

# Searching for top squarks from the string landscape at HL-LHC

Juhi Dutta  
University of Oklahoma

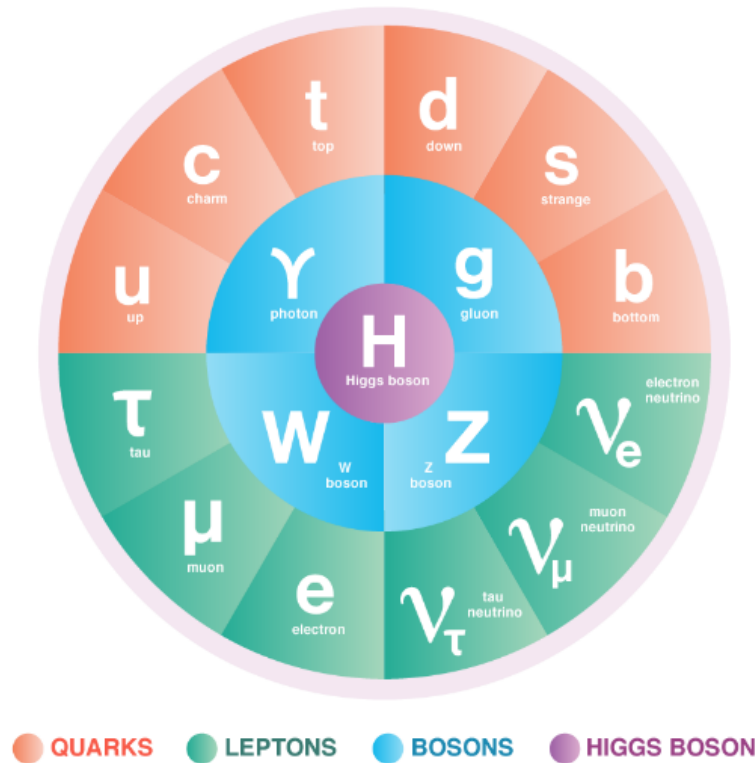
based on *Phys.Rev.D* 108 (2023) 7, 075027  
with H.Baer, V.Barger, D.Sengupta, K.Zhang.

*DPF-PHENO 2024 University of Pittsburgh / Carnegie Mellon University*



May 13, 2024

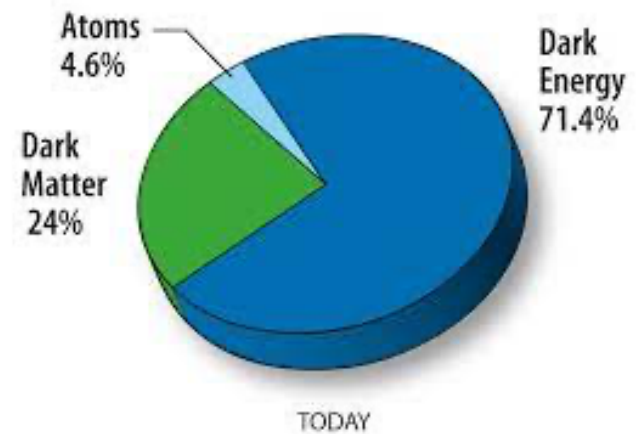
# The Standard Model of Particle Physics



The Standard Model of Particle Physics, FERMILAB

# Why New Physics?

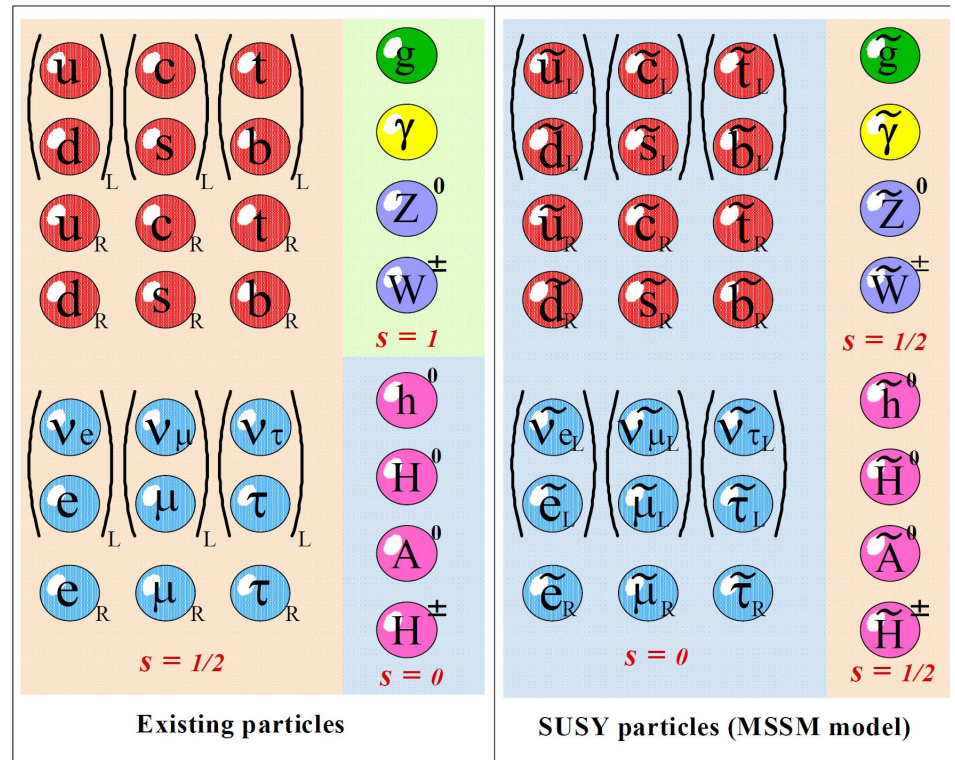
- Gauge hierarchy problem
- Dark Matter
- Dark Energy
- Non-zero neutrino masses
- Matter-antimatter asymmetry
- Gravity



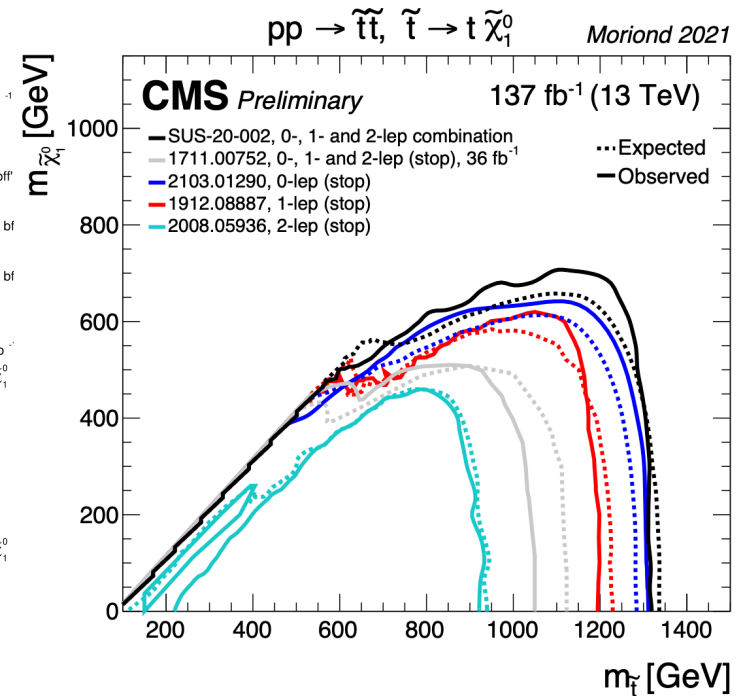
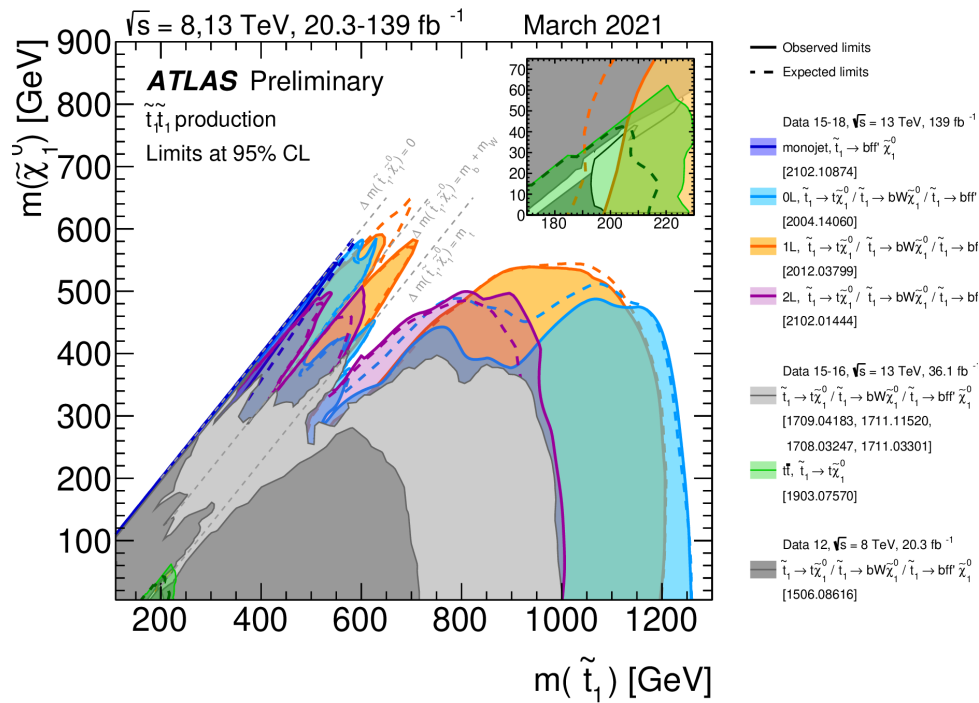
NASA.

# Supersymmetry

- Special symmetry providing a super partner to SM particles differing by spin-1/2 => Minimal Supersymmetric Standard Model (MSSM).
- Resolves big Hierarchy problem.
- R-parity, a discrete symmetry, when conserved provides a dark matter candidate.
- Gauge coupling unification.
- Minimal extensions also address neutrino masses.



# Current Status from LHC: lightest stop



Limits of  $m_{\tilde{t}_1} > 1.3 \text{ TeV}$  for massless  $\tilde{\chi}_1^0$  from simplified model analyses assuming 100% decay modes.

## Expectations from Naturalness measures for SUSY

- Light stops are a lucrative target for high luminosity LHC.

- Early estimates of naturalness such as by Barbieri-Gudice,  $\Delta_{BG} = \max_i \left| \frac{p_i \partial m_Z^2}{m_Z^2 \partial p_i} \right|$ . R. Barbieri and G. F. Giudice, Nucl. Phys. B 306 (1988)
- Stringent upper limits on third generation squarks around ~450 GeV. ↑  $p_i$  independent soft terms at high scale.

- Assumes soft terms are independent, however not true for eg, in UV complete scenario such as SUGRA.

- More conservative measures of naturalness:  $\Delta_{EW}$ , the ratio of the largest term on the right-hand-side (below) to  $m_Z^2/2$ :

$$\frac{m_Z^2}{2} = \frac{m_{H_d}^2 + \Sigma_d^d - (m_{H_u}^2 + \Sigma_u^u) \tan^2 \beta}{\tan^2 \beta - 1} - \mu^2$$

$$\Delta_{EW} = \frac{\max \left| \frac{m_{H_d}^2 + \Sigma_d^d - (m_{H_u}^2 + \Sigma_u^u) \tan^2 \beta}{\tan^2 \beta - 1} - \mu^2 \right|}{m_Z^2/2}$$

Limits on particles for 3%  $\Delta_{BG}$  ( $\Delta_{BG} < 30$ ) fine-tuning and for  $\Delta_{EW} < 30$ .

Mass	BG/DG	$\Delta_{EW}$
$\mu$	<350 GeV	<350 GeV
$m_{\tilde{g}}$	<400 – 600 GeV	<6 TeV
$m_{\tilde{t}_1}$	<450 GeV	<3 TeV
$m_{\tilde{q}, \tilde{\ell}}$	<550 – 700 GeV	<10 – 30 TeV

Baer, et.al, 2202.11578

# The Minimal Supersymmetric Standard Model from the landscape

- Motivated by Weinberg's anthropic solution to the cosmological constant, one tries to address the origin of the SUSY breaking scale in the string landscape where  $10^{500}$  vacua solutions arise from compactification from 10 to 4 spacetime dimensions.

Weinberg, Phys. Rev. Lett. 59, 2607

- Supersymmetric models with low electroweak fine-tuning are expected to be more prevalent on the string landscape than fine-tuned models.

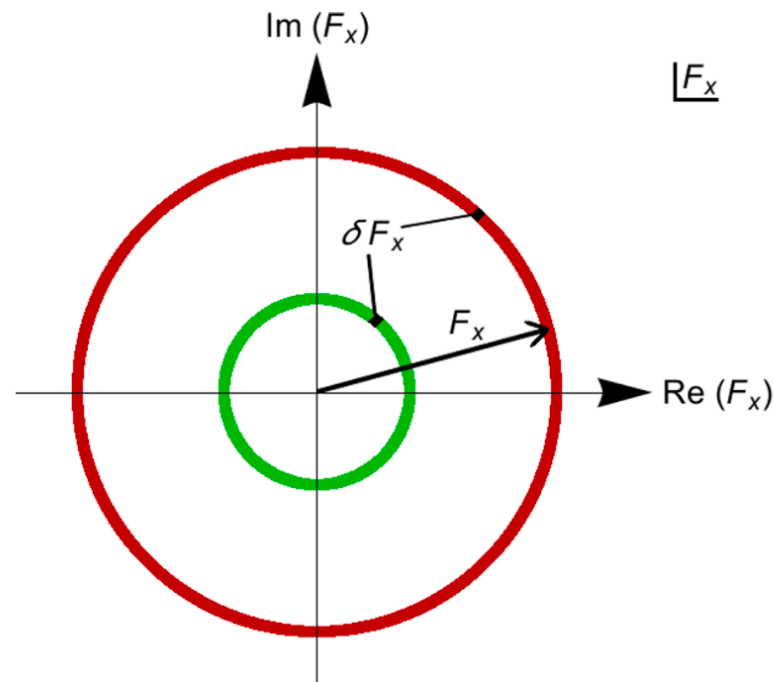
Baer, et.al, 2202.11578

- Douglas et.al, proposed a probabilistic view of naturalness, *stringy naturalness*, by identifying the statistical trends for the observables over the many landscape vacua solution we are likely to be in.

Douglas, Comptes Rendus  
Physique 5 (2004) 965–977

- Stringy naturalness: an observable  $\mathcal{O}_2$  is more natural than  $\mathcal{O}_1$  if more phenomenological vacua lead to  $\mathcal{O}_2$  than  $\mathcal{O}_1$ .

- For  $n_F = 1$ ,  $n_D = 0$ , linear statistical draw to soft terms.



Baer, et.al, 2202.11578

**Annuli of the complex  $F_x$  plane giving rise to linearly increasing selection of soft SUSY breaking terms.**



- This stringy statistical draw must be compensated for by requiring that the **derived value for the weak scale in each pocket universe of the multiverse be not too far from our measured value**, so that complex nuclei and hence atoms arise in any anthropically allowed pocket universe => ABDS window.

Agrawal et.al, Phys. Rev. D 57 (1998) 5480–5492

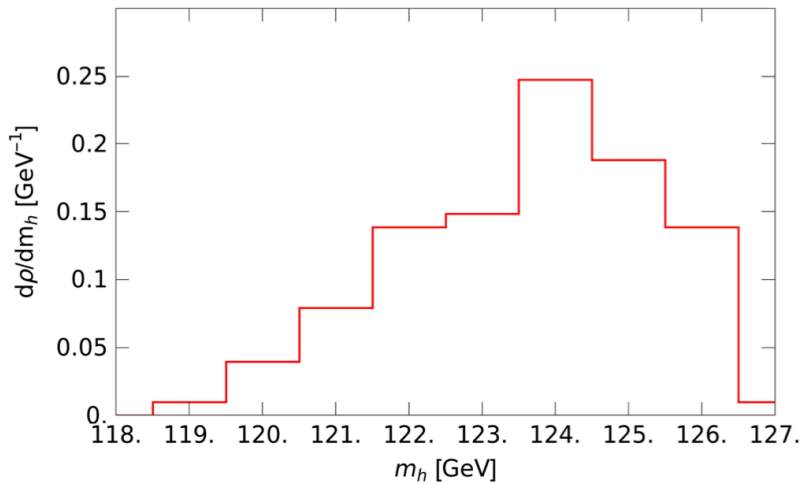
- We restrict ourselves to landscape vacua with MSSM as the low energy theory.
- Assume a linear draw for the soft terms  $m_0(1,2)$ ,  $m_0(3)$ ,  $A_0$ ,  $m_A$  which scan independently in the string landscape.
- The soft terms scan as  $f_{SUSY} \sim m_{soft}^n$ , where  $n = 2n_F + n_D - 1$ .

# Non Universal Higgs Model (NUHM(3)): Global Scan

- Free parameters:  $m_0(1,2), m_0(3), \tan \beta, A_0, \mu, m_A$
- Global scan:
  - $m_0(1,2) : 0.1 - 45 \text{ TeV}$ ,
  - $m_0(3) : 0.1 - 10 \text{ TeV}$ ,
  - $m_{1/2} : 0.5 - 3 \text{ TeV}$ ,
  - $A_0 : 0 - (-20) \text{ TeV}$ ,
  - $\tan \beta : 3 - 60 \text{ TeV}$  (uniform scan),
  - $m_A : 0.3 - 10 \text{ TeV}$ ,

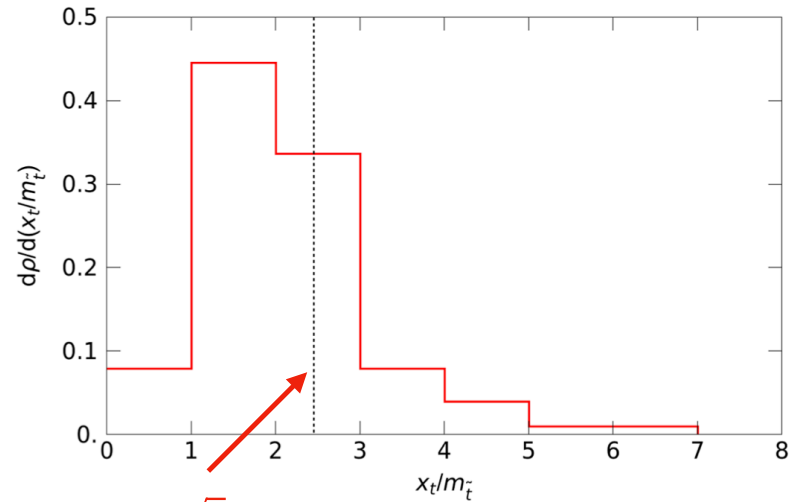
with fixed  $\mu = 200 \text{ GeV}$  and the upper bounds on the parameters are set beyond the upper bounds derived from the ABDS window.

# Probability distributions: Higgs



Probability distribution for lightest higgs mass for n=1.

$$m_h^2 \simeq m_Z^2 \cos^2 2\beta + \frac{3g^2}{8\pi^2} \frac{m_t^4}{m_W^2} \left[ \ln \frac{m_t^2}{m_{\tilde{t}}^2} + \frac{x_t^2}{m_{\tilde{t}}^2} \left( 1 - \frac{x_t^2}{12m_{\tilde{t}}^2} \right) \right]$$



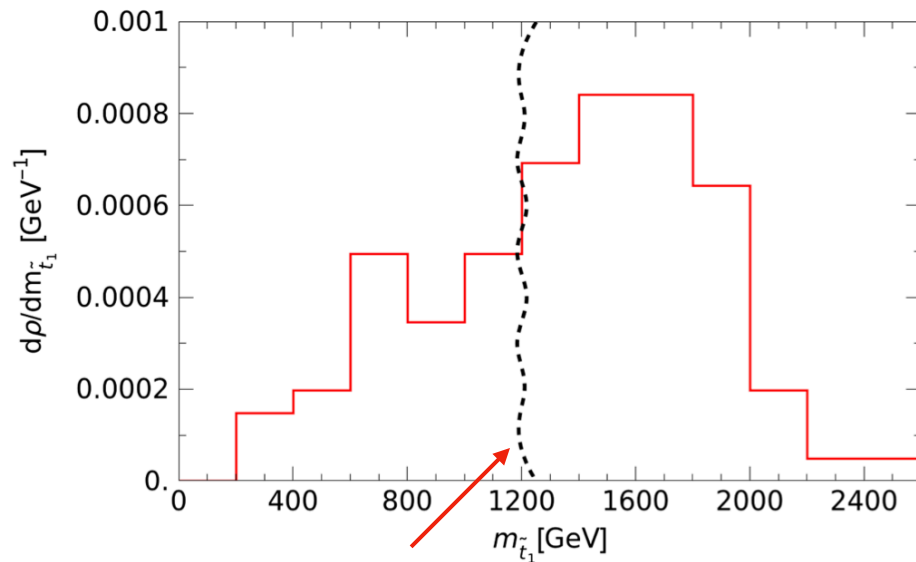
$$x_t = \sqrt{6}m_{\tilde{t}}$$

Probability distribution of  $x_t/m_{\tilde{t}_1}$  for n=1.

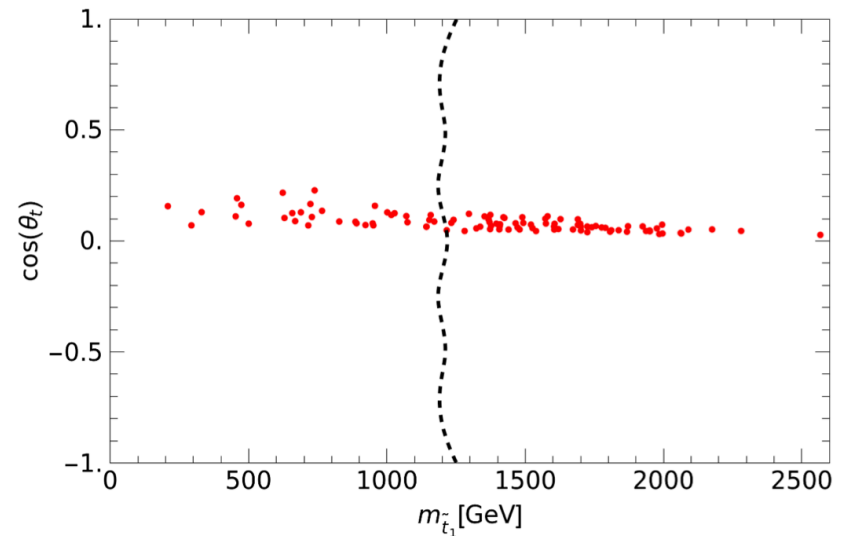
$$x_t = A_t - \mu \cot \beta$$

$$m_{\tilde{t}}^2 = m_{Q_3} m_{U_3}$$

# Properties of the Stop Sector from the landscape



Limit from LHC ~ 1.1-1.3 TeV.



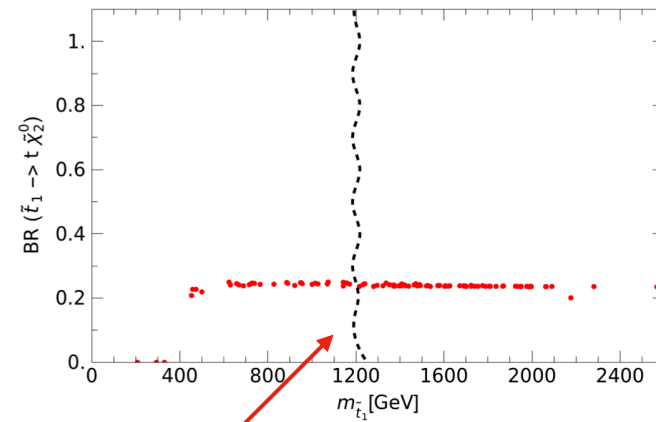
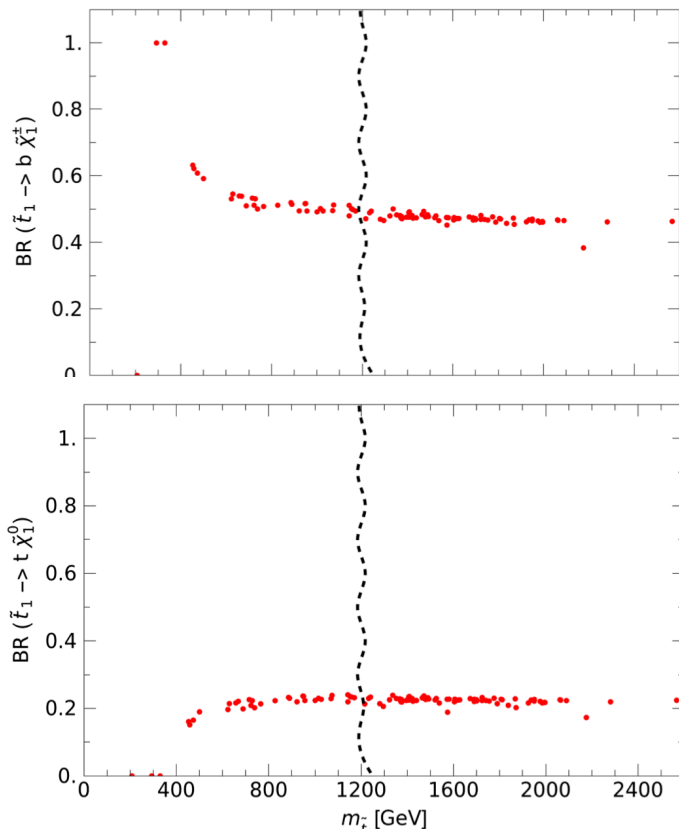
The lightest stop mass = 1-2.5 TeV and mostly right-handed.

Also consistent with the lightest CP-even Higgs mass ~125 GeV.

$$\tilde{t}_1 = \cos \theta_t \tilde{t}_L - \sin \theta_t \tilde{t}_R$$

# Decays of the stops

- Lightest stop, mostly right-handed, hence decays to  $b\tilde{\chi}_1^\pm, t\tilde{\chi}_1^0, t\tilde{\chi}_2^0$ : 2:1:1



Limit from LHC

# Benchmark point

Natural spectra characterized by

- gluinos = 2-6 TeV
- light stops = 1-2.5 TeV,
- higgsinos = 100-400 GeV

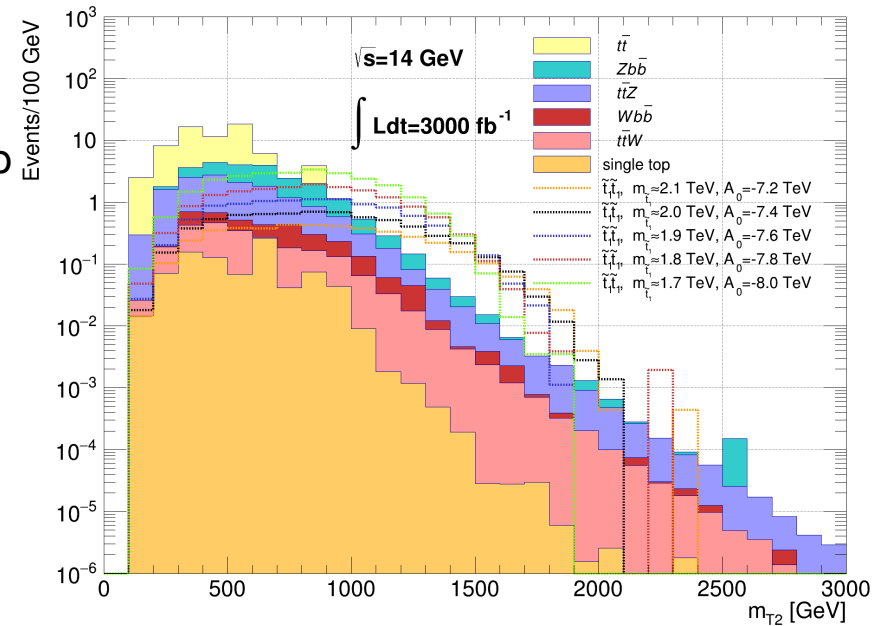
consistent with  $\Delta_{EW} \leq 30$ .

Parameter	Stringy natural BM point
$m_0$	5 TeV
$m_{1/2}$	1.2 TeV
$A_0$	-8 TeV
$\tan \beta$	10
$\mu$	250 GeV
$m_A$	2 TeV
$m_{\tilde{g}}$	2830 GeV
$m_{\tilde{u}_L}$	5440 GeV
$m_{\tilde{u}_R}$	5561 GeV
$m_{\tilde{e}_R}$	4822 GeV
$m_{\tilde{t}_1}$	1714 GeV
$m_{\tilde{t}_2}$	3915 GeV
$m_{\tilde{b}_1}$	3949 GeV
$m_{\tilde{b}_2}$	5287 GeV
$m_{\tilde{\tau}_1}$	4746 GeV
$m_{\tilde{\tau}_2}$	5110 GeV
$m_{\tilde{\nu}_\tau}$	5107 GeV
$m_{\tilde{\chi}_1^\pm}$	261.7 GeV
$m_{\tilde{\chi}_2^\pm}$	1020.6 GeV
$m_{\tilde{\chi}_1^0}$	248.1 GeV
$m_{\tilde{\chi}_2^0}$	259.2 GeV
$m_{\tilde{\chi}_3^0}$	541.0 GeV
$m_{\tilde{\chi}_4^0}$	1033.9 GeV
$m_h$	124.7 GeV
$\Omega_{\tilde{\chi}_1^0}^{\text{std}} h^2$	0.016
$\text{BR}(b \rightarrow s\gamma) \times 10^4$	3.1
$\text{BR}(B_s \rightarrow \mu^+\mu^-) \times 10^9$	3.8
$\sigma^{\text{SI}}(\tilde{\chi}_1^0, p)$ (pb)	$2.2 \times 10^{-9}$
$\sigma^{\text{SD}}(\tilde{\chi}_1^0, p)$ (pb)	$2.9 \times 10^{-5}$
$\langle \sigma v \rangle _{v \rightarrow 0}$ (cm <sup>3</sup> /sec)	$1.3 \times 10^{-25}$
$\Delta_{EW}$	22

# Phenomenology

- Signal topologies:  $tb + E_T, tt + E_T, bb + E_T$
- Key SM backgrounds:  $b\bar{b}Z, t\bar{t}Z, t\bar{t}W, t\bar{t}$ , single top suppressed using highly boosted top-jets:
- For  $tb + E_T$ :  $p_T(\text{top} - \text{jet}) > 400$  GeV,  $R=1.5$ .
- Cuts:  $E_T > 400$  GeV,
  - $H_T > 1.4$  TeV,  $L_T > 1.8$  TeV,
  - $\min(m_T(b_1, E_T), m_T(b_2, E_T)) > 175$  GeV,
  - $\Delta\Phi(b, E_T) > 40^\circ, \Delta\Phi(J, E_T) > 30^\circ$ .

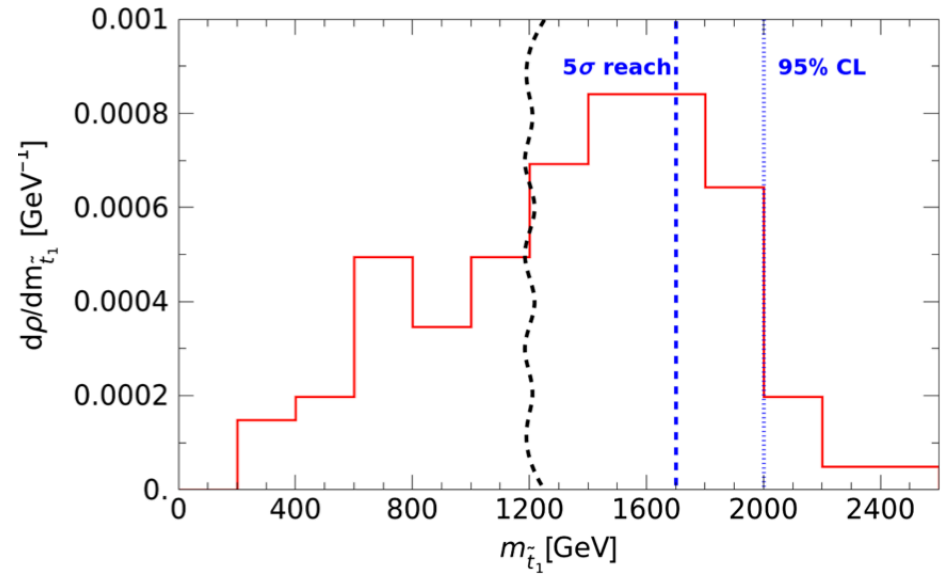
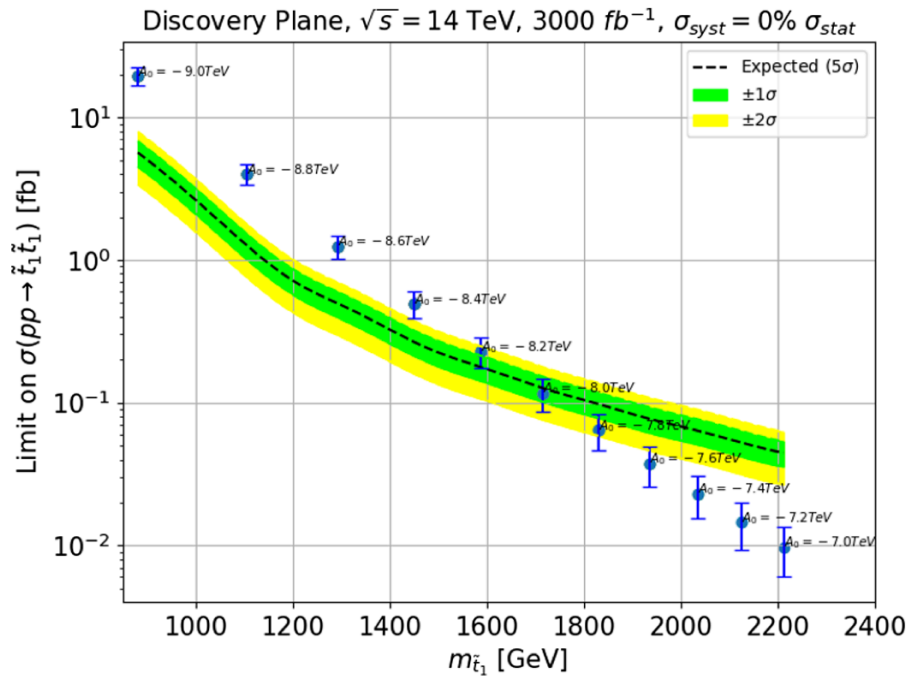
Similar cuts for  $tt + E_T$  and  $bb + E_T$ .



Distribution for  $m_{T_2}$  in the  $tb + \text{MET}$  final state.

Key kinematic variable for discrimination between signal and background  $m_{T_2}$  after cuts.

# Results



Probability distribution of the lightest stop with the predicted reach at HL-LHC.

Combined reach of stops  $\sim 1.65 \text{ TeV}$  at  $5\sigma$  and  $\sim 1.95 \text{ TeV}$ , at 95% C.L., covering most of the region allowed from the string landscape!



# Summary

- We examined top squark masses and other properties are expected from the string landscape where a power-law draw to large soft terms is expected.
- The derived value of the weak scale must lie within the ABDS window in order to allow for complex nuclei (and hence atoms) in each anthropically-allowed pocket universe.
- Under the stringy naturalness requirement, we find  $m_{\tilde{t}_1} = 1\text{-}2.5$  TeV with large mixing. The large mixing helps boost  $m_h$  to 125 GeV while minimizing the top squark contributions to the weak scale.
- In spite of the large mixing, the lightest top-squark is mainly right-handed, thus leading to mixed final states of  $tb + E_T, tt + E_T, bb + E_T$ .
- The top-squark pair production is revealed as an enhancement in the  $m_{T_2}$  distribution at high values of  $m_{T_2}$ .
- A combined reach of all channels at HL-LHC lead to a reach of  $\sim 1.65$  TeV at  $5\sigma$  and  $\sim 1.95$  TeV at 95%CL.

**Thank You!**