

An Area Law for Entanglement Entropy in Particle Scattering

Based on 2405.08056 and 240X.XXXXX with Ian Low

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Introduction

- Entanglement is a key phenomenon in quantum theories; characterized by entropy
- Area laws exist for entropy; may be related to holography
- We consider entanglement entropy in $2 \rightarrow 2$ scattering

Wave packets

Ingredient 1: Distinguishable particles

Consider the elastic scattering of $AB \rightarrow AB$. The initial state:

$$|\text{in}\rangle = |\psi_A\rangle \otimes |\psi_B\rangle$$

The density matrix: $\rho^i = |\text{in}\rangle\langle\text{in}|$: We need $\text{Tr} \rho^i = 1$

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$$|\psi\rangle = \int_p \psi(p) |p\rangle, \quad \langle\psi|\psi\rangle = \int \frac{d^3\vec{p}}{(2\pi)^3} |\psi(p)|^2 = 1$$

with $\langle p|q\rangle = (2\pi)^3 2E_p \delta^3(\vec{p} - \vec{q})$.

The final state

Ingredient 2: Unitary S -matrix

- $|\text{out}\rangle = S|\text{in}\rangle$, $S = 1 + iT$. $S^\dagger S = 1$, $2\text{Im} T = T^\dagger T$.

$$\langle \{k_f\} | T | \{k_i\} \rangle = (2\pi)^4 \delta^4 \left(\sum k_f - \sum k_i \right) M(\{k_i\}; \{k_f\})$$

- The “complete” final state density matrix is $|\text{out}\rangle\langle\text{out}|$

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Probabilities:

$$\mathcal{P}_{\text{tot}} = \langle \text{in} | T^\dagger T | \text{in} \rangle, \quad \mathcal{P}_{\text{el}} = \langle \text{in} | T^\dagger P_{\text{AB}} T | \text{in} \rangle, \quad \mathcal{P}_{\text{inel}} = \langle \text{in} | T^\dagger (1 - P_{\text{AB}}) T | \text{in} \rangle$$

The plane wave limit

Ingredient 3: The plane wave limit

We require $\psi_{A/B}(p)$ to

- Be centered around $\vec{k}_{A/B}$. In the CoM frame, $\vec{k}_A = -\vec{k}_B \equiv \vec{k}$
- Have a width of δ_p in momentum space
- Take the plane wave limit: $\delta_p \ll |\vec{k}|$, and expand around small $\delta_p/|\vec{k}|$

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Probabilities (“ \doteq ” means keeping the leading piece in $\delta_p/|\vec{k}|$):

$$\mathcal{P}_{\text{tot}} \doteq I_0(|\vec{k}|)\sigma_{\text{tot}}, \quad \mathcal{P}_{\text{el}} \doteq I_0(|\vec{k}|)\sigma_{\text{el}}, \quad \mathcal{P}_{\text{inel}} \doteq I_0(|\vec{k}|)\sigma_{\text{inel}}$$

where

$$I_0(|\vec{k}|) = 4|\vec{k}|\sqrt{s} \int_{p_1, p_2, q_1, q_2} \psi_A(p_1)\psi_B(p_2)\psi_A^*(q_1)\psi_B^*(q_2) \\ \times (2\pi)^4 \delta^4(q_1 + q_2 - p_1 - p_2)$$

The uniform wave packet

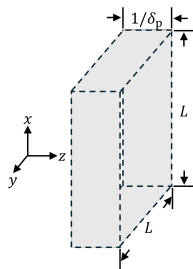
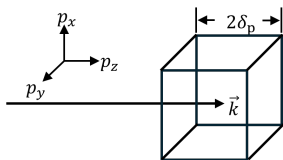
$\mathcal{P} \doteq I_0(|\vec{k}|)\sigma$ is small in the plane wave limit

- $I_0(|\vec{k}|) \sim \delta_p^2$ has the dimension of area⁻¹
- The size of $I_0(|\vec{k}|)$ depends on the overlap of the wave functions in the transverse directions in position space

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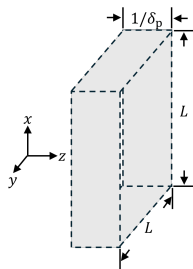
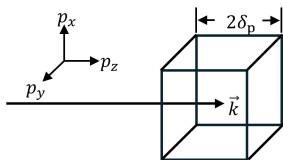
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The uniform wave packet

$\mathcal{P} \doteq I_0(|\vec{k}|)\sigma$ is small in the plane wave limit

- $I_0(|\vec{k}|) \sim \delta_p^2$ has the dimension of area^{-1}
- The size of $I_0(|\vec{k}|)$ depends on the overlap of the wave functions in the transverse directions in position space



- Roughly uniform in the transverse direction: $L \gg 1/\delta_p$, with $\mathcal{O}(1/(\delta_p L))$ corrections
- For head-on collisions, $I_0(|\vec{k}|) \doteq 1/L^2$, $\mathcal{P} \doteq \sigma/L^2$

Final state entropy

For the elastic scattering $AB \rightarrow AB$, $|\text{out}\rangle_{\text{el}} = P_{AB}|\text{out}\rangle$

- $\text{Tr}(|\text{out}\rangle\langle\text{out}|) = \langle\text{in}|S^\dagger S|\text{in}\rangle = 1$
- Properly normalized ρ^f :

$$\rho^f = \frac{1}{1 - \mathcal{P}_{\text{inel}}} |\text{out}\rangle_{\text{el}} \langle\text{out}|$$

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The key result

For the linear entropy $\mathcal{E}_2(\rho) = 1 - \text{Tr} \rho^2$, the final entanglement entropy is

$$\mathcal{E}_2^f \doteq 2I_0(|\vec{k}|) \sigma_{\text{el}} \doteq 2\mathcal{P}_{\text{el}}$$

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For uniform wave packets,

$$\mathcal{E}_2^f \doteq 2\frac{\sigma_{\text{el}}}{L^2}$$

For Tsallis and Rényi entropies

$$\mathcal{E}_n^f \doteq \frac{n}{n-1} I_0(|\vec{k}|)\sigma_{\text{el}} \doteq \frac{n}{n-1} \mathcal{P}_{\text{el}}$$

Generalization

Consider 3 type of quantum numbers:

- What distinguishes A and B, which we will call “charge”
- Kinematics, described by wave packets
- Other discrete quantum numbers $f_{A/B}$ (“flavors”), labelled by i and \bar{i}

Generalization

Consider 3 type of quantum numbers:

- What distinguishes A and B, which we will call “charge”
- Kinematics, described by wave packets
- Other discrete quantum numbers $f_{A/B}$ (“flavors”), labelled by i and \bar{i}

The initial state:

$$|\text{in}\rangle = |\psi_A\rangle \otimes |i\rangle \otimes |\psi_B\rangle \otimes |\bar{i}\rangle,$$

with $\langle i|j\rangle = \delta^{ij}$, $\langle \bar{i}|\bar{j}\rangle = \delta^{\bar{i}\bar{j}}$

Cutting the cake

	A	B
p	p_A	p_B
f	f_A	f_B

Cutting the cake

$$\begin{array}{cc} & \begin{array}{cc} \text{A} & \text{B} \end{array} \\ \begin{array}{c} p \\ f \end{array} & \begin{array}{|cc|} \hline p_A & p_B \\ \hline f_A & f_B \\ \hline \end{array} \end{array}$$

$$\begin{aligned} \mathcal{E}_{2,A}^f &\doteq 2\mathcal{P}_{\text{el}} \\ &\doteq 2I_0(|\vec{k}|)\sigma_{\text{el}} \\ &\rightarrow 2\frac{\sigma_{\text{el}}}{L^2} \end{aligned}$$

i, \bar{i} to any j, \bar{j}

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i, \bar{i} to any j, \bar{j}

i, \bar{i} to any $j \neq i, \bar{j} \neq \bar{i}$

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Cutting the cake

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Ingredients

- 1 Distinguishable particles
 - 2 Unitary S -matrix
 - 3 The plane wave limit
- Entanglement entropy as cross section/probability
 - General and non-perturbative (in coupling)
 - An area law for a 2-body system: What is the boundary? Holography?
 - Other entropies: Specific momenta? Other kinds of final states?
 - Beyond the plane wave limit: Constraints from entropy?