{\mu}$) [CMS monojet]	(0.098, 1.58)	(0.065, 2.33)	(0.049, 1.15)	(0.044, 1.40)
($p, \hat{\mu}$) [ATLAS monojet]	(0.277, 1.21)	(0.163, 2.44)	(0.127, 1.53)	(0.277, 0.879)
M_1 [GeV]	248.0	254.6	269.9	238.2
M_2 [GeV]	241.7	251.3	254.0	228.6
m_h [GeV]	127.0	126.5	126.8	126.8
$m_{\tilde{\chi}_1^0}$ [GeV]	256.8	263.7	279.5	246.7
$m_{\tilde{\chi}_2^0}$ [GeV]	273.1	283.7	286.8	258.4
$m_{\tilde{\chi}_1^\pm}$ [GeV]	273.3	283.9	287.0	258.6
($N_{11}, N_{12}, N_{13}, N_{14}$)	(0.9995, -0.0211, 0.0232, -0.0038)	(0.9996, -0.0175, 0.0231, -0.0039)	(0.9984, -0.0501, 0.02443, -0.0043)	(0.9993, -0.0284, 0.0235, -0.0038)
($N_{21}, N_{22}, N_{23}, N_{24}$)	(0.0220, 0.9990, -0.0392, 0.0066)	(0.0184, 0.9990, -0.0394, 0.0068)	(0.0511, 0.9979, -0.0386, 0.0068)	(0.0293, 0.9988, -0.0390, 0.0063)

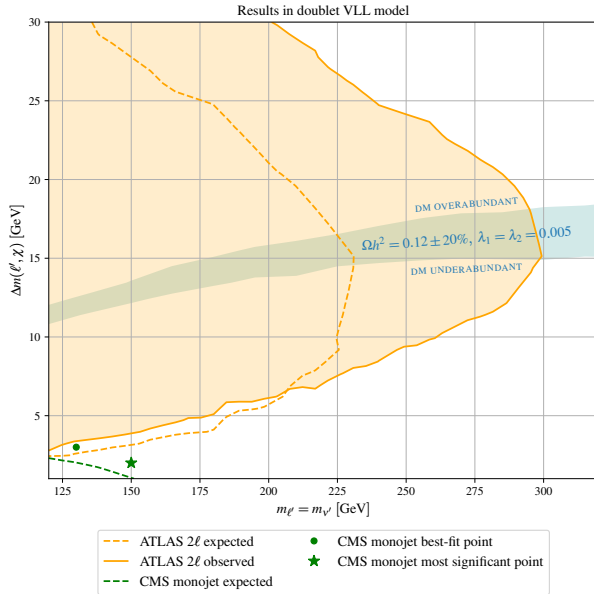
- 2ℓ best fit ● shows modest overlap with CMS monojet
- Most significant ★ in vicinity of correct Ωh^2
- Monojet best fit ● with $\Delta m < 10$ GeV invisible to 2ℓ analysis



NMSSM SIGNIFICANT POINTS



Point ($m_{\tilde{\chi}_2^0}, \Delta m$) [GeV]	★ ATLAS 2ℓ best (197, 17.4)	● ATLAS 2ℓ second best (184, 15.5)	● CMS monojet best fit (205, 6.66)	★ CMS monojet most significant (179, 16.3)
($p, \hat{\mu}$) [ATLAS 2ℓ]	(0.041, 0.97)	(0.044, 0.80)	(0.435, < 0.1)	(0.071, 0.64)
($p, \hat{\mu}$) [ATLAS 3ℓ]	(0.446, 0.12)	(> 0.5, 0.12)	(> 0.5, 0.71)	(> 0.5, 0.13)
($p, \hat{\mu}$) [CMS monojet]	(0.132, 3.00)	(0.129, 2.65)	(0.712, 1.91)	(0.051, 3.79)
($p, \hat{\mu}$) [ATLAS monojet]	(0.277, 2.44)	(0.277, 2.02)	(0.127, 2.96)	(0.277, 2.08)
μ_{eff} [GeV]	189.3	177.0	199.1	172.6
κ	0.0157	0.0108	0.0025	0.0146
λ	0.0330	0.0226	0.0050	0.0309
$\tan \beta$	19.71	25.94	10.70	12.82
M_t^2 [GeV ²]	8.06×10^7	7.20×10^7	3.42×10^7	9.12×10^7
A_t [GeV]	2.61×10^3	-1.28×10^3	2.07×10^3	-2.64×10^3
A_λ [GeV]	-34.60	-92.77	189.4	192.0
A_κ [GeV]	-43.01	-8.771	-161.3	-55.91
m_h [GeV]	124.0	123.4	123.0	122.6
$m_{\tilde{\chi}_1^0}$ [GeV]	179.6	168.5	198.5	162.7
$m_{\tilde{\chi}_2^0}$ [GeV]	197.0	183.9	205.2	179.0
$m_{\tilde{\chi}_1^\pm}$ [GeV]	198.1	185.5	207.1	180.3
$m_{\tilde{\chi}_3^0}$ [GeV]	199.9	187.3	209.0	182.1
($N_{11}, N_{12}, N_{13}, N_{14}, N_{15}$)	(0.0042, -0.0070, 0.1547, -0.1683, 0.9735)	(0.0032, -0.0053, 0.1201, -0.1299, 0.9841)	(0.0016, -0.0026, 0.0580, -0.0597, 0.9965)	(0.0041, -0.0069, 0.1479, -0.1622, 0.9756)
($N_{21}, N_{22}, N_{23}, N_{24}, N_{25}$)	(-0.0173, 0.0289, -0.6932, 0.6827, 0.2284)	(-0.0172, 0.0287, -0.7004, 0.6907, 0.1767)	(-0.0184, 0.0312, -0.7077, 0.7006, 0.0832)	(-0.0176, 0.0294, -0.6951, 0.6838, 0.0219)
Im ($N_{31}, N_{32}, N_{33}, N_{34}, N_{35}$)	(-0.0134, 0.0226, 0.7039, 0.7097, 0.0110)	(-0.0137, 0.0231, 0.7036, 0.7101, 0.0080)	(-0.0126, 0.0216, 0.7041, 0.7097, 0.0016)	(-0.0131, 0.0220, 0.7035, 0.7100, 0.0116)



NON-SUSY SIGNIFICANT POINTS

Point ($m_{\ell'} = m_{\nu'}$, $\Delta m(\ell', \chi)$) [GeV]	● CMS monojet best fit (273, 16.3)	★ CMS monojet most significant (284, 20.0)
Significance (p -value)	0.047	0.041
Signal strength $\hat{\mu}$	1.12	1.26

- VLL model has no points with $2/3\ell$ discovery potential: 2ℓ observed limits stronger than expected
- Monojet points correspond to underabundant DM
- Type-II seesaw unconstrained by all searches:

	CL_s^{exp}	CL_s^{obs}	p -value	$\hat{\mu}$
ATLAS 2ℓ	0.387	0.620	0.050	2.06
ATLAS 3ℓ	0.611	0.615	≥ 0.5	0.62
CMS monojet	0.270	0.395	0.081	1.17
ATLAS monojet	0.553	0.399	0.277	0.72

- ($m_{S^\pm}, \Delta m(S^\pm, S)$) = (95, 5) GeV
- Excesses still visible ($CL_s^{\text{obs}} > CL_s^{\text{exp}}$)

OUTLOOK

- We explored excesses in the soft lepton + E_T^{miss} channel for four well motivated models
- Excesses might have some overlap with monojets for pure higgsinos [8]; we now find modest compatibility in bino-wino MSSM and singlino-higgsino NMSSM, but poor fits for two non-SUSY models
- With this puzzle unresolved and a workflow worked out, we have choices for future work:
 - Detailed scans of other popular models (go ahead, call them out)
 - Cook up model(s) specifically for these excesses
 - Improved statistical treatments (combined likelihood, allowed signal strengths)

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Thank you for your attention

I am happy to answer questions if we have time

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Bonus material

ATLAS-SUSY-2018-16 YIELDS



		SR $m_{\ell\ell}$ bin [GeV]							
		[1,2]	[2,3]	[3,2,5]	[5,10]	[10,20]	[20,30]	[30,40]	[40,60]
SR-E high ee	Observed			1	16	13	8	8	18
	Fitted SM events			0.7 ± 0.4	10.3 ± 2.5	12.1 ± 2.2	10.1 ± 1.7	10.4 ± 1.7	19.3 ± 2.5
SR-E high $\mu\mu$	Observed	5	5	0	9	23	3	5	20
	Fitted SM events	3.4 ± 1.2	3.5 ± 1.3	3.9 ± 1.3	11.0 ± 2.0	17.8 ± 2.7	8.3 ± 1.4	10.1 ± 1.5	19.6 ± 2.3
SR-E med ee	Observed			0	4	11	4		
	Fitted SM events			0.11 ± 0.08	5.1 ± 1.6	7.3 ± 1.9	2.2 ± 0.9		
SR-E med $\mu\mu$	Observed	16	8	6	41	59	21		
	Fitted SM events	14.6 ± 2.9	6.9 ± 2.1	6.2 ± 1.9	34 ± 4	52 ± 6	18.5 ± 3.2		
SR-E low ee	Observed			7	11	16	16	10	9
	Fitted SM events			5.3 ± 1.5	8.6 ± 1.8	16.7 ± 2.5	15.5 ± 2.6	12.9 ± 2.1	18.8 ± 2.2
SR-E low $\mu\mu$	Observed	9	7	7	12	17	18	16	44
	Fitted SM events	15.4 ± 2.4	8.0 ± 1.7	6.5 ± 1.6	11.3 ± 1.9	15.6 ± 2.3	16.7 ± 2.3	15.3 ± 2.0	35.9 ± 3.3

NEUTRALINO DECAYS TO LEPTONS

- Full differential decay rate, $\tilde{\chi}_2^0 \rightarrow Z^* + \tilde{\chi}_1^0$:

$$\frac{d\Gamma}{dm_{\ell\ell}} = C m_{\ell\ell} \frac{\{m_{\ell\ell}^4 - m_{\ell\ell}^2[(\Delta m)^2 + M^2] + (M\Delta m)^2\}^{1/2}}{(m_{\ell\ell}^2 - M_Z^2)^2} \times \{-2m_{\ell\ell}^4 + m_{\ell\ell}^2[2M^2 - (\Delta m)^2] + (M\Delta m)^2\}$$

$M, \Delta m = m_{\tilde{\chi}_2^0} \pm m_{\tilde{\chi}_1^0}$: $2M^2 - (\Delta m)^2$ sensitive to signed $\tilde{\chi}^0$ masses

- “Flat” decay rate used by PYTHIA 8:

$$\frac{d\Gamma^{\text{flat}}}{dm_{\ell\ell}} = C^{\text{flat}} m_{\ell\ell} \{m_{\ell\ell}^4 - m_{\ell\ell}^2[(\Delta m)^2 + M^2] + (M\Delta m)^2\}^{1/2}$$

FLAT VS. ACCURATE $m_{\ell\ell}$

