

## **Timing RPCs: 25 years**

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### **Outlook**

The early (good old) days

tRPCs in the world

tRPC physics: what is done and what remains to be done

#### **ALICE TOF requirements**

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Requirements (ALICE TDP CERN/LHCC/95–71):
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 $\sim$ 150 m<sup>2</sup> area granularity (determined by occupancy)  $\sim$ 9 cm<sup>2</sup> (160000 channels) system time resolution  $\leq$  100 ps

#### **Existing (and discarded by 1998) technologies**



### **ALICE TOF candidate technologies (autumn 1998 – spring 1999)**

- Double Parallel Plate Chamber (PPC)
- Being studied for calorimetry since the early 1990s' mainly by the ITEP group. This was the baseline solution and had been under development for a long time(see V.A.Akimov et al., Instrum. and Exp. Tech. 45:4 (2002) 493)Time resolution ~200 ps (already a big step forward). Prone to sparking...
- Melamine or glass multigap RPC (MRPC) Based on the muon trigger RPCs that were being studied by the LAA group for ATLAS/CMS since 1996.
- Metal-glass symmetric multigap RPC Proposed by the Coimbra group. Inspired by the Pestov spark counter and by both detectors mentioned above.

There were two beamtimes: autumn 1998 and spring 1999.

**The spring 1999 beamtime was supposed to demonstrate the "1m2" prototype as the final step for the elaboration of the TOF TDR, which was due in July 1999**.

The ALICE TOF R&D effort was coordinated by François Piuz and W.Klempt.

#### **Double PPC**



32 channels prototype



#### Chromium-coated ceramic plates



## **Melamine/glass MRPC**Autumn 1998 Spring 1999



### **Metal-glass symmetric MRPC**



The metal plates were borrowed from the ceramic PPCs.

P.Fonte et al., NIM A (2000) 201



#### 32 channelsprototype



Read out at ground potential



7Designed to allow to place shielding between the cells to minimize crosstalk.New dedicated electronics was produced.A CERN/Coimbra/ITEP collaboration

100

98

96

94

92

90

TOF efficiency (%)

#### **Results of the autumn 1998 beamtime**

Nuclear Instruments and Methods in Physics Research A 443 (2000) 201-204

Letter to the Editor The following colleagues, from our home insti-A new high-resolution TOF technology tutes, contributed to this work: A. Akindinov, E. Cerron-Zeballos, R. Ferreira-Marques, V. Go-P. Fontea,b,\*,1, A. Smirnitski<sup>c</sup>, M.C.S. Williams<sup>a,d</sup> lovine, D. Hatzifotiadou, J. Lamas-Valverde, A. Martemianov, V. Petrov, F. Piuz, A. Policarpo, <sup>a</sup>CERN, EP Division, 1211 Geneva 23, Switzerland K. Voloshin. <sup>b</sup>LIP-Coimbra, Coimbra, Portugal <sup>c</sup>ITEP. Moscow. Russia <sup>d</sup>INFN Bologna, Bologna, Italy Received 24 March 1999; received in revised form 27 August 1999; accepted 2 September 1999 180 B 160  $\circ$ 8 Ō 140 TOF sigma (ps) 120 100 Conductive Spacers 80  $(0.3$ mm $)$ glass 60  $\bullet$  TOF sigma 40 O TOF efficiency 20 Metallisation Electrical Ceramic  $\Omega$ plate Contact 1600 2000 2400 0 400 800 1200 Counting rate (Hz/cm<sup>2</sup>)

> Fig. 5. Timing resolution and efficiency as a function of the counting rate per unit area. For counting rates below 800 Hz/cm<sup>2</sup> a resolution better than 120 ps sigma was achieved with efficiency above 98%.

EUROPEAN LABORATORY FOR PARTICLE PHYSICS (CERN)

**CERN-EP 99-166** 28 October 1999

### A four-gap glass-RPC time-of-flight array with 90 ps time resolution

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Figure 5: One eight-channel prototype readout card mounted on the prototype. The card carlifiers made with discrete components, discrimination logics and voltage regulation







#### 10





A. Akindinov et al., CERN EP 99-166 (28 Oct. 1999); IEEE Trans. Nucl. Sci 48, 5 (2001) 1658



A. Akindinov et al., CERN EP 99-166 (28 Oct. 1999); IEEE Trans. Nucl. Sci 48, 5 (2001) 1658





Figure 2.23: TOF module test results (10 PPCs shown):

- $a$  TOF resolutions;
- $b$  efficiency plateau widths;
- $c$  efficiencies in the plateau centres.



#### (study of single gaps)**Results of the spring 1999 beamtime**Recent developments in very high time resolution RPCs **CERN 1999 P.Fonte RECENT DEVELOPMENTS IN** Event by event ratio of fast/total charge **VERY HIGH TIME RESOLUTION RPCs**  $(fast charge = electron component)$ Standard mixture P. Fonte Methane - 0.1 mm Isobutane - 0.1 mm  $0.1$  mm **CERN-EP**  $\alpha(Q)$ d  $0.8$  $0.8$  $\alpha d$ 0.8  $0.6$ and  $0.6$  $0.6$ LIP-Coimbra, Portugal  $0.4$  $0.4$  $0.4$  $0.2$  $0.2$  $0.2$ 0.5 0.5  $0.5$ Methane - 0.3 mm Isobutane - 0.3 mm Standard mixture Fast charge (unid.arb.)  $0.3$  mm  $2.5$ 1.5  $\Omega$  $\Omega$ **EP Detector Seminar** Total charge  $(pC)$ **CERN 27/9/99** Fast charge vs. Total charge • Electronic charge/total charge should be a constant  $[1/(\alpha d)]$  but for most mixtures (except methane in a

0.1 mm gap) there is a strong non-linearity  $\Rightarrow$  space-charge effect.

Later published asP. Fonte, V. Peskov. NIM A 477 (2002) 17

#### **ALICE TOF 2**×**8 pads MRPC prototype**



Figure 2.48: Artist's view of strip detector used in the November 1999 test beam.



8x2 cell strip detector 12.5 kV

Figure 2.52: Efficiency and resolution of the 16 cells of the strip chamber.



igure 2.49: Cross section of the strip detector used in the November 1999 test beam.

Crosstalk not measuredFinal ALICE TOF design was symmetric double-stack

#### **The HARP experiment**



Fig. 3. Cross-section through the glass stack.



Fig. 4. Layout of the readout electrodes; eight pads are connected to the same prear

Symmetric double-gap MRPCs $4 \times 0.3$  mm gaps All electrodes glass368 pads of 29 $\times$ 106 mm<sup>2</sup> = 30.7 cm<sup>2</sup> readout on one side onlyTotal area  $\sim$ 6 m2

The first experiment to use timing RPCs.

Built in 5 months from December 2000!



#### **The HARP experiment**



Fig. 7. Time resolution and efficiency as function of the impact point of the beam. The  $x$  direction is along the RPC, one readout channel covers 240 mm; the boundary between the pads read out by two different preamplifiers is at  $x = 750$  mm.

#### The first experiment to use timing RPCs.



Fig. 8. Time-charge relation for an RPC exposed to a  $-8$  GeV/c beam.

#### Resolution 140 to 170 ps.

Later there was some controversy about the performance of the detector in the experiment.



## **Extension to large area/channel** (inspired by the needs of the HARP experiment)



A.Blanco et al., arXiv:physics/0103086 (26 March, 2001)A.Blanco et al.,. NIM A 485 (2002) 328



## **Extension to large area/channel** (inspired by the needs of the HARP experiment)



No degradation when the area/channel was doubled (800 cm<sup>2</sup>/channel)

A.Blanco et al., arXiv:physics/0103086 (26 March, 2001)A.Blanco et al.,. NIM A 485 (2002) 328

#### **The FOPI experiment**

The FOPI experiment (GSI) had been already for several years pursuing the Pestovspark counter technology, but quickly adopted the tRPC approach owing to its enormous practical advantages over the very technically challenging Pestov counters.



Final resolution 70 to 80 psLimited multihit capability

#### **Use of tRPCs in physics experiments**



Except for HARP, in all cases the time resolution is typically between 50 and 100 ps  $\sigma$ 

Clearly the future directions are:

- large count rate density
- super resolution ~20ps
- simultaneous accurate position resolution (TOF-tracker)

(please suggest corrections/additions to this list for the paper)

#### **Overview of detector structures - chambers**





#### Triple-stack (MPD TOF)Honeycom panel (5 mm) PET scrue Outer PCB (1.5 mm) الها Mylar  $(100 \mu m)$  $HV-$ Outer HV glass  $(400 \mu m)$ Inner glass  $(280 \mu m)$ Spacer (monofilment 200 µm) HV+ Inner PCB with strips (1.5 mm) HV-HV+ Differential preamp  $\left| \frac{10}{10} \right|$  (pitch - 12.5 mm) 300

Typically 4 to 10 gaps with width in the range 0.2 to 0.3 mm.

#### **Overview of detector structures - readout**Multistrip (most – OK for low occupancy)



#### Single shielded cells (HADES)



Strips tend to be <4 cm wide. No strong length limitation (transmission lines).





muon tomography, RPC-PET, etc. 25 Also useful for

# **tRPC <sup>p</sup>hysics - efficiency**

An efficiency of 75% has been measured on single 0.3 mm gaps. How to understand this?



The probability that no charge is created on the efficient part of the gap is P(0)= $e^{-\lambda g^*}$ =0.25 where  $\lambda$  is the primary ionization density <~9/mm, so  $g^*$ =0.15 mm: about half of the gap.

The problem is that for an electron to generate sufficient charge on half the gap  $(z_0 = g - g^* \sim g/2)$ , lets say, modestly,  $10^5$  e<sup>-</sup> (16 fC), then an electron released from the cathode would generate  $(10^5)^2\!\!=\!10^{10}$  electrons!

Such avalanches were never seen. The famous Raether limit is  $10^8$  electrons and that's for wide ~cm gaps. For small gaps it will be less.

# **tRPC <sup>p</sup>hysics – space charge**

As the avalanche grows the first Townsend coefficient gets smaller and this limits the final size of the larger avalanches.Standard mixture

There is rather direct evidence of this:



Fast charge vs. Total charge Sometimes this same effect generates a streamer and that defines the gain limit of the (avalanche mode) RPC

P. Fonte, V. Peskov. NIM A 477 (2002) 17



 (theory developed by several people over ~ a decade)**tRPC <sup>p</sup>hysics – timing** How can the excellent time resolution be understood on a gaseous detector where the transit time of the electrons across the gap is on the order of 3 ns? **1st order explanation: the time is determined already during the progression of the avalanches, so it doesn't depend on the position where they started.**



Progression distance or time

A variable number of primary electrons is created in each gap (Poisson distribution): primary statistics.

Each primary electron generates an avalanche whose charge is noisy in its preliminary stages (exponential distribution in the worst case): avalanche statistics

Avalanches started too close to the cathode don't grow enough to contribute to the current at the level of the timing threshold, so there is an effective gap width  $g^*(-g/2)$ .

Above ~100 e<sup>-</sup> each avalanche starts to behave deterministically.

. See P.Fonte, JINST 8 (2013) P11001 and refs therein 28 All these effects can be summarized by extrapolating the deterministic part back to time 0 and growingdeterministically from there (red dashed lines).

analytically:

#### (theory developed by several people over ~ a decade)**tRPC <sup>p</sup>hysics – timing** The consequences of these (rather straightforward) assumptions have been worked out

Single-electron avalanche of $\mathbb{F}_{\mathcal{A}+\mathcal{D}}(ST) = ue^{-u} \frac{\sqrt{r\lambda g^*}I_1\left(2\sqrt{r\lambda g^*u}\right)}{\left(e^{r\lambda g^*}-1\right)\sqrt{u}}, u = e^{\ln(r)+ST_u - ST_u}$ 

fixed standard deviation  $= 1.28$ 

 $\alpha^*$  = effective ionization density in the avalanche

*λ=* primary ionization density

*r =* probability of multiplication vs. attachment.

 $S = \alpha^* v_e$  is the ionization rate in the avalanche. This is the basic resolution-scaling parameter. Therefore the famous formula for the time resolution  $\sigma = 1,28/(\alpha *v_e)$ , corresponding to the single-electron limit.

A remarkable feature of this distribution is that its shape depends only on *r λ <sup>g</sup>\** = number of primary charges in the effective gap region that have not been captured by the electronegative gas (effective primary charge)  $\Rightarrow$  related to the intrinsic inefficiency of the detector intrinsic inefficiency of the detector.

*I1* is the modified Bessel function of first order.



 (theory developed by several people over ~ a decade)**tRPC <sup>p</sup>hysics – timing**

The variance of this distribution is also known analytically (in series form)



 $\left( g^{\ast }\,/\,g\right)$  $\left( \alpha ^{*}g\right)$ \*  $\alpha^*\nu_e$   $\sqrt{r\lambda\left(g^*/g\right)\over \frac{Q^*}{\lambda^{1/2}}}$   $\frac{(\alpha^*)}{\text{limited by}}$ ~1/2 limited by (will increase $e$  streamer onset for smaller gaps) (will decrease for smaller gaps)2 1  $2g$  1 /*T* $r\lambda g^* \alpha^* v_e \qquad \boxed{r\lambda\ \left( g^*/\,g\right)} \qquad \boxed{\alpha^*}$ *e* | (0 0) (0 *e g* $V_e$   $\left[\right.$   $\left.\left\vert r\lambda \right. \right. \left. \left. \left( \begin{matrix} g & g \\ g & g \end{matrix} \right) \right.$   $\left( \begin{matrix} \alpha & g \\ g & g \end{matrix} \right)$   $\left. \right.$   $V_e$ σ≃= $\overbrace{\hspace{2.5cm}}^{\hspace{2.5cm} \frown\hspace{2.8cm}}$ 

But beware that this  $\alpha^*$  is the one unaffected by space charge. It is much larger than what  $v_e$  can be inferred from the observed charge.

In very thin gaps at some point electronics and mechanics will start to dominate.

30See D.González-Diáz et al., JINST 12 (2017) C03029In MRPC the currents induced from all N gaps add analogically, so this is equivalent to replacing  $\lambda \rightarrow N \lambda$  and all the rest will be the same.

# **tRPC <sup>p</sup>hysics – what is more or less done**

Many other aspects of tRPC physics have been worked out analytically and/or numerically in more or less detail over the last 25 years:

- Progression of avalanches, space-charge regime and streamer onset
- Deterministic and stochastic voltage drops on the electrodes
- Shape of the charge distribution (in Townsend regime only)
- Signal induction in the presence of conductive materials

see W. Riegler and P. Windischhofer, NIM A 980 (2020) 164471

• Signal propagation in multi-conductor transmission lines

# **tRPC <sup>p</sup>hysics – what remains to be done**

A comprehensive simulation of RPC behavior, benchmarked with reality. This is a long term objective of DRD1 WG4. Anyone interested to contribute is welcome.

All analytical models have been formulated in the Townsend regime…

 There are 3 empirical analytical models of space charge. It has been shown that for the practical ranges of interest they are almost equivalent. But the comparison with reality or even numerical simulations hasn't been done.

The time resolution seems to be unaffected when the timing threshold lies in the space charge region. Both analytical and numerical models agree on this. Why?

The comparison between analytical predictions for time resolution and experiment is difficult owing to the large influence of technical factors, It lacks a fundamental measurement, free as possible of such complications.

Is the time-charge correction correcting something more than the amplifier rise time?Analytical calculation of the charge distribution in the space-charge regime.

Some prediction, even approximate, of where the avalanche-streamer transition will occur. This is a practical limitation to  $\alpha^*$  and therefore to the time resolution.

Clarification of the role of  $SF_6$  in streamer suppression and of Ar in streamer enhancement.

### **Conclusion**

Timing RPCs have been invented in the framework of the ALICE TOF R&D program in 1998/99 and opened a new era of large-area TOF systems for nuclear and HEP.

Have been or are being used in 10 physics experiments, with a total active area of 590 m2 readout by almost 200000 channels. Typical time resolutions range from 50 to 100 ps.

Are being proposed for 7 future experiments, some with requirements beyond the current baseline technology:

- large count rate density ( $\sim$  tens of kHz/cm<sup>2</sup>)
- super time resolution (~20ps)
- simultaneous accurate (<1 mm) position resolution (TOFtracker)

The advent of the SiPM has brought the scintillator technology into the range of<br>options. It remains to be seen a large area detector with resolution <100 ps, but it options. It remains to be seen a large area detector with resolution <100 ps, but it is the baseline technology for the SHiP TD  $(50 \text{ m}^2)$ .

 Considerable progress has been made on the understanding and modeling of tRPCphysics, but a lot still remains to be done. The DRD1 Work Group 4 welcomes contributors on these matters.