

# Phenomenology of an unusual Composite Higgs Model

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Based on G. Cacciapaglia, T. Flacke, M. Kunkel and WP, JHEP **02** (2022), 208 (arXiv:2112.00019)  
G. Cacciapaglia, T. Flacke, M. Kunkel, WP and L. Schwarze, JHEP **12** (2022), 087 (arXiv:2210.01826)

10 June 2024

# Generic Composite Higgs set-up

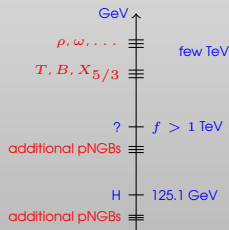
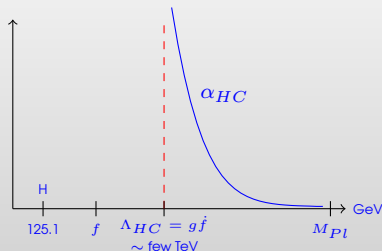
Possible solution to hierarchy problem

- ▶ Generate a scale  $\Lambda_{HC} \ll M_{Pl}$  through a new confining gauge group
- ▶ Interpret Higgs as a pseudo-Nambu-Goldstone boson (pNGB) of a spontaneously broken global symmetry of the new strong sector

(Georgi, Kaplan, PLB **136** (1984), 136)

'Price' to pay

- ▶ additional resonances at the scale  $\Lambda_{HC}$  (spin-1 resonances, vector-like fermions, scalars)
- ▶ additional light pNGBs/ extended scalar sector
- ▶ deviations of the Higgs couplings from their SM values of  $O(v/f)$



# Towards underlying models

A wish list to construct and classify candidate models:

Gerghetta et al (2015), Ferretti et al. PLB (2014), PRD 94 (2016), JHEP 1701.094

Underlying models of a composite Higgs should

- ▶ contain no elementary scalars (otherwise there would be again a hierarchy problem)
- ▶ have a simple hyper-color group
- ▶ have a Higgs candidate amongst the pNGBs of the bound states
- ▶ have a top-partner amongst its bound states (for top mass via partial compositeness)
- ▶ satisfy further 'standard' consistency conditions (asymptotic freedom, no gauge anomalies)

The resulting models have several common features:

- ▶ All models predict pNGBs beyond the Higgs multiplet
- ▶ All models contain several top partner multiplets

can be extended to include neutrino masses and dark matter, e.g. G. Cacciapaglia, M. Rosenlyst, JHEP **09** (2021), 167

## List of "minimal" CHM UV embeddings

$G_{\text{HC}}$	$\psi$	$\chi$	Restrictions	$-q_x/q_\psi$	$Y_x$	Non Conformal	Model Name
	Real	Real	$SU(5)/SO(5) \times SU(6)/SO(6)$				
$SO(N_{\text{HC}})$	$5 \times \mathbf{S}_2$	$6 \times \mathbf{F}$	$N_{\text{HC}} \geq 55$	$\frac{5(N_{\text{HC}}+2)}{6}$	1/3	/	
$SO(N_{\text{HC}})$	$5 \times \mathbf{Ad}$	$6 \times \mathbf{F}$	$N_{\text{HC}} \geq 15$	$\frac{5(N_{\text{HC}}-2)}{6}$	1/3	/	
$SO(N_{\text{HC}})$	$5 \times \mathbf{F}$	$6 \times \mathbf{Spin}$	$N_{\text{HC}} = 7, 9$	$\frac{5}{6}, \frac{5}{12}$	1/3	$N_{\text{HC}} = 7, 9$	M1, M2
$SO(N_{\text{HC}})$	$5 \times \mathbf{Spin}$	$6 \times \mathbf{F}$	$N_{\text{HC}} = 7, 9$	$\frac{5}{6}, \frac{5}{3}$	2/3	$N_{\text{HC}} = 7, 9$	M3, M4
	Real	Pseudo-Real	$SU(5)/SO(5) \times SU(6)/Sp(6)$				
$Sp(2N_{\text{HC}})$	$5 \times \mathbf{Ad}$	$6 \times \mathbf{F}$	$2N_{\text{HC}} \geq 12$	$\frac{5(N_{\text{HC}}+1)}{3}$	1/3	/	
$Sp(2N_{\text{HC}})$	$5 \times \mathbf{A}_2$	$6 \times \mathbf{F}$	$2N_{\text{HC}} \geq 4$	$\frac{5(N_{\text{HC}}-1)}{3}$	1/3	$2N_{\text{HC}} = 4$	M5
$SO(N_{\text{HC}})$	$5 \times \mathbf{F}$	$6 \times \mathbf{Spin}$	$N_{\text{HC}} = 11, 13$	$\frac{5}{24}, \frac{5}{48}$	1/3	/	
	Real	Complex	$SU(5)/SO(5) \times SU(3)^2/SU(3)$				
$SU(N_{\text{HC}})$	$5 \times \mathbf{A}_2$	$3 \times (\mathbf{F}, \bar{\mathbf{F}})$	$N_{\text{HC}} = 4$	$\frac{5}{3}$	1/3	$N_{\text{HC}} = 4$	M6
$SO(N_{\text{HC}})$	$5 \times \mathbf{F}$	$3 \times (\mathbf{Spin}, \bar{\mathbf{Spin}})$	$N_{\text{HC}} = 10, 14$	$\frac{5}{12}, \frac{5}{48}$	1/3	$N_{\text{HC}} = 10$	M7
	Pseudo-Real	Real	$SU(4)/Sp(4) \times SU(6)/SO(6)$				
$Sp(2N_{\text{HC}})$	$4 \times \mathbf{F}$	$6 \times \mathbf{A}_2$	$2N_{\text{HC}} \leq 36$	$\frac{1}{3(N_{\text{HC}}-1)}$	2/3	$2N_{\text{HC}} = 4$	M8
$SO(N_{\text{HC}})$	$4 \times \mathbf{Spin}$	$6 \times \mathbf{F}$	$N_{\text{HC}} = 11, 13$	$\frac{8}{3}, \frac{16}{3}$	2/3	$N_{\text{HC}} = 11$	M9
	Complex	Real	$SU(4)^2/SU(4) \times SU(6)/SO(6)$				
$SO(N_{\text{HC}})$	$4 \times (\mathbf{Spin}, \bar{\mathbf{Spin}})$	$6 \times \mathbf{F}$	$N_{\text{HC}} = 10$	$\frac{8}{3}$	2/3	$N_{\text{HC}} = 10$	M10
$SU(N_{\text{HC}})$	$4 \times (\mathbf{F}, \bar{\mathbf{F}})$	$6 \times \mathbf{A}_2$	$N_{\text{HC}} = 4$	$\frac{2}{3}$	2/3	$N_{\text{HC}} = 4$	M11
	Complex	Complex	$SU(4)^2/SU(4) \times SU(3)^2/SU(3)$				
$SU(N_{\text{HC}})$	$4 \times (\mathbf{F}, \bar{\mathbf{F}})$	$3 \times (\mathbf{A}_2, \bar{\mathbf{A}}_2)$	$N_{\text{HC}} \geq 5$	$\frac{4}{3(N_{\text{HC}}-2)}$	2/3	$N_{\text{HC}} = 5$	M12
$SU(N_{\text{HC}})$	$4 \times (\mathbf{F}, \bar{\mathbf{F}})$	$3 \times (\mathbf{S}_2, \bar{\mathbf{S}}_2)$	$N_{\text{HC}} \geq 5$	$\frac{4}{3(N_{\text{HC}}+2)}$	2/3	/	

G. Ferretti, JHEP **06** (2016), 107; A. Belyaev et al. JHEP **01** (2017), 094

# M5: $HC = Sp(4), SU(5) \times SU(6)/SO(5) \times Sp(6)$

## pNGBs:

electroweak:	$SO(5)$ 14	$SU(2)_L \times SU(2)_R$ (1,1) + (2,2) + (3,3)	states $\eta, H, \eta_1^0, \eta_3^{+,0,-}, \eta_5^{+,+,0,-,--}$ $(S_i^0 = \eta, \eta_{1,3,5}^0, S_i^+ = \eta_{3,5}^+, S^{++} = \eta_5^{++})$
strong:	$Sp(6)$ 14	$SU(3)_C \times U(1)_{em}$ $3_{2/3} + \bar{3}_{-2/3} + 8_0$	states $\pi_3, \pi_3^*, \pi_8$

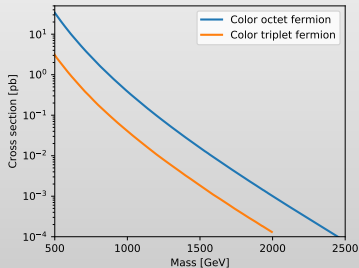
## fermionic bound states:

$SO(5) \times Sp(6)$	$SU(3)_L \times SU(2)_L \times U(1)_Y$ / names				
$(\mathbf{5}, \mathbf{14})$	$(3, 2)_{7/6}$	$(3, 2)_{1/6}$	$(8, 2)_{1/2}$	$(3, 1)_{2/3}$	$(8, 1)_0$
	$(X_{5/3}, X_{3,2})$	$(T_L, B_L)$	$(\tilde{G}^+, \tilde{G}^0)$	$T_R$	$\tilde{g}$
$(\mathbf{5}, \mathbf{1})$	$(1, 2)_{1/2}$	$(1, 1)_0$			
	$(\tilde{H}^+, \tilde{H}^0)$	$\tilde{B}$			

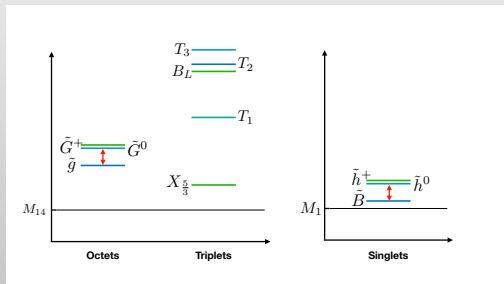
$\tilde{g}$  and  $\tilde{B}$  are Majorana fermions, all other are Dirac fermions

accidental global symmetry: 'baryon' number

# Hyper-baryons (top-partners)



**3 @ NLO, 8 @ NNLO<sub>approx</sub> + NNLL**  
 G. Cacciapaglia *et al.*,  
 arXiv:2112.00019



**Assumption:** 1) fermions within an  $SO(5) \times Sp(6)$  multiplet have about the same mass  
mass splitting due to SM gauge interactions

2)  $\tilde{B}$  is stable

$\Rightarrow$  **LHC:** 1) fermionic color octets have largest cross section  
2) events with large missing  $p_T$

Possible decays:

$$\begin{array}{l} \tilde{g} \rightarrow t \pi_3^*, \bar{t} \pi_3 \\ \rightarrow \tilde{B} \pi_8 \end{array} \quad \left| \quad \begin{array}{l} \tilde{G}^0 \rightarrow \bar{t} \pi_3 \\ \rightarrow \tilde{H}^0 \pi_8 \end{array} \quad \left| \quad \begin{array}{l} \tilde{G}^+ \rightarrow \bar{b} \pi_3 \\ \rightarrow \tilde{H}^+ \pi_8 \end{array} \right.$$

$\tilde{H}^+ \rightarrow \pi^+ \tilde{B}, \tilde{H}^0 \rightarrow \pi^0 \tilde{B}$  with very soft pions

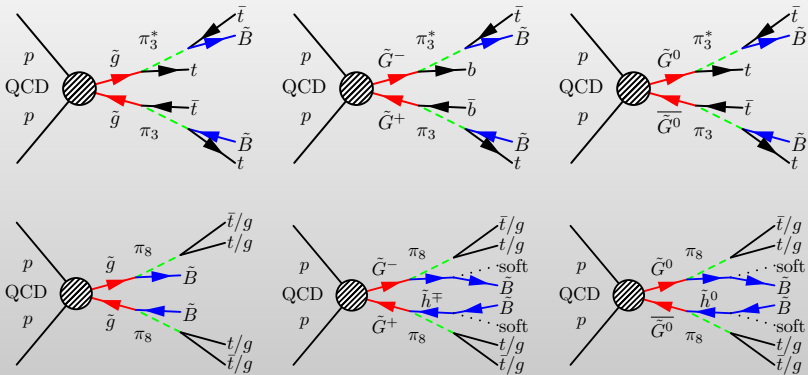
$$\begin{array}{l} \pi_3 \rightarrow t \tilde{B} \\ (\rightarrow t \nu) \\ (\rightarrow \bar{s} \bar{d}) \end{array} \quad \left| \quad \begin{array}{l} \pi_8 \rightarrow g g \\ \rightarrow t \bar{t} \\ (\rightarrow q \bar{q}, q = u, d, s, c, b) \end{array}$$

Bounds on  $\pi_3$ :  $\tilde{t}_R$  searches,  $\simeq 1.3 \text{ TeV}^\dagger$

$\pi_8$ :  $\simeq 1.1 \text{ TeV}^*$

$^\dagger$  (ATLAS, arXiv:2102.01444 (hep-ex); CMS, arXiv:2107.10892 (hep-ex))

\* G. Cacciapaglia et al., arXiv:2002.01474 (hep-ph)



for later use:

$$Q_8 = \{\tilde{g}, \tilde{G}^0, \tilde{G}^\pm\}$$



# Recast of existing LHC analyses

LHC signatures:

- ▶  $4 t + \text{missing } p_T$
- ▶  $3 t + j + \text{missing } p_T$
- ▶  $2 t + 2 j + \text{missing } p_T$
- ▶  $t + 3 j + \text{missing } p_T$
- ▶  $4 j + \text{missing } p_T$

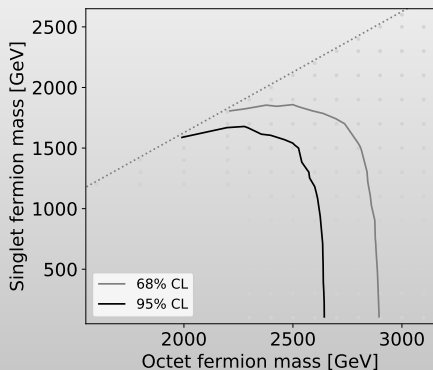
In all cases: additional soft pions possible.

We used here and in the following

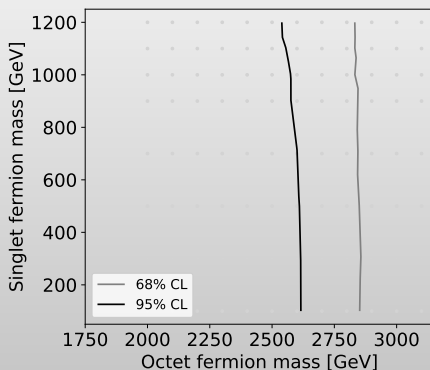
- ▶ generated  $10^5$  events per data point using `MadGraph5_aMC@NLO`, hadronized with `Pythia8`
- ▶ recast tools
  - ▶ `MadAnalysis5`, mainly SUSY searches, E. Conte et al., arXiv:1206.1599, arXiv:1808.00480
  - ▶ `CheckMate`, SUSY searches, M. Drees et al., arXiv:1312.2591; D. Dercks et al., arXiv:1611.09856
  - ▶ `Contur`, based on SM measurements implemented in `Rivet`, J. Butterworth et al., arXiv:1606.05296, arXiv:1902.03067
- ▶ check for each data point which tool gives the best constraint

Cross sections: NNLOapprox + NNLL, from <https://twiki.cern.ch/twiki/bin/view/LHCPhysics/SUSYCrossSections13TeVgluglu>

[//twiki.cern.ch/twiki/bin/view/LHCPhysics/SUSYCrossSections13TeVgluglu](https://twiki.cern.ch/twiki/bin/view/LHCPhysics/SUSYCrossSections13TeVgluglu)

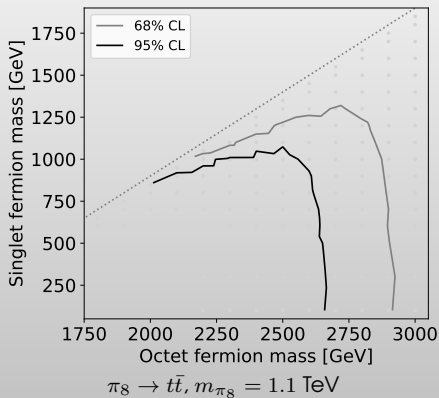
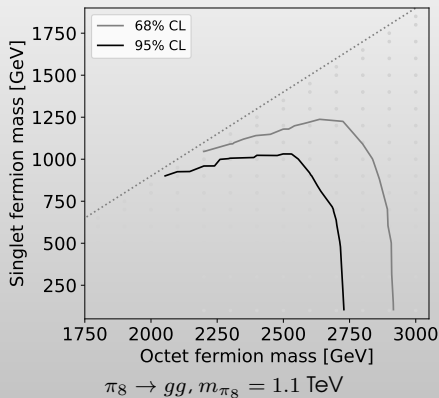
Octet decays with 100% decays into  $\pi_3$ 

$$m_{Q_8} - m_{\pi_3} = 200 \text{ GeV}$$



$$m_{\pi_3} = 1.4 \text{ TeV}$$

G. Cacciapaglia *et al.*, arXiv:2112.00019

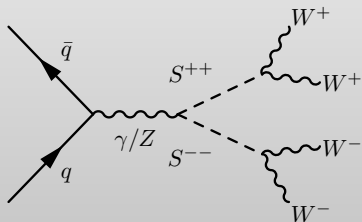
Octet decays with 100% decays into  $\pi_8$ G. Cacciapaglia *et al.*, arXiv:2112.00019

# Electroweak pNGBs

$$pp \rightarrow S_i^{\pm\pm} S_j^{\mp}, S_i^{\pm} S_j^0, S_i^{++} S_j^{--}, S_i^+ S_j^-, S_i^0 S_j^0$$

two limiting scenarios

fermiophobic

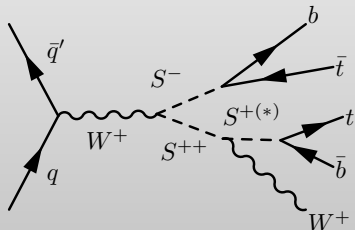


$$S_i^{++} \rightarrow W^+ W^+$$

$$S_i^+ \rightarrow W^+ \gamma, W^+ Z$$

$$S_i^0 \rightarrow W^+ W^-, \gamma \gamma, \gamma Z, Z Z.$$

fermiophilic



$$S^{++} \rightarrow W^+ t \bar{b},$$

$$S^+ \rightarrow t \bar{b},$$

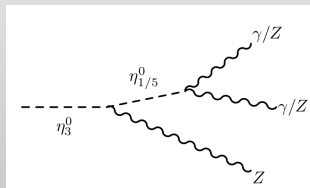
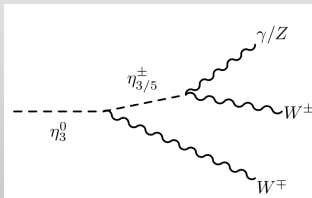
$$S^0 \rightarrow t \bar{t}, b \bar{b}.$$

# Electroweak pNGBs

fermiophilic scenario: only weak bounds in a very small region of parameter space

⇒ focus on fermiophobic scenario

- ▶ assume custodial multiplets are mass-degenerate
- ▶ lights multiplet decays only via anomaly terms, except  $\eta_3^0$  which does not couple to the anomaly, but

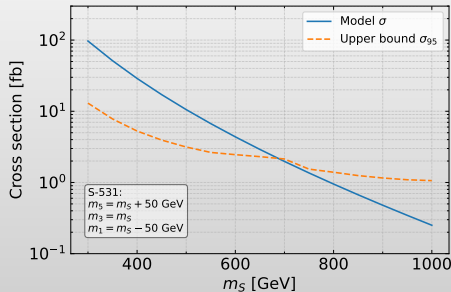
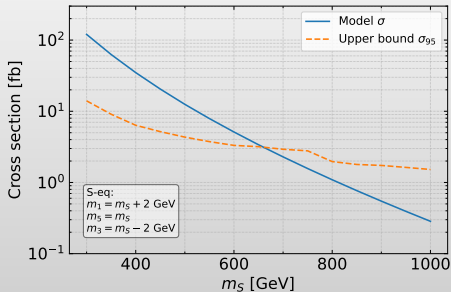


- ▶ for the heavier custodial multiplets: decays into (off-shell) vector bosons + lighter multiplet, e.g.

$$\eta_3^+ \rightarrow \eta_5^{++} W^{-(*)}, \eta_5^+ Z^{(*)}, \eta_5^0 W^{+(*)}, \eta_1^0 W^{+(*)};$$

$$\eta_3^0 \rightarrow \eta_5^{\pm} W^{\mp(*)}, \eta_5^0 Z^{(*)}, \eta_1^0 Z^{(*)}.$$

# Bounds on $\eta_1, \eta_3, \eta_5$

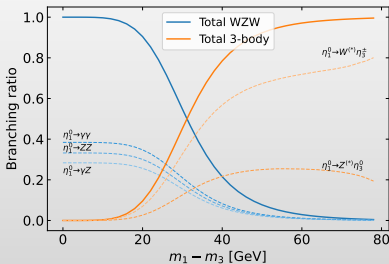


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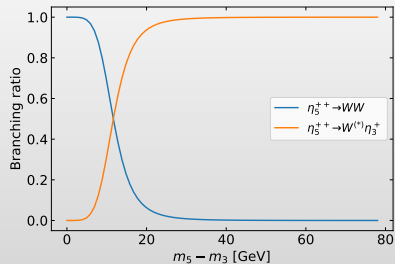
## Conclusions:

- ▶ Composite Higgs models provide a viable solution to the hierarchy problem but they still provide many challenges and room for exploration in theory and model-building.
- ▶ In general:
  - ▶ several pNGBs, also in the strongly interacting sector
  - ▶ fermionic bound states: not only color triplets, but also for example octets and singlets
  - ▶ **bounds depend strongly on possible decay modes**
- ▶ example, M5-model:
  - ▶ mass bounds on electroweak pNGBs:  
fermiophilic scenario: only very weak bounds in small part of parameter space  
fermiophobic scenario:  $\sim 400\text{-}650$  GeV depending on mass splittings
  - ▶ color octets among the top-partners: bounds of up to 2.8 TeV on their masses

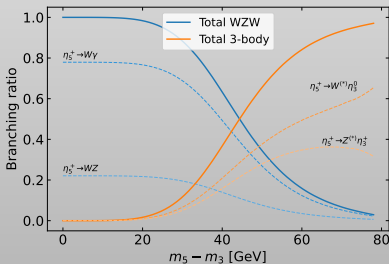
# Branching ratios 1, M5-model



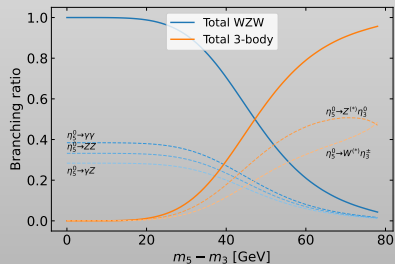
$$\eta_1^0, m_1 = 600 \text{ GeV} > m_3$$



$$\eta_5^{++}, m_5 = 600 \text{ GeV} > m_3$$



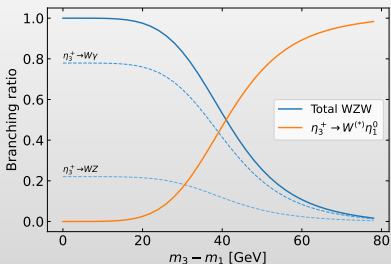
$$\eta_5^+, m_5 = 600 \text{ GeV} > m_3$$



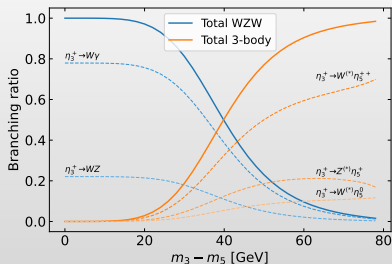
$$\eta_5^0, m_5 = 600 \text{ GeV} > m_3$$



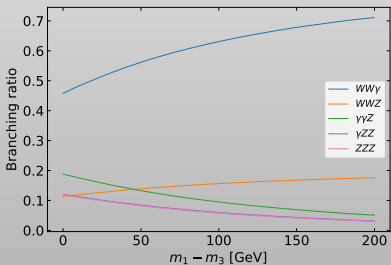
# Branching ratios 2, M5-model



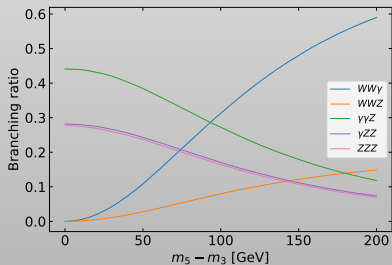
$$\eta_3^+, m_5 \gg m_3 = 600 \text{ GeV} > m_1$$



$$\eta_3^+, m_1 \gg m_3 = 600 \text{ GeV} > m_5$$

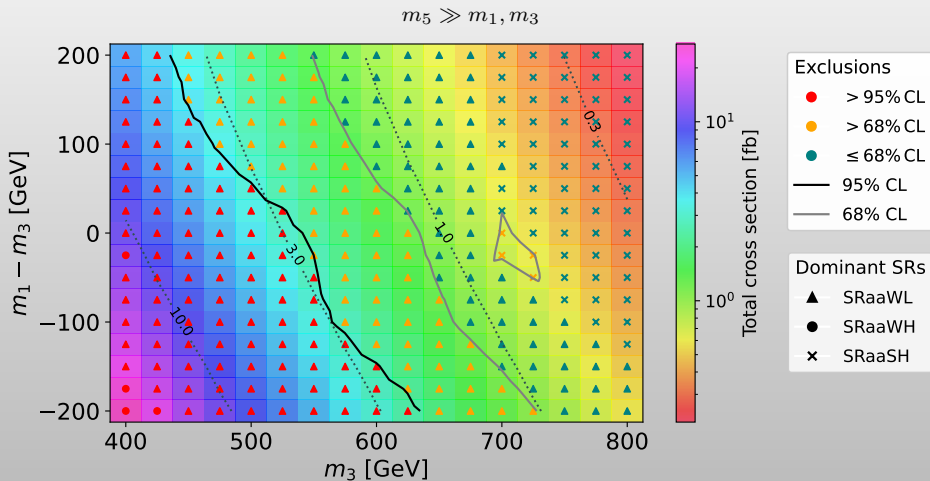


$$\eta_3^0, m_5 \gg m_1 > m_3 = 600 \text{ GeV}$$

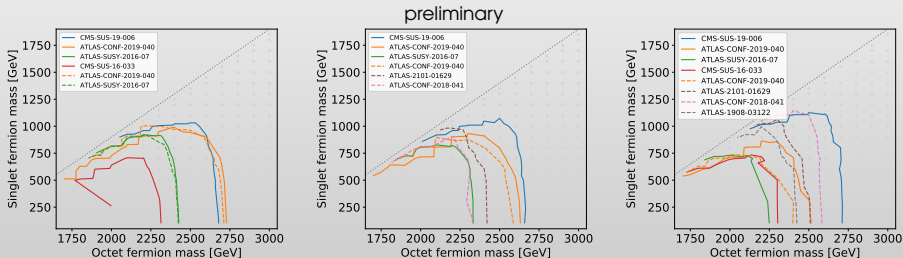


$$\eta_3^0, m_1 \gg m_5 > m_3 = 600 \text{ GeV}$$

# Bounds on $\eta_1, \eta_3, \eta_5$



## Contribution of different searches for color octets



Comparison of the bounds at 95% CL obtained from different searches implemented in MADANALYSIS 5 (solid lines) and CHECKMATE 2 (dashed lines).