

# Status of sub-GeV dark matter

Tomás Gonzalo

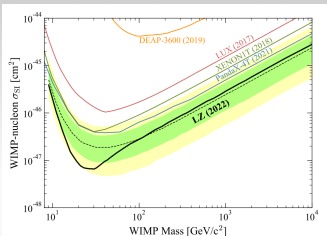
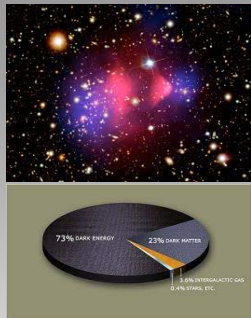
Karlsruhe Institute for Technology

SUSY 2024,  
10 June 2024

[S. Balan et al, [arXiv:2405.17548](https://arxiv.org/abs/2405.17548)]

# Dark Matter

- Plenty of evidence for DM from astrophysics (e.g bullet cluster) and cosmology (e.g CMB)
- If DM is a particle and if interacts then we should be able to detect it
- Most popular DM models are WIMPs
  - EW-scale mass, accesible at colliders
  - Just right RD through freeze-out
  - Form part of complete models (e.g. MSSM)



- No evidence of WIMPs
  - Very strong constraints from experimental searches (e.g LZ)
  - Many WIMP models in trouble, only survive in fine-tuned scenarios
- What if DM was not a WIMP?

# Sub-GeV DM

- Most DD experiments threshold 1 GeV  $\rightarrow$  sub-GeV DM avoids DD
- Sub-GeV DM (scalar or fermion) with dark photon mediator

$$\mathcal{L}_\Phi = |\partial_\mu \Phi|^2 - m_{\text{DM}}^2 |\Phi|^2 + ig_{\text{DM}} A'^\mu [\Phi^* (\partial_\mu \Phi) - (\partial_\mu \Phi^*) \Phi] - g_{\text{DM}}^2 A'_\mu A'^\mu |\Phi|^2,$$

$$\mathcal{L}_\psi = \bar{\psi} (i \not{\partial} - m_{\text{DM}} \psi + g_{\text{DM}} A'^\mu \bar{\psi} \gamma_\mu \psi).$$

- Dark photon mixes with SM photon

$$\mathcal{L}_{A'} = -\frac{1}{2} m_{A'}^2 A'^\mu A'_\mu - \frac{1}{4} A'^{\mu\nu} A'_{\mu\nu} - \kappa e A'^\mu \sum_f q_f \bar{f} \gamma_\mu f$$

- We consider only  $m_{A'} \geq 2m_{\text{DM}}$  so that  $\text{BR}(A' \rightarrow \chi\bar{\chi}) \sim 1$
- Strongly constrained annihilation cross section (CMB & X-rays)

$\rightarrow$  Resonant enhancement

$$\epsilon_R = (m_{A'}^2 - 4m_{\text{DM}}^2) / (4m_{\text{DM}}^2)$$

$\rightarrow$  Particle-antiparticle asymmetry

$$\eta_{\text{DM}} = (n_\xi - n_{\bar{\chi}}) / s$$

$\rightarrow$  Underabundant DM

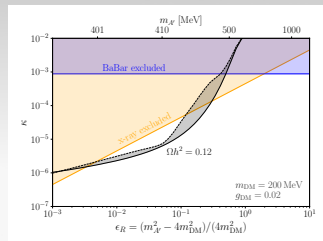
$$f_{\text{DM}} = \Omega_{\text{DM}} / \Omega_{\text{DM,obs}} < 1$$

# Resonant enhancement

- Resonant enhancement of ann at freezeout and suppression of ID
- Resonant parameter  $\epsilon_R = \frac{m_{A'}^2 - 4m_{DM}^2}{4m_{DM}^2}$
- The kinetic energy available in an ann process  $\epsilon = \frac{s - 4m_{DM}^2}{4m_{DM}^2}$  which is around  $\epsilon \sim 0.1$  at freezeout and  $\epsilon \sim 10^{-6}$  in the MW
- In the non-relativistic limit,  $\epsilon = v_{DM}^2$ , so the propagator of  $A'$  is

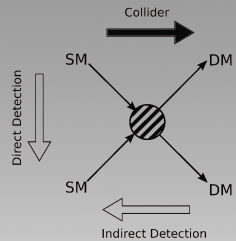
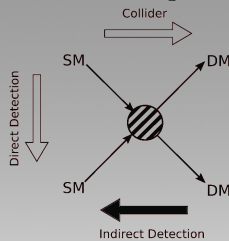
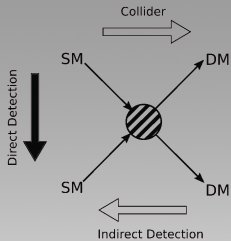
$$\frac{1}{(s - m_{A'}^2)^2 + m_{A'}^2 \Gamma_{A'}^2} = \frac{1}{16m_{DM}^4 (\epsilon - \epsilon_R)^2 + m_{A'}^2 \Gamma_{A'}^2}$$

- So a value of  $\epsilon_R \sim 0.1$  enhances ann at freeze-out but not today
- Optimal range  $\epsilon_R \in [10^{-3}, 0.3]$



# Constraints on sub-GeV DM

- Constraints change a lot with respect to GeV-scale WIMPs



→ Nuclear (CRESSTIII)

→ Migdal (DarkSide-50, XENON1T, PandaX4T)

→ Electron  
(XENON1T, SENSEI, DarkSide-50, PandaX4T, DAMIC, SuperCDMS)

→ X-rays (INTEGRAL)

→ Bullet Cluster  
 $\sigma_0/m_{\text{DM}} < 1.4 \text{ cm}^2 \text{g}^{-1}$

→ CMB  $E$  injection

→  $N_{\text{eff}}$  at BBN

→ RD of asym DM  
 $\Omega_{\text{DM}} h^2 \leq 0.120 \pm 0.001$

→ Beam dumps:

LSND, MiniBooNE

$\pi^0, \eta \rightarrow \gamma A'$

→ Fixed target: NA64

$e^- Z \rightarrow e^- Z A'$

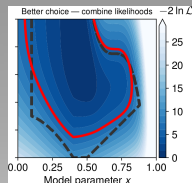
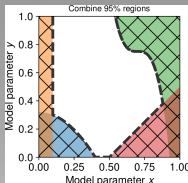
→ Single- $\gamma$  search:

BaBar  $e^+ e^- \rightarrow \gamma A'$

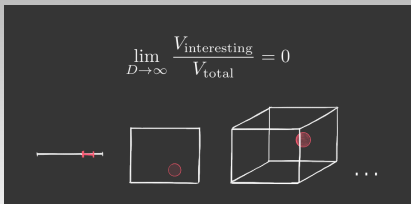
# Global fits of DM models

- Multitude of constraints
- Exclusion regions do not properly represent the model predictions
- Composite likelihood

$$\mathcal{L} = \mathcal{L}_{Direct} \mathcal{L}_{Indirect} \mathcal{L}_{Collider} \mathcal{L}_{Astro} \dots$$



[arXiv:2012.09874 [hep-ph]]



- Multitude of parameters
- Hard to find interesting regions
- Random methods are inefficient
- Need smart sampling strategies (differential, nested, genetic, ...)

- Rigorous statistical interpretations (frequentist / Bayesian)
- Parameter estimation, goodness-of-fit, model comparison, ...

## GAMBIT: The Global And Modular BSM Inference Tool

[gambit.hepforge.org](http://gambit.hepforge.org)

[github.com/GambitBSM](https://github.com/GambitBSM)

EPJC 77 (2017) 784

arXiv:1705.07908

- Extensive model database, beyond SUSY
- Fast definition of new datasets, theories
- Extensive observable/data libraries
- Plug&play scanning/physics/likelihood pack
- Various statistical options (frequentist /Bayesian)
- Fast LHC likelihood calculator
- Massively parallel
- Fully open-source



**Members of:** ATLAS, Belle-II, CLIC, CMS, CTA, Fermi-LAT, DARWIN, IceCube, LHCb, SHiP, XENON

**Authors of:** BubbleProfiler, Capt'n General, Contur, DarkAges, DarkSUSY, DDCalc, DirectDM, Diver, EasyScanHEP, ExoCLASS, FlexibleSUSY, gamLike, GM2Calc, HEPLike, IsaTools, MARTY, nuLike, PhaseTracer, PolyChord, Rivet, SOFTSUSY, SuperIso, SUSY-AI, xsec, Vevacious, WIMPSim

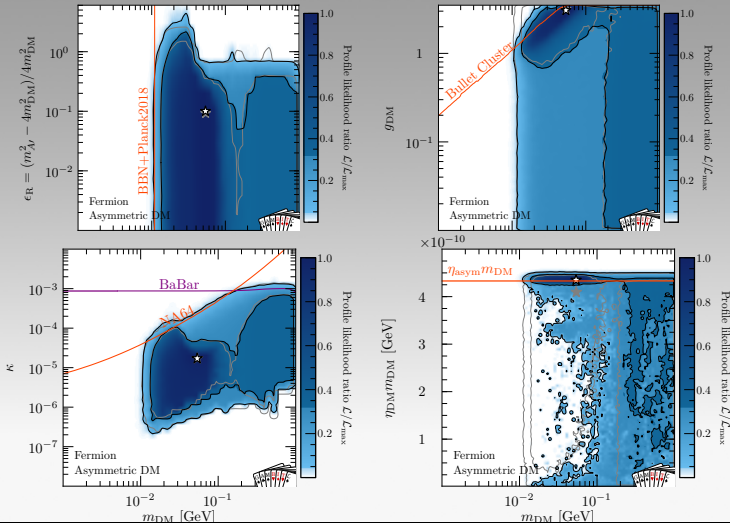
**Recent collaborators:** V Ananyev, P Athron, N Avis-Kozar, C Balázs, A Beniwal, LL Braseth, T Bringmann, A Buckley, J Butterworth, JE Camargo-Molina, C Chang, J Cornell, M Danninger, A Fowlie, T Gonzalo, W Handley, S Hoof, A Jueid, F Kahlhoefer, A Kvellestad, M Lacroq, C Lin, M Lucente, FN Mahmoudi, DJE Marsh, G Martinez, H Pacey, MT Prim, T Procter, F Rajec, A Raklev, R Ruiz, A Scaffidi, P Scott, W Shorrock, C Sierra, P Stöcker, W Su, J Van den Abeele, A Vincent, M White, A Woodcock, Y Zhang ++

70+ participants in many experiments and numerous major theory codes

- Global fits of BSM models: DM, ALPs, SUSY,  $\nu$ s, flavour, ...
- Other applications: nuclear physics, COVID spread models, ...

# Fermion asymmetric DM

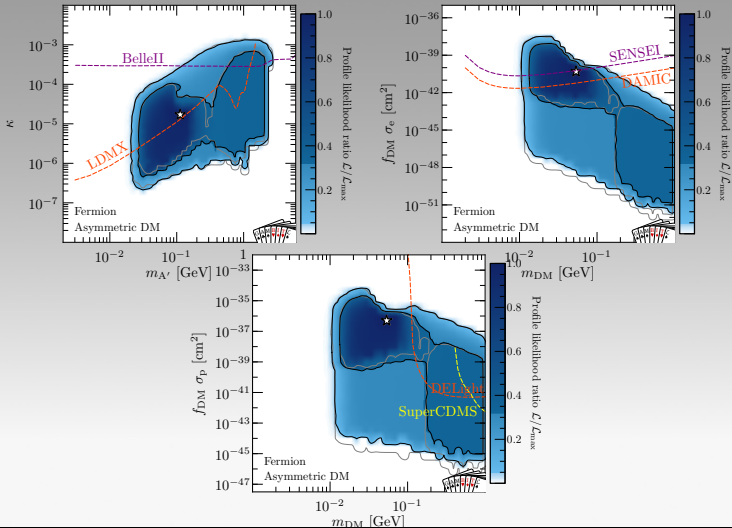
- Frequentist results





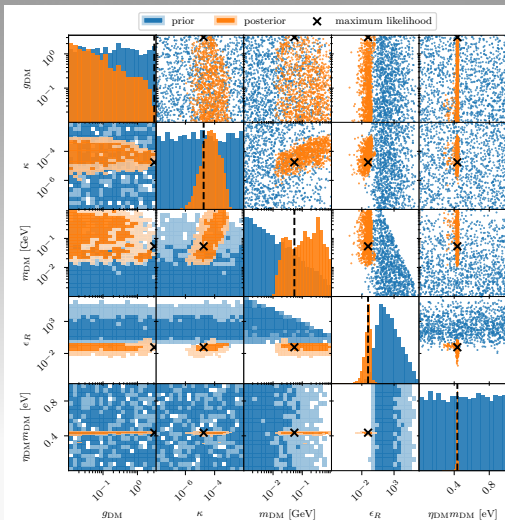
# Fermion asymmetric DM

- Future sensitivities



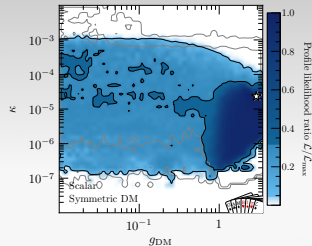
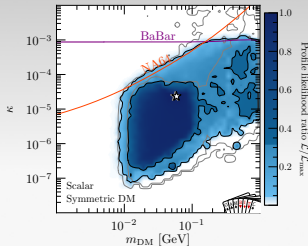
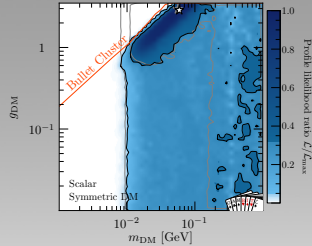
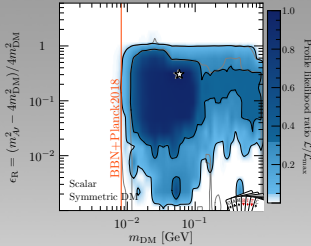
# Fermion asymmetric DM

- Bayesian results



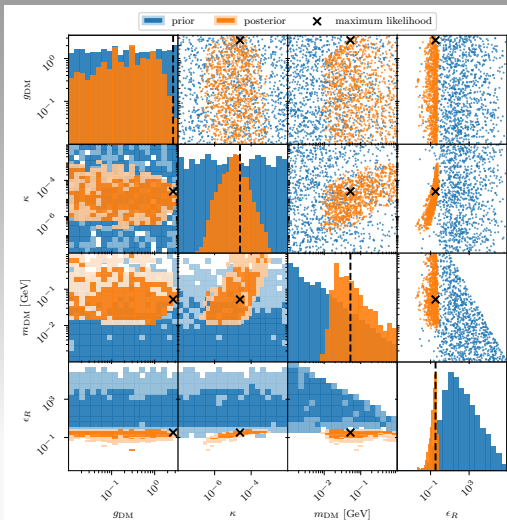
# Scalar symmetric DM

## • Frequentist results



# Scalar symmetric DM

- Bayesian results



# Bayesian model comparison

- Bayesian evidence

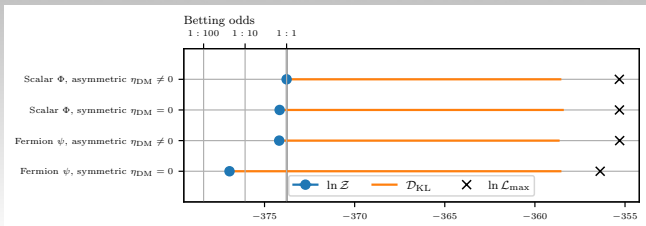
$$\mathcal{Z} = \int \mathcal{L}(\theta)\pi(\theta)d\theta \quad \rightarrow \quad \log \mathcal{Z} = -\langle \log \mathcal{L} \rangle_{\mathcal{P}} - \mathcal{D}_{\text{KL}}$$

- Posterior-weighted log-likelihood

$$\langle \log \mathcal{L} \rangle_{\mathcal{P}} = \int \mathcal{P}(\theta) \log \mathcal{L}(\theta) d\theta$$

- Kullback-Leibler divergence

$$\mathcal{D}_{\text{KL}} = \int \mathcal{P}(\theta) \log \frac{\mathcal{P}(\theta)}{\pi(\theta)} d\theta$$



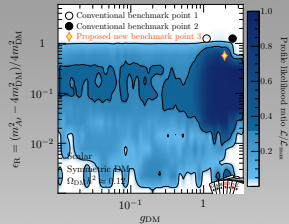
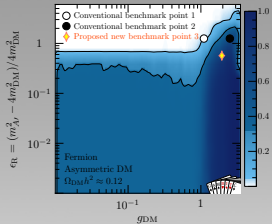
- $\mathcal{Z}_{\text{asym}}/\mathcal{Z}_{\text{sym}} = 15.6$

# Benchmark points

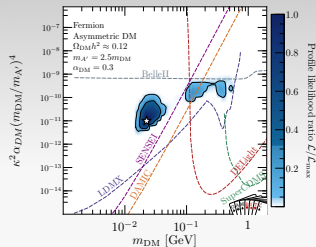
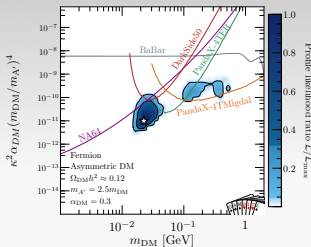
- Past benchmark points are excluded by current constraints

$$m_{A'} / m_{\text{DM}} = 3$$

$$g_{\text{DM}} = 1.1 \text{ or } 2.5$$



- We propose new BP:  $m_{A'} = 5/2 m_{\text{DM}}$ ,  $g_{\text{DM}} = 1.94$



# Conclusions and outlook

- GeV-scale WIMPs might not be the right answer → sub-GeV DM
- There are many models of DM constrained by multitude of constraints from different sources
  - Global studies the only way to give definitive status of models
- Fermionic DM survives either on the resonance  $m_{A'} \sim 2m_{\text{DM}}$ , or in the case of maximum asymmetry  $\eta_{\text{DM}}m_{\text{DM}} \sim 4 \times 10^{-10}$ 
  - Bayesian evidence prefers asymmetric case  $\mathcal{Z}_{\text{asym}}/\mathcal{Z}_{\text{sym}} = 15.6$
- Scalar DM does not need either extreme resonance or asymmetry
  - No significant Bayesian preference for either
- Old benchmarks are (mostly) excluded with recent data
  - New benchmarks can be discovered in the next gen of searches

$$m_{A'} = \frac{5}{2}m_{\text{DM}} \quad \text{or} \quad \epsilon_R = \frac{9}{16}, \quad \alpha_{\text{DM}} = 0.3 \quad \text{or} \quad g_{\text{DM}} = 1.94$$

Thanks!

# Backup



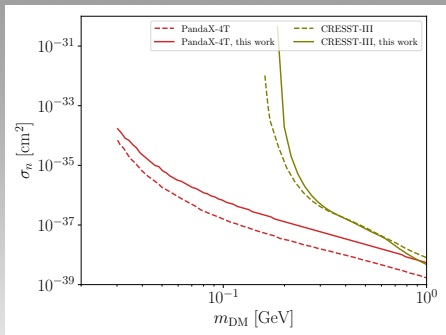
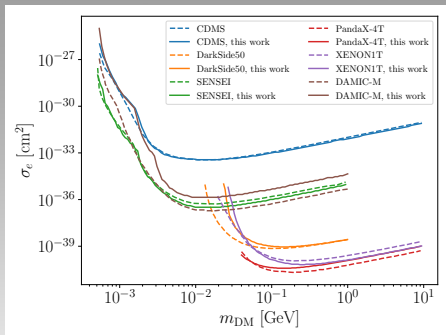
# Sub-GeV DM

- Parameter ranges and priors

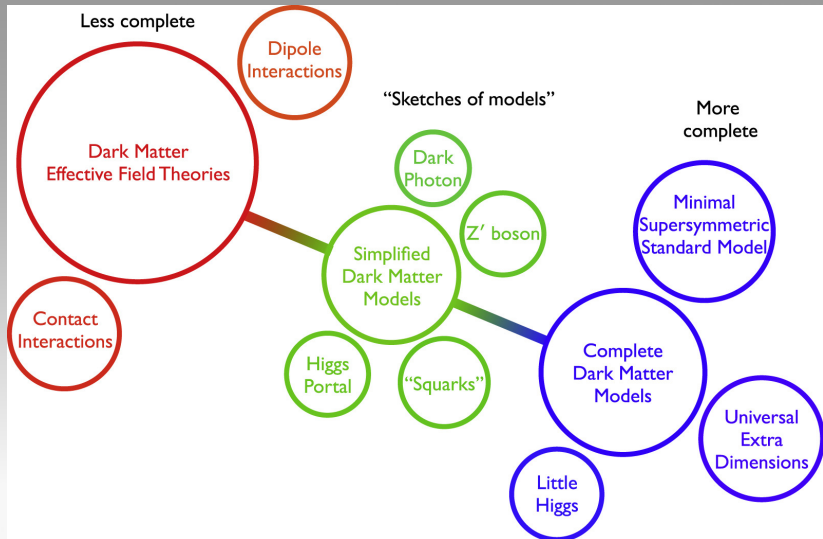
Parameter name	Symbol	Unit	Range	Prior
Kinetic mixing	$\kappa$	–	$[10^{-8}, 10^{-2}]$	logarithmic
Dark sector coupling	$g_{\text{DM}}$	–	$[10^{-2}, \sqrt{4\pi}]$	logarithmic
Asymmetry parameter	$\eta_{\text{DM}}$	–	$[0, 10^{-9} \text{ GeV}/m_{\text{DM}}]$	linear
Dark matter mass	$m_{\text{DM}}$	MeV	$[1, 1000]$	logarithmic
Dark photon mass <i>or</i>	$m_{A'}$	MeV	$[2, 6000]$ with $m_{A'} \geq 2m_{\text{DM}}$	logarithmic
Resonance parameter	$\epsilon_R$	–	$[10^{-3}, 8]$	logarithmic

# Sub-GeV DM

- Reproduction of the DD results (ER, NR and Migdal)

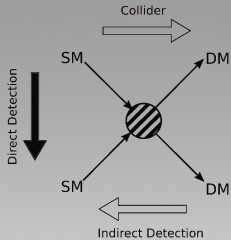


# Dark Matter

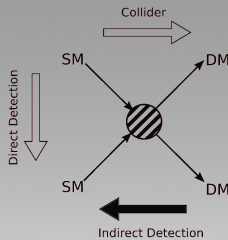


# Dark Matter

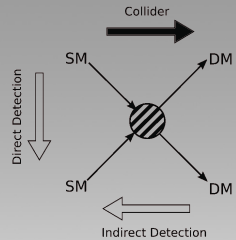
- Searches for DM in particle physics, astrophysics and cosmology



- DM interacting with nuclei
- LZ, XENON1T, PandaX, LUX, CDMSlite, CRESST, PICO-60, DarkSide-50



- DM annihilates into SM particles
- $\gamma$  rays,  $\nu$ s,  $\bar{p}$ , ...
- Fermi-LAT, IceCube, AMS02
- BBN and CMB
- $\Omega_{\text{DM}} h^2 \leq 0.120 \pm 0.001$



- LHC searches for large  $\cancel{E}_T$
- Mono-X (jet, ...)
- $pp \rightarrow \chi\chi j \rightarrow j + \cancel{E}_T$
- Mediator searches (e.g.  $\Gamma_{H \rightarrow \text{inv}}$ , dijets)

# Higgs portal DM

- Scalar DM ( $S$ )

[GAMBIT, Eur.Phys.J.C 77 (2017) 8, 568]

[S.Balan et al, arXiv:2303.07362 [hep-ph]]

$$\mathcal{L}_S = \frac{1}{2}\mu_S^2 S^2 + \frac{1}{2}\lambda_{hS} S^2 |H|^2 + \frac{1}{4}\lambda_S S^4 + \frac{1}{2}\partial_\mu S \partial^\mu S,$$

$$m_S^2 = \mu_S^2 + \frac{1}{2}\lambda_{hS} v^2$$

- Vector DM ( $V_\mu$ )

[GAMBIT, Eur.Phys.J.C 79 (2019) 1, 38]

$$\mathcal{L}_V = -\frac{1}{4}W_{\mu\nu}W^{\mu\nu} + \frac{1}{2}\mu_V^2 V_\mu V^\mu - \frac{1}{4!}\lambda_V (V_\mu V^\mu)^2 + \frac{1}{2}\lambda_{hV} V_\mu V^\mu H^\dagger H$$

$$m_V^2 = \mu_V^2 + \frac{1}{2}\lambda_{hV}^2 v^2$$

- Fermionic DM (Dirac,  $\psi$ )

[GAMBIT, Eur.Phys.J.C 79 (2019) 1, 38]

$$\mathcal{L}_\psi = \bar{\psi}(i\not{\partial} - m_\psi)\psi - \frac{\lambda_{h\psi}}{\Lambda_\psi} (\cos \xi \bar{\psi}\psi + \sin \xi \bar{\psi}i\gamma_5\psi)(vh + \frac{1}{2}h^2)$$

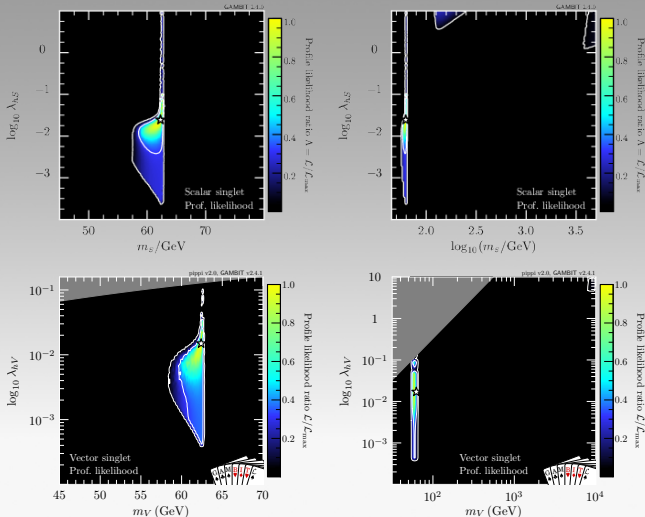
- Fermionic DM (Majorana,  $\chi$ )

[GAMBIT, Eur.Phys.J.C 79 (2019) 1, 38]

$$\mathcal{L}_\chi = \frac{1}{2}\bar{\chi}(i\not{\partial} - m_\chi)\chi - \frac{1}{2}\frac{\lambda_{h\chi}}{\Lambda_\chi} (\cos \xi \bar{\chi}\chi + \sin \xi \bar{\chi}i\gamma_5\chi)(vh + \frac{1}{2}h^2)$$

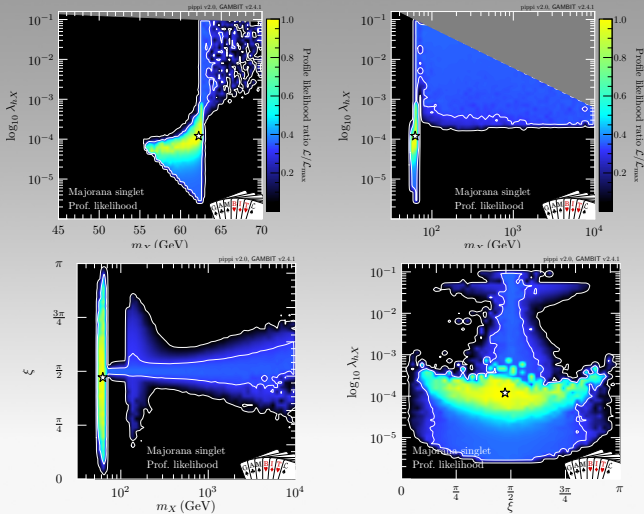
# Higgs portal DM

- Bosonic DM (scalar and vector)



# Higgs portal DM

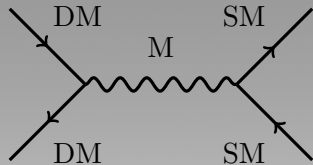
- Majorana fermion DM ( $\approx$  Dirac DM)



# Simplified DM models

- Singlet DM candidate plus vector mediator that couples to SM particles (quarks)

$$\mathcal{L}_V = -\frac{1}{4}F'_{\mu\nu}F'^{\mu\nu} - \frac{1}{2}m_M^2 V_\mu V^\mu + g_q V_\mu \bar{q}\gamma^\mu q$$



- DM can be a scalar ( $\phi$ ), a fermion ( $\psi$  or  $\chi$ ) or a vector ( $X_\mu$ )

[C.Chang et al, Eur.Phys.J.C 83 (2023) 3, 249]

$$\mathcal{L}_\phi = \partial_\mu \phi^\dagger \partial^\mu \phi - m_{\text{DM}}^2 \phi^\dagger \phi + i g_{\text{DM}}^V V_\mu \left( \phi^\dagger (\partial^\mu \phi) - (\partial^\mu \phi^\dagger) \phi \right),$$

$$\mathcal{L}_\chi = i \bar{\chi} \gamma^\mu \partial_\mu \chi - m_{\text{DM}} \bar{\chi} \chi + V_\mu \bar{\chi} (g_{\text{DM}}^V + g_{\text{DM}}^A \gamma^5) \gamma^\mu \chi,$$

$$\mathcal{L}_\psi = \frac{1}{2} i \bar{\psi} \gamma^\mu \partial_\mu \psi - \frac{1}{2} m_{\text{DM}} \bar{\psi} \psi + \frac{1}{2} g_{\text{DM}}^A V_\mu \bar{\psi} \gamma^5 \gamma^\mu \psi$$

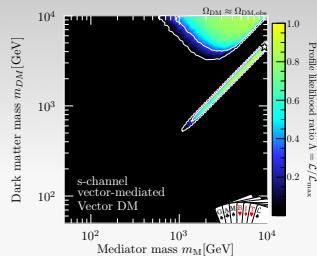
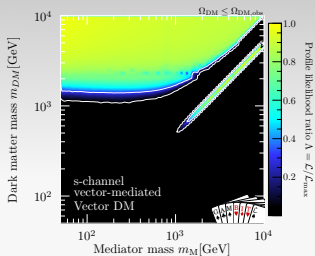
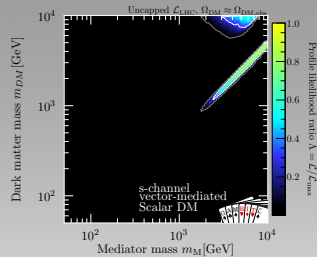
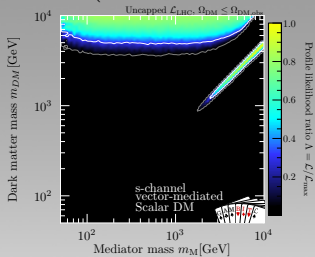
[C.Chang et al, arXiv:2303.08351 [hep-ph]]

$$\mathcal{L}_X = \frac{1}{2} X_{\mu\nu}^\dagger X^{\mu\nu} + m_{\text{DM}}^2 X_\mu^\dagger X^\mu - i g_{\text{DM}} \left( X_\nu^\dagger \partial_\mu X^\nu - (\partial_\mu X^{\dagger\nu}) X_\nu \right) V^\mu$$



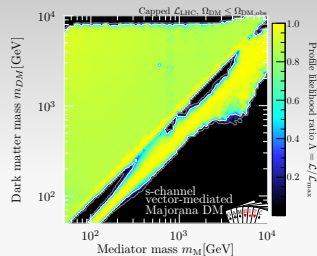
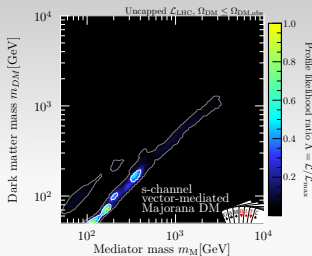
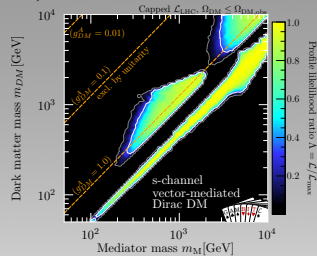
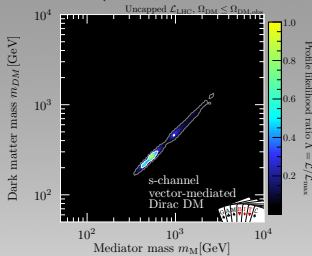
# Simplified DM models

- Bosonic DM (scalar and vector)



# Simplified DM models

- Fermion DM (Dirac and Majorana)



- Dirac fermionic DM  $\chi$ :  $\mathcal{L} = \mathcal{L}_{\text{SM}} + \mathcal{L}_{\text{int}} + \bar{\chi} (i\not{\partial} - m_\chi) \chi$
- Effective interactions (quarks/gluons):  $\mathcal{L}_{\text{int}} = \sum_{a,d} \frac{\mathcal{C}_a^{(d)}}{\Lambda^{d-4}} \mathcal{Q}_a^{(d)}$

$$\mathcal{Q}_1^{(5)} = \frac{e}{8\pi^2} (\bar{\chi} \sigma_{\mu\nu} \chi) F^{\mu\nu},$$

$$\mathcal{Q}_2^{(5)} = \frac{e}{8\pi^2} (\bar{\chi} i \sigma_{\mu\nu} \gamma_5 \chi) F^{\mu\nu}$$

$$\mathcal{Q}_{1,q}^{(6)} = (\bar{\chi} \gamma_\mu \chi) (\bar{q} \gamma^\mu q),$$

$$\mathcal{Q}_{2,q}^{(6)} = (\bar{\chi} \gamma_\mu \gamma_5 \chi) (\bar{q} \gamma^\mu q),$$

$$\mathcal{Q}_{3,q}^{(6)} = (\bar{\chi} \gamma_\mu \chi) (\bar{q} \gamma^\mu \gamma_5 q),$$

$$\mathcal{Q}_{4,q}^{(6)} = (\bar{\chi} \gamma_\mu \gamma_5 \chi) (\bar{q} \gamma^\mu \gamma_5 q).$$

$$\mathcal{Q}_1^{(7)} = \frac{\alpha_s}{12\pi} (\bar{\chi} \chi) G^{a\mu\nu} G_{\mu\nu}^a,$$

$$\mathcal{Q}_2^{(7)} = \frac{\alpha_s}{12\pi} (\bar{\chi} i \gamma_5 \chi) G^{a\mu\nu} G_{\mu\nu}^a,$$

$$\mathcal{Q}_3^{(7)} = \frac{\alpha_s}{8\pi} (\bar{\chi} \chi) G^{a\mu\nu} \tilde{G}_{\mu\nu}^a,$$

$$\mathcal{Q}_4^{(7)} = \frac{\alpha_s}{8\pi} (\bar{\chi} i \gamma_5 \chi) G^{a\mu\nu} \tilde{G}_{\mu\nu}^a,$$

$$\mathcal{Q}_{5,q}^{(7)} = m_q (\bar{\chi} \chi) (\bar{q} q),$$

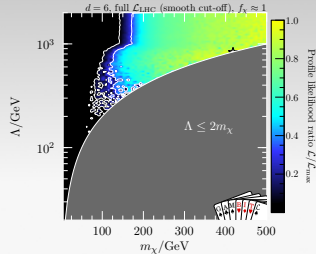
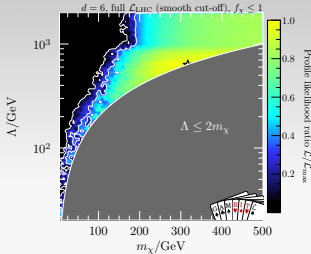
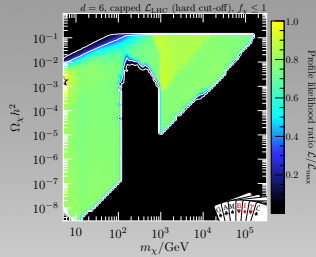
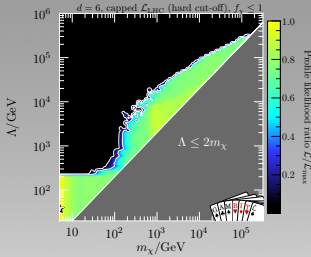
$$\mathcal{Q}_{6,q}^{(7)} = m_q (\bar{\chi} i \gamma_5 \chi) (\bar{q} q),$$

$$\mathcal{Q}_{7,q}^{(7)} = m_q (\bar{\chi} \chi) (\bar{q} i \gamma_5 q),$$

$$\mathcal{Q}_{8,q}^{(7)} = m_q (\bar{\chi} i \gamma_5 \chi) (\bar{q} i \gamma_5 q),$$

$$\mathcal{Q}_{9,q}^{(7)} = m_q (\bar{\chi} \sigma^{\mu\nu} \chi) (\bar{q} \sigma_{\mu\nu} q),$$

$$\mathcal{Q}_{10,q}^{(7)} = m_q (\bar{\chi} i \sigma^{\mu\nu} \gamma_5 \chi) (\bar{q} \sigma_{\mu\nu} q).$$



# DM EFT

- Running and mixing

→ For direct detection WCs are needed at  $\mu = 2 \text{ GeV}$  (DirectDM)

→ For  $\Lambda > m_t(m_t)$ :

$$\mathcal{C}_{1,2}^{(5)} = -4 \frac{m_t(m_t)^2}{\Lambda^2} \log \frac{\Lambda^2}{m_t(m_t)^2} \mathcal{C}_{9,10}^{(7)}$$

$$\Delta \mathcal{C}_i^{(7)} = -\mathcal{C}_{i+4,q}^{(7)} \quad (i = 1, 2)$$

$$\Delta \mathcal{C}_i^{(7)} = \mathcal{C}_{i+4,q}^{(7)} \quad (i = 3, 4)$$

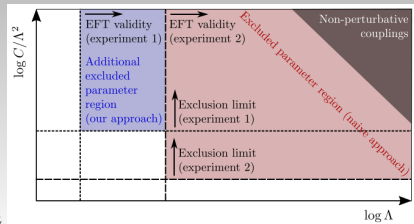
- EFT validity,  $\Lambda$  free parameter

→ DD requires  $\Lambda > 2 \text{ GeV}$

→ Annihilation processes (ID/RD) require  $\Lambda > 2m_\chi$

→ Collider searches  $\Lambda > \cancel{E}_T$

$$\Lambda < \cancel{E}_T \quad \left\{ \begin{array}{l} \frac{d\sigma}{d\cancel{E}_T} = 0 \\ \frac{d\sigma}{d\cancel{E}_T} \rightarrow \frac{d\sigma}{d\cancel{E}_T} \left( \frac{\cancel{E}_T}{\Lambda} \right)^{-a} \end{array} \right.$$



# Likelihoods

- Direct Detection

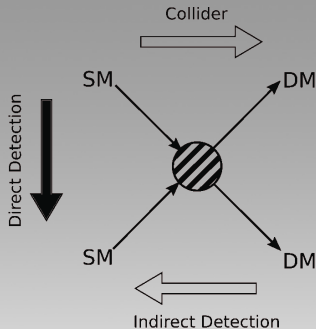
$$\frac{dR}{dE_R} = \frac{\rho}{m_T m_\chi} \int_{v_{\min}}^{\infty} v f(v) \frac{d\sigma}{dE_R} d^3v$$

$$v_{\min}(E_R) = \sqrt{\frac{m_T E_R}{2 \mu^2}}$$

→ Non-relativistic operators

$$\mathcal{L}_{\text{NR}} = \sum_{i,N} c_i^N (q^2) \mathcal{O}_i^N,$$

→ XENON1T, LUX 2016, PandaX 2016-17, CDMSlite, CRESST-II, CRESST-III, PICO-60 2017-19, and DarkSide-50



- Relic abundance  $\frac{dn_\chi}{dt} + 3Hn_\chi = -\langle\sigma v_{\text{rel}}\rangle (n_\chi n_{\bar{\chi}} - n_{\chi,\text{eq}} n_{\bar{\chi},\text{eq}})$   
 → Planck 2018:  $\Omega_{\text{DM}} h^2 \leq 0.120 \pm 0.001$

# Likelihoods

- Indirect detection with  $\gamma$ -rays
  - $\gamma$ -rays from DM annihilation in dSphs

$$\ln \mathcal{L}_{\text{dwarfs}}^{\text{prof.}} = \ln \mathcal{L}_{ki}(\Phi_i \cdot J_k) + \ln \mathcal{L}_J$$

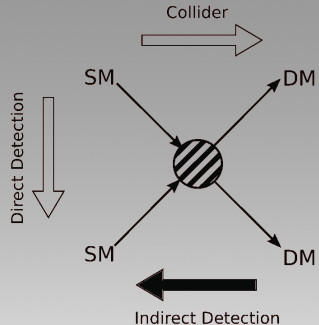
- Pass-8 combined of 15 dSphs from *Fermi*-LAT data

- Indirect detection with  $\nu$ s
  - Solar capture of DM leads to very high energy  $\nu$ s  $>$  solar  $\nu$ s
  - 79-string IceCube search

- Indirect detection constraints from CMB

- Injected energy ( $\gamma, e^\pm$ ) changes reion history and optical depth  $\tau$
- CMB is sensitive to energy deposition efficiency  $f_{\text{eff}}$  via combination

$$p_{\text{ann}} = f_\chi f_{\text{eff}} \frac{\langle \sigma v \rangle}{m_\chi}$$



# Likelihoods

- Collider constraints

→ Many signatures for DM searches

$$pp \rightarrow \chi\chi j \rightarrow j + \cancel{E}_T$$

→ MadGraph\_aMC@NLO  $\rightsquigarrow$  Pythia

→ Interpolated grids for  $\sigma$  and  $\epsilon A$

→ Events per  $\cancel{E}_T$  bin (signal regions)

$$N = L \times \sigma \times (\epsilon A)$$

→ ATLAS  $139\text{fb}^{-1}$  mono-jet

$\rightsquigarrow$  SR with best significance

$$\rightsquigarrow \mathcal{L}_{\text{ATLAS}}(s_i) \equiv \mathcal{L}_{\text{ATLAS}}(s_i, \hat{\gamma}_i)$$

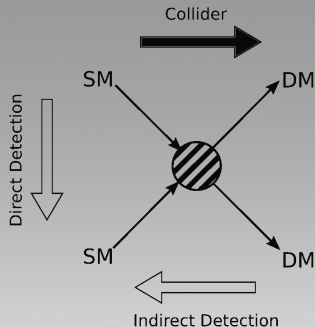
→ *Capped* likelihood

$$\mathcal{L}_{\text{cap}}(\mathbf{s}) = \min[\mathcal{L}_{\text{LHC}}(\mathbf{s}), \mathcal{L}_{\text{LHC}}(\mathbf{s} = \mathbf{0})]$$

→ CMS  $36\text{fb}^{-1}$  mono-jet

$\rightsquigarrow$  Profile over systematics

$$\rightsquigarrow \mathcal{L}_{\text{CMS}}(\mathbf{s}) \equiv \mathcal{L}_{\text{CMS}}(\mathbf{s}, \hat{\gamma})$$





# Scan framework

- Model parameters

DM mass	$m_\chi$
New physics scale	$\Lambda$
Wilson coefficients	$C_a^{(d)}$

- Nuisance parameters

Local DM density	$\rho_0$
Most probable speed	$v_{\text{peak}}$
Galactic escape speed	$v_{\text{esc}}$
<hr/>	
Running top mass ( $\overline{\text{MS}}$ scheme)	$m_t(m_t)$
<hr/>	
Pion-nucleon sigma term	$\sigma_{\pi N}$
s-quark contrib. to nucleon spin	$\Delta s$
s-quark nuclear tensor charge	$g_T^s$
s-quark charge radius of the proton	$r_s^2$

- Needs smart sampling to efficiently scan over all parameters and explore interference effects among WCs

## GAMBIT: The Global And Modular BSM Inference Tool

[gambit.hepforge.org](https://gambit.hepforge.org)

[github.com/GambitBSM](https://github.com/GambitBSM)

EPIC 77 (2017) 784

arXiv:1705.07908

- Extensive model database, beyond SUSY
- Fast definition of new datasets, theories
- Extensive observable/data libraries
- Plug&play scanning/physics/likelihood packs
- Various statistical options (frequentist /Bayesian)
- Fast LHC likelihood calculator
- Massively parallel
- Fully open-source



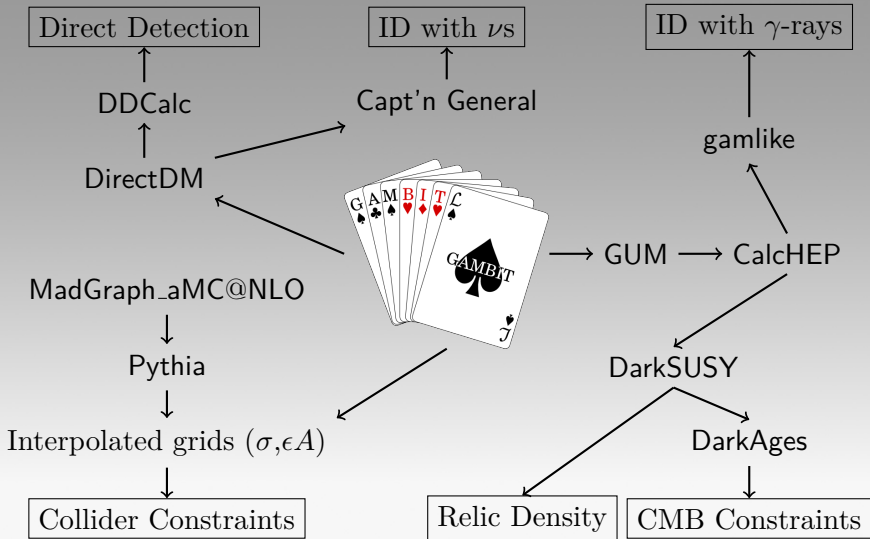
**Members of:** ATLAS, Belle-II, CLIC, CMS, CTA, Fermi-LAT, DARWIN, IceCube, LHCb, SHIP, XENON

**Authors of:** BubbleProfiler, Capt'n General, Contur, DarkAges, DarkSUSY, DDCalc, DirectDM, Diver, EasyScanHEP, ExoCLASS, FlexibleSUSY, gamLike, GM2Calc, HEPLike, IsaTools, MARTY, nuLike, PhaseTracer, PolyChord, Rivet, SOFTSUSY, SuperIso, SUSY-AI, xsec, Vevacious, WIMPSim

**Recent collaborators:** V Ananyev, P Athron, N Avis-Kozar, C Balázs, A Beniwal, LL Braseth, T Bringmann, A Buckley, J Butterworth, JE Camargo-Molina, C Chang, J Cornell, M Danninger, A Fowlie, T Gonzalo, W Handley, S Hoof, A Jueid, F Kahlhoefer, A Kvellestad, M Lecroq, C Lin, M Lucente, FN Mahmoudi, DJE Marsh, G Martinez, H Pacey, MT Prim, T Procter, F Rajec, A Raklev, R Ruiz, A Scaffidi, P Scott, W Shorrock, C Sierra, P Stöcker, W Su, J Van den Abeele, A Vincent, M White, A Woodcock, Y Zhang ++

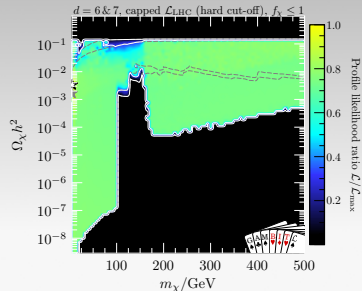
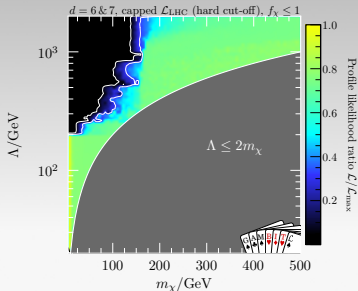
**70+ participants in many experiments and numerous major theory codes**

# Scan framework



# Results

- Include dim-7 operators,  $\Omega_{\text{DM}} h^2$  upper limit, LHC loglike *capped*
  - No change on large  $\Lambda$  - small  $m_\chi$  region
  - Neither  $Q_{1-4}^{(7)}$  (LHC) nor  $Q_{5-10,q}^{(7)}$  (suppressed) contribute to ann xsec
  - However, RD can be saturated for  $m_\chi < 100$  GeV (and small  $\Lambda$ )
  - $Q_3^{(7)}$  and  $Q_{7,q}^{(7)}$  give unconstrained signals in DD and ID
  - Similar fits to LHC excesses, even when dim-6 ops are zero

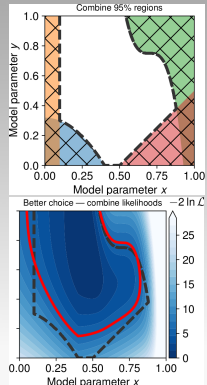


# Global fits

- Combine all constraints into a **composite likelihood**

$$\mathcal{L} = \mathcal{L}_{\text{Collider}} \mathcal{L}_{\text{Higgs}} \mathcal{L}_{\text{DM}} \mathcal{L}_{\text{Flavour}} \dots$$

- Perform an extensive **parameter scan**
  - Old-school sampling methods (random, grid) are inefficient
  - Harder to make statement about statistics
  - Need **smart sampling strategies** (differential, nested, genetic, ...)
  - **Rigorous** statistical interpretation (frequentist/Bayesian)
    - Goodness-of-fit
    - Parameter estimation
    - Model comparison



[arXiv:2012.09874 [hep-ph]]

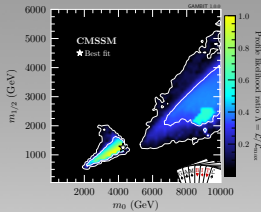
# Modules (Bits)

- **Physics Modules**
  - **ColliderBit**: collider searches [Eur.Phys.J. C77 (2017) no.11, 795]
  - **DarkBit**: relic density, dd,... [Eur.Phys.J. C77 (2017) no.12, 831]
  - **FlavBit**: flavour observables [Eur.Phys.J. C77 (2017) no.11, 786]
  - **SpecBit**: spectra, RGE running [Eur.Phys.J. C78 (2018) no.1, 22]
  - **DecayBit**: decay widths [Eur.Phys.J. C78 (2018) no.1, 22]
  - **PrecisionBit**: precision tests [Eur.Phys.J. C78 (2018) no.1, 22]
  - **NeutrinoBit**: neutrino likelihoods [Eur.Phys.J.C 80 (2020) no.6, 569]
  - **CosmoBit**: cosmological constraints [JCAP 02 (2021) 022]
- **ScannerBit** : stats and sampling [Eur.Phys.J. C77 (2017) no.11, 761]
  - Diver, GreAT, Multinest, Polychord, ...
- **Models**: hierarchical model database
- **Core** : dependency resolution [Eur.Phys.J. C78 (2018) no.2, 98]
- **Backends** : External tools to calculate observables
- **GUM**: Autogeneration of code [S. Bloor, TG, P. Scott et. al., soon]

# Examples

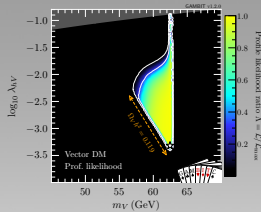
## CMSSM

[Eur.Phys.J.C 77 (2017) 12,824]



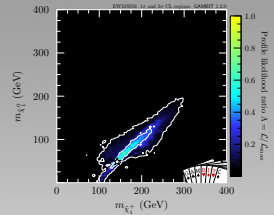
## Higgs-portal DM

[Eur.Phys.J.C 79 (2019) 1, 38]



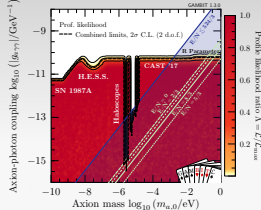
## MSSM-EW

[Eur.Phys.J.C 79 (2019) 5, 395]



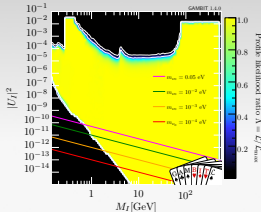
## QCD axions

[JHEP 03 (2019) 191]



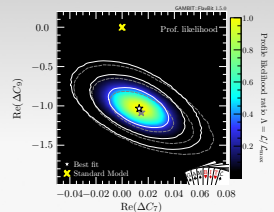
## Right-Handed Neutrinos

[Eur.Phys.J.C 80 (2020) 6, 569]



## Flavour EFT

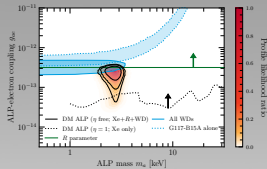
[Eur.Phys.J.C 81 (2021) 12, 1076]



# Examples

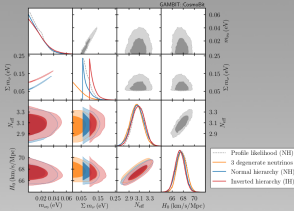
## DM ALPs

[JHEP 05 (2021) 159]



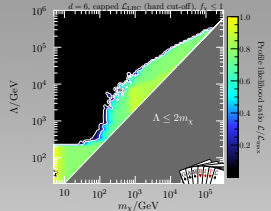
## Neutrino Masses

[Phys.Rev.D 103(2021)12,123508]



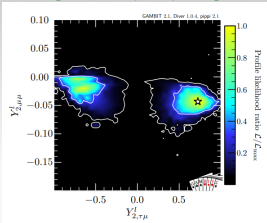
## DMEFT

[Eur.Phys.J.C 81 (2021) 11,992]



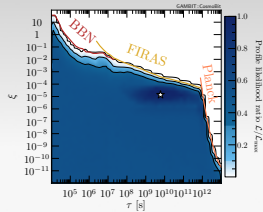
## THDM-III

[JHEP 01 (2022) 037]



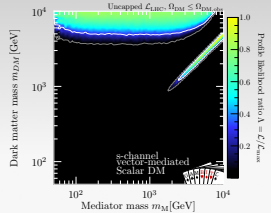
## Cosmo ALPs

[arXiv:2205.13549[astro-ph.CO]]



## S-channel DM

[arXiv:2209.13266 [hep-ph]]





# Core

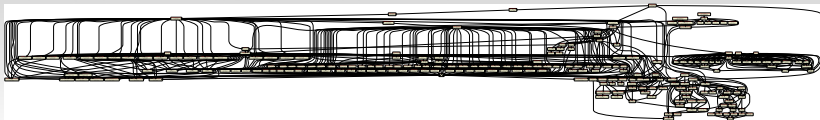
- Each module contains a collection of module functions
- Module functions provide a *capability*
- They have dependencies and backend requirements
- Allowed for specific models
- At run time a dependency tree is generated and resolved

```
// SM-like Higgs mass with theoretical uncertainties
#define CAPABILITY prec_nh
START_CAPABILITY

#define FUNCTION FH_HiggsMass
START_FUNCTION(triplet<double>)
DEPENDENCY(unImproved_MSSM_spectrum, Spectrum)
DEPENDENCY(FH_HiggsMasses, fh_HiggsMassObs)
ALLOW_MODELS(MSSM63atQ, MSSM63atMGUT)
#undef FUNCTION

#define FUNCTION SHD_HiggsMass
START_FUNCTION(triplet<double>)
DEPENDENCY(unImproved_MSSM_spectrum, Spectrum)
BACKEND_REQ(SUSYHD_MHiggs, (), MReal, (const MList<MReal>&))
BACKEND_REQ(SUSYHD_DeltaMHiggs, (), MReal, (const MList<MReal>&))
ALLOW_MODELS(MSSM63atQ, MSSM63atMGUT)
#undef FUNCTION

#undef CAPABILITY
```



# Models

- Extensive model database

## SUSY

CMSSM  
 NUHM1,2  
 MSSM63atQ

## DM

Scalar Singlet  
 Fermionic Singlet  
 Vector Singlet  
 Axions

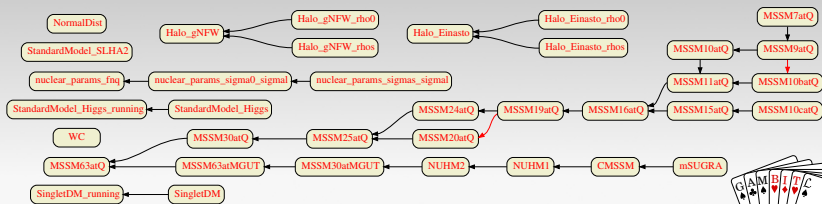
## Cosmo

$\Lambda$ CDM  
 $\Delta N_{\text{eff}}$   
 Power-law  
 inflation

## Others

SM  
 RH neutrinos  
 WC  
 nuisance models

- Parent-daughter hierarchy
- Module functions are activated for each model



# Backends

- C, Fortran  $\rightsquigarrow$  POSIX d1
- C++  $\rightsquigarrow$  BOSS + POSIX d1
- Mathematica  $\rightsquigarrow$  WSTP
- Python  $\rightsquigarrow$  pybind11

## CosmoBit

AlterBBN 2.2  
 DarkAges 1.2.0  
 MontePythonLike 3.3.0  
 MultiModeCode 2.0.0  
 classy 2.9.4  
 plc 3.0

## DarkBit

CaptnGeneral 1.0  
 DDCalc 2.2.0  
 DarkSUSY 6.2.2  
 MicrOmegas 3.6.9.2  
 gamLike 1.0.1  
 nulike 1.0.9

## ColliderBit

HiggsBounds 4.3.1  
 HiggsSignals 1.4  
 Pythia 8.212  
 nulike 1.0.9

## PrecisionBit

FeynHiggs 2.12.0  
 SUSYHD 1.0.2  
 gm2calc 1.3.0

## SpecBit

*FlexibleSUSY 2.0.1*  
 SPheno 4.0.3

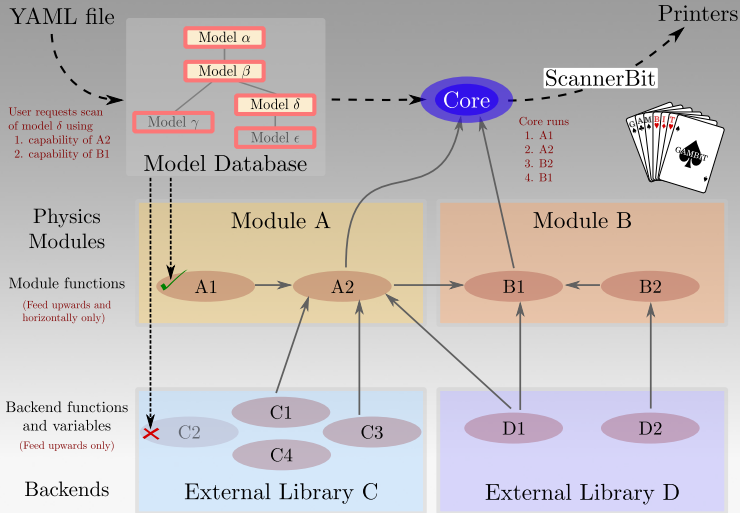
## FlavBit

SuperISO 3.6

## DecayBit

SUSY\_HIT 1.5

# An example run



# Operators

	SI scattering	SD scattering	Annihilations
$\mathcal{Q}_{1,q}^{(6)} = (\bar{\chi}\gamma_{\mu}\chi)(\bar{q}\gamma^{\mu}q)$	unsuppressed	—	s-wave
$\mathcal{Q}_{2,q}^{(6)} = (\bar{\chi}\gamma_{\mu}\gamma_5\chi)(\bar{q}\gamma^{\mu}q)$	suppressed	—	p-wave
$\mathcal{Q}_{3,q}^{(6)} = (\bar{\chi}\gamma_{\mu}\chi)(\bar{q}\gamma^{\mu}\gamma_5q)$	—	suppressed	s-wave
$\mathcal{Q}_{4,q}^{(6)} = (\bar{\chi}\gamma_{\mu}\gamma_5\chi)(\bar{q}\gamma^{\mu}\gamma_5q)$	—	unsuppressed	s-wave $\propto m_q^2/m_{\chi}^2$
$\mathcal{Q}_1^{(7)} = \frac{\alpha_s}{12\pi}(\bar{\chi}\chi)G^{a\mu\nu}G_{\mu\nu}^a$	unsuppressed	—	p-wave
$\mathcal{Q}_2^{(7)} = \frac{\alpha_s}{12\pi}(\bar{\chi}i\gamma_5\chi)G^{a\mu\nu}G_{\mu\nu}^a$	suppressed	—	s-wave
$\mathcal{Q}_3^{(7)} = \frac{\alpha_s}{8\pi}(\bar{\chi}\chi)G^{a\mu\nu}\tilde{G}_{\mu\nu}^a$	—	suppressed	p-wave
$\mathcal{Q}_4^{(7)} = \frac{\alpha_s}{8\pi}(\bar{\chi}i\gamma_5\chi)G^{a\mu\nu}\tilde{G}_{\mu\nu}^a$	—	suppressed	s-wave
$\mathcal{Q}_{5,q}^{(7)} = m_q(\bar{\chi}\chi)(\bar{q}q)$	unsuppressed	—	p-wave $\propto m_q^2/m_{\chi}^2$
$\mathcal{Q}_{6,q}^{(7)} = m_q(\bar{\chi}i\gamma_5\chi)(\bar{q}q)$	suppressed	—	s-wave $\propto m_q^2/m_{\chi}^2$
$\mathcal{Q}_{7,q}^{(7)} = m_q(\bar{\chi}\chi)(\bar{q}i\gamma_5q)$	—	suppressed	p-wave $\propto m_q^2/m_{\chi}^2$
$\mathcal{Q}_{8,q}^{(7)} = m_q(\bar{\chi}i\gamma_5\chi)(\bar{q}i\gamma_5q)$	—	suppressed	s-wave $\propto m_q^2/m_{\chi}^2$
$\mathcal{Q}_{9,q}^{(7)} = m_q(\bar{\chi}\sigma^{\mu\nu}\chi)(\bar{q}\sigma_{\mu\nu}q)$	loop-induced	unsuppressed	s-wave $\propto m_q^2/m_{\chi}^2$
$\mathcal{Q}_{10,q}^{(7)} = m_q(\bar{\chi}i\sigma^{\mu\nu}\gamma_5\chi)(\bar{q}\sigma_{\mu\nu}q)$	loop-induced	suppressed	s-wave $\propto m_q^2/m_{\chi}^2$

# Hadronic input parameters

Parameter	Value	Parameter	Value
$\sigma_{\pi N}$	50(15) MeV [1]	$\mu_p$	2.793 - [2]
$Bc_5(m_d - m_u)$	-0.51(8) MeV [3]	$\mu_n$	-1.913 [2]
$g_A$	1.2756(13) [2]	$\mu_s$	-0.036(21) [4]
$m_G$	836(17) MeV [1]	$g_T^u$	0.784(30) [5]
$\sigma_s$	52.9(7.0) MeV [6]	$g_T^d$	-0.204(15) [5]
$\Delta u + \Delta d$	0.440(44) [7]	$g_T^s$	$-27(16) \cdot 10^{-3}$ [5]
$\Delta s$	-0.035(9) [7]	$B_{T,10}^{u/p}$	3.0(1.5) [8]
$B_0 m_u$	0.0058(5) GeV <sup>2</sup> [9]	$B_{T,10}^{d/p}$	0.24(12) [8]
$B_0 m_d$	0.0124(5) GeV <sup>2</sup> [9]	$B_{T,10}^{s/p}$	0.0(2) [8]
$B_0 m_s$	0.249(9) GeV <sup>2</sup> [9]	$r_s^2$	-0.115(35) GeV <sup>-2</sup> [4]

[1][F. Bishara et. al., JHEP 11 (2017) 059] [2][PDG 2020] [3][A. Crivellin et. al., Phys. Rev. D 89 (2014) 054021] [4][D. Djukanovic et. al., Phys. Rev. Lett. 123 (2019) 212001, R. S. Sufian et. al., Phys. Rev. Lett. 118 (2017) 042001] [5][R. Gupta, et. al., Phys. Rev. D 98 (2018) 091501] [6][S. Aoki et. al., Eur. Phys. J. C 80 (2020) 113] [7][J. Liang et. al., Phys. Rev. D 98 (2018) 074505] [8][B. Pasquini et. al., Phys. Rev. D 72 (2005) 094029] [9][F. Bishara et. al., arXiv:1708.02678.]

# Nuisance parameters

Nuisance parameter		Value ( $\pm 3\sigma$ range)
Local DM density	$\rho_0$	0.2–0.8 GeV cm <sup>-3</sup>
Most probable speed	$v_{\text{peak}}$	240 (24) km s <sup>-1</sup>
Galactic escape speed	$v_{\text{esc}}$	528 (75) km s <sup>-1</sup>
Running top mass ( $\overline{\text{MS}}$ scheme)	$m_t(m_t)$	162.9 (6.0) GeV
Pion-nucleon sigma term	$\sigma_{\pi N}$	50 (45) MeV
Strange quark contrib. to nucleon spin	$\Delta s$	-0.035 (0.027)
Strange quark nuclear tensor charge	$g_T^s$	-0.027 (0.048)
Strange quark charge radius of the proton	$r_s^2$	-0.115 (0.105) GeV <sup>-2</sup>

# Collider Likelihoods

- ATLAS, Poisson loglike marginalised over nuisance  $\xi =$  relative signal/bkg uncertainties

$$\mathcal{L}_{\text{marg}}(n|p) = \int_0^\infty \frac{[\xi p]^n e^{-\xi p}}{n!} \times \frac{1}{\sqrt{2\pi}\sigma_\xi} \frac{1}{\xi} \exp\left[-\frac{1}{2} \left(\frac{\ln \xi}{\sigma_\xi}\right)^2\right] d\xi.$$

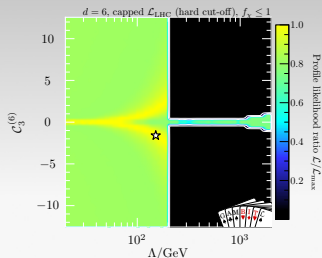
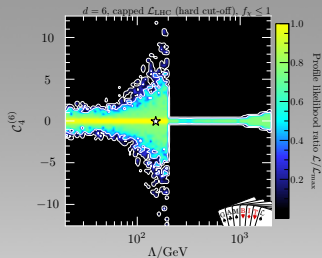
- CMS, convolved Poisson-Gaussian, profiled over systematic uncertainties  $\gamma$  on expected background yields with covariance matrix  $\Sigma$

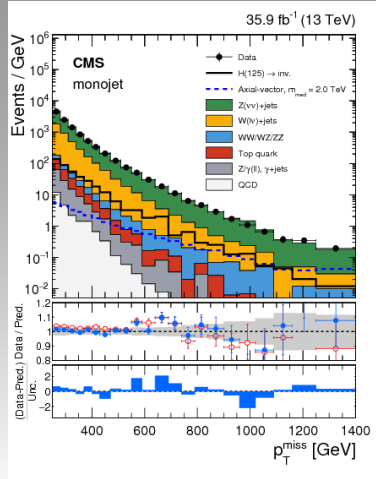
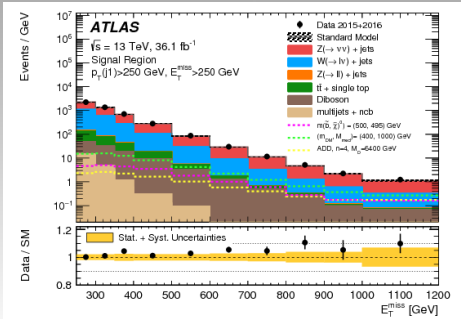
$$\mathcal{L}(\mathbf{s}, \gamma) = \prod_i^{N_{\text{bin}}} \left[ \frac{(s_i + b_i + \gamma_i)^{n_i} e^{-(s_i + b_i + \gamma_i)}}{n_i!} \right] \times \frac{1}{\sqrt{\det 2\pi\Sigma}} e^{-\frac{1}{2}\gamma^T \Sigma^{-1} \gamma}.$$



# Results

- $\mathcal{C}_1^{(6)}$ 
  - spin-independent scattering
  - strongly constrained  $\rightsquigarrow$  very small
- $\mathcal{C}_2^{(6)}$ 
  - momentum-dependent scattering
  - $\Lambda < 250$  GeV DD constrained
  - $\Lambda > 250$  GeV LHC constrained
- $\mathcal{C}_3^{(6)}$ 
  - *both* SD and MD scattering
  - $\Lambda < 250$  GeV weak DD constraints
  - Main contribution to *Fermi – LAT*
  - $\Lambda > 250$  GeV LHC constrained
- $\mathcal{C}_4^{(6)}$ 
  - spin-dependent scattering
  - identical to  $\mathcal{C}_2^{(6)}$





# But...

*How do I use GAMBIT with my favourite model?*

↪ Adding a model

↪ Sorting out hierarchy

↪ Making physics computations work with that model

*How do I add a new physical observable or likelihood?*

↪ Create capabilities

↪ Declare dependencies

↪ and models

↪ and backend requirements

## 1. Add the model to the **model hierarchy**:

- Choose a model name, and declare any **parent model**
- Declare the model's parameters
- Declare any **translation function** to the parent model

```
#define MODEL NUHM1
#define PARENT NUHM2
START_MODEL
DEF THEPARS(NO_W12,mf,AS,TanBeta,SigMa)
DECLAREPAR AS_PAREN_FUNCTION(NUHM1_to_NUHM2)
#undef PARENT
#undef MODEL
```

## 2. Write the translation function as a standard C++ function:

```
void MODEL_NAMESPACE::NUHM1_to_NUHM2 (const ModelParameters &mpP, ModelParameters &targetP)
{
    // Set NO, W12, AS, TanBeta and SigMa in the NUHM2 to the same values as in the NUHM1
    targetP.setValues(mpP,false);
    // Set the values of mf and mf in the NUHM2 to the value of mf in the NUHM1
    targetP.setValues("mf", mpP["mf"]);
    targetP.setValues("mf2", mpP["mf"]);
}
```

## 3. If needed, declare that existing module functions work with the new model, or add new functions that do.

Adding a new module function is easy:

## 1. Declare the function to GAMBIT in a module's **rollcall header**

- Choose a capability
- Declare any **backend requirements**
- Declare any **dependencies**
- Declare any specific **allowed models**
- other more advanced declarations also available

```
#define MODULE FlavBit // A tasty GAMBIT module.
START_MODULE

#define CAPABILITY Flav // Observable: BR(K->nu nu)/BR(pi->nu nu)
START_CAPABILITY
#define FUNCTION SI_Flav // Name of a function that can compute Flav
FUNCTION(double) // Function computes a double precision result
BACKEND(Superlu, ylmom, (ey_tag), double, (const parameters*)) // Needs function from a backend
BACKEND_OPTION( (Superlu, 8.0), (ey_tag) ) // Backend must be Superlu 3.6
DEPENDENCY(Superlu_modelinfo, parameters) // Needs another function to calculate Superlu info
ALLOW_MODEL(NUHM2b0, NUHM2b0D07) // Works with weak/DP-scale NUHM and descendants
#undef FUNCTION
#undef CAPABILITY
```

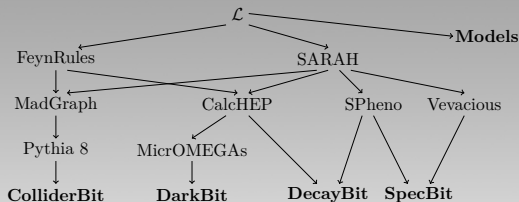
## 2. Write the function as a standard C++ function (one argument: the result)

# Solution

The **G**AMBIT **U**niversal **M**odel Machine



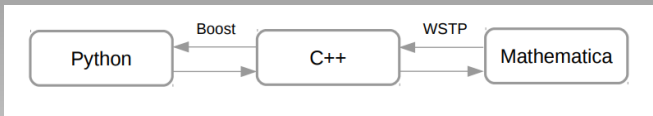
- GUM interfaces LLT SARAH and FeynRules with GAMBIT
- Uses existing HEP toolchains



- GAMBIT-compatible outputs from GUM

Generated output	FeynRules	SARAH	Usage in GAMBIT
CalcHEP	✓	✓	Decays, cross-sections
micrOMEGAs (via CalcHEP)	✓	✓	DM observables
Pythia (via MadGraph)	✓	✓	Collider physics
SPheno	✗	✓	Particle mass spectra, decay widths
Vevacious	✗	✓	Vacuum stability

- Primarily written in Python, with interface to Mathematica via Boost and WSTP



- Automatically generates GAMBIT code
  - Particles → particle database and parameters → Models
  - Module functions for ColliderBit, DarkBit, DecayBit and SpecBit
  - Writes interfaces to requested backends
- GUM will release with GAMBIT 2.0 **VERY SOON**

# An example

- Majorana DM  $\chi$  with scalar mediator  $Y$

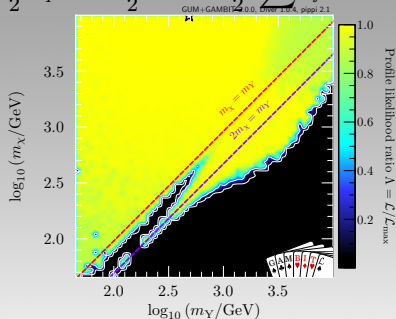
$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \frac{1}{2} \bar{\chi} (i\not{\partial} - m_\chi) \chi + \frac{1}{2} \partial_\mu Y \partial^\mu Y - \frac{1}{2} m_Y^2 Y^2 - \frac{g_\chi}{2} \bar{\chi} \chi Y - \frac{c_Y}{2} \sum_f y_f \bar{f} f Y.$$

```

math:
# Choose FeynRules
package: feynrules
# Name of the model
model: MDMSM
# Model builds on the Standard Model FeynRules file
base_model: SM
# The Lagrangian is defined by the DM sector (LDM),
# defined in MDMSM.fr, plus the SM Lagrangian (LSM)
# imported from the 'base model', SM.fr
Lagrangian: LDM + LSM
# Make CKM matrix = identity to simplify output
restriction: DiagonalCKM

# PDG code of the annihilating DM candidate in
# FeynRules file
wimp_candidate: 52

# Select outputs for DM physics.
# Collider physics is not as important in this model.
output:
  pythia: false
  calchep: true
  micromegas: true
  
```



→ Follow Sanjay's tutorial  
3pm Room A