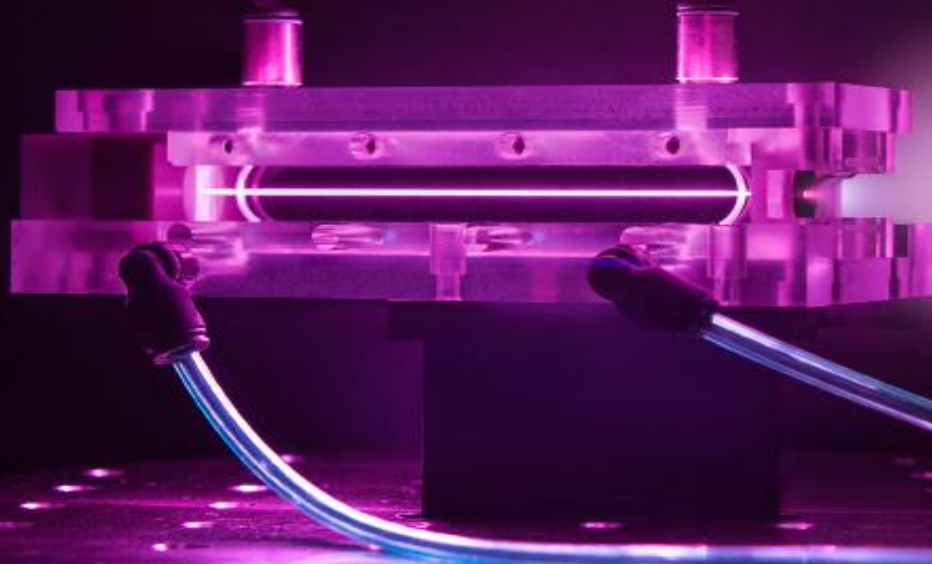


Advanced Accelerator Concepts I

Massimo.Ferrario@lnf.infn.it



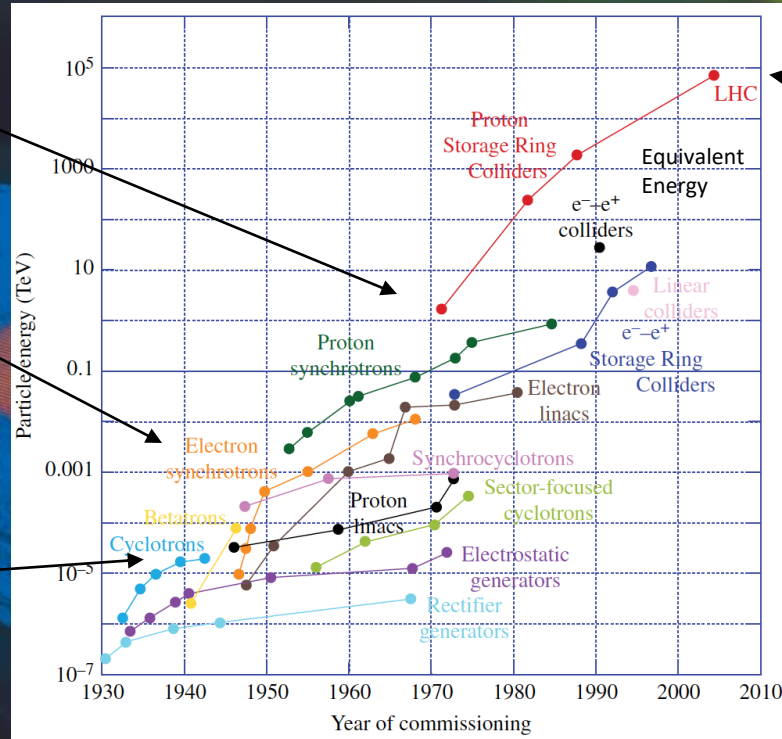
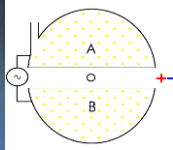
Introduction to Accelerator Physics

S. Susanna – 30 September 2024

Livingstone Diagram



$$\begin{cases} \omega_L = \frac{qB_y}{\gamma m_0} \\ p_z = qB_y R \end{cases}$$



Energy of colliders is plotted in terms of the laboratory energy of particles colliding with a proton at rest to reach the same center of mass energy.

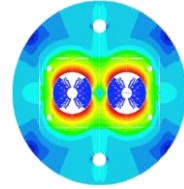
Options towards higher energies

Hadron (p) circular collider

$$p = e \cdot R \cdot B_y$$

Increase bending field
SC bend magnet work (FCC-hh)

Increase radius = size (FCC-hh)



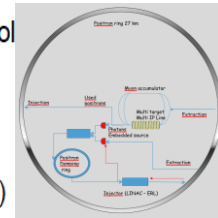
Lepton (e-,e+) circular collider

$$p \propto E_0 \cdot \sqrt[4]{\rho \cdot U_0}$$

Increase supplied RF vol
(FCC-ee)

Increase mass of acc. particle (muon)

Increase radius = size (FCC-ee)



Lepton (e-,e+) linear collider

$$p = L \cdot G_{acc}$$

Increase length (ILC, CLIC)

Compact , Cost
Effective and
Improved
Sustainability

High Gradient Options

Metallic accelerating structures =>

$$100 \text{ MV/m} < E_{\text{acc}} < 1 \text{ GV/m}$$

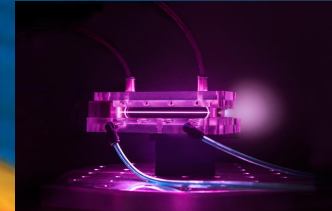
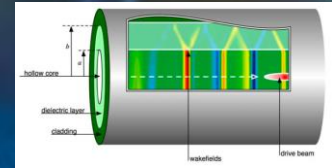
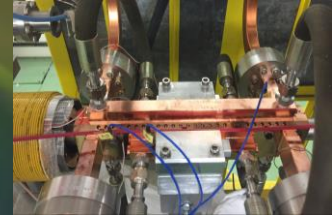
Dielectric structures, laser or particle driven =>

$$E_{\text{acc}} < 10 \text{ GV/m}$$

Plasma accelerator, laser or particle driven =>

$$E_{\text{acc}} < 100 \text{ GV/m}$$

Related Issues: Power Sources and Efficiency, Stability, Reliability, Staging, Synchronization, Rep. Rate and short (fs) bunches with small (μm) spot to match high gradients



Beam Quality Requirements

Future accelerators will require also high quality beams :

==> High Luminosity & High Brightness,

==> High Energy & Low Energy Spread



$$L = \frac{N_{e^+} N_{e^-} f_r}{4 \rho S_x S_y}$$



$$B_n \gg \frac{2I}{e_n^2}$$



-N of particles per pulse
=> 10^9

-High rep. rate f_r =>
bunch trains

-Small spot size => low
emittance

-Short pulse (ps => fs)

-Little spread in
transverse momentum and
angle => low emittance

Laser Electron Accelerator

T. Tajima and J. M. Dawson

Department of Physics, University of California, Los Angeles, California 90024

(Received 9 March 1979)

An intense electromagnetic pulse can create a weak of plasma oscillations through the action of the nonlinear ponderomotive force. Electrons trapped in the wake can be accelerated to high energy. Existing glass lasers of power density $10^{18}\text{W}/\text{cm}^2$ shone on plasmas of densities 10^{18}cm^{-3} can yield gigaelectronvolts of electron energy per centimeter of acceleration distance. This acceleration mechanism is demonstrated through computer simulation. Applications to accelerators and pulsers are examined.

Acceleration of Electrons by the Interaction of a Bunched Electron Beam with a Plasma

Pisin Chen^(a)

Stanford Linear Accelerator Center, Stanford University, Stanford, California 94305

and

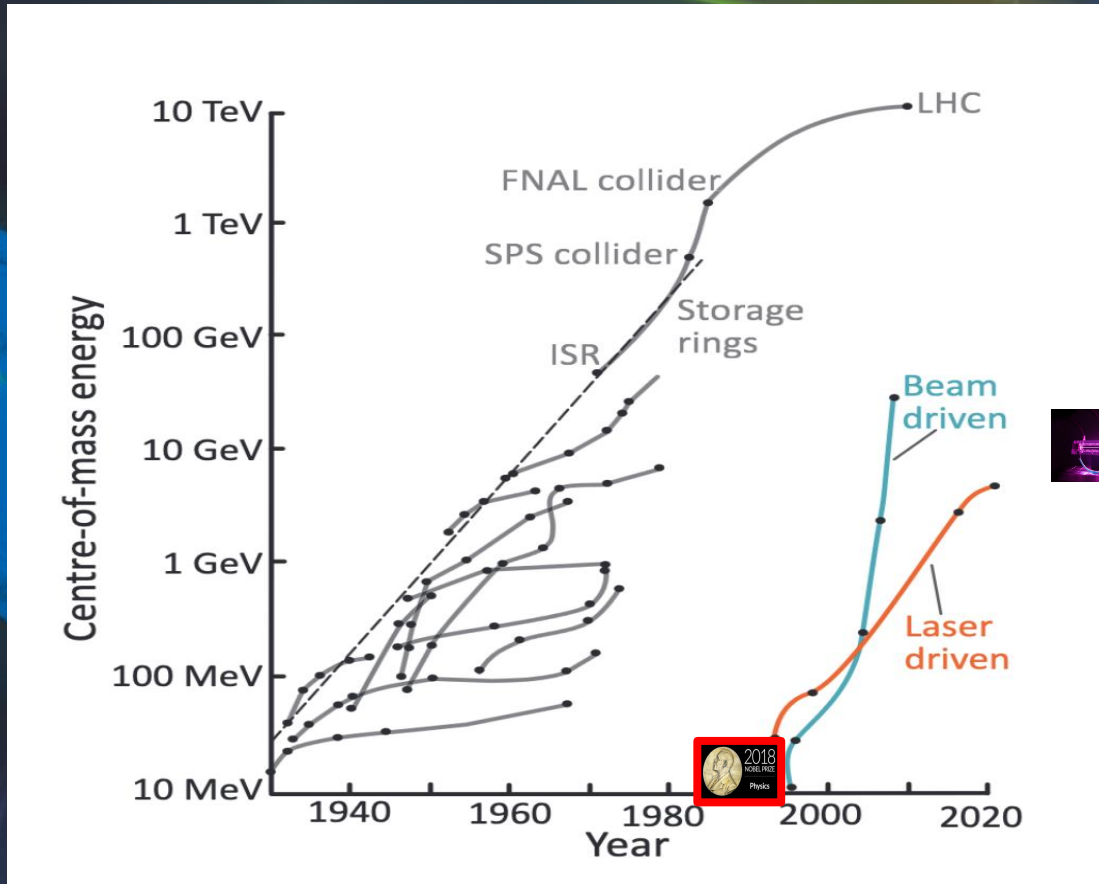
J. M. Dawson, Robert W. Huff, and T. Katsouleas

Department of Physics, University of California, Los Angeles, California 90024

(Received 20 December 1984)

A new scheme for accelerating electrons, employing a bunched relativistic electron beam in a cold plasma, is analyzed. We show that energy gradients can exceed $1\text{ GeV}/\text{m}$ and that the driven electrons can be accelerated from $\gamma_0 mc^2$ to $3\gamma_0 mc^2$ before the driving beam slows down enough to degrade the plasma wave. If the driving electrons are removed before they cause the collapse of the plasma wave, energies up to $4\gamma_0 mc^2$ are possible. A noncollinear injection scheme is suggested in order that the driving electrons can be removed.

Livingstone Diagram with PWFAs



Principles of plasma physics

Definition of Plasma: a quasi-neutral gas of charged particles showing collective behaviour

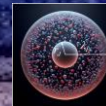
=> a plasma responds to external forces as a single entity

The Debye length is a fundamental property of nearly all plasmas of interest and depends equally on its temperature

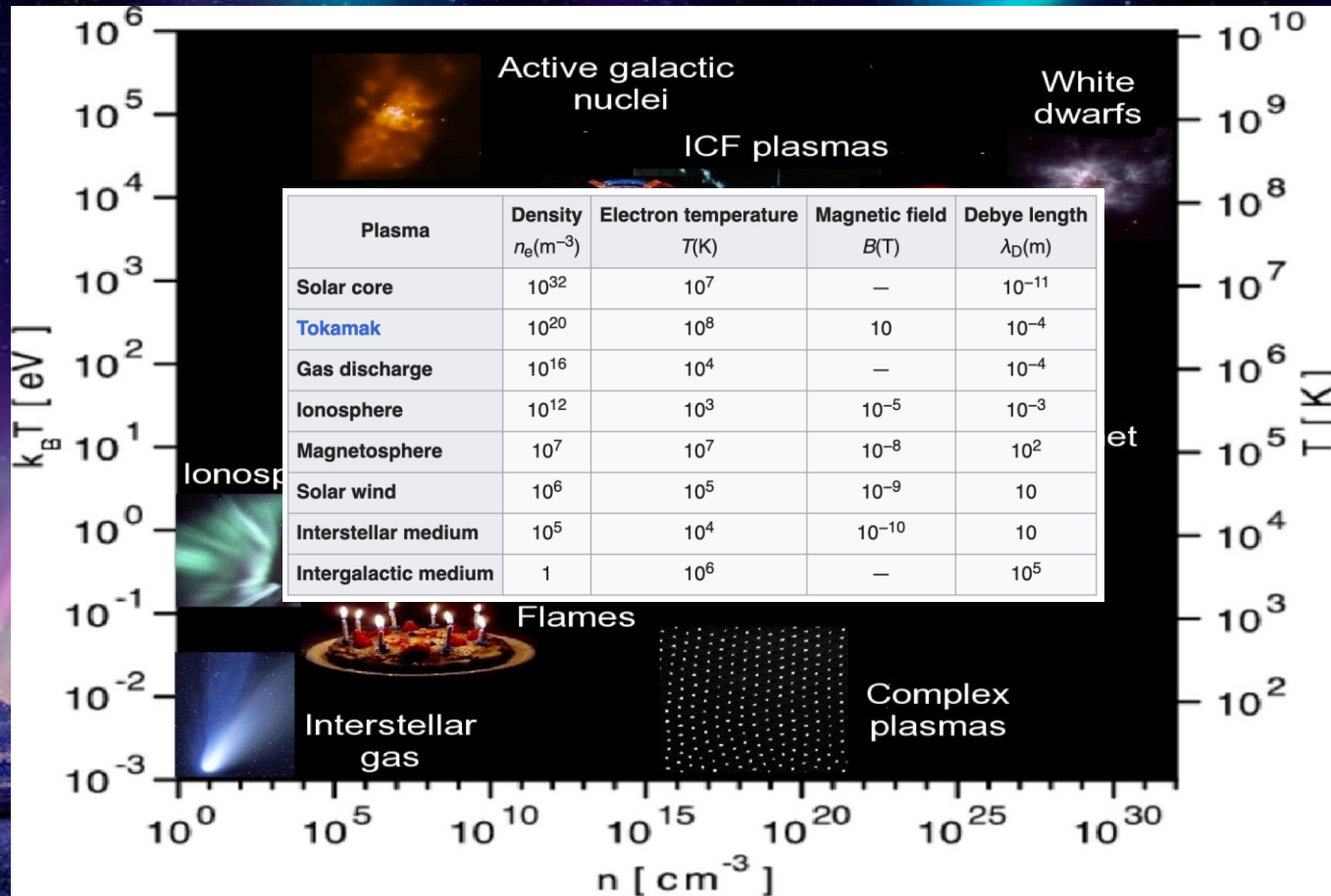
$$\lambda_D = \left(\frac{\epsilon_0 k_B T_e}{e^2 n_e} \right)^{1/2} = 743 \left(\frac{T_e}{\text{eV}} \right)^{1/2} \left(\frac{n_e}{\text{cm}^{-3}} \right)^{-1/2} \text{ cm.}$$

=> An ideal plasma has many particles per Debye sphere, a prerequisite for the collective behaviour:

$$N_D \equiv n_e \frac{4\pi}{3} \lambda_D^3 \gg 1$$

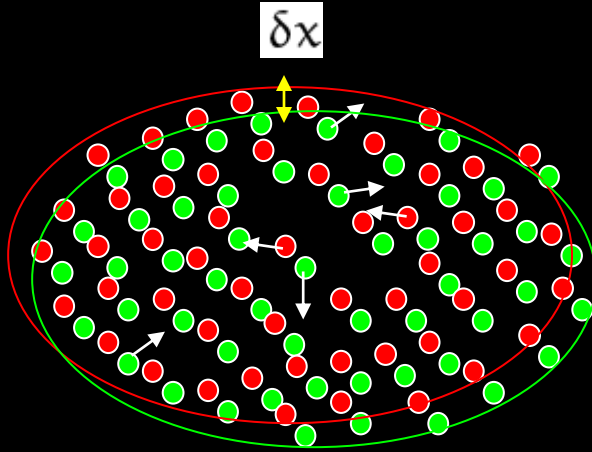


Plasma Temperature and Density



Surface charge density

$$\sigma = e n \delta x$$



Surface electric field

$$E_x = -\sigma/\epsilon_0 = -e n \delta x/\epsilon_0$$

Restoring force

$$m \frac{d^2 \delta x}{dt^2} = e E_x = -m \omega_p^2 \delta x$$

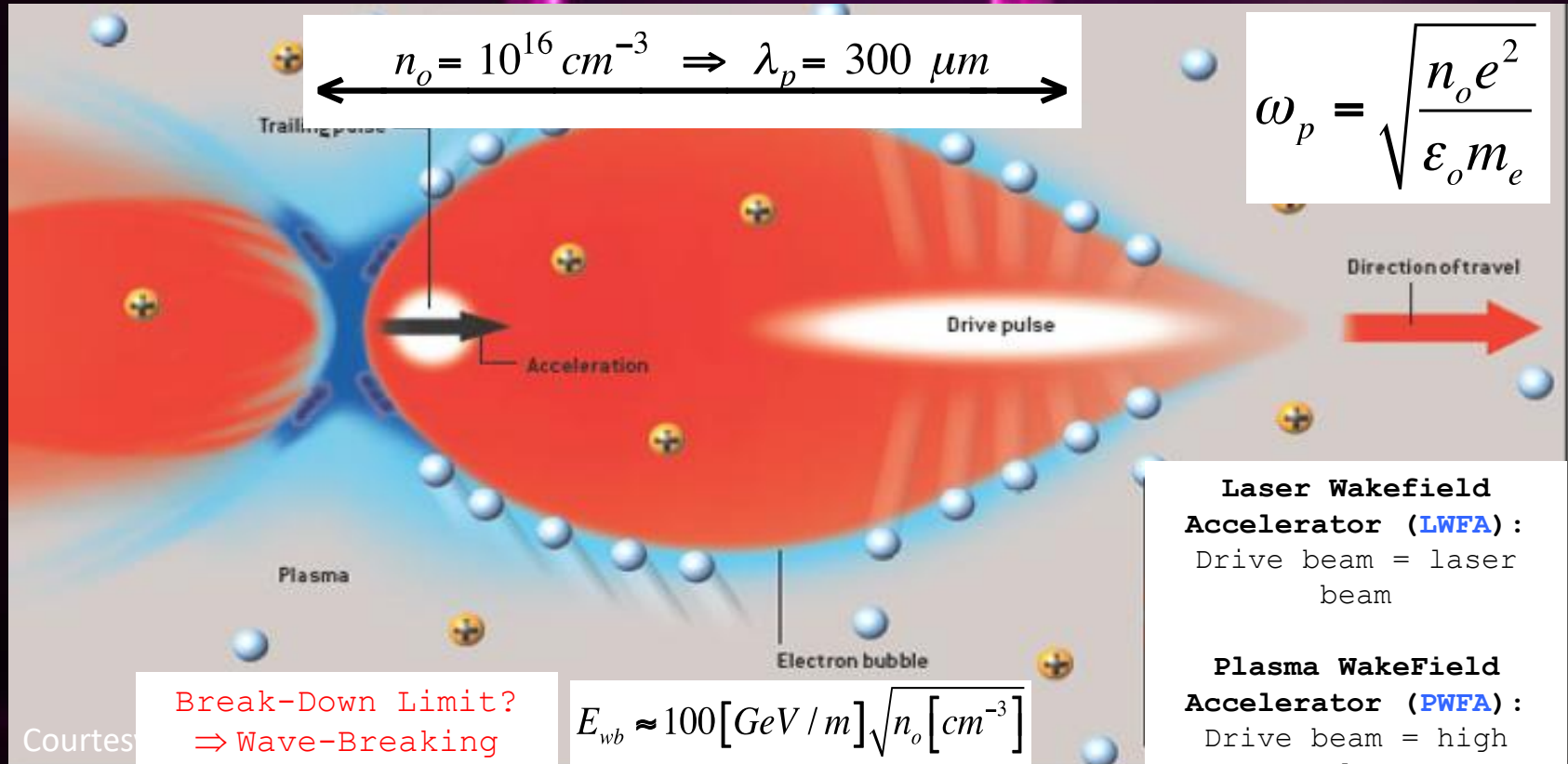
Plasma frequency

$$\omega_p^2 = \frac{n e^2}{\epsilon_0 m}$$

Plasma oscillations

$$\delta x = (\delta x)_0 \cos(\omega_p t)$$

Principle of plasma acceleration



$$n_o = 10^{16} \text{ cm}^{-3} \Rightarrow \lambda_p = 300 \mu\text{m}$$

$$\omega_p = \sqrt{\frac{n_o e^2}{\epsilon_o m_e}}$$

Laser Wakefield Accelerator (LWFA):

Drive beam = laser beam

Plasma WakeField Accelerator (PWFA):

Drive beam = high energy electron or proton beam

Break-Down Limit?
 ⇒ Wave-Breaking field:

$$E_{wb} \approx 100 [\text{GeV} / \text{m}] \sqrt{n_o [\text{cm}^{-3}]}$$

Courtesy

Principle of plasma acceleration

From Maxwell's equations, the electric field in a (positively) charged sphere with uniform density n_i at location r is

$$\vec{E}(r) = \frac{q_i n_i}{3 \epsilon_0} r$$

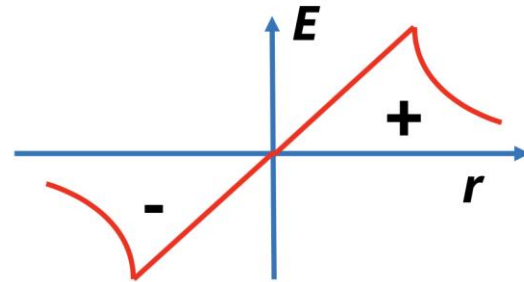
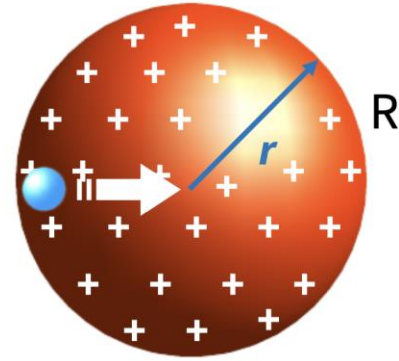
The field is **increasing** inside the sphere

Let's put some numbers

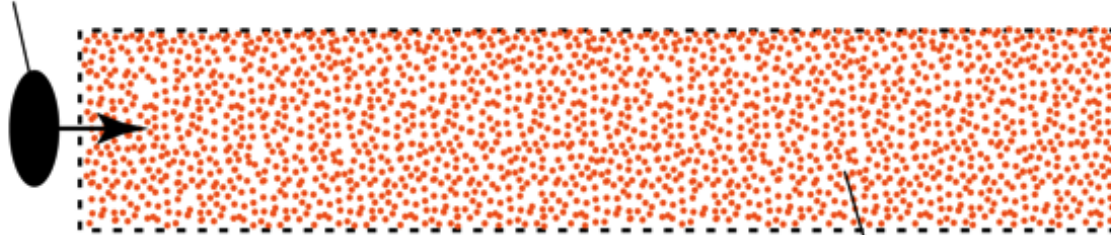
$$n_i = 10^{16} \text{ cm}^{-3}$$

$$R = 0.5$$

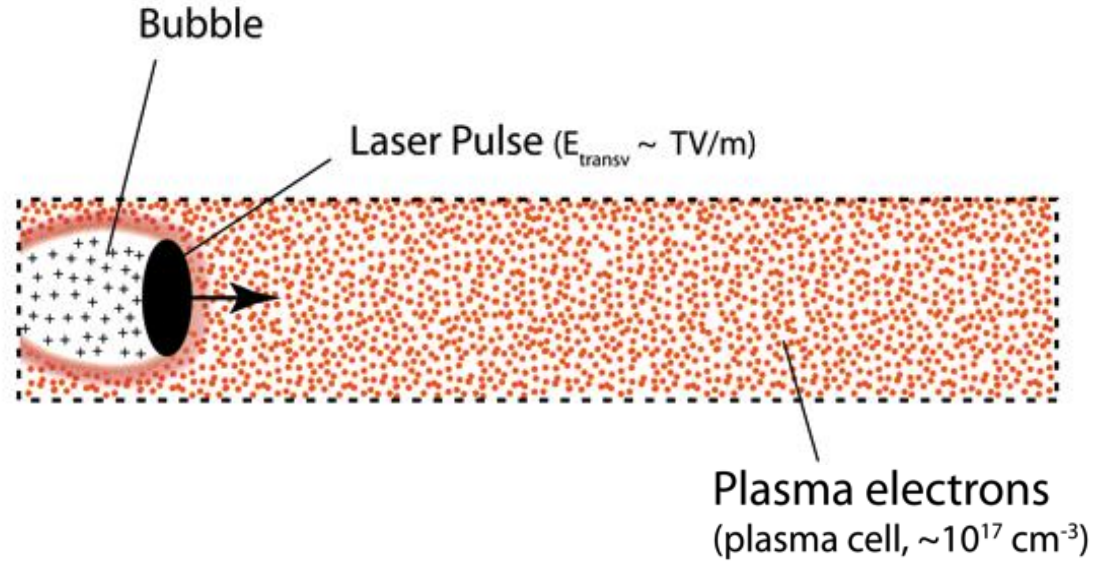
$$\Rightarrow E \approx 10 \frac{\text{GV}}{\text{m}}$$

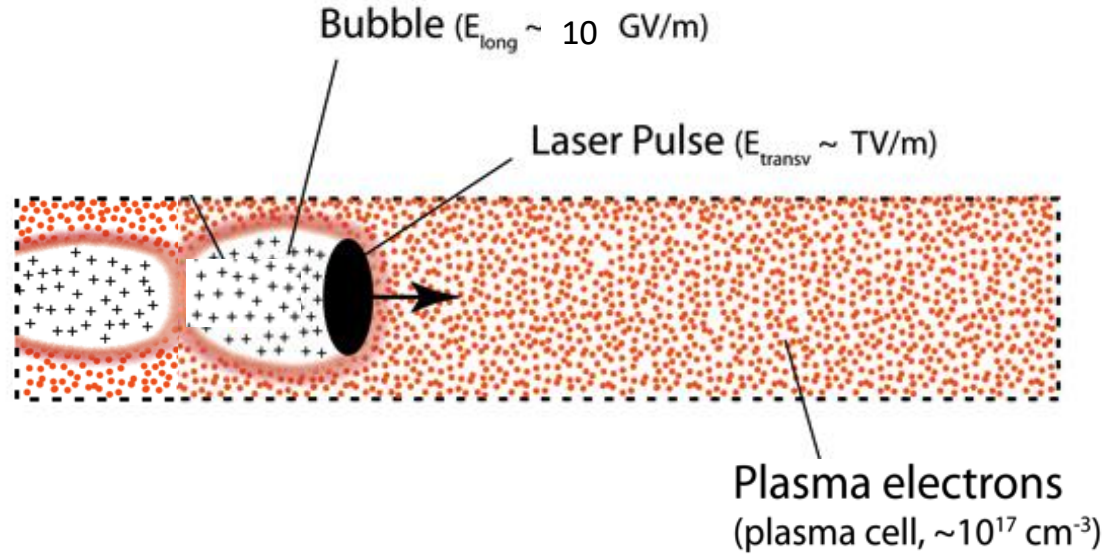


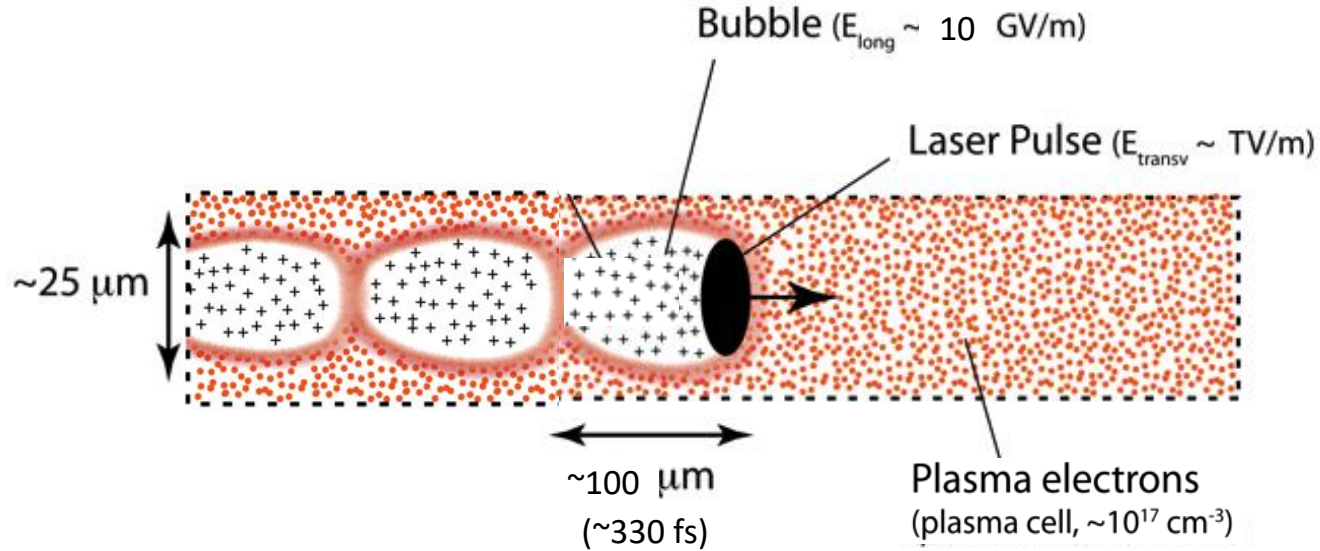
Laser Pulse (200 TW, ~ 30 fs, $E_{\text{transv}} \sim \text{TV/m}$)



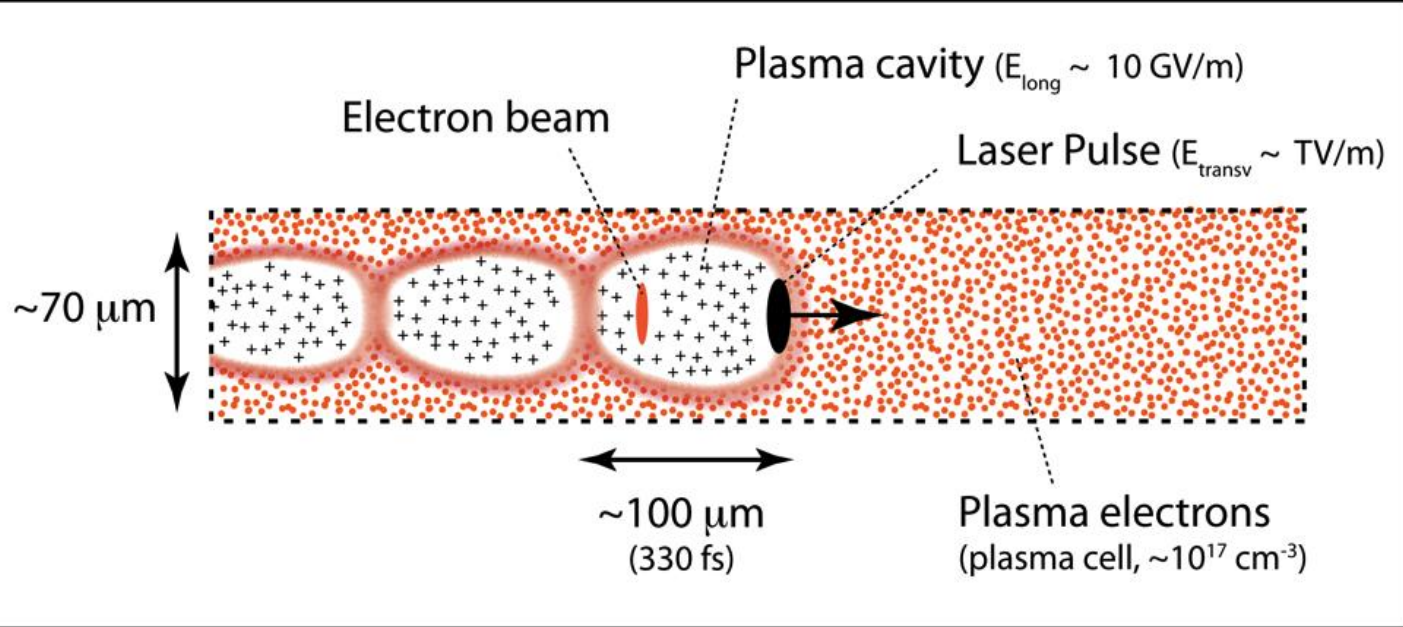
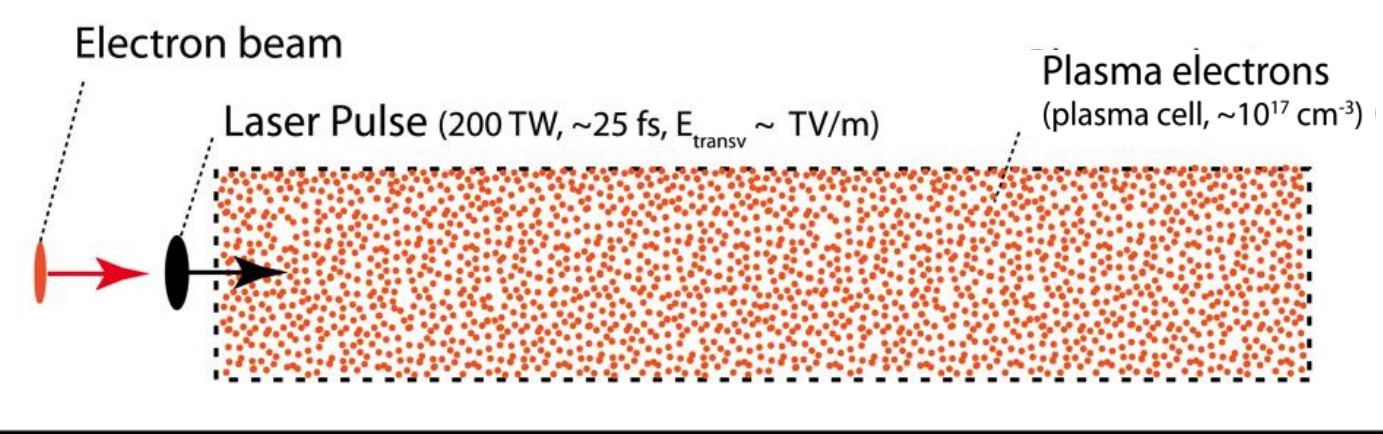
Plasma electrons
(plasma cell, $\sim 10^{17} \text{ cm}^{-3}$)



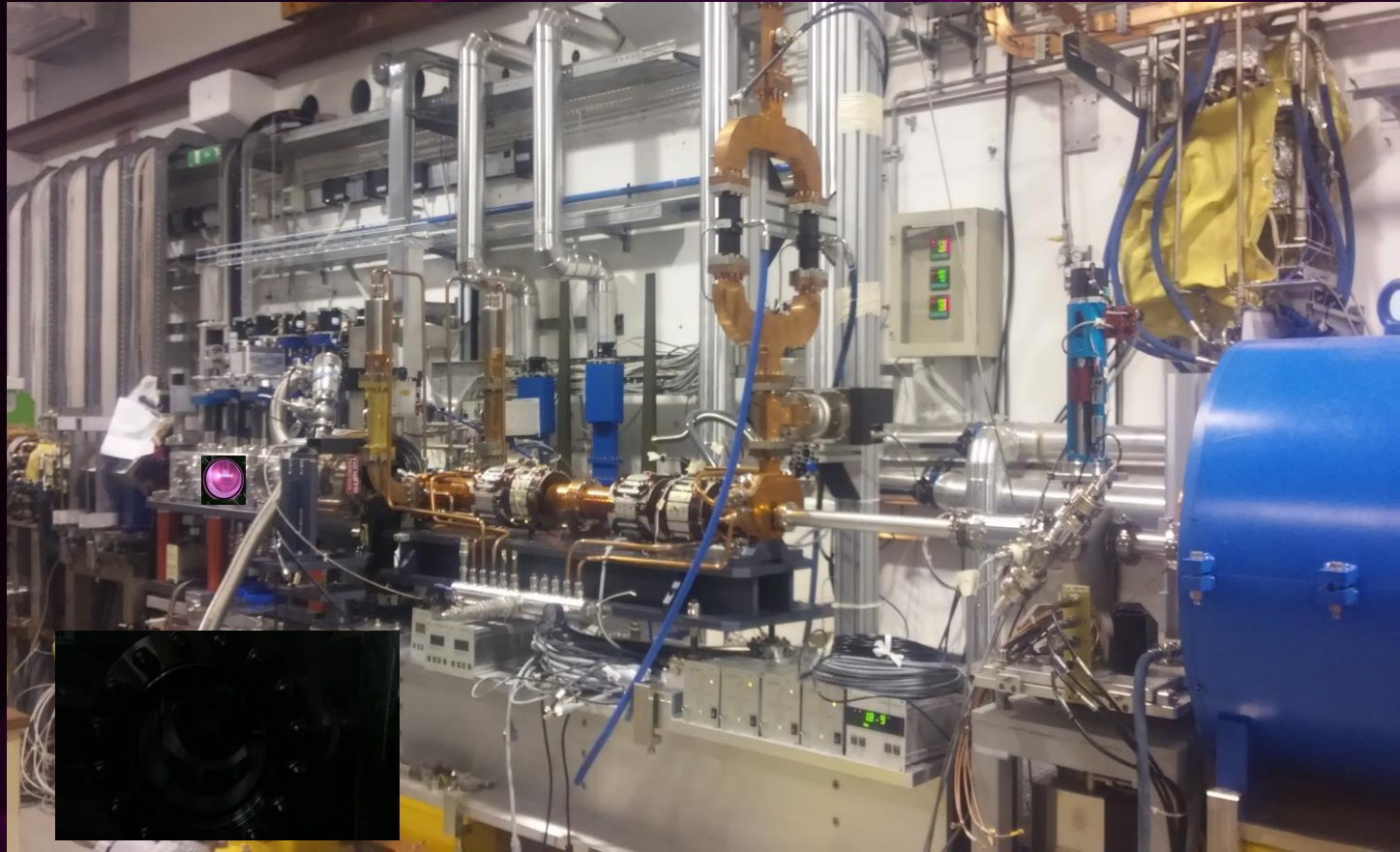


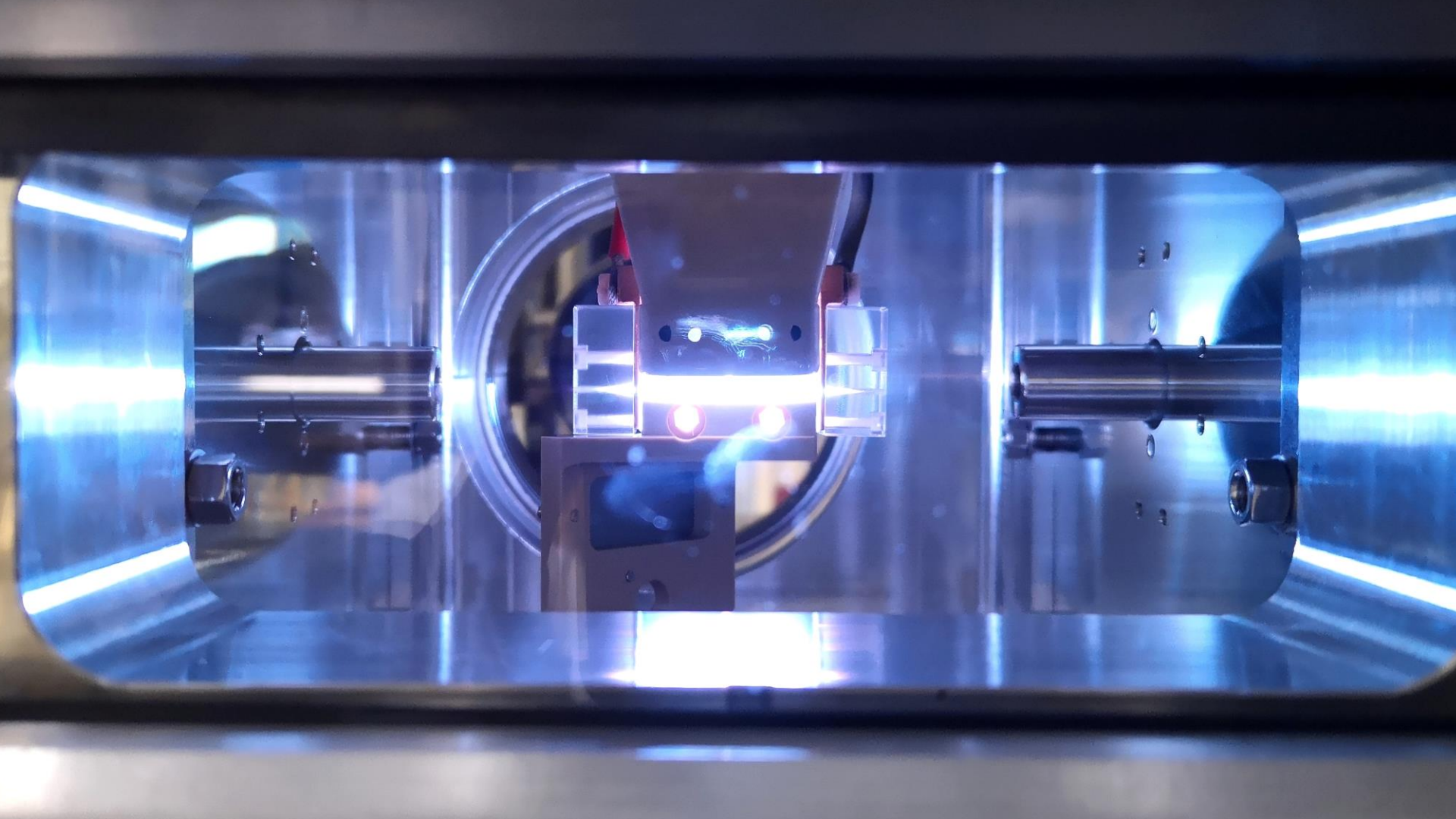


This accelerator fits into a human hair!



PWFA beam line at SPARC_LAB





Looking for a plasma target

He



Ne



Ar



Kr

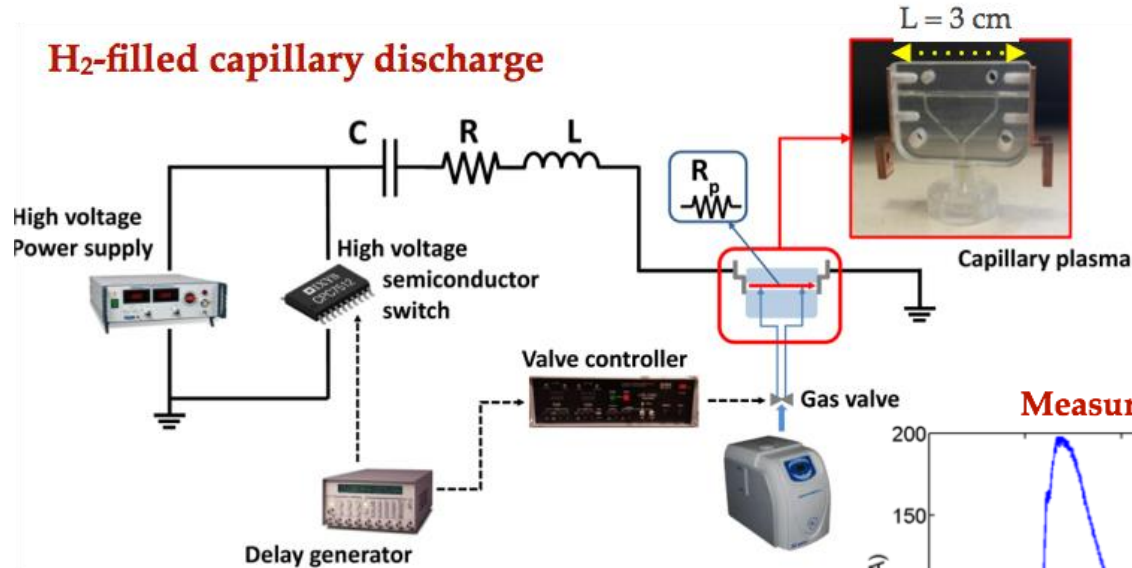


Xe

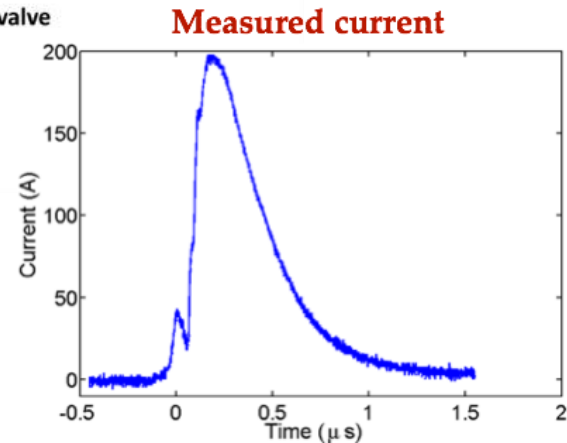


Plasma Source

H₂-filled capillary discharge

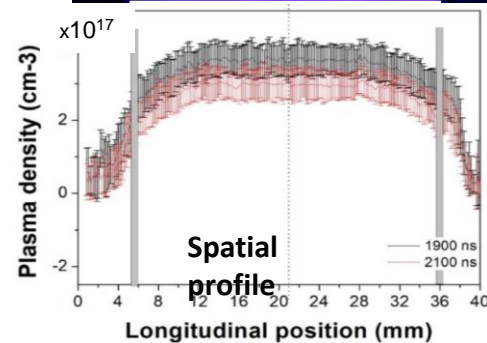
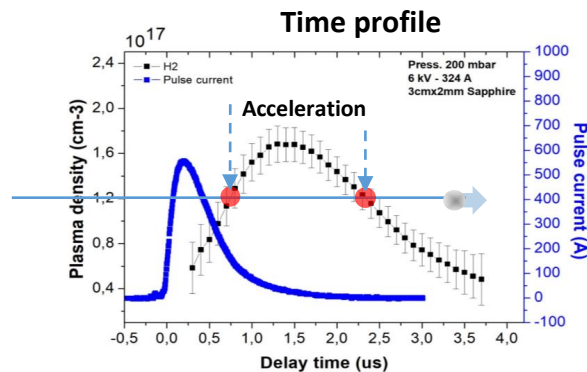
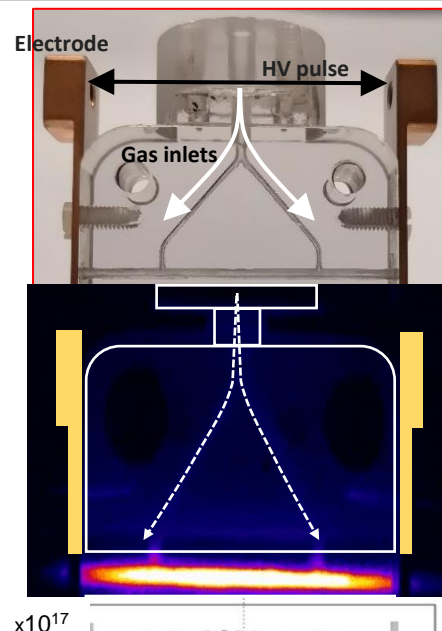
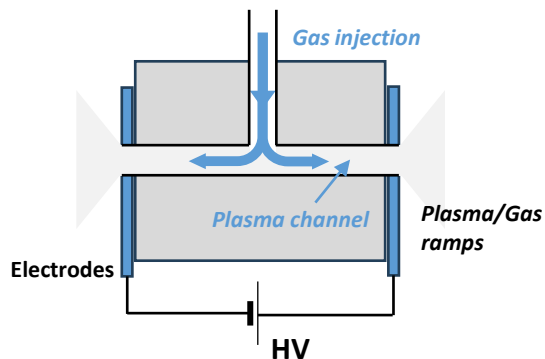


$P_{H_2} = 10$ mbar
 Total discharge duration: 800 ns
 Voltage: 20 kV
 Peak current: 200 A
 Capacitor: 6 nF

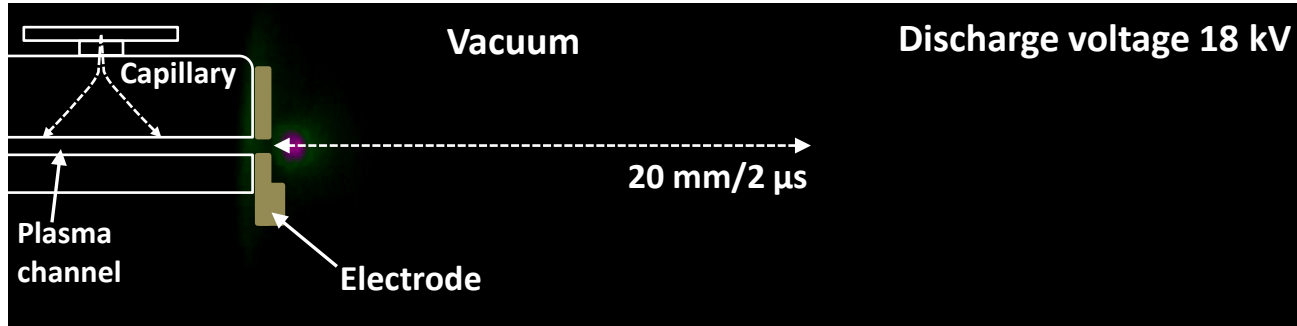


Discharge Plasma sources properties

- Region of uniform neutral gas contained by differential pumping at two input/output extremities and realized by a thin capillary
- Modulation of longitudinal gas distribution by special geometric shapes
- Small gas flow to operate in steady state regime with continuous flow ($P < 1$ bar)
- There is no movable parts
- Small dimension of the HV source to produce plasma
- Plume of gas from front and back of cell, plasma ramps to be controlled
- Density easily adjusted by controlling gas flow, geometry or voltage in the range: $10^{15} - 10^{19}$ cm^{-3}
- Erosion of the plasma channel due to the electrical discharges to form plasma



- 20 images separated by 100 ns, so 2 μ s of total observation time of the plasma plumes
- The ICCD camera area is 1024 x 256 pixel

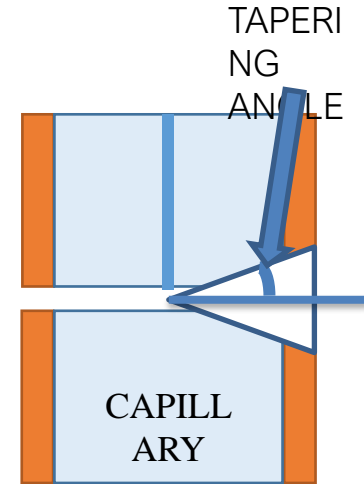
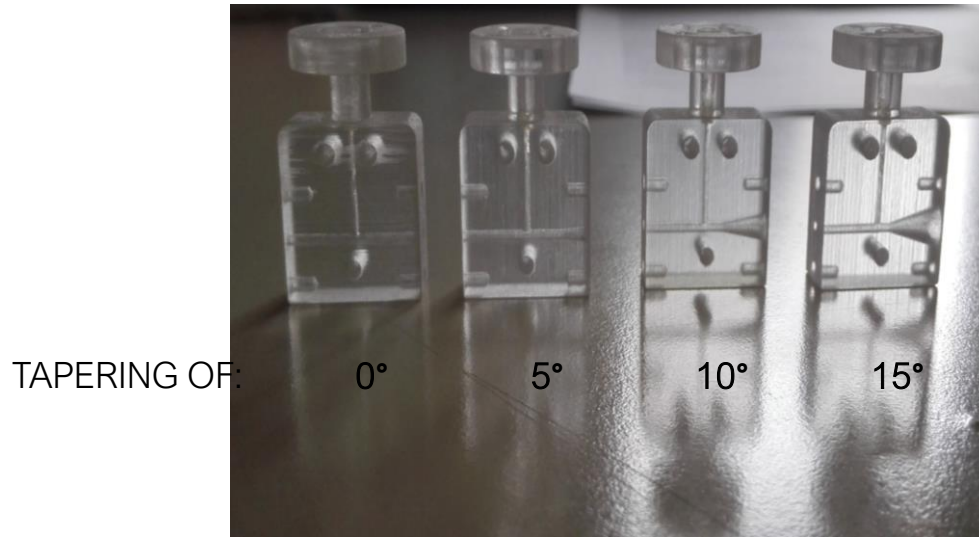


- Both plasma plumes can reach a total expansion length around 40 mm (20 mm each one) that is comparable with the channel length of 30 mm, so they can strongly affect the beam properties that pass through the capillary
- Temperature, pressure and plasma density, inside and outside the gas-filled capillary plasma source, represent essential parameters that have to be investigated to understand the plasma evolution and how it can affect the electron beam.

Tapered capillaries

Local control of the plasma density is required to match the laser/electron beam into the plasma.

Tapering the capillary diameter is the easiest way to change locally the density.



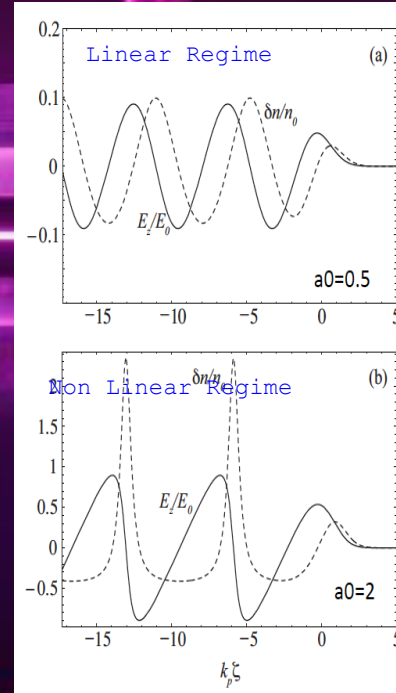
Principle of plasma acceleration

Driven by Radiation Pressure

$$\left(\frac{\partial^2}{\partial t^2} + \omega_p^2\right) \frac{n}{n_o} = c^2 \nabla^2 \frac{a^2}{2}$$
$$a = \frac{eA}{mc^2} \propto \lambda J^{1/2}$$

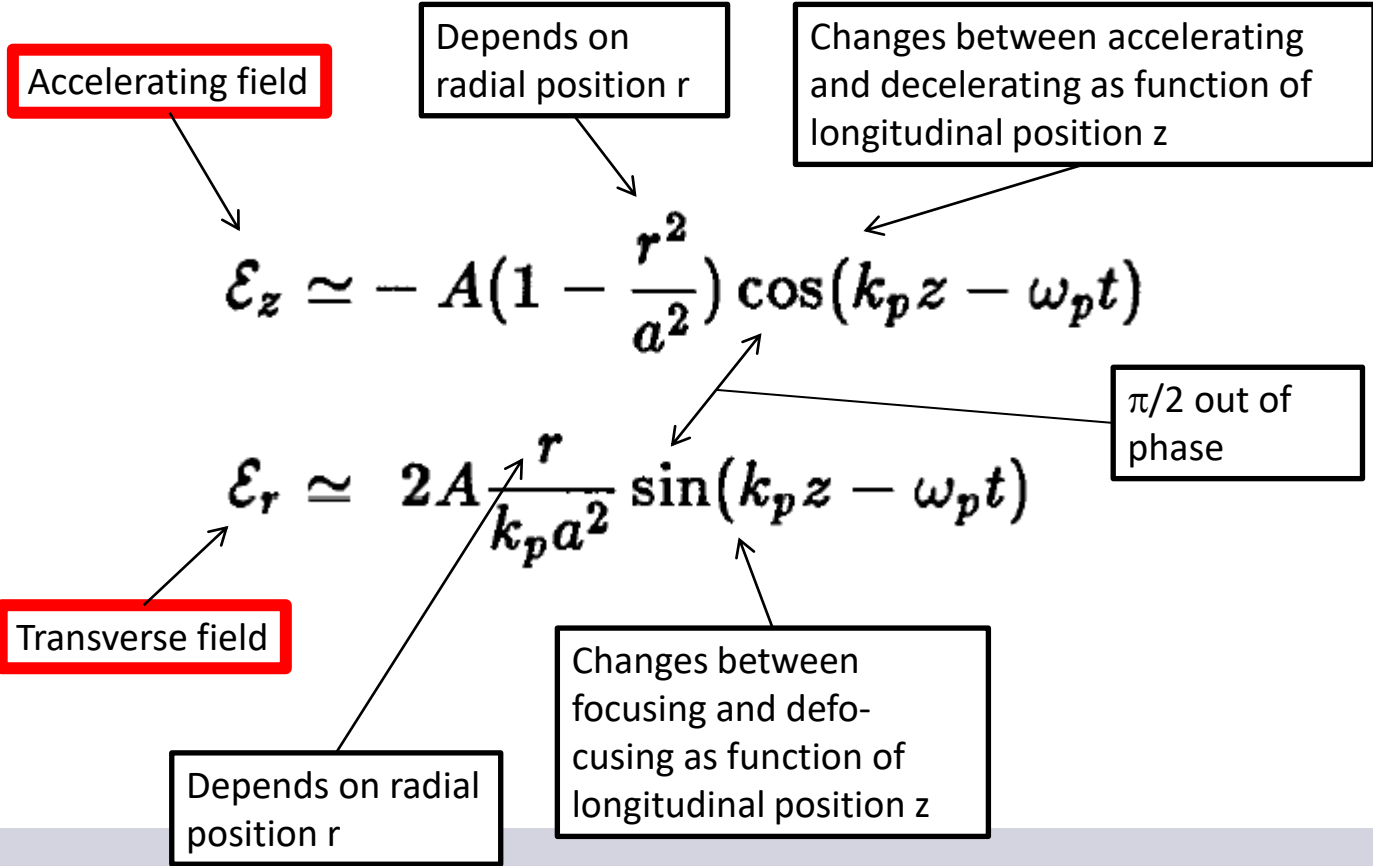
Driven by Space Charge

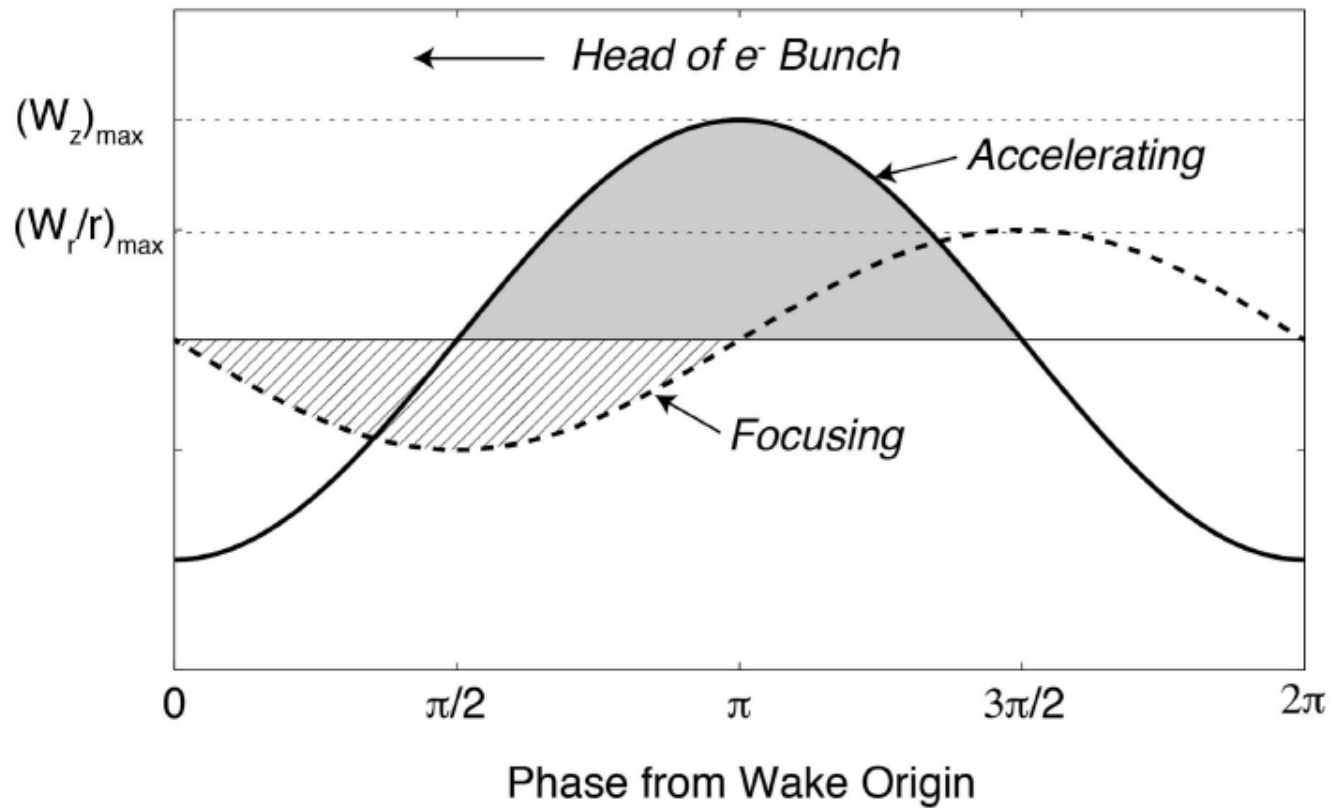
$$\left(\frac{\partial^2}{\partial t^2} + \omega_p^2\right) \frac{n}{n_o} = -\omega_p^2 \frac{n_{beam}}{n_o}$$
$$n_{beam} = \frac{N}{\sqrt{(2\pi)^3 \sigma_r^2 \sigma_z}}$$

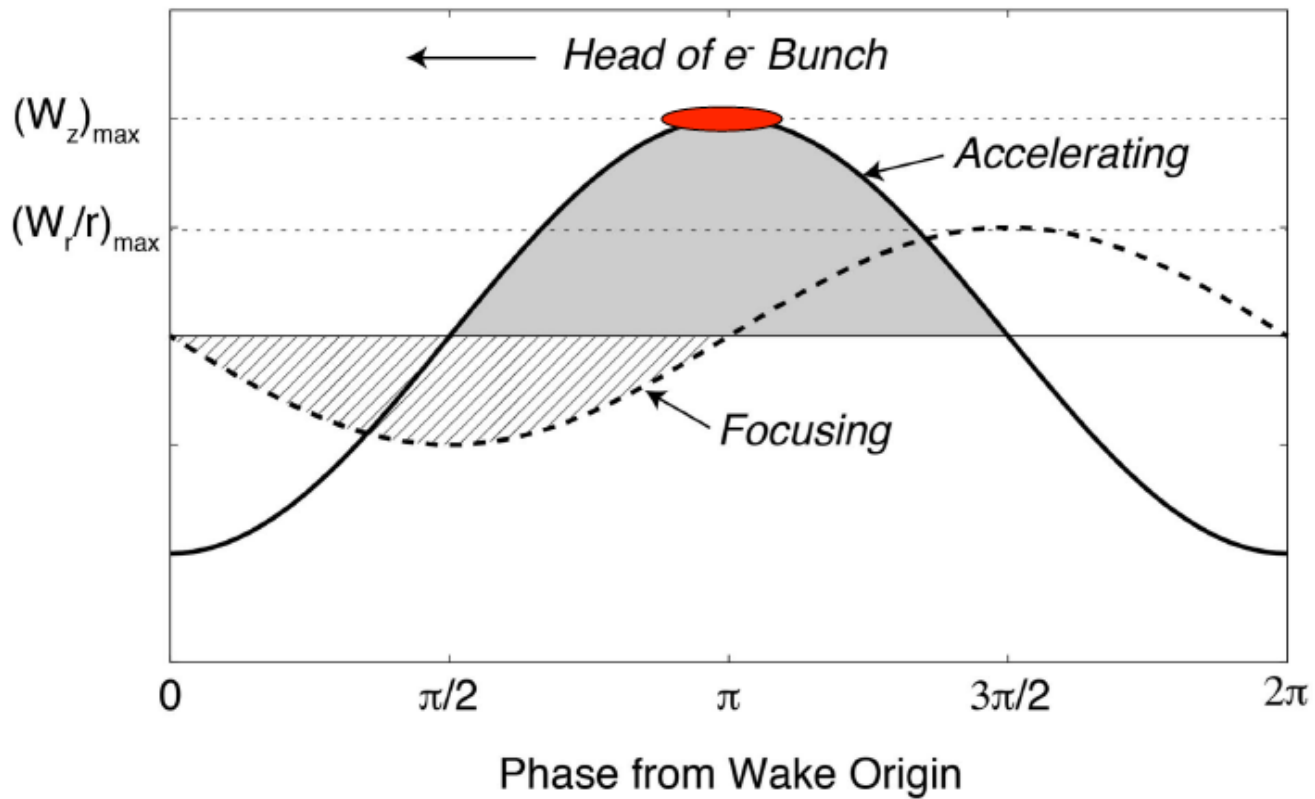


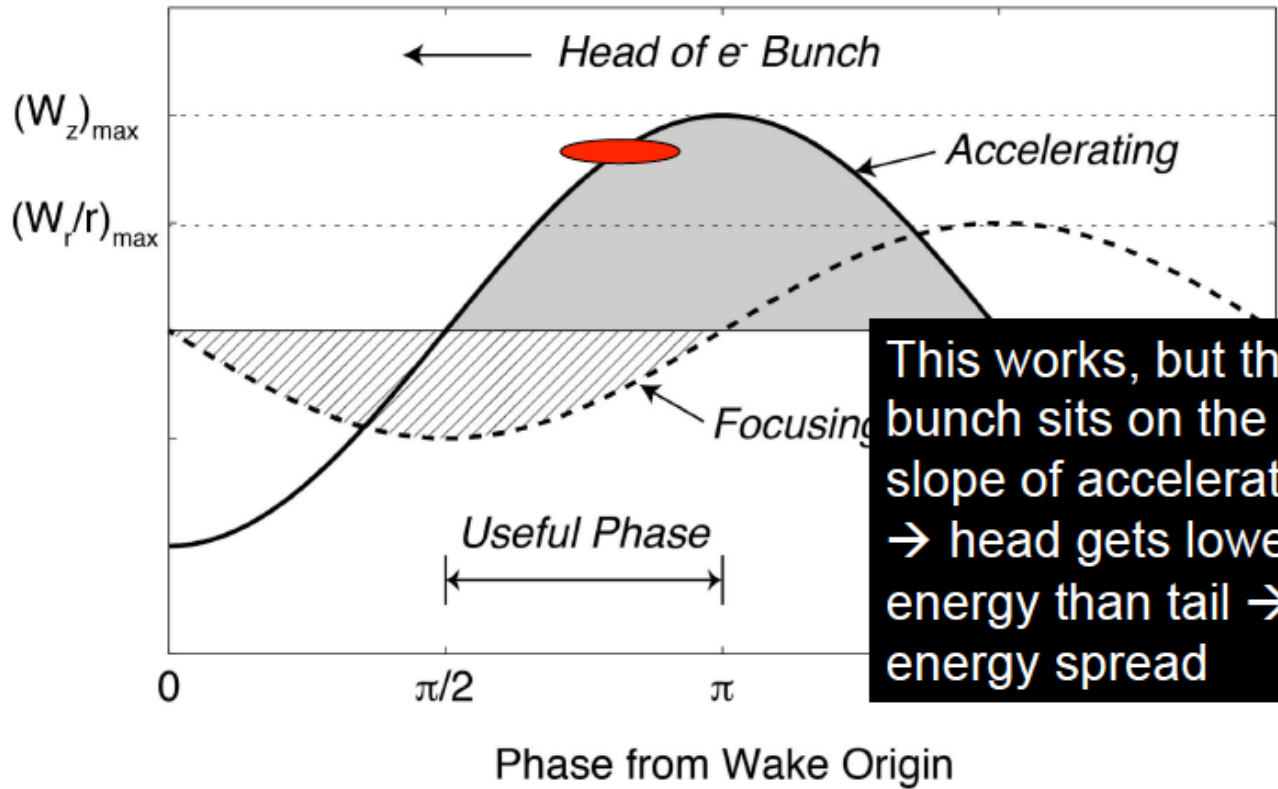
LWFA limitations: Diffraction, Dephasing, Depletion

PWFA limitations: Head Erosion, Hose





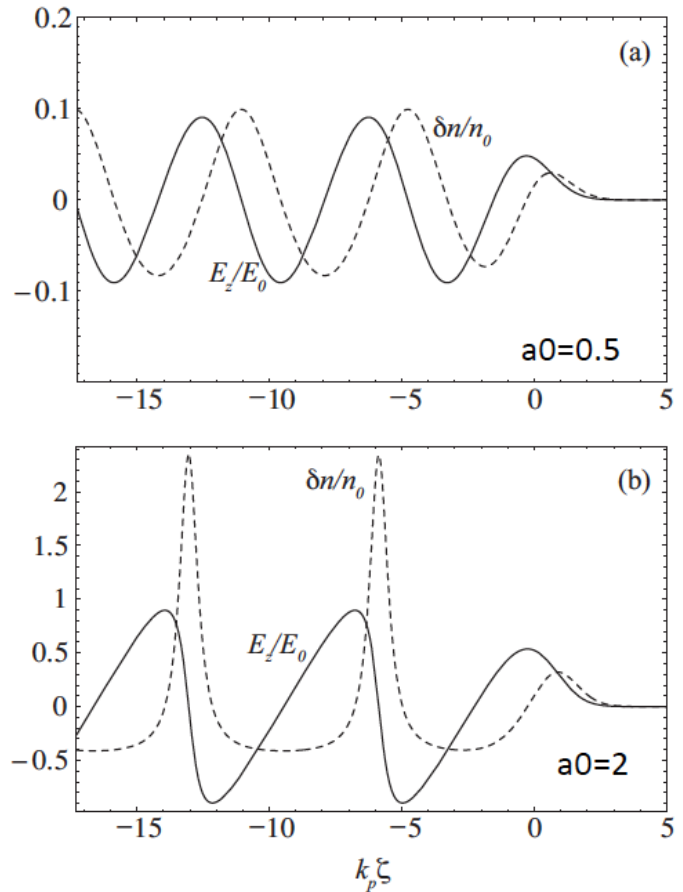




This works, but the bunch sits on the slope of acceleration → head gets lower energy than tail → energy spread



Regimes: Linear & Non-Linear



Linear



FIG. 8. Time-averaged density variation $\delta n/n_0$ (dashed curve) and axial electric field E_z/E_0 (solid curve) in an LWFA driven by a Gaussian laser pulse (pulse is moving to the right, centered at $k_p \zeta=0$ with rms intensity length $L_{\text{rms}}=k_p^{-1}$) for (a) $a_0=0.5$ and (b) $a_0=2.0$.

Non-Linear

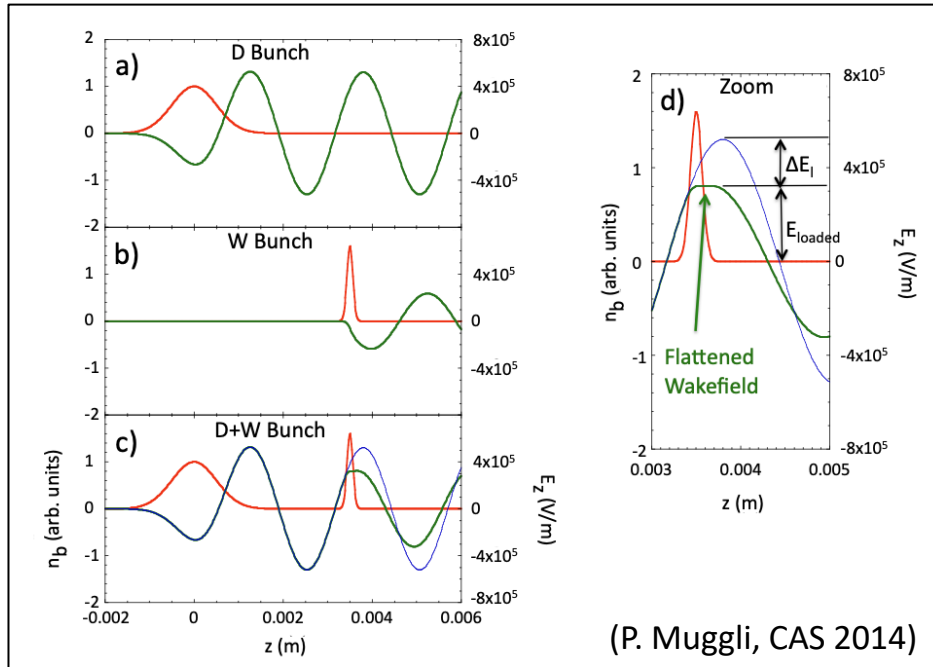


Beam Loading

→ BEAM LOADING:

The presence of the witness bunch affects the wakefields

Linear regime



Non-linear regime

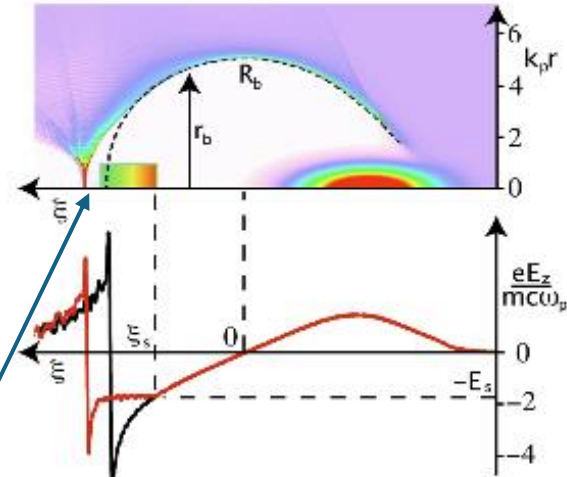
PRL **101**, 145002 (2008)

PHYSICAL REVIEW LETTERS

week ending
3 OCTOBER 2008

Beam Loading in the Nonlinear Regime of Plasma-Based Acceleration

M. Tzoufras,¹ W. Lu,¹ F. S. Tsung,² C. Huang,² W. B. Mori,^{1,2} T. Katsouleas,³ J. Vieira,⁴ R. A. Fonseca,⁴ and L. O. Silva⁴



The bubble closes later

Non Linear Regime – Ellipsoidal Bubble Model

$$\alpha = \frac{n_b}{n_p} \geq 1$$

$$\begin{cases} X = 2\sqrt{\alpha}\sigma_{x,d} \\ Y = 2\sqrt{\alpha}\sigma_{y,d} \\ Z = \frac{\lambda_p}{2} \end{cases}$$

$$E_z(\xi) = An_p\sqrt{I_d}\xi$$

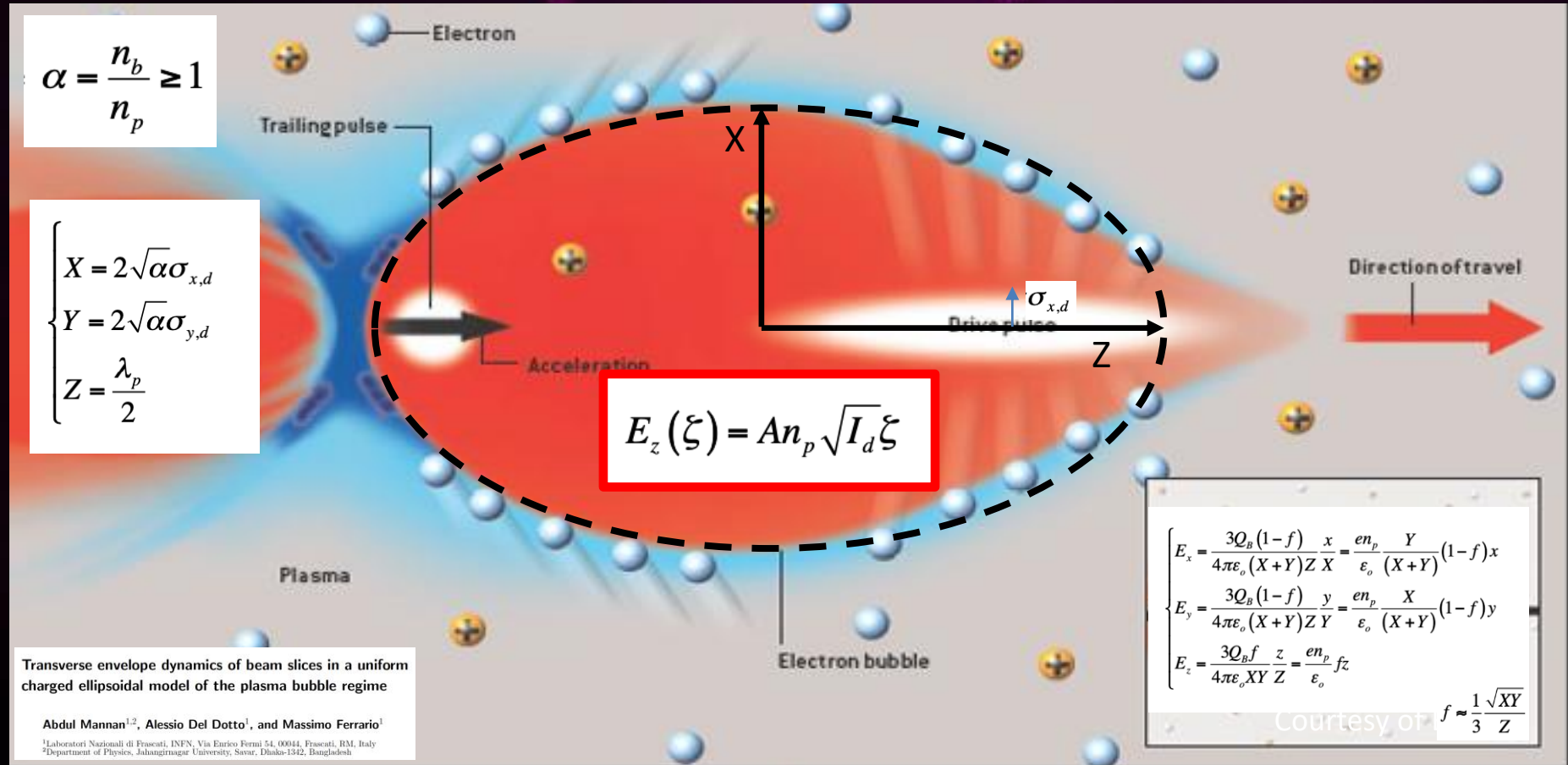
$$\begin{cases} E_x = \frac{3Q_B(1-f)}{4\pi\epsilon_0(X+Y)Z} \frac{x}{X} = \frac{en_p}{\epsilon_0} \frac{Y}{(X+Y)}(1-f) \frac{x}{X} \\ E_y = \frac{3Q_B(1-f)}{4\pi\epsilon_0(X+Y)Z} \frac{y}{Y} = \frac{en_p}{\epsilon_0} \frac{X}{(X+Y)}(1-f) \frac{y}{Y} \\ E_z = \frac{3Q_B f}{4\pi\epsilon_0 XY Z} \frac{z}{Z} = \frac{en_p f}{\epsilon_0} \frac{z}{Z} \end{cases}$$

$$f \approx \frac{1}{3} \frac{\sqrt{XY}}{Z}$$

Transverse envelope dynamics of beam slices in a uniform charged ellipsoidal model of the plasma bubble regime

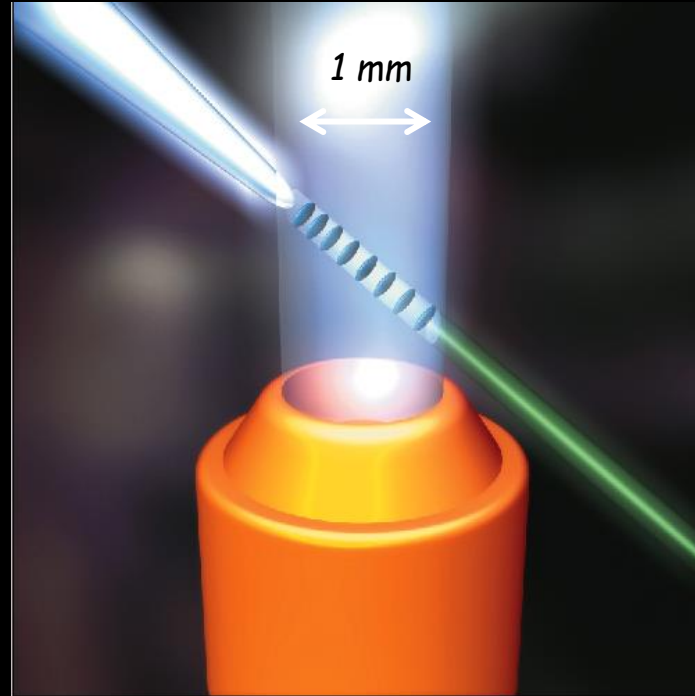
Abdul Mannan^{1,2}, Alessio Del Dotto¹, and Massimo Ferrario¹

¹Laboratori Nazionali di Frascati, INFN, Via Enrico Fermi 54, 00044, Frascati, RM, Italy
²Department of Physics, Jahangirnagar University, Savar, Dhaka-1342, Bangladesh



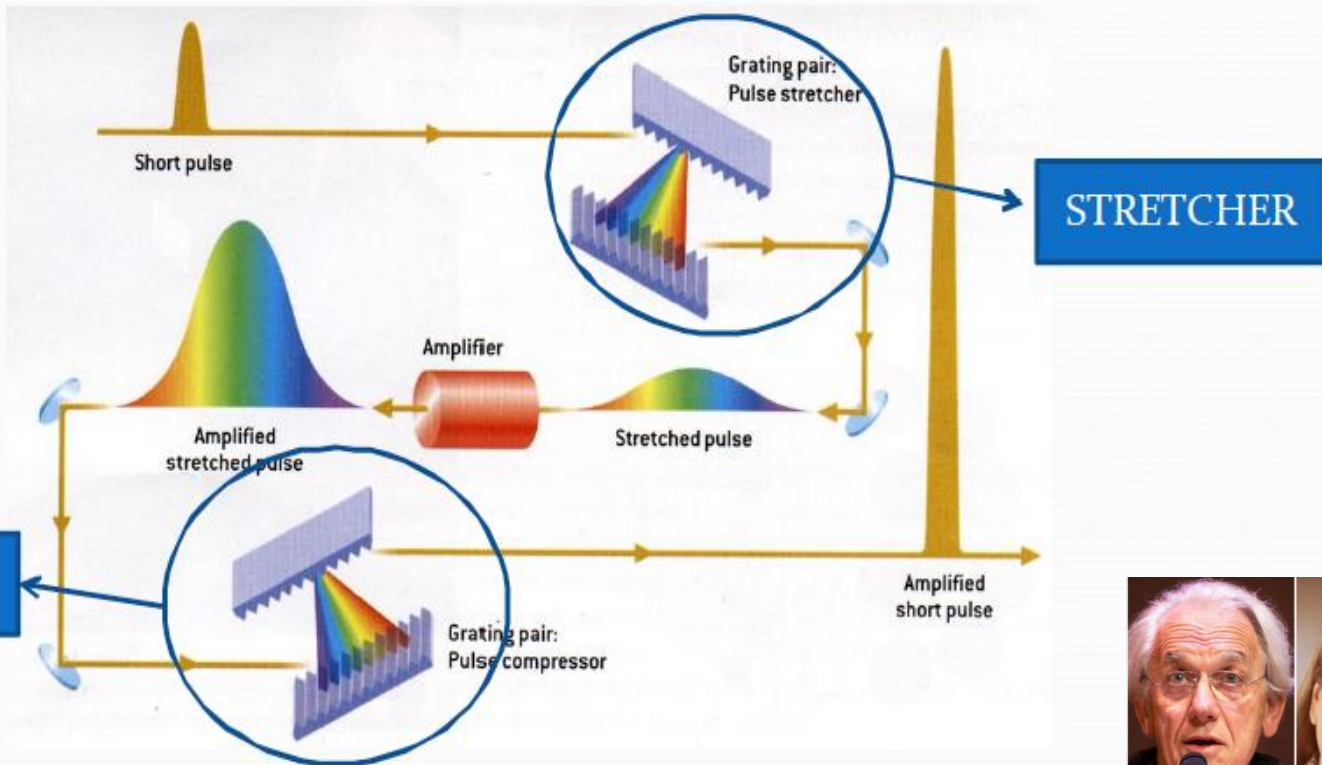
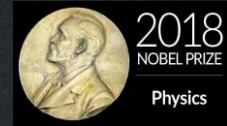
Laser Driven LWFA

Direct production of e-beam

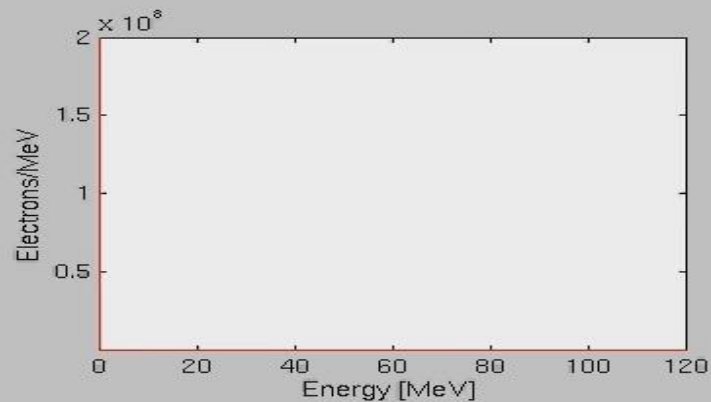
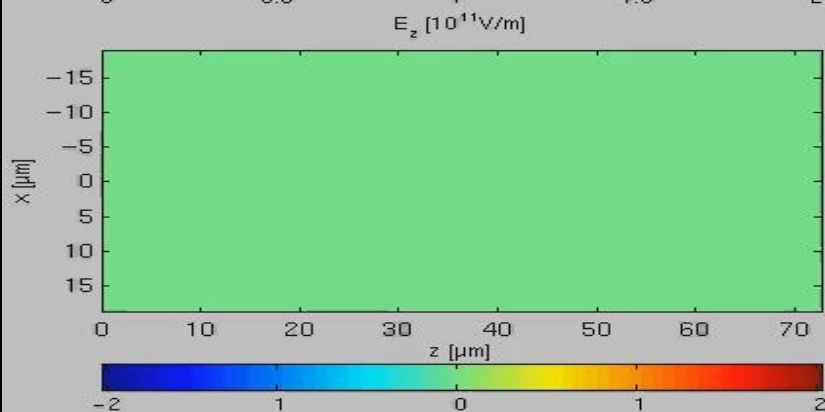
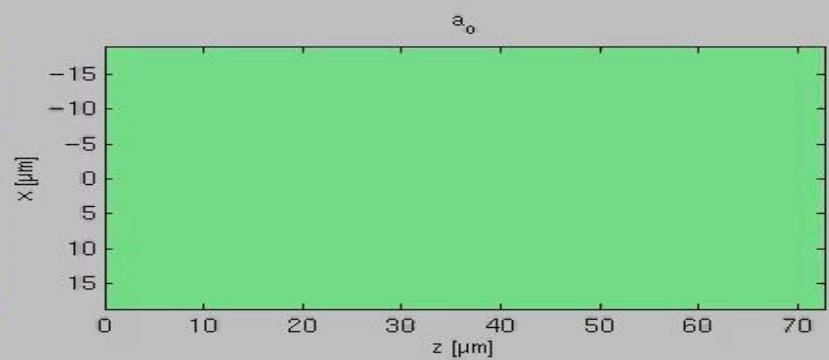
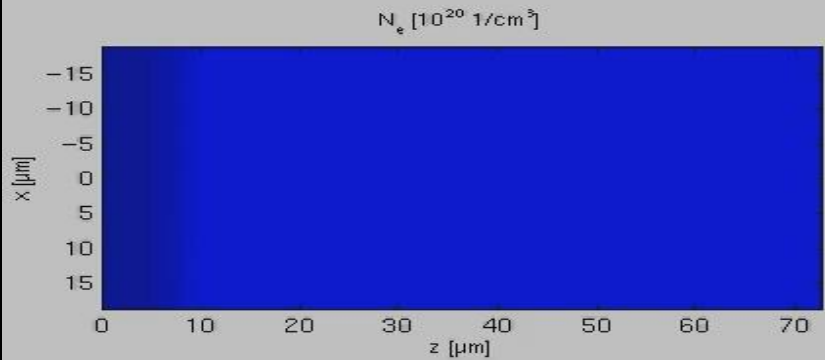


Electron beam

Chirped Pulse



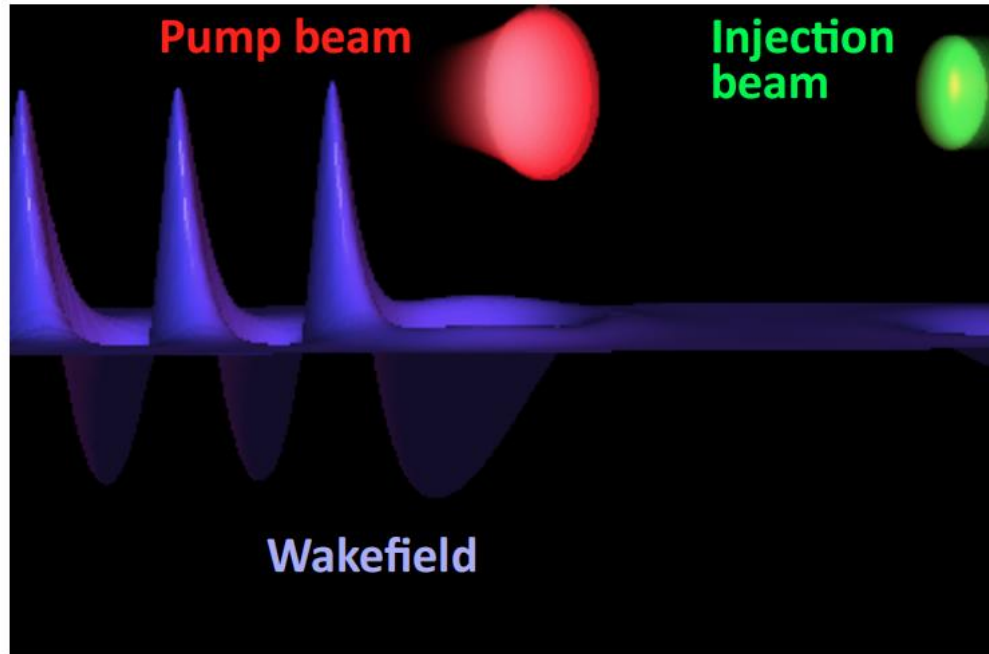
Diffraction - Self injection - Dephasing - Depletion



Colliding Laser Pulses Scheme



The first laser creates the accelerating structure, a second laser beam is used to heat electrons



Theory : E. Esarey *et al.*, PRL **79**, 2682 (1997), H. Kotaki *et al.*, PoP **11** (2004)
Experiments : J. Faure *et al.*, Nature **444**, 737 (2006)



<http://loa.ensta.fr/>

1st European Advanced Accelerator Concepts Workshop, La Biodola, Isola d'Elba - Italy, June 2-7 (2013)

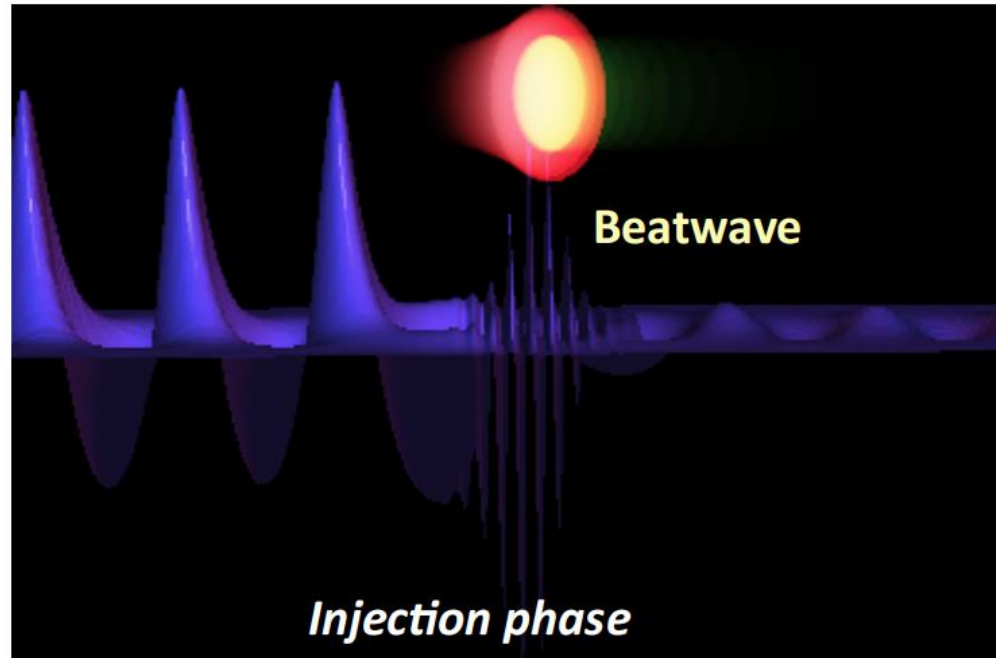


UMR 7639



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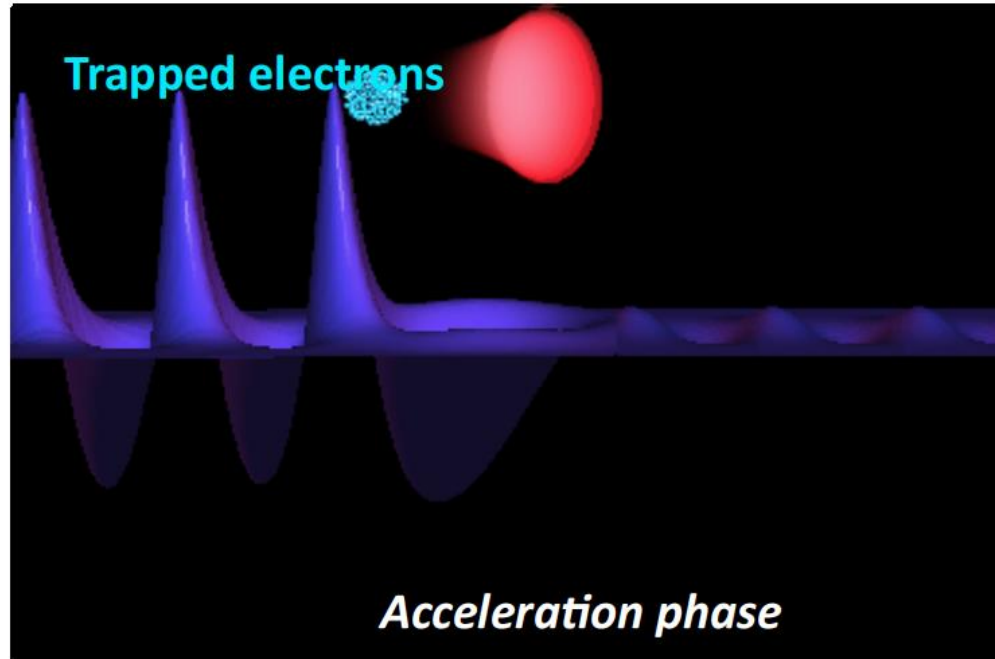
UMR 7639



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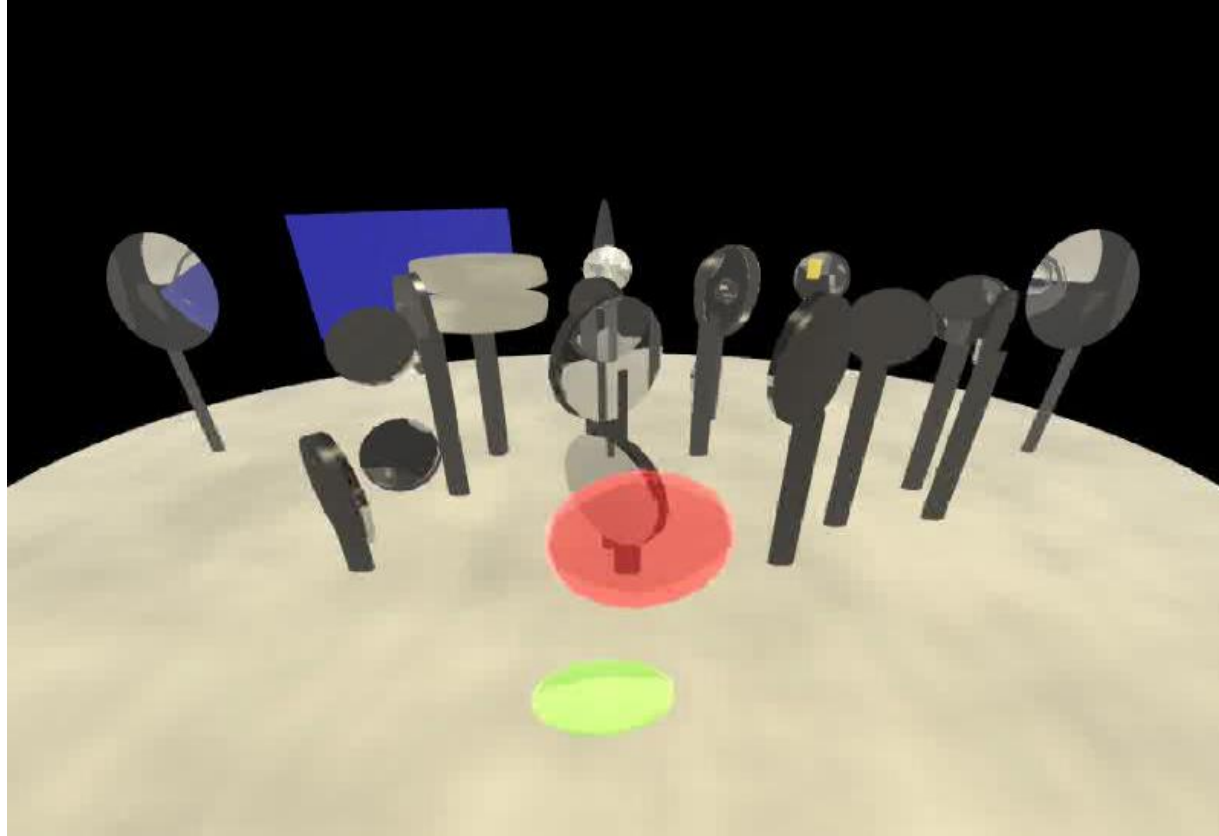
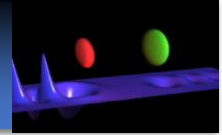


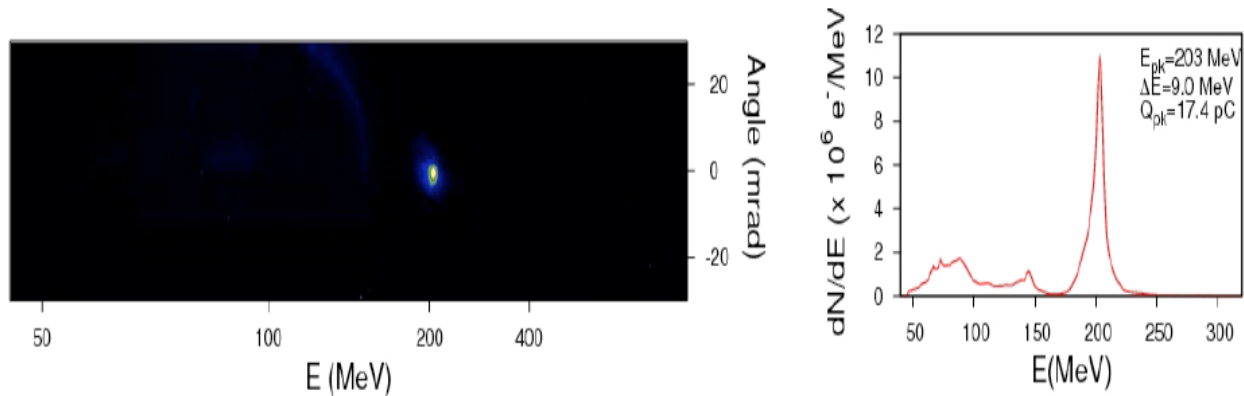
<http://loa.ensta.fr/>

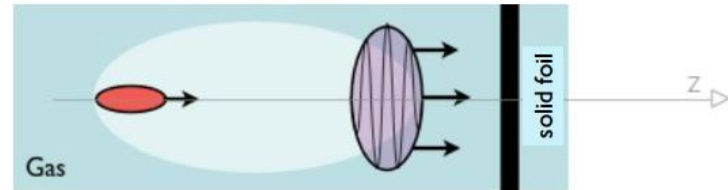
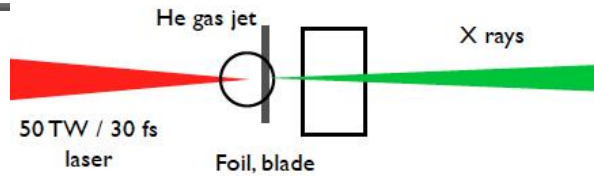
UMR 7639



The colliding of two laser pulses

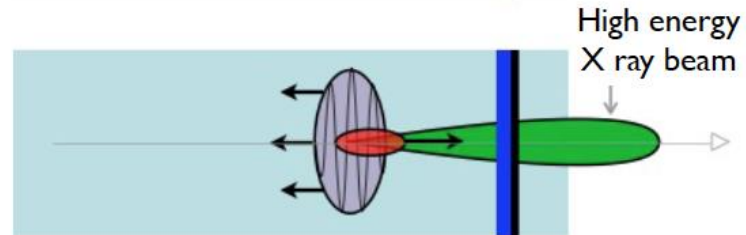
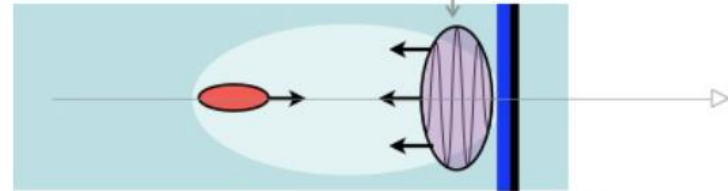






Back reflected laser pulse

Plasma mirror



A single laser pulse

A plasma mirror reflects the laser beam

The back reflected laser collides with the accelerated electrons

No alignment : the laser and the electron beams naturally overlap

Save the laser energy !

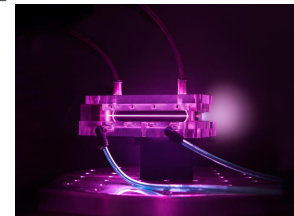
BELLA: BERkeley Lab Laser Accelerator

BELLA Facility: state-of-the-art 1.3 PW-laser for laser accelerator science:
>42 J in <40 fs (> 1PW) at 1 Hz laser and supporting infrastructure at LBNL

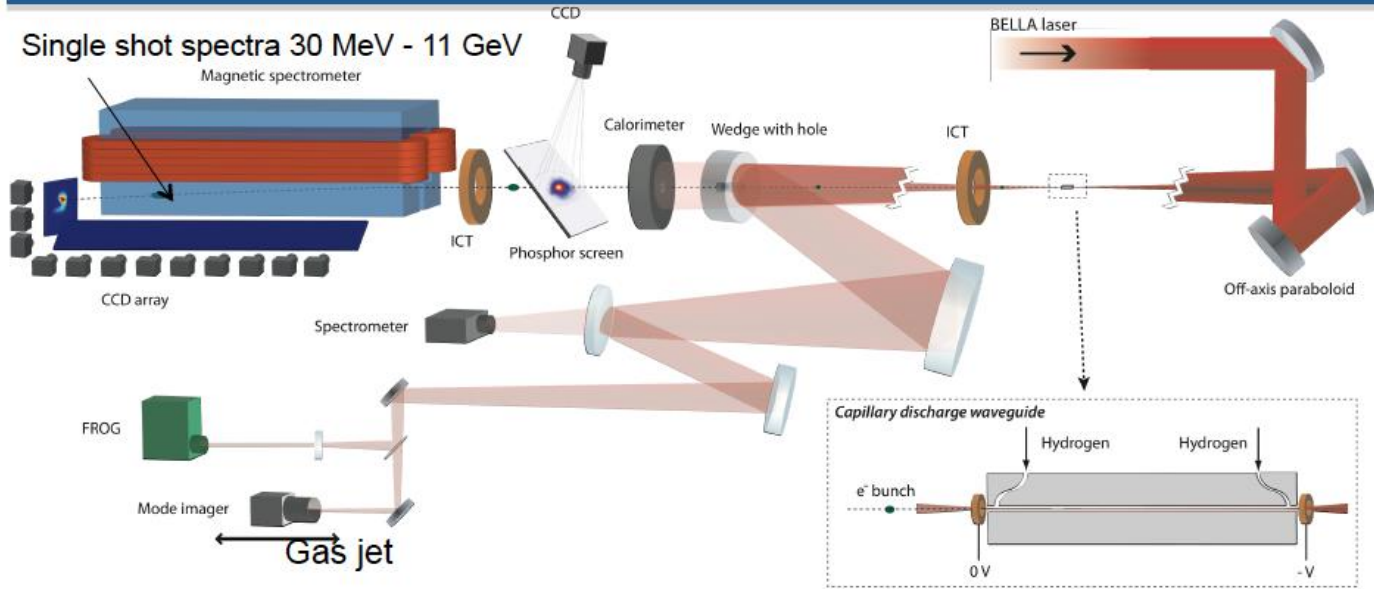


Critical HEP experiments:

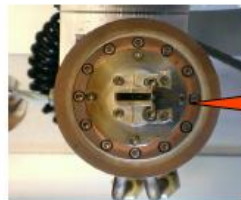
- 10 GeV electron beam from <1 m LPA
- Staging LPAs
- Positron acceleration



Experiments at LBNL use the BELLA laser focused by a 14 m focal length off-axis paraboloid onto gas jet or capillary discharge targets



Capillary discharge

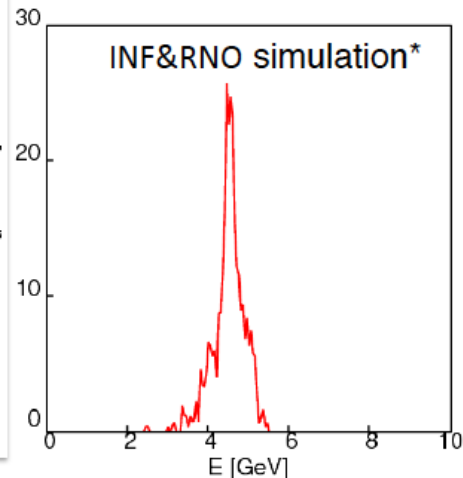
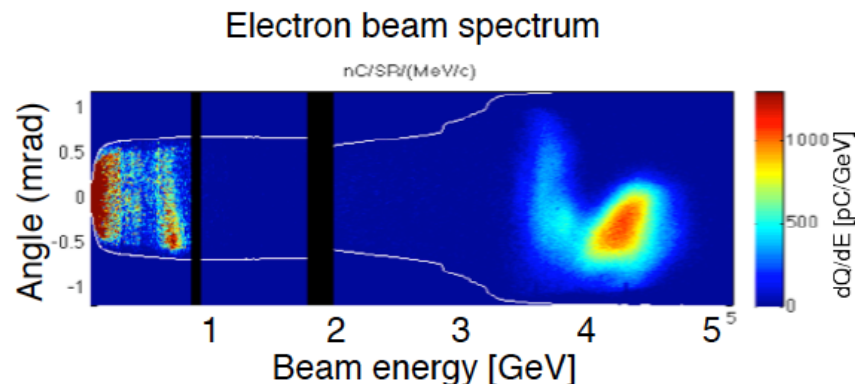


Big Laser In



4.25 GeV beams have been obtained from 9 cm plasma channel powered by 310 TW laser pulses (15 J)

*C. Benedetti et al., proceedings of AAC2010, proceedings of ICAP2012



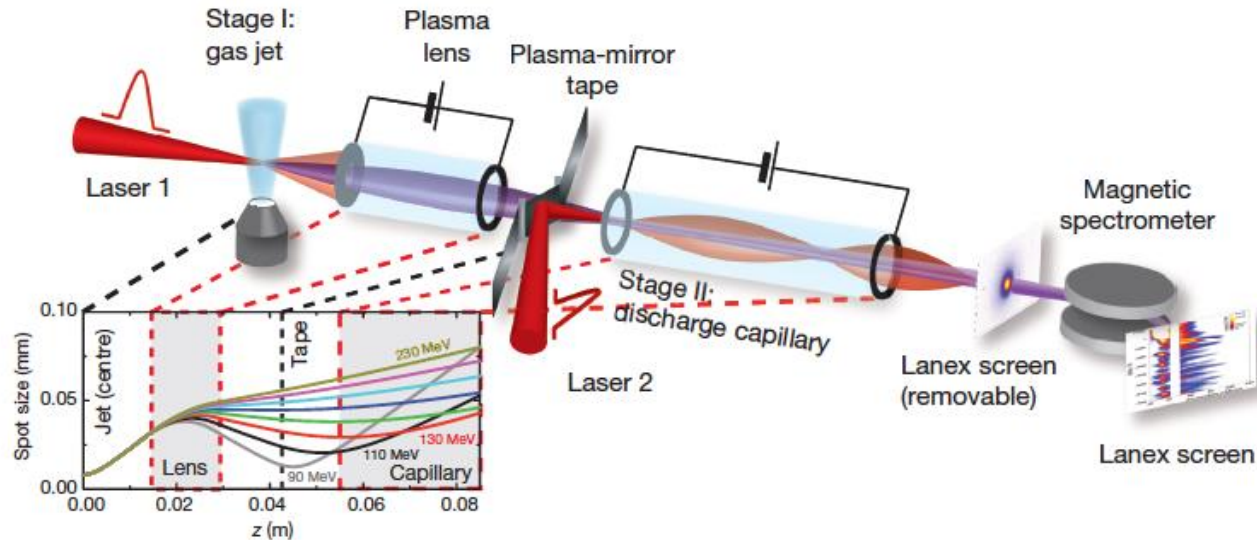
- **Laser** (E=15 J):
 - Measured) longitudinal profile ($T_0 = 40$ fs)
 - Measured far field mode ($w_0 = 53$ μm)
- **Plasma**: parabolic plasma channel (length 9 cm, $n_0 \sim 6-7 \times 10^{17}$ cm^{-3})

	Exp.	Sim.
Energy	4.25 GeV	4.5 GeV
$\Delta E/E$	5%	3.2%
Charge	~ 20 pC	23 pC
Divergence	0.3 mrad	0.6 mrad

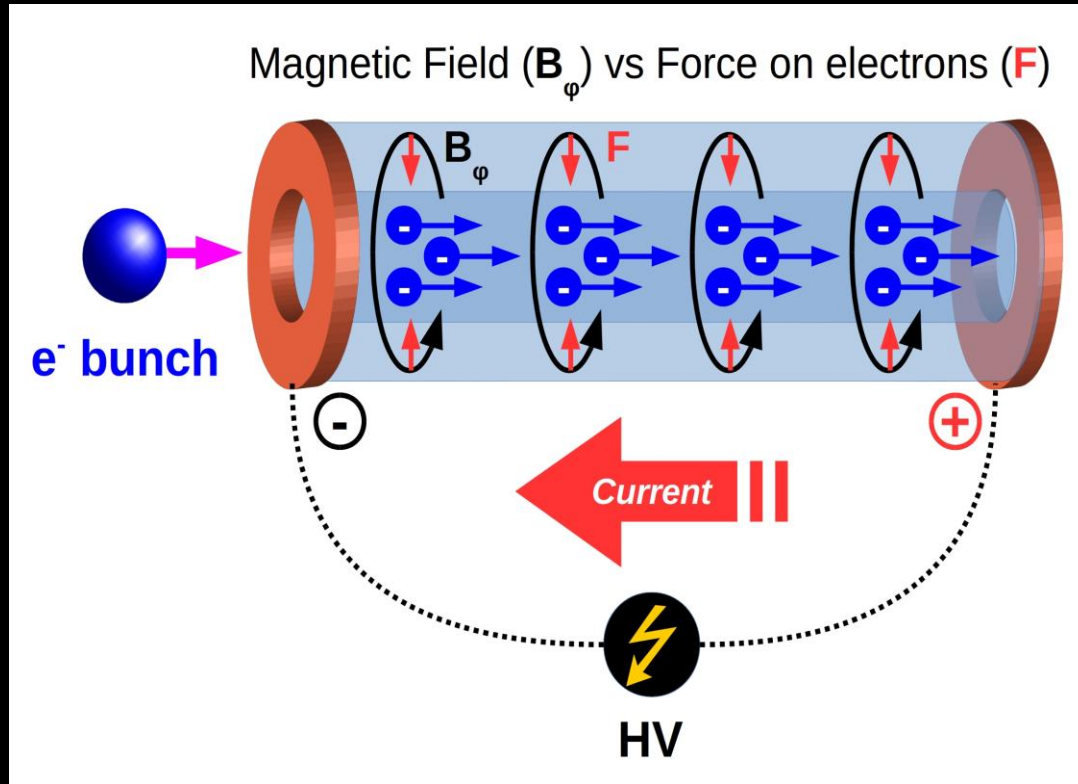
W.P. Leemans et al., PRL 2014

Multistage coupling of independent laser-plasma accelerators

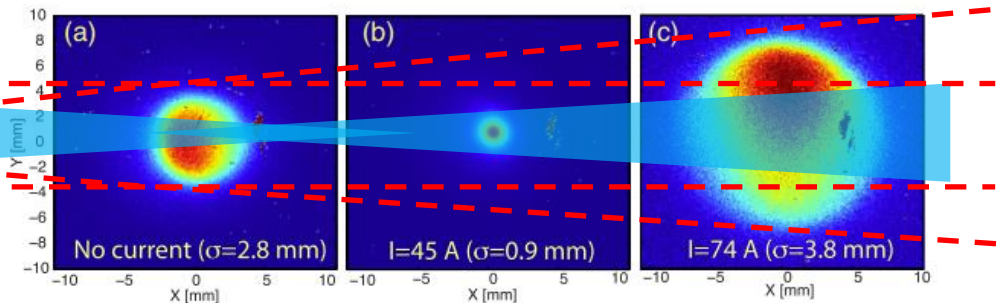
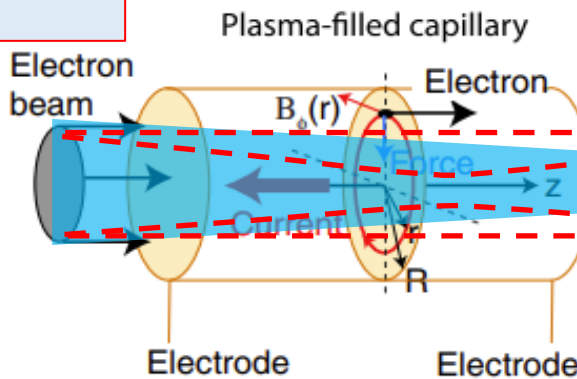
S. Steinke¹, J. van Tilborg¹, C. Benedetti¹, C. G. R. Geddes¹, C. B. Schroeder¹, J. Daniels^{1,3}, K. K. Swanson^{1,2}, A. J. Gonsalves¹, K. Nakamura¹, N. H. Matlis¹, B. H. Shaw^{1,2}, E. Esarey¹ & W. P. Leemans^{1,2}



Active Plasma Lens

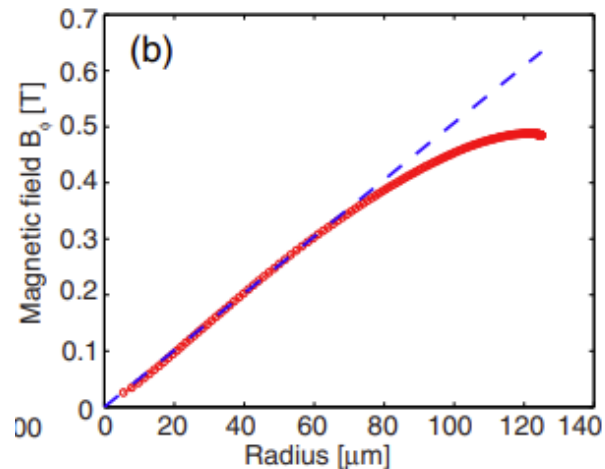
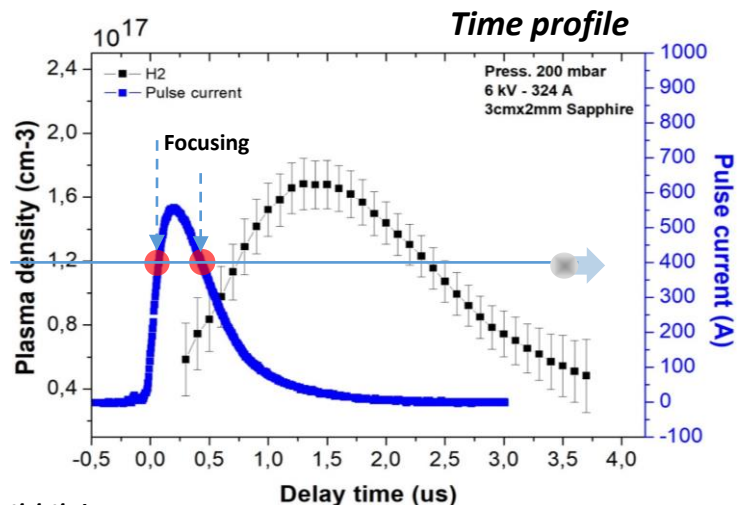


Active Plasma lenses
(APL)



$$\frac{\partial B_\phi}{\partial r} = \mu_0 I_0 / (2\pi R^2)$$

- R capillary radius
- I_0 peak current



Beam Quality Preservation requires proper matching – Focusing

→ Ion column provides linear focusing force

Radial electric field: $E_r(r) = \frac{en_{pe}}{2\epsilon_0} r$ (Gauss' law on cylinder of ions)

→ Plug it in envelope equation: $\sigma_r''(z) + \sigma_r(z) \left(K - \frac{\epsilon_g^2}{\sigma_r^4(z)} \right) = 0$

=0:
Matching condition

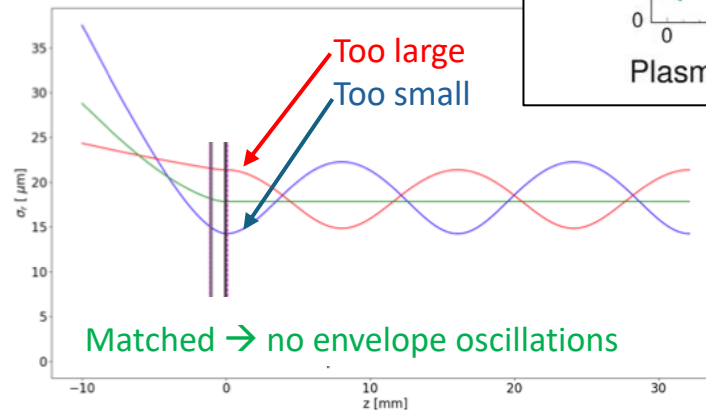
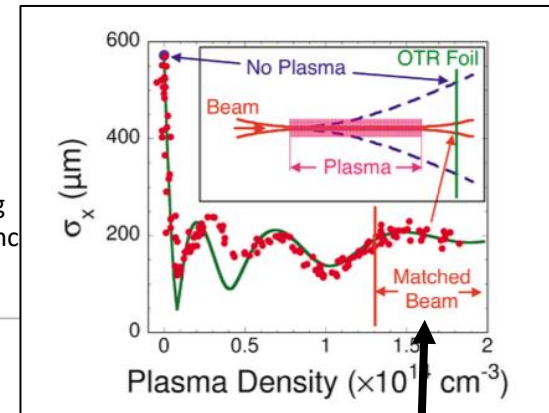
Meter-Scale Plasma-Wakefield Accelerator Driven by a Matched Electron Beam

P. Muggli,¹ B. E. Blue,² C. E. Clayton,² S. Deng,¹ F.-J. Decker,³ M. J. Hogan,³ C. Huang,² R. Iverson,³ C. Joshi,² T. C. Katsouleas,¹ S. Lee,¹ W. Lu,² K. A. Marsh,² W. B. Mori,² C. L. O'Connell,³ P. Raimondi,³ R. Siemann,³ and D. Walz³

Equilibrium between focusing force and emittance

- $\beta = \frac{\sigma^2(0)}{\epsilon_g} = \sqrt{\frac{2\epsilon_0 m_e c^2 \gamma}{n_{pe} e^2}}$
- Injection at waist: $\sigma'(z=0) = 0$

Note: the matching condition is given by β !
→ Stronger focusing required for higher density



Beam leaves the plasma at waist
→ matched → constant size

Beam Quality Preservation requires proper matching – Focusing

→ Ion column provides linear focusing force

Radial electric field: $E_r(r) = \frac{en_{pe}}{2\epsilon_0} r$ (Gauss' law on cylinder of ions)

→ Plug it in envelope equation: $\sigma_r''(z) + \sigma_r(z) \left(K - \frac{\epsilon_g^2}{\sigma_r^4(z)} \right) = 0$

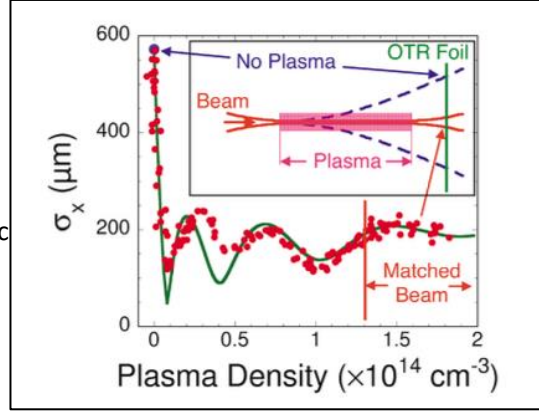
=0:
Matching condition

VOLUME 93, NUMBER 1 PHYSICAL REVIEW LETTERS week ending 2 JULY 2004

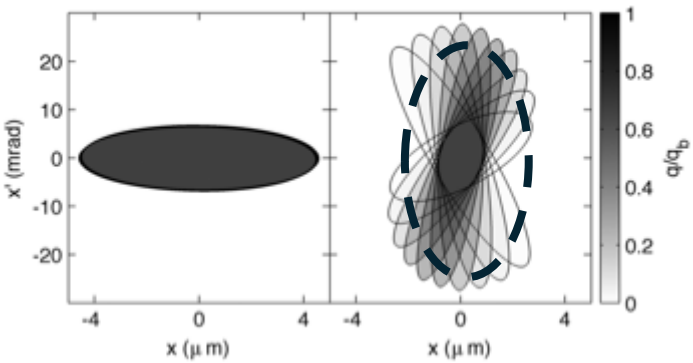
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Equilibrium between focusing force and emittance

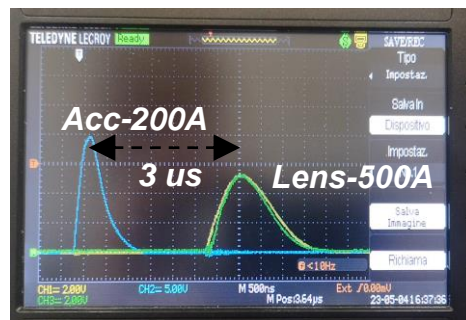
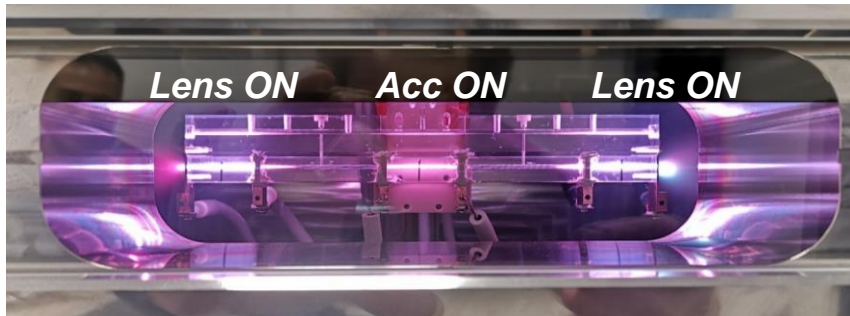


Q: WHY IS IMPORTANT TO MATCH THE BEAM ENVELOPE?

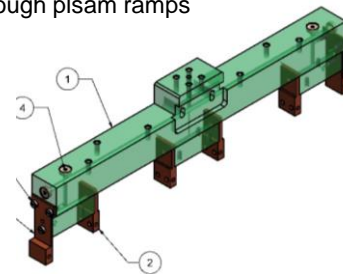


- Finite energy spread:
 - different energy slices rotate at different rates in transverse phase space
 - slice emittance preserved (linear focusing)
 - projected (i.e., overall) emittance grows!

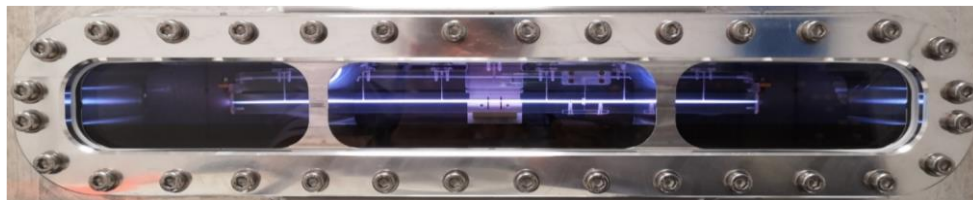
Integrated capillary



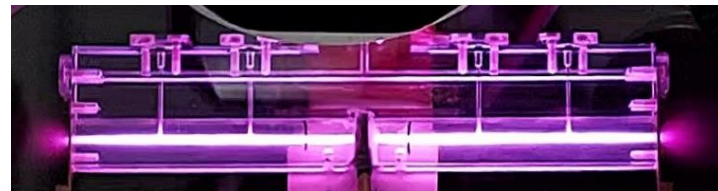
Studies on cross-talk effects:
Design of electrodes and HV-circuits to
reduce the interaction among discharges
through plasm ramps



Very long capillary

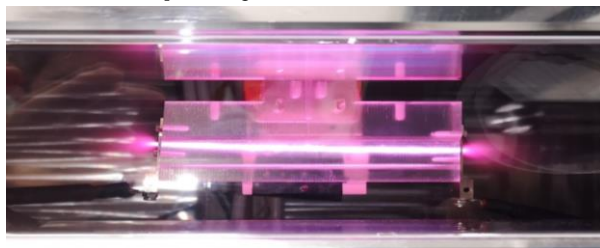


Design of m-scale capillaries for EuPRAXIA project by using segmented
capillaries: design of HV-voltage circuits and discharge synchronization



Segmented capillary

Curved capillary for APD



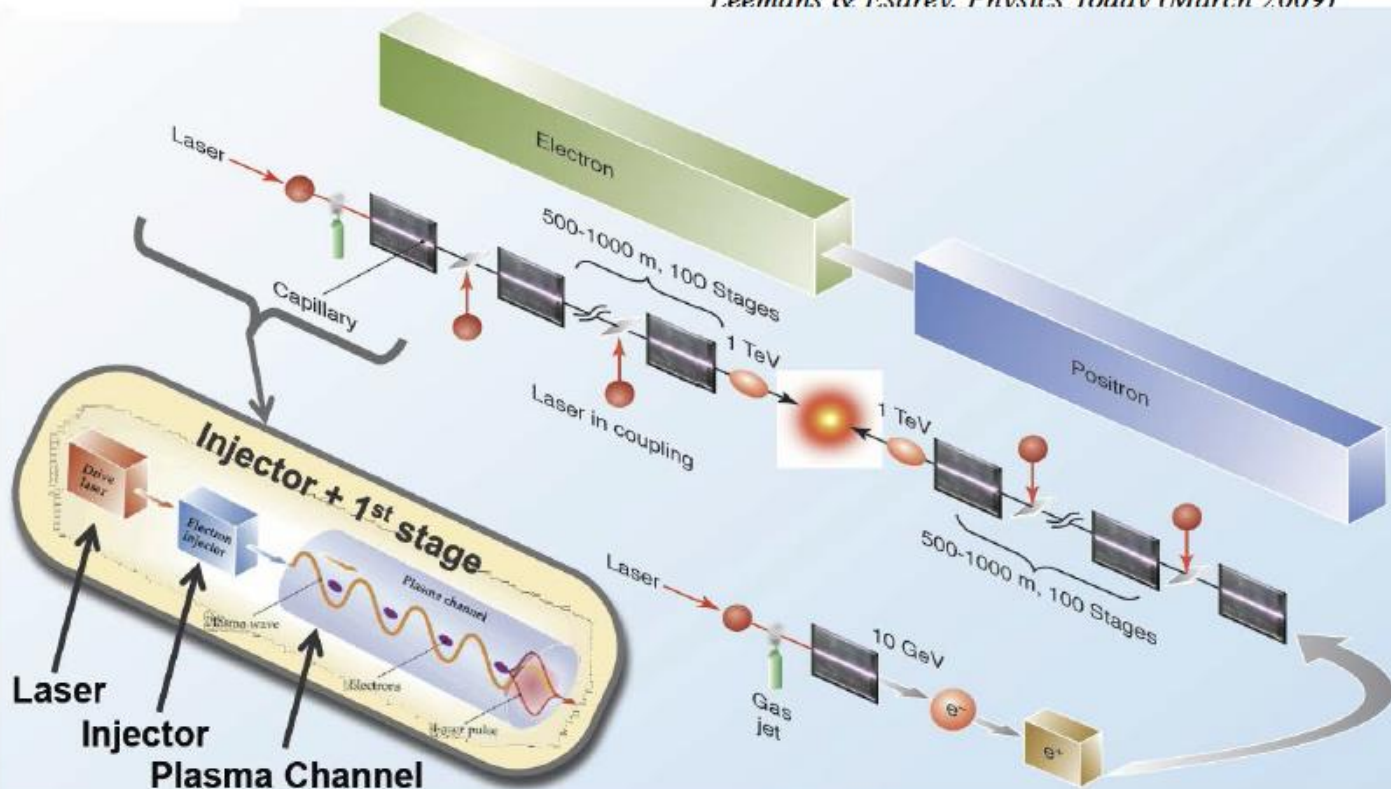
Design of new geometries for curved
channels: HV-circuits to allow high current
pulses

Pompili, Riccardi, et al, Guiding of charged particle beams in curved capillary-discharge
waveguides, AIP Advances 8.1 (2018)



Laser-Plasma-Accelerator LC

Leemans & Esarev. Physics Today (March 2009)





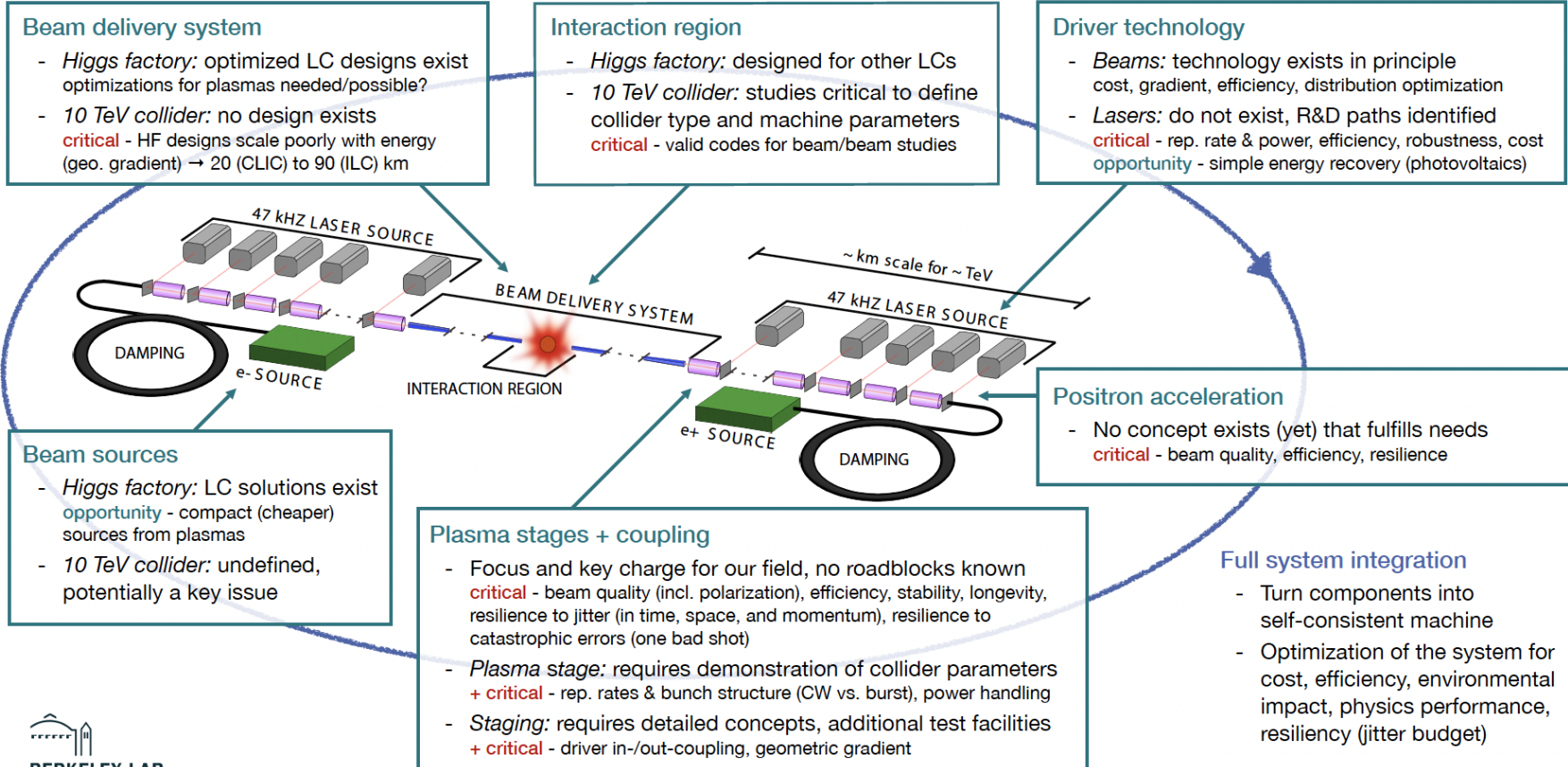
Parameter Set for LPWA LC

Case: CoM Energy (Plasma density)	1 TeV (10^{17} cm^{-3})	1 TeV ($2 \times 10^{15} \text{ cm}^{-3}$)	10 TeV (10^{17} cm^{-3})	10 TeV ($2 \times 10^{15} \text{ cm}^{-3}$)
Energy per beam (TeV)	0.5	0.5	5	5
Luminosity ($10^{34} \text{ cm}^{-2} \text{ s}^{-1}$)	2	2	200	200
Electrons per bunch ($\times 10^{10}$)	0.4	2.8	0.4	2.8
Bunch repetition rate (kHz)	15	0.3	15	0.3
Horizontal emittance $\gamma \varepsilon_x$ (nm-rad)	100	100	50	50
Vertical emittance $\gamma \varepsilon_y$ (nm-rad)	100	100	50	50
β^* (mm)	1	1	0.2	0.2
Horizontal beam size at IP σ_x^* (nm)	10	10	1	1
Vertical beam size at IP σ_y^* (nm)	10	10	1	1
Disruption parameter	0.12	5.6	1.2	56
Bunch length σ_z (μm)	1	7	1	7
Beamstrahlung parameter Υ	180	180	18,000	18,000
Beamstrahlung photons per e, n_γ	1.4	10	3.2	22
Beamstrahlung energy loss δ_E (%)	42	100	95	100
Accelerating gradient (GV/m)	10	1.4	10	1.4
Average beam power (MW)	5	0.7	50	7
Wall plug to beam efficiency (%)	6	6	10	10
One linac length (km)	0.1	0.5	1.0	5



×2+FF

Plasma collider challenges

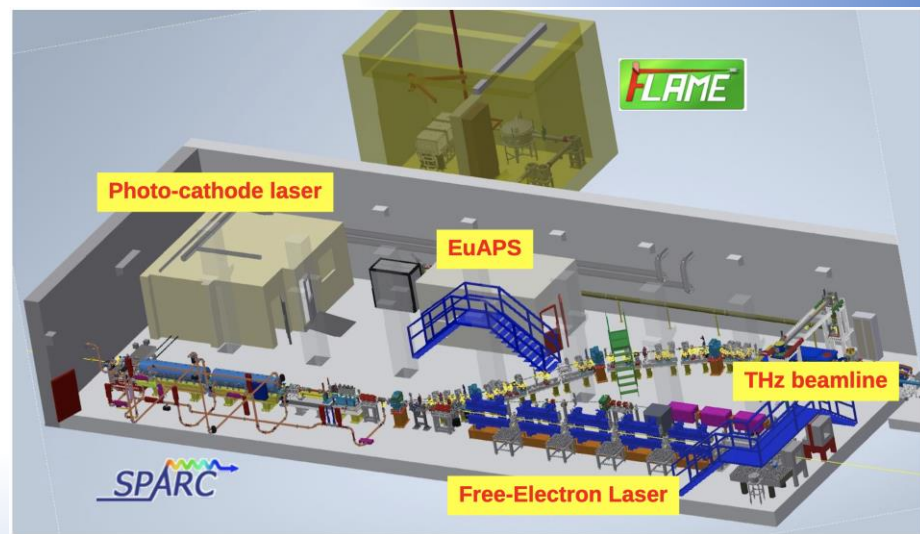


https://jobs.dsi.infn.it/dettagli_job.php?id=4180

DETTAGLIO DEL BANDO

Numero :	27077
Anno :	2024
Titolo:	n. 20 POST-DOCTORAL SENIOR LEVEL 3 RESEARCH GRANT IN EXPERIMENTAL PHYSICS
Numero Posti:	20
Tipo:	Assegno di ricerca
Data bando:	12-09-2024
Data scadenza:	15-11-2024
Bando:	27077.pdf

Dead line: November 15



LPAW 2025 Laser and Plasma Accelerators Workshop 2025 14-18 April 2025, Ischia Island, Italy



<https://agenda.infn.it/event/42311/>

The **Laser and Plasma Accelerators Workshop 2025 (LPAW 2025)** will be held at **Hotel Continental Ischia**, in the **Ischia Island (Campania, Italy)**, from **Monday 14 to Friday 18 April 2025**.

The Laser and Plasma Accelerators Workshop (LPAW) series is one of the leading workshops in the field of plasma-based acceleration and radiation generation.

The following scientific topics will be the main focus of the conference:

- Plasma-based lepton acceleration (experiments, simulations, theory, diagnostics...).
- Plasma-based ion acceleration (experiments, simulations, theory, diagnostics...).
- Secondary radiation generation and applications (experiments, simulations, theory, diagnostics...).

John Dawson Thesis Prize

“John Dawson Thesis Prize” is awarded on a biannual basis to the best PhD thesis in the area of plasma accelerators driven by laser or particle beams. The prize will be awarded for fundamental (theoretical or experimental) or applied aspects.

Each prize winner will receive a certificate of merit, up to 500 Euros, and financial support to attend the “Laser and Plasma Accelerators Workshop,” where the prize will be awarded.

Beam Driven PWFA