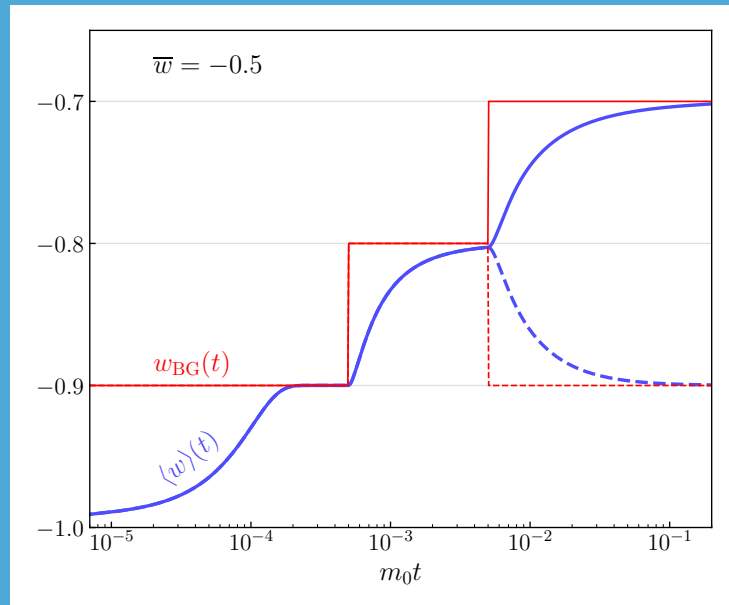


Cosmological Stasis from Dynamical Scalars



Brooks Thomas
LAFAYETTE
COLLEGE

Work supported
in part by



Based on work done in collaboration with:
Keith R. Dienes, Fei Huang, Lucien Heurtier, and
Timothy M. P. Tait [arXiv:2405.xxxxx]
DPF/PHENO, University of Pittsburgh, May 15th, 2024

Cosmological Stasis

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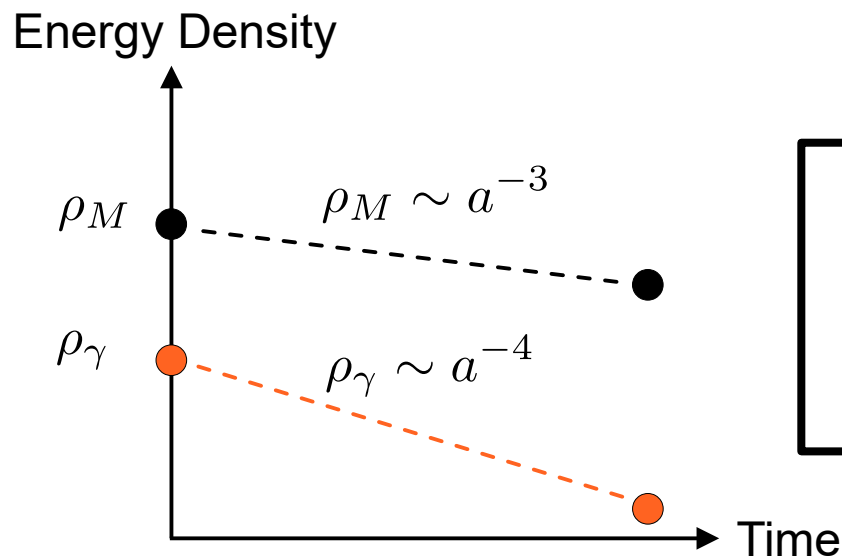
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Example: Matter ($w_M = 0$) and Radiation ($w_\gamma = 1/3$)



Boltzmann Equations

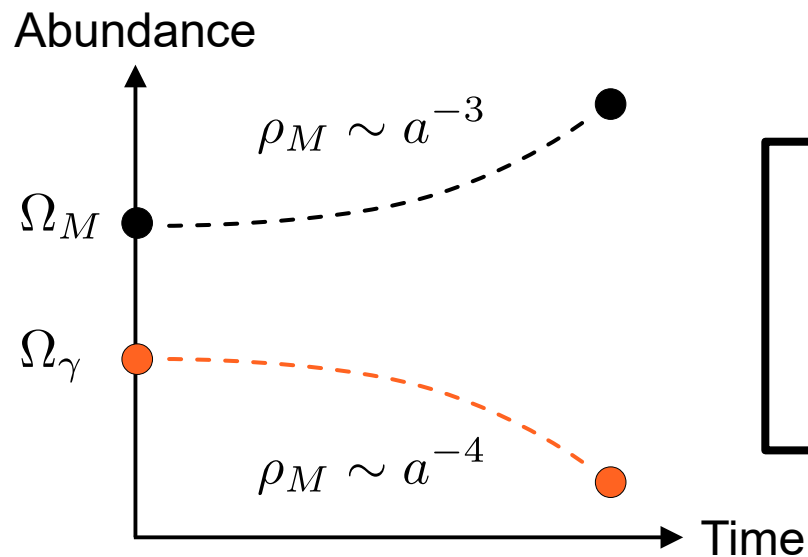
$$\frac{d\rho_M}{dt} = -3H\rho_M$$

$$\frac{d\rho_\gamma}{dt} = -4H\rho_\gamma$$

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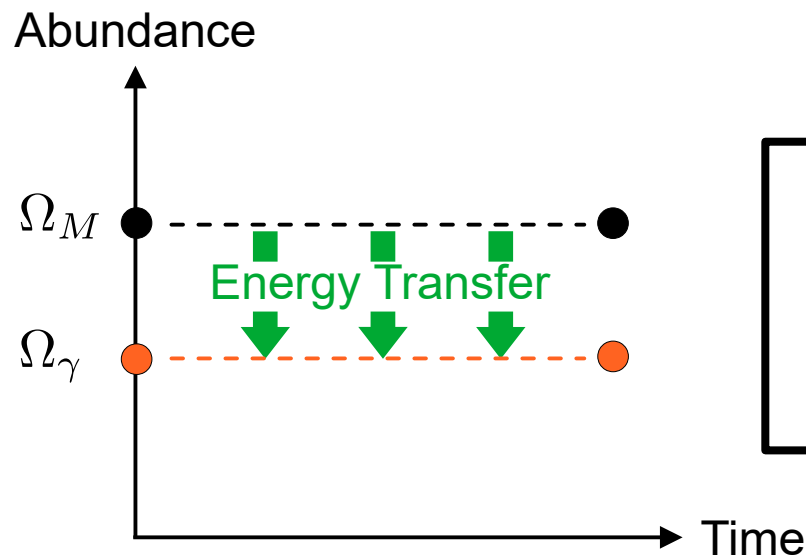
$$\frac{d\Omega_M}{dt} = H\Omega_M\Omega_\gamma$$

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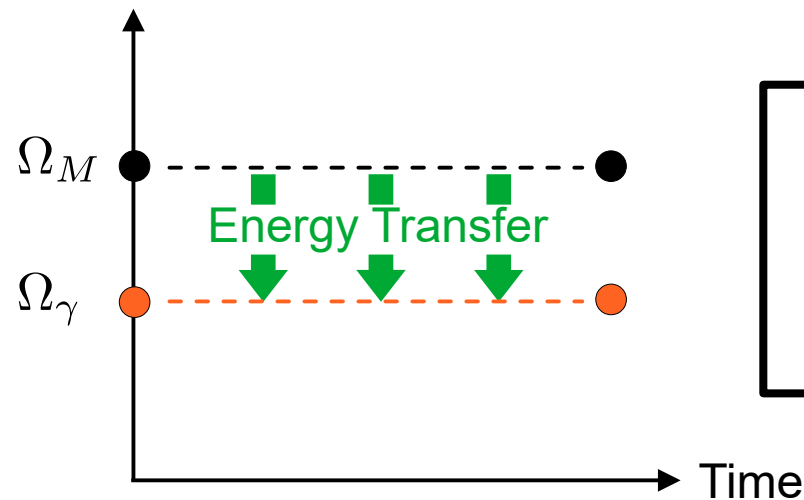
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Abundance



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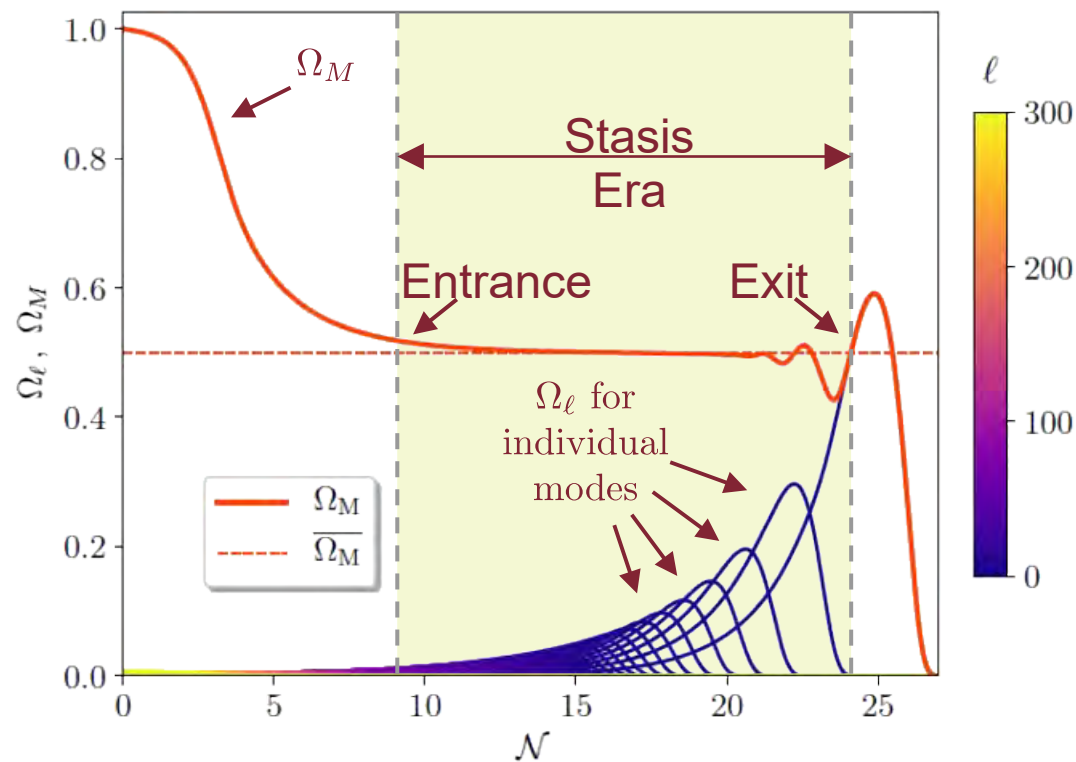
$\propto t^{-1}$, so $P(t)$ should be as well.

“Pump” from component with higher to lower w .

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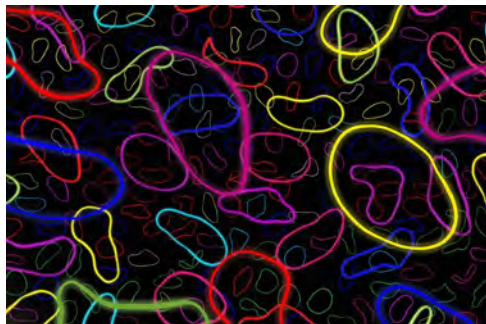


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- Pump terms with the right time-dependence for stasis emerge naturally in scenarios involving towers of states with broad spectra of masses, cosmological abundances, lifetimes, etc.

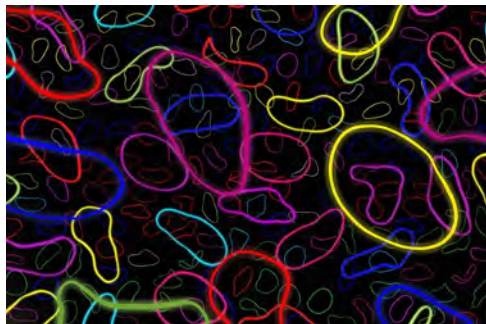
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- Such towers are a facet of numerous BSM-physics scenarios including...
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 - Theories with extra spacetime dimensions (KK towers)
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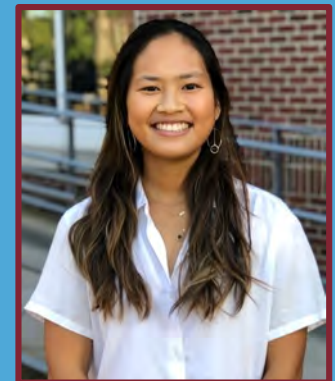
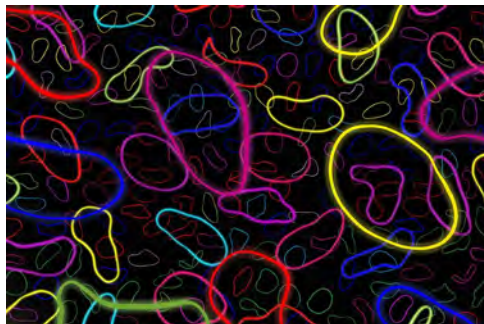
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- When they do emerge, stasis is typically a **global attractor**: the universe will evolve toward stasis regardless of initial conditions.
- The modified cosmological histories associated with stasis can affect the evolution of **scalar and tensor perturbations**.



See Anna Paulsen's talk
(directly after this one)

Stasis from Dynamical Scalar Fields

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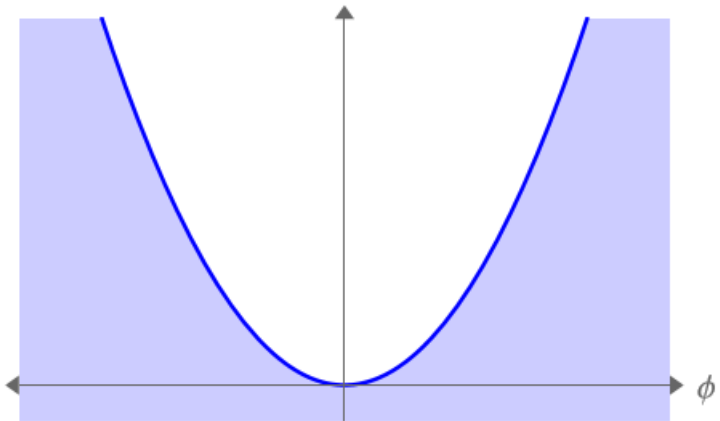
Is it possible to achieve a prolonged epoch of cosmological stasis from a tower of such scalars?

- Such a stasis, as we'll see, would be characterized by an effective equation-of-state parameter between that of vacuum energy ($w_\Lambda = -1$) and matter ($w_\Lambda = 0$)
- Moreover, stases involving dynamical scalars give rise to some phenomena not seen in other realizations of stasis which could potentially be useful for addressing fundamental questions in cosmology.

Warm-Up: A Single Scalar

- To set the stage, let's recall how the homogeneous zero-mode of a **single scalar field** ϕ of mass m with a quadratic potential $V(\phi)$ evolves in a flat FRW universe.

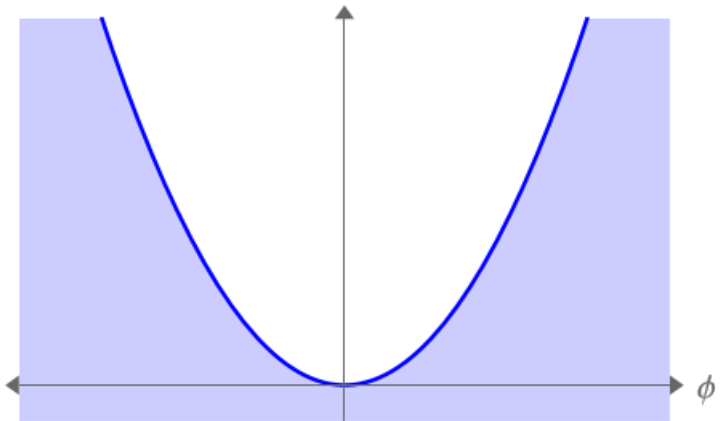
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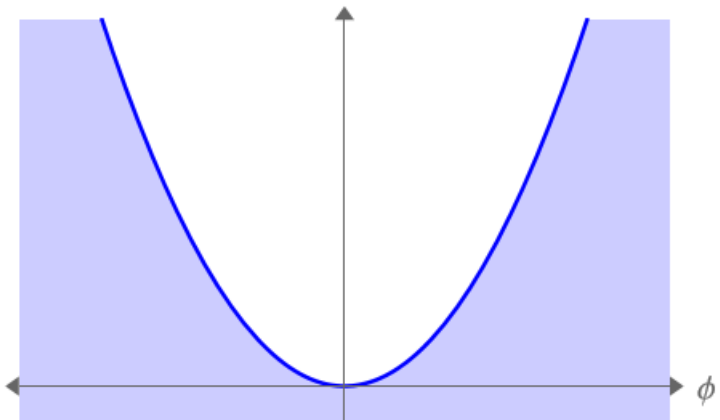
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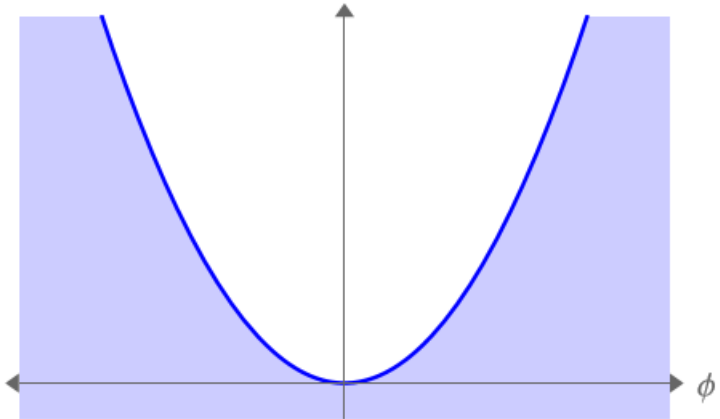
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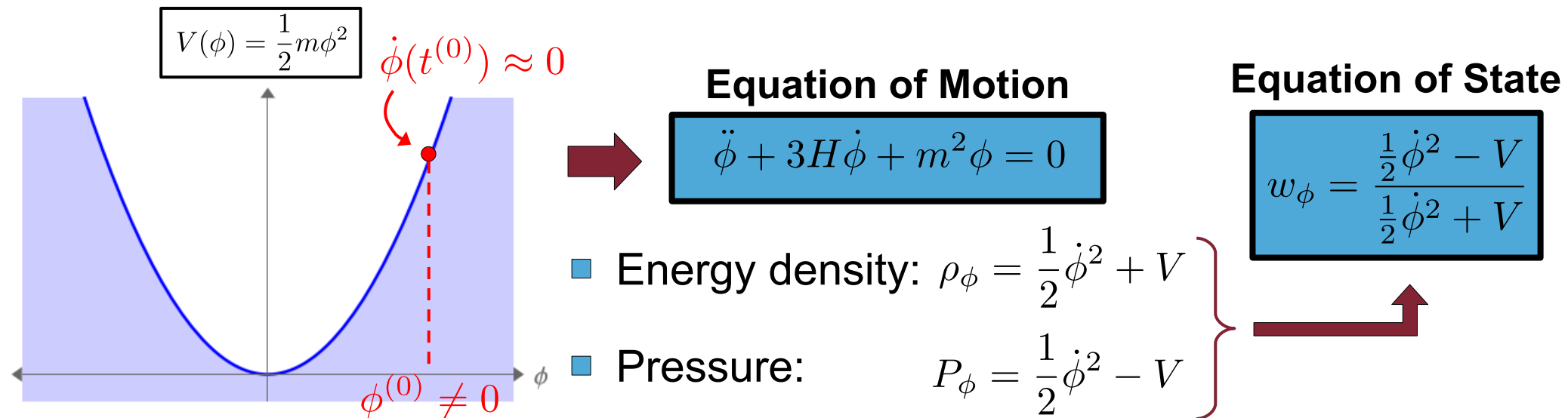
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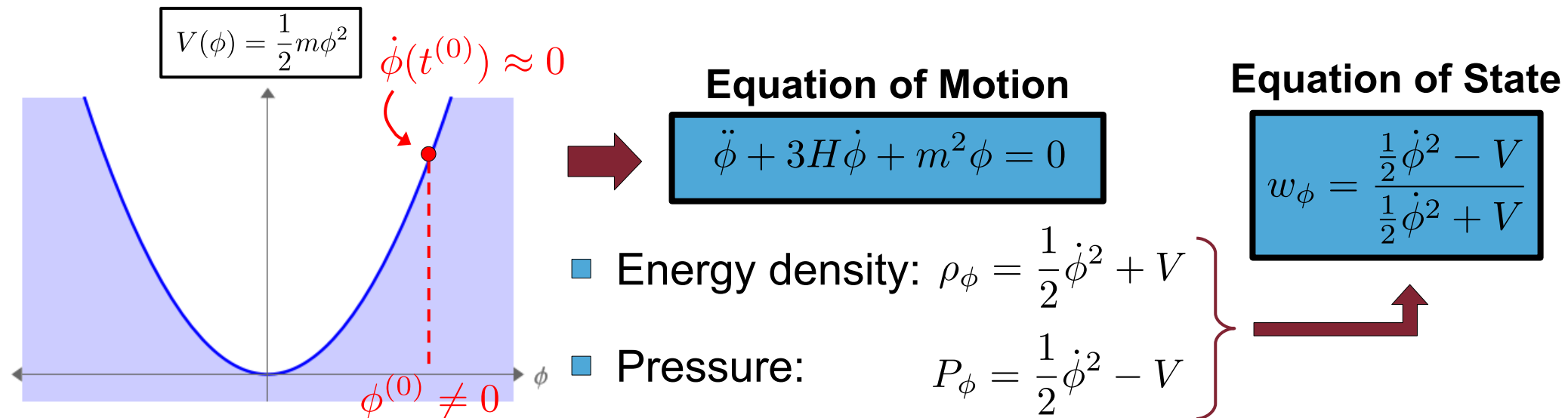
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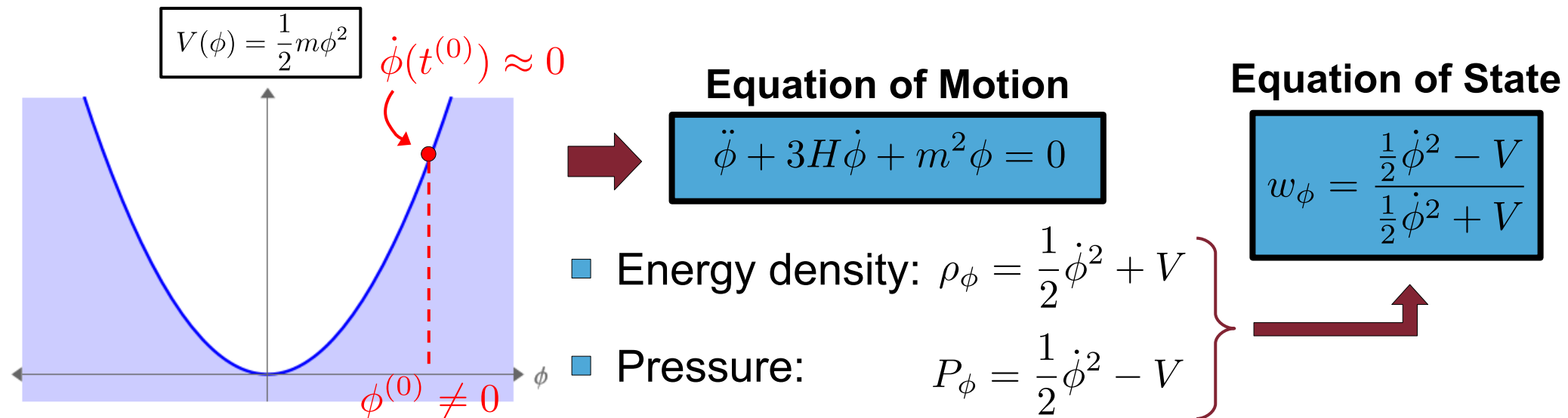
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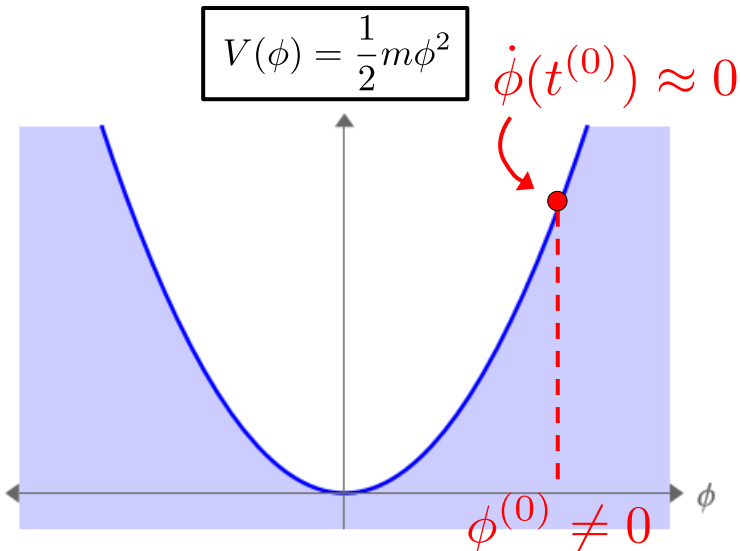


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$$H \approx \frac{\kappa}{3t}, \quad \text{where } \kappa \equiv \frac{2}{1+w}$$

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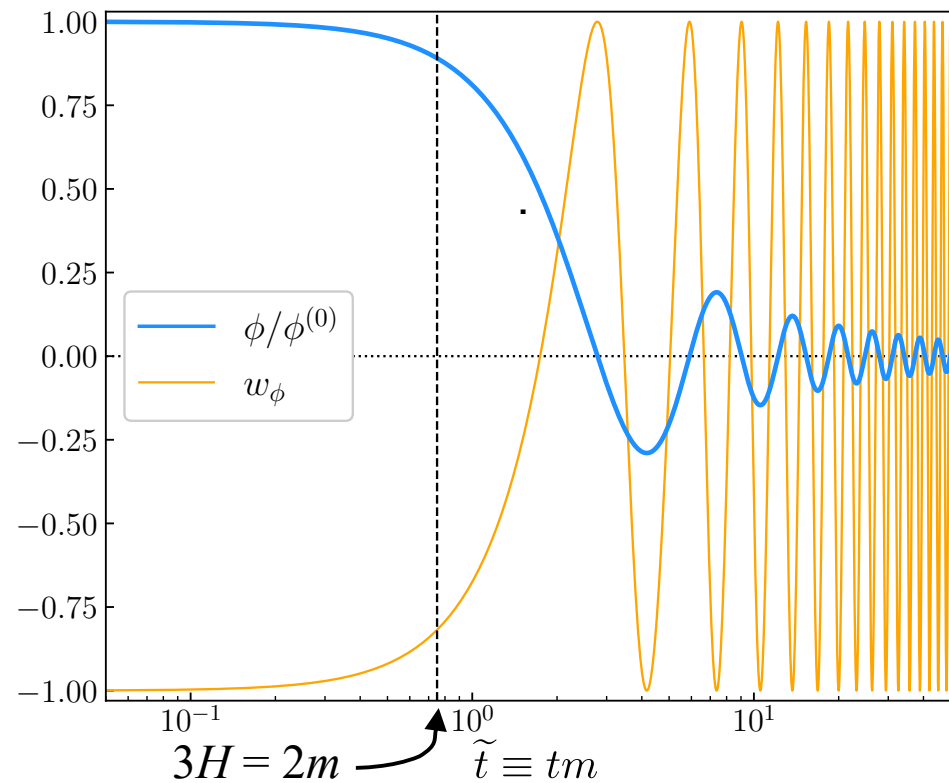
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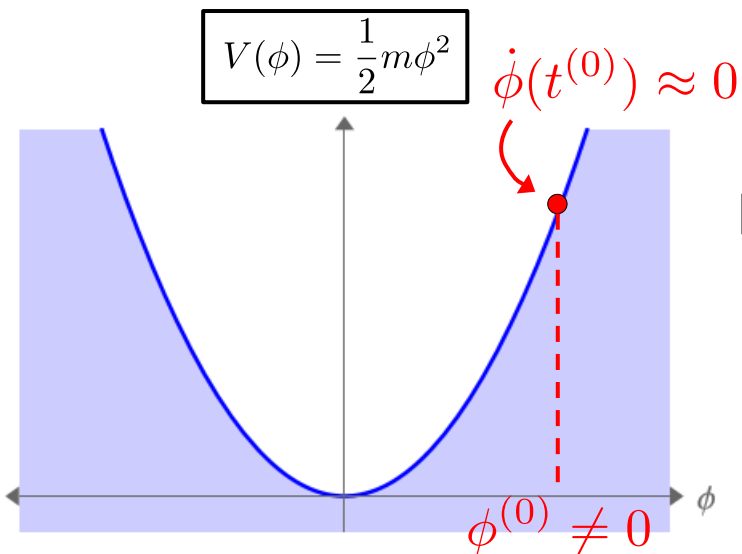
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Scalar in a Fixed Background ($w = 1/3$)



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Approximate Solution

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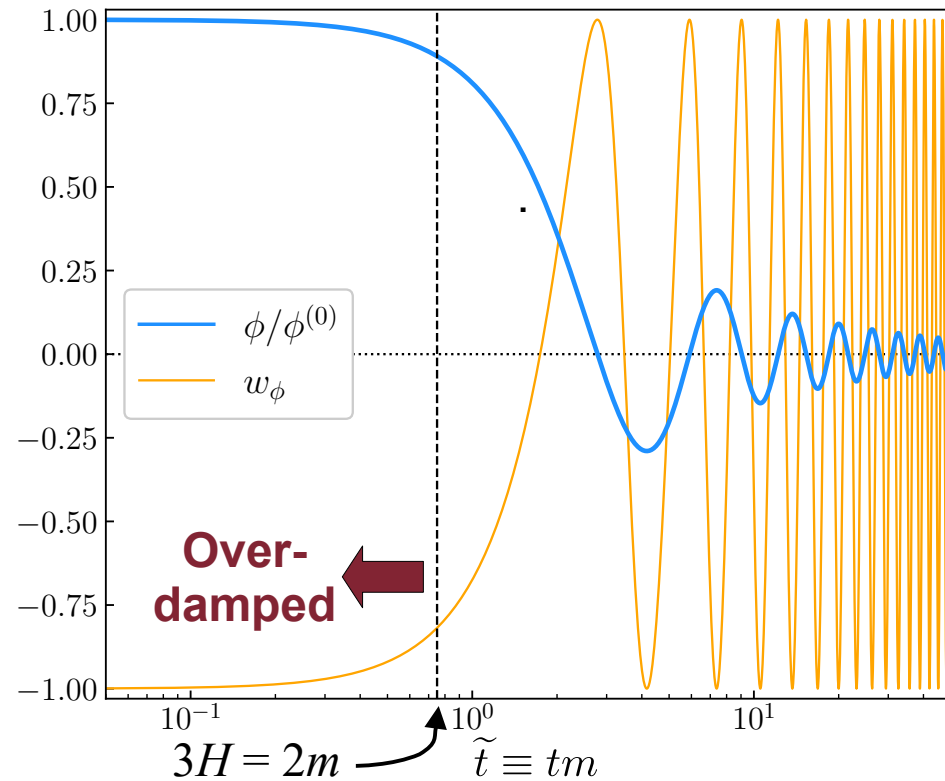
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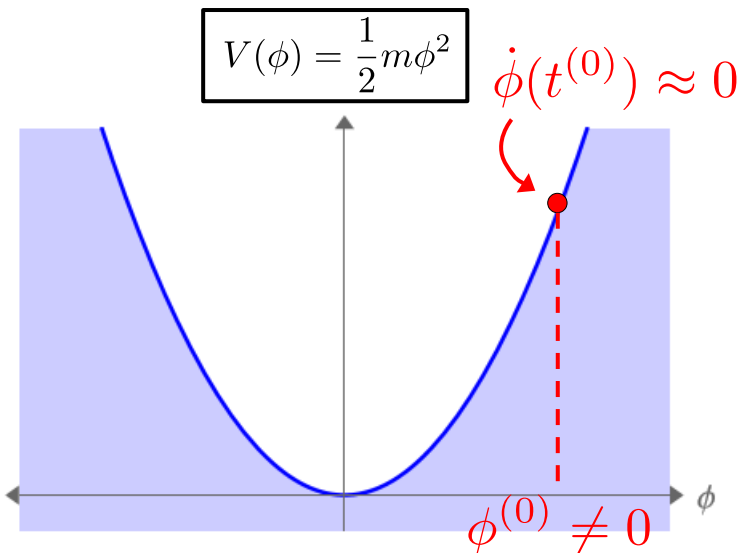
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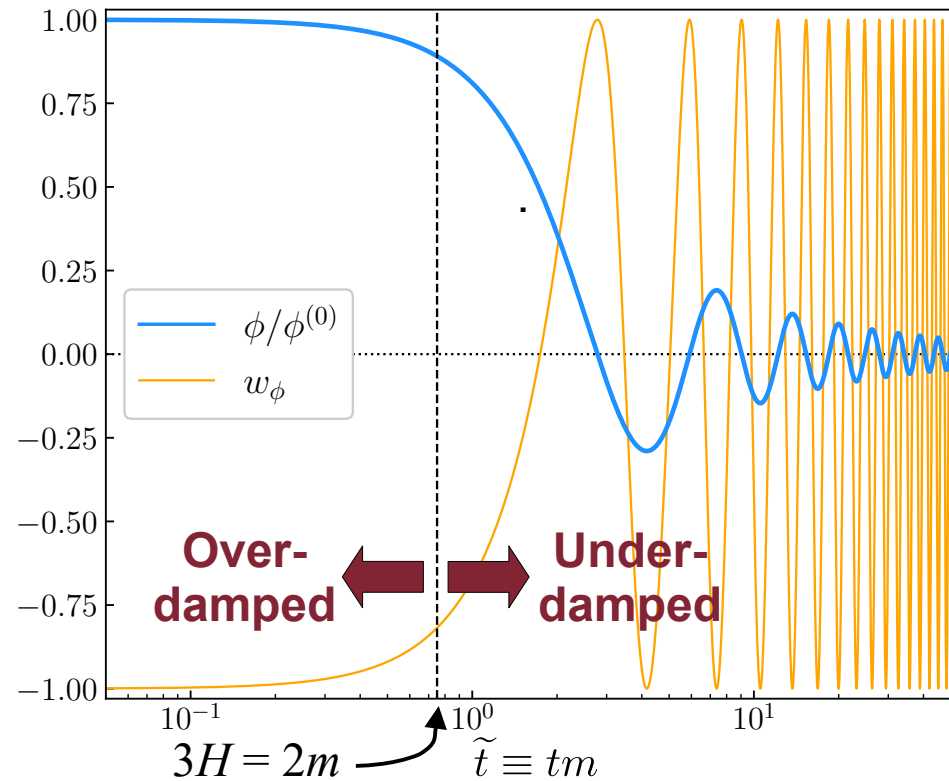
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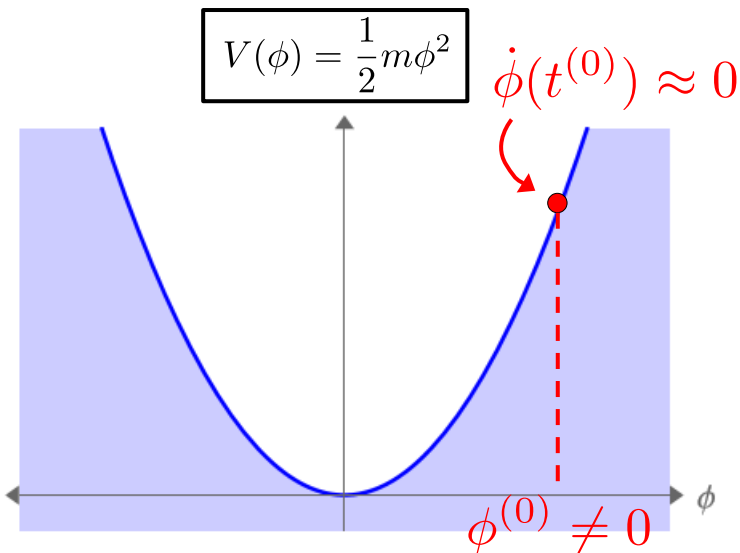
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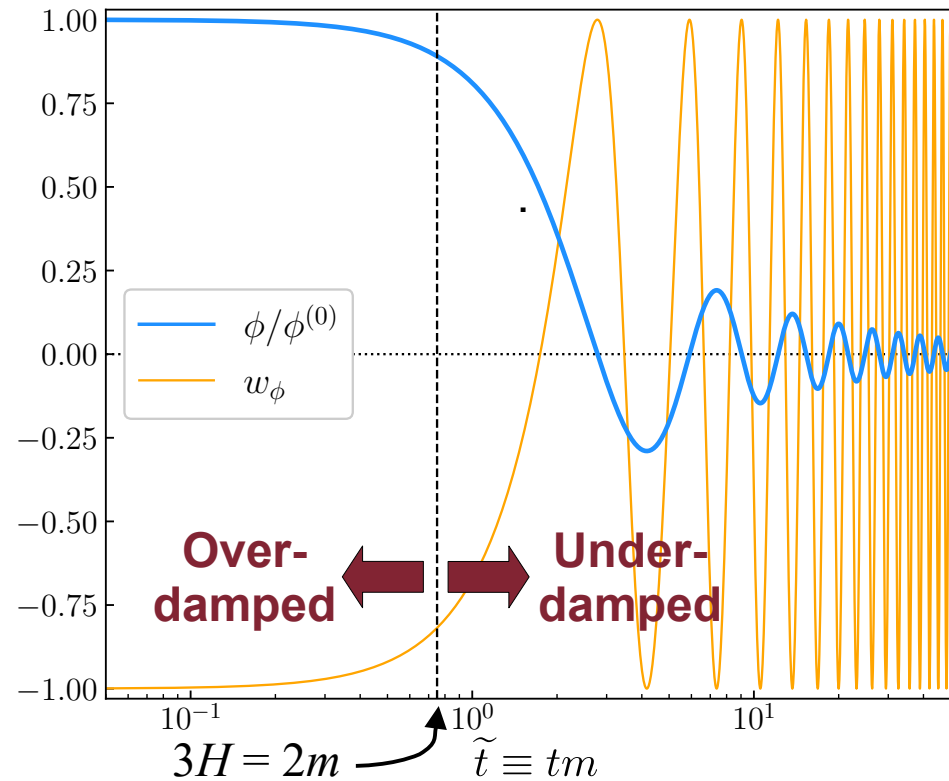
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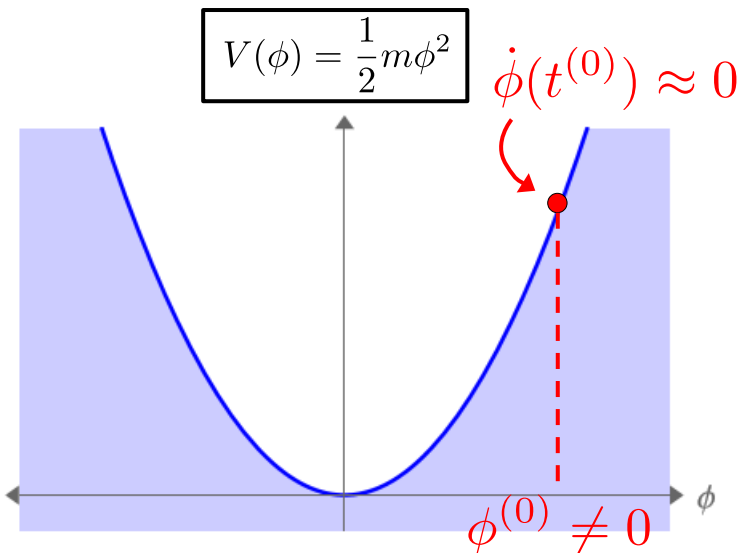
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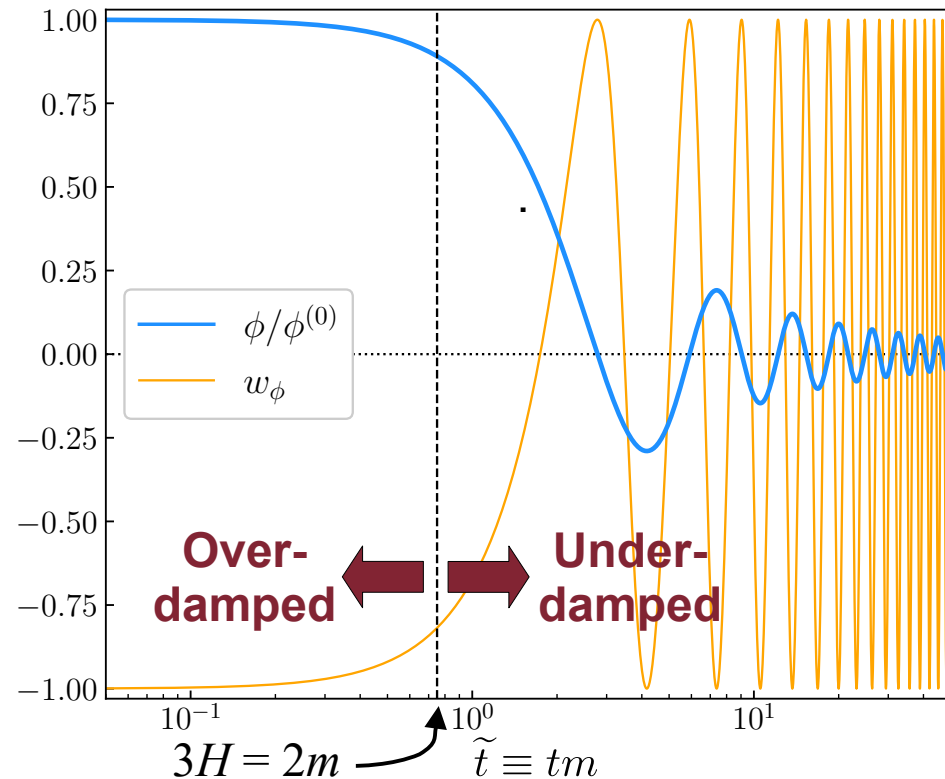
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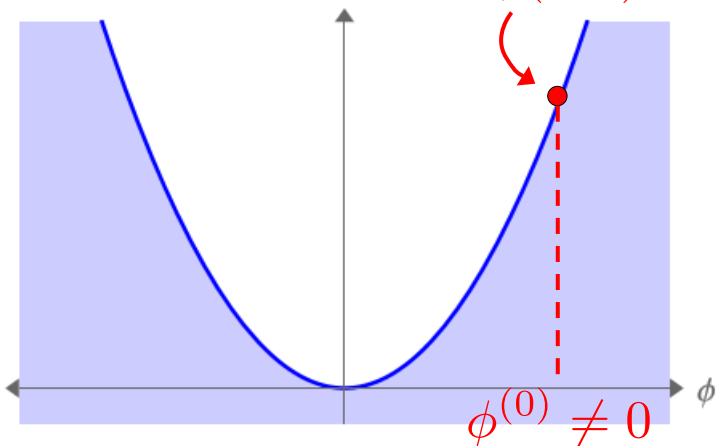
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➔ **Behaves like massive matter**

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Tower of Scalars

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Tower of Scalars

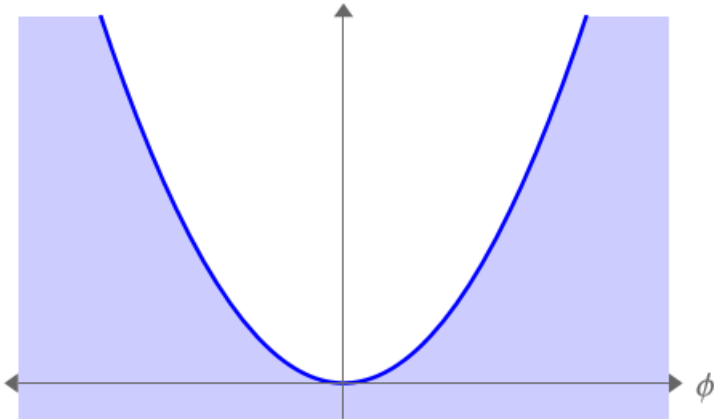
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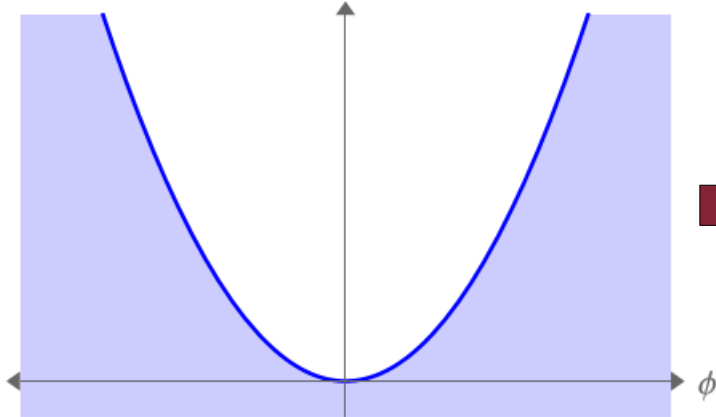
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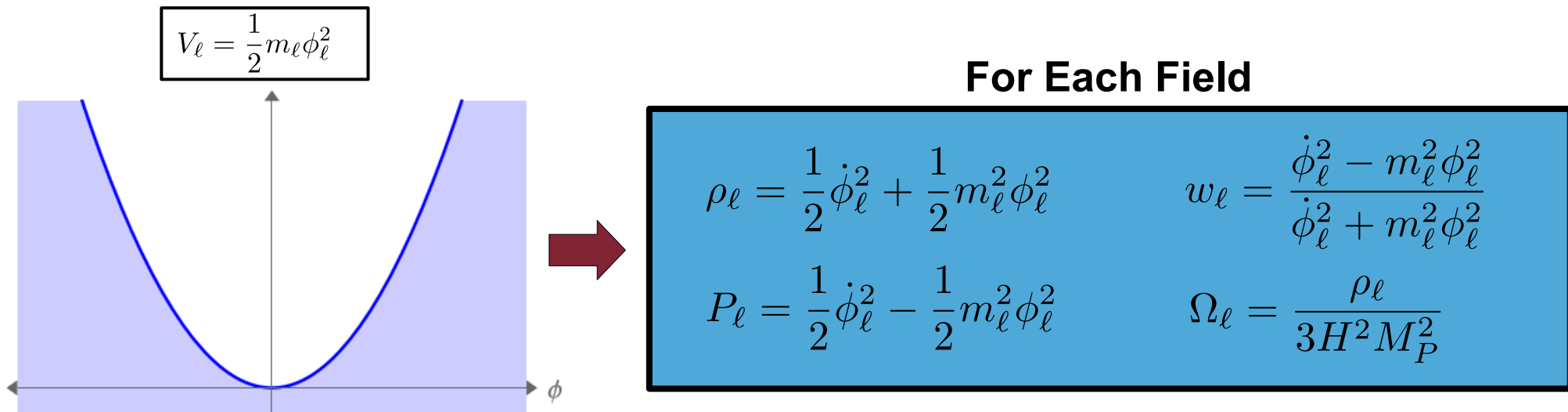
For Each Field

$$\rho_\ell = \frac{1}{2}\dot{\phi}_\ell^2 + \frac{1}{2}m_\ell^2\phi_\ell^2 \quad w_\ell = \frac{\dot{\phi}_\ell^2 - m_\ell^2\phi_\ell^2}{\dot{\phi}_\ell^2 + m_\ell^2\phi_\ell^2}$$

$$P_\ell = \frac{1}{2}\dot{\phi}_\ell^2 - \frac{1}{2}m_\ell^2\phi_\ell^2 \quad \Omega_\ell = \frac{\rho_\ell}{3H^2 M_P^2}$$

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- We'll also assume (for the moment) that there's **no background energy component**: the collective energy density of the ϕ_ℓ dominates the universe.

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- The system of (coupled) Boltzmann and Friedmann equations that describes the evolution of the ϕ_ℓ and H in this case is therefore...

N "copies" of this

$$\ddot{\phi}_\ell + 3H\dot{\phi}_\ell + m_\ell^2\phi_\ell = 0$$

$$H = \frac{1}{3M_P^2} \sum_{\ell=0}^{N-1} \rho_\ell$$

Tower of Scalars

- Now let's consider the case in which the universe comprises a **tower** of N such scalars ϕ_ℓ , where the index $\ell = 0, 1, 2, \dots, N-1$ labels these states in order of increasing mass.
- Consider a **mass spectrum** (motivated by KK towers): $m_\ell = m_0 + (\Delta m)\ell^\delta$
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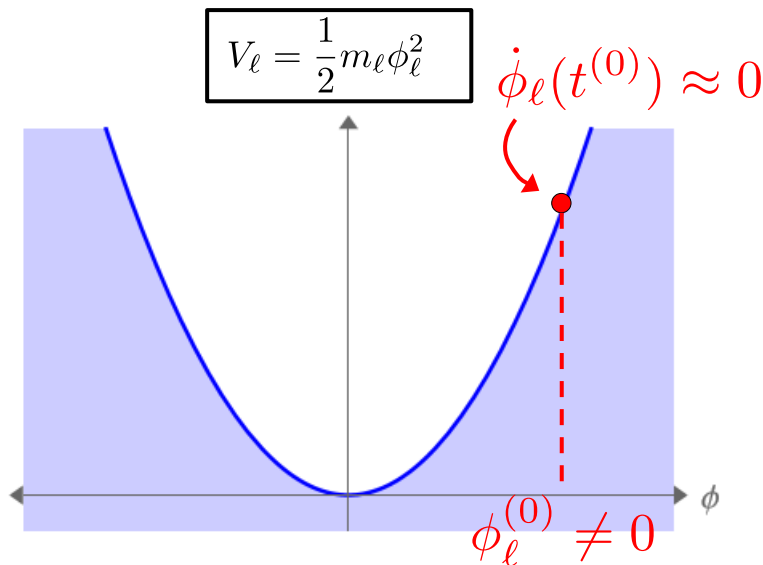
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Let's see what the cosmology of such a tower of scalar-field zero modes looks like!

Initial Conditions

- All that now remains is to specify the initial conditions for our scalars.
- For simplicity (and because it's consistent with many standard abundance-generation mechanisms for fields of this sort – e.g., vacuum misalignment), we once again take $\dot{\phi}_\ell(t^{(0)}) \approx 0$ for all of the ϕ_ℓ .
- However, we still need both an **overall mass scale** for the displacements and to know **how they scale** with ℓ across the tower.
- We assume a power-law scaling for the initial abundances of the form

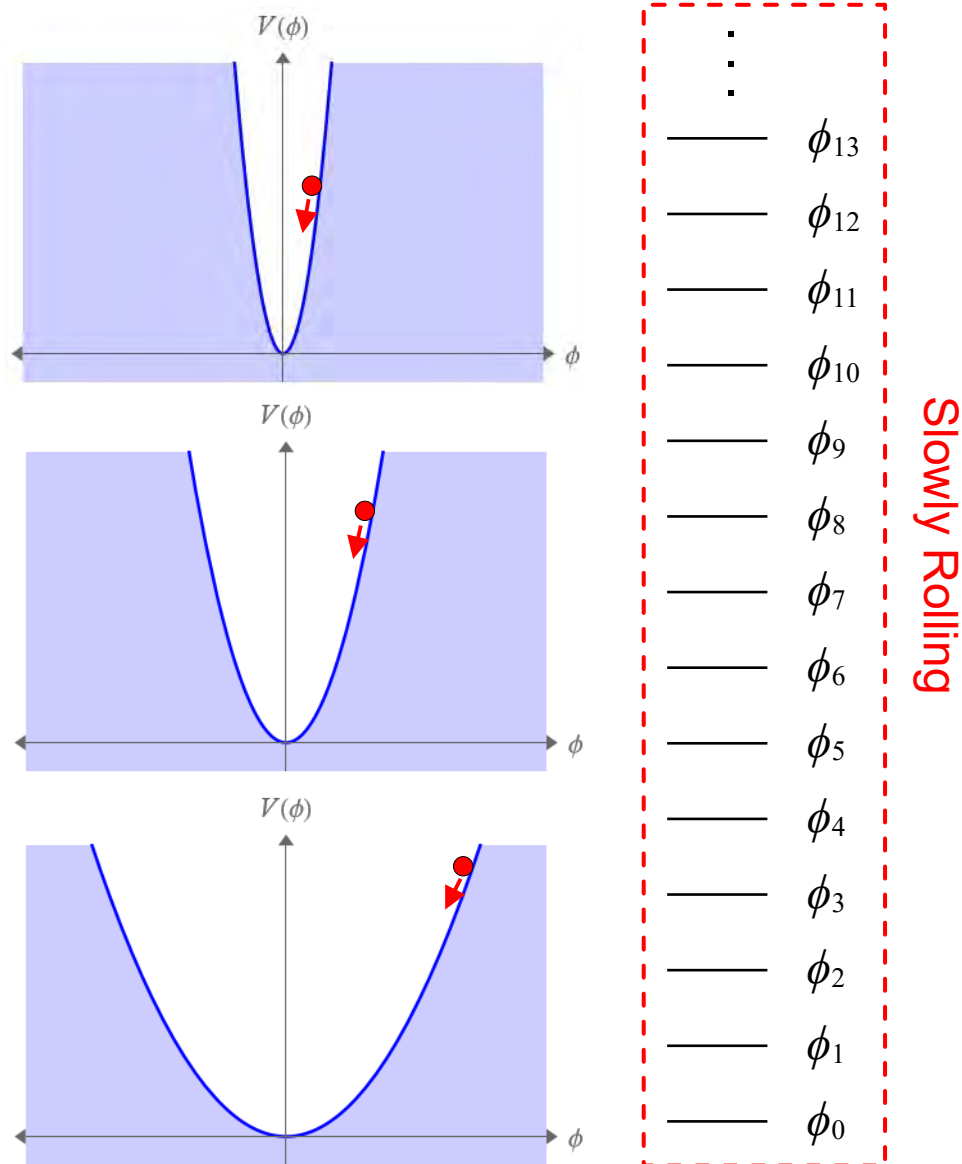
$$\Omega_\ell^{(0)} = \Omega_0^{(0)} \left(\frac{m_\ell}{m_0} \right)^\alpha$$



- For a given mass spectrum, the overall scale of the abundances can be parameterized by the ratio $\phi_0^{(0)}/M_P$, or, equivalently, by the ratio $H^{(0)}/m_{N-1}$.

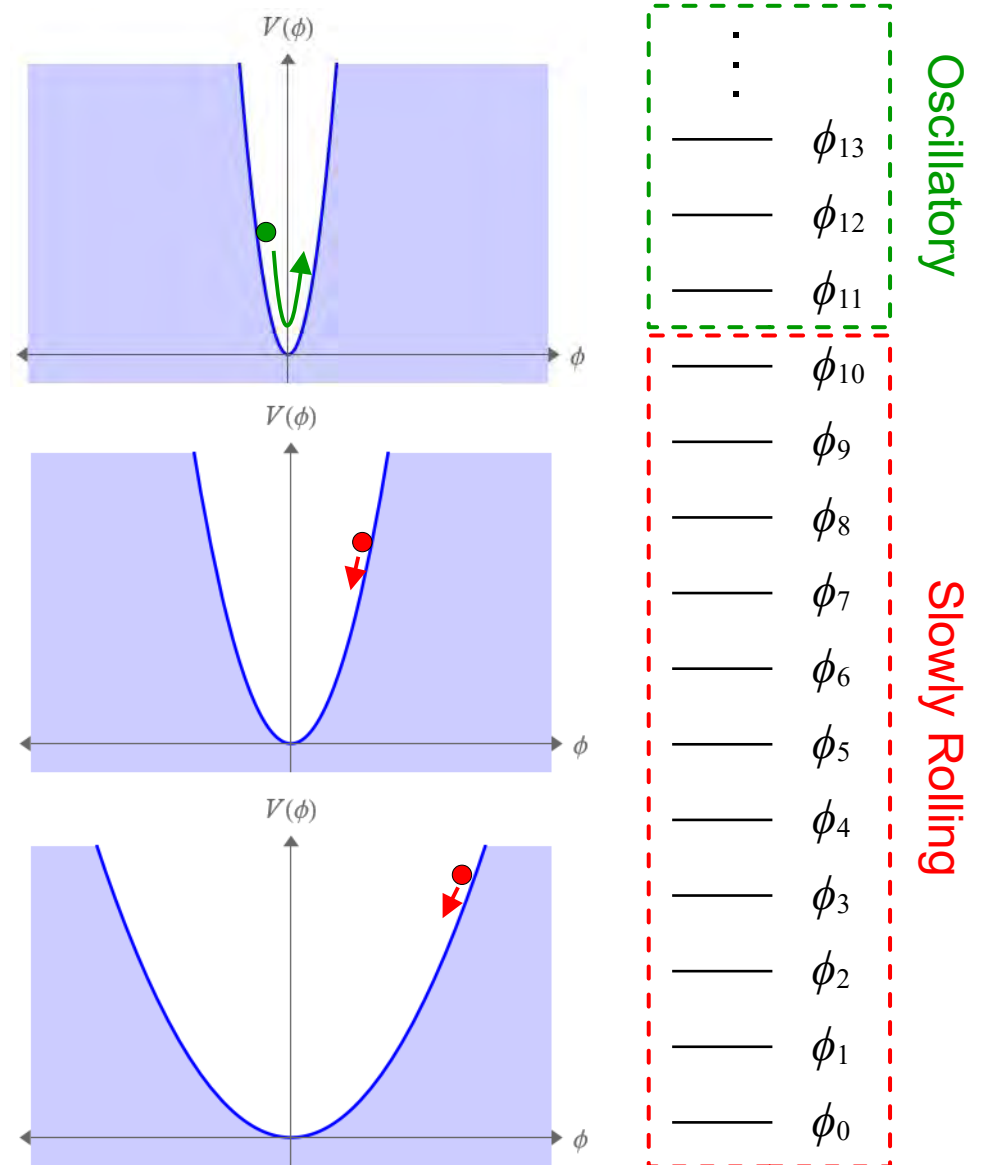
Dynamical Evolution

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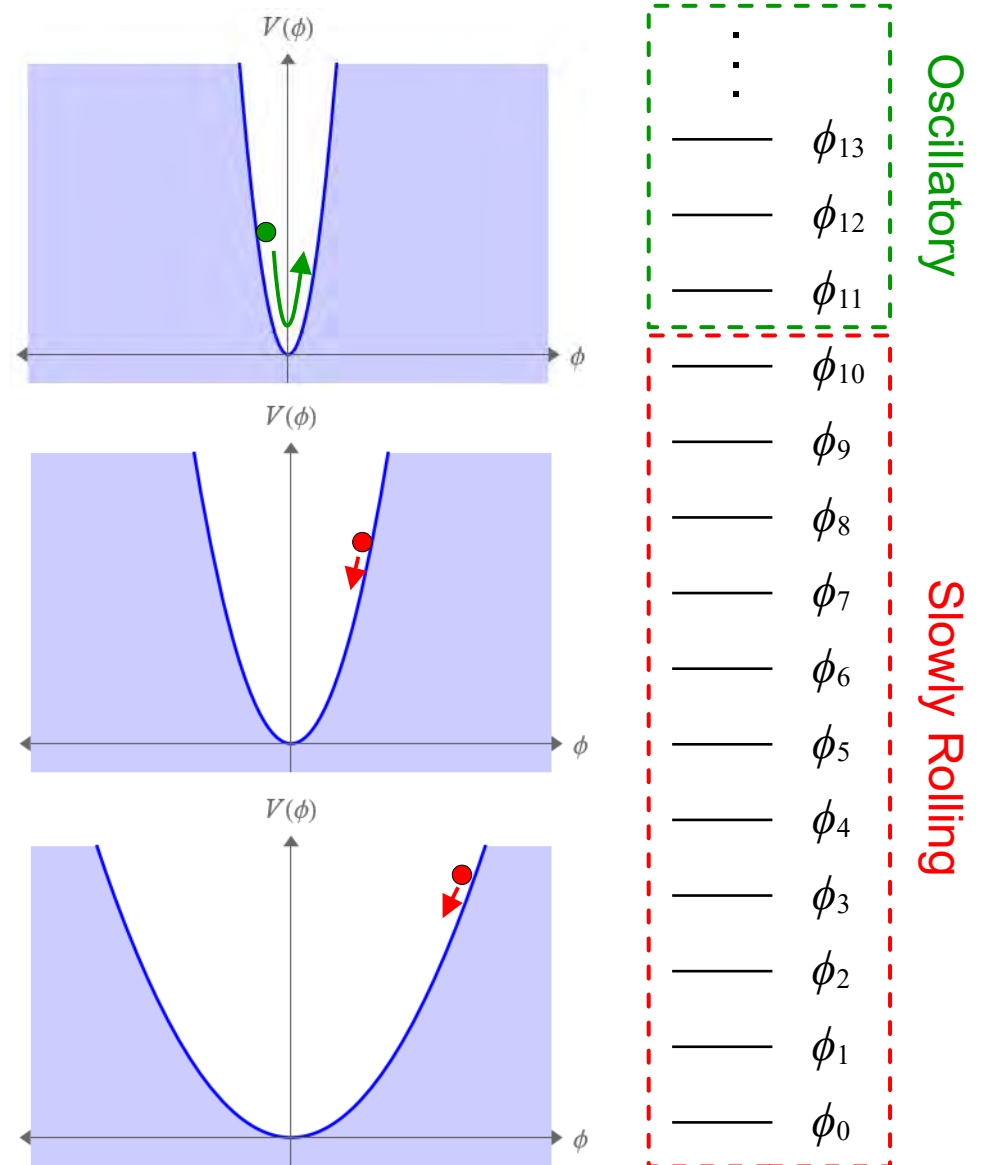
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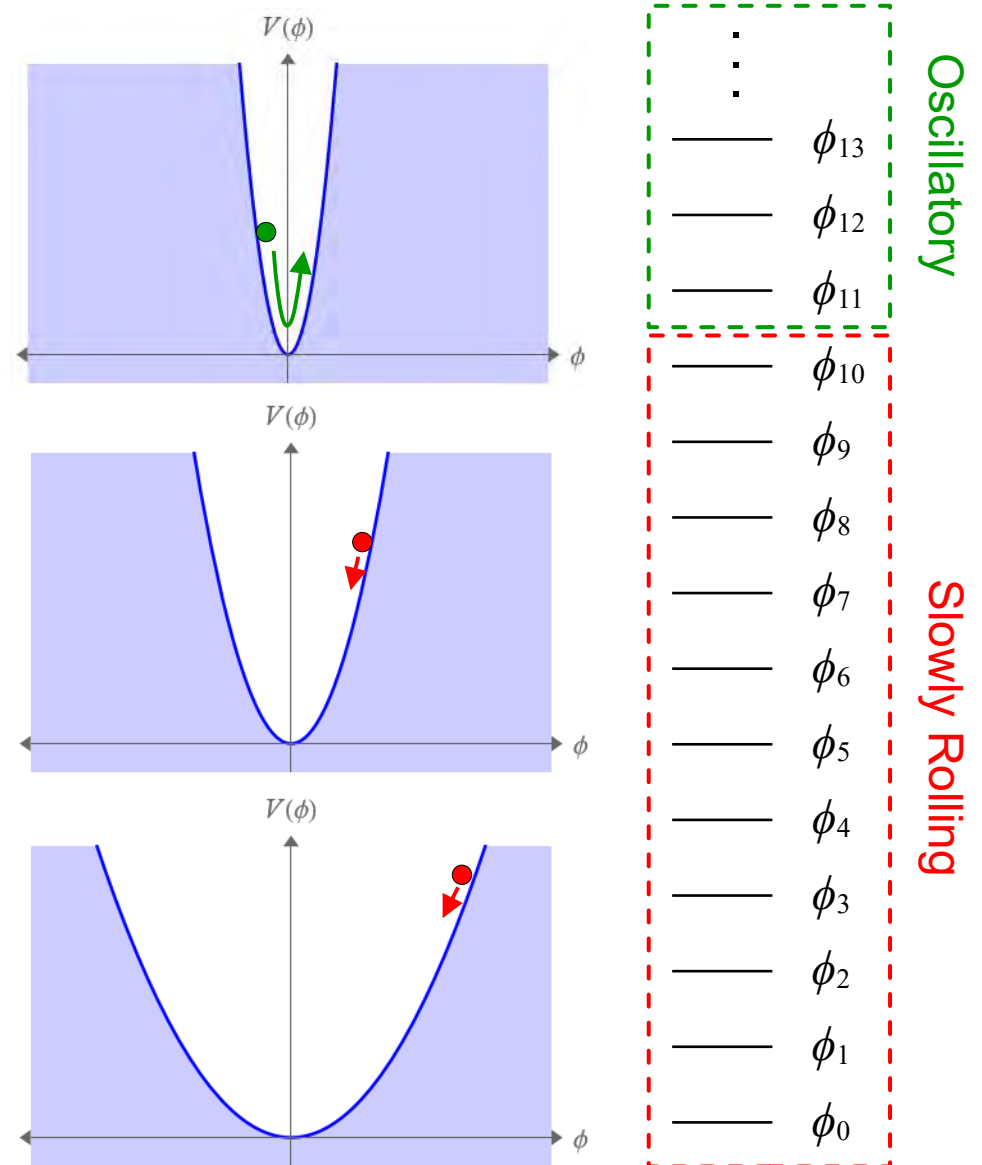
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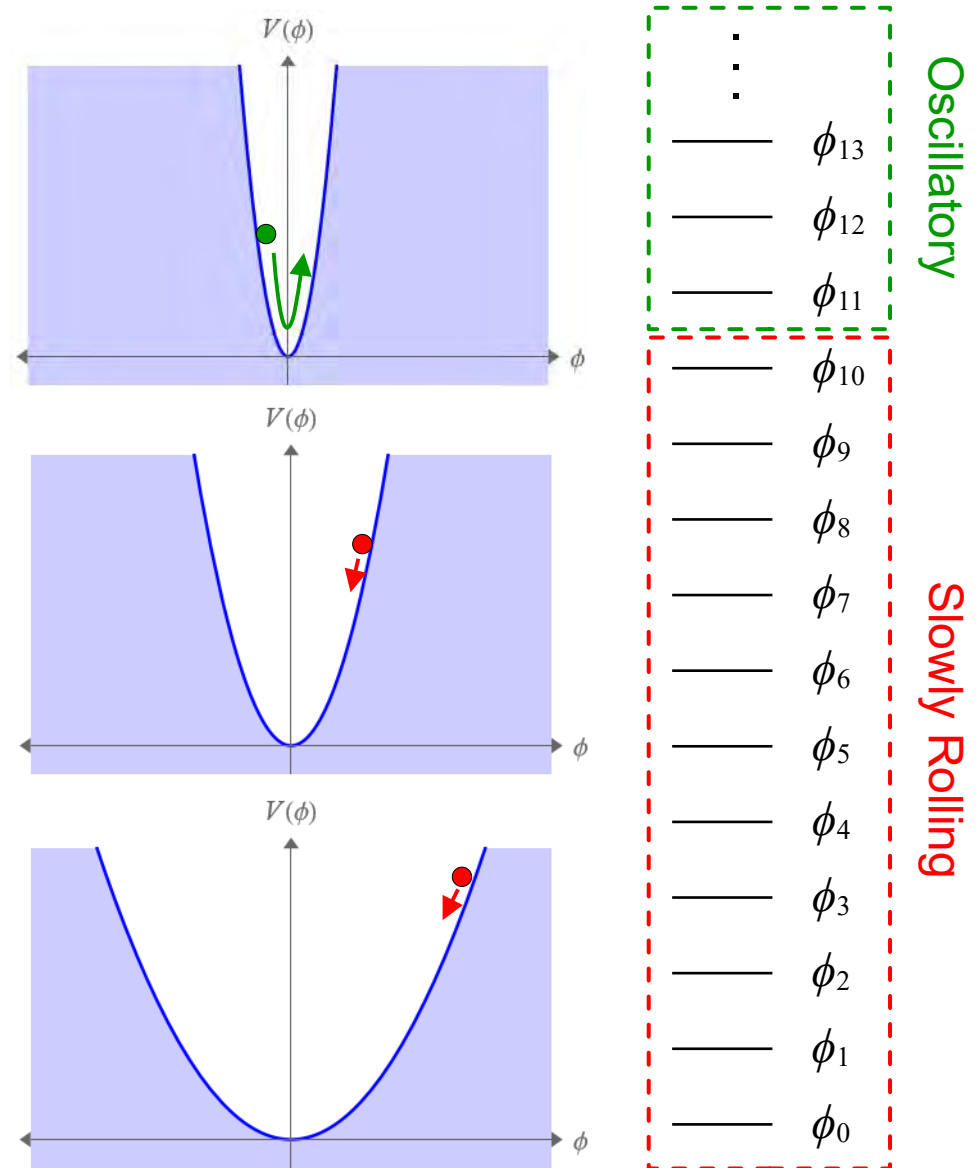
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states with $3H(t) \geq 2m_\ell$.

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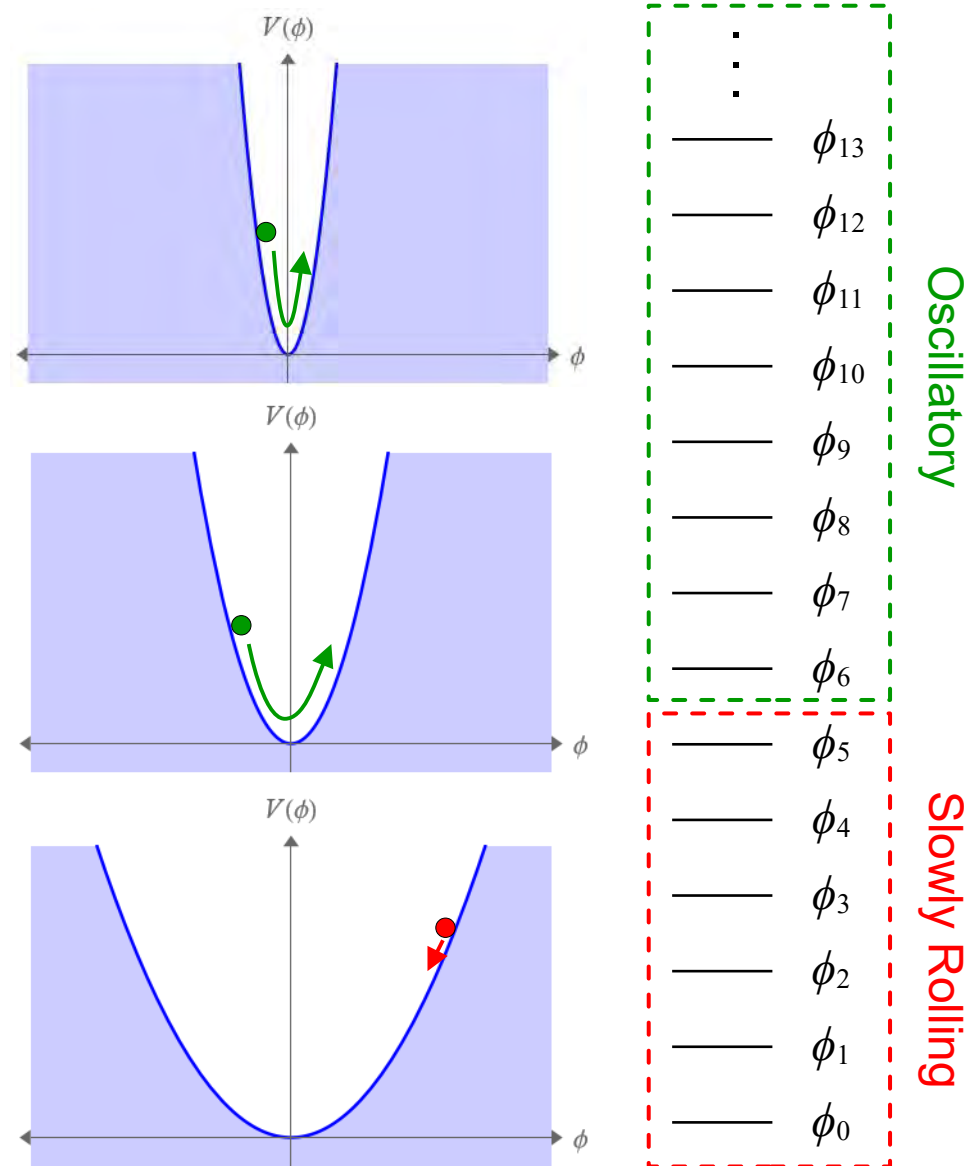
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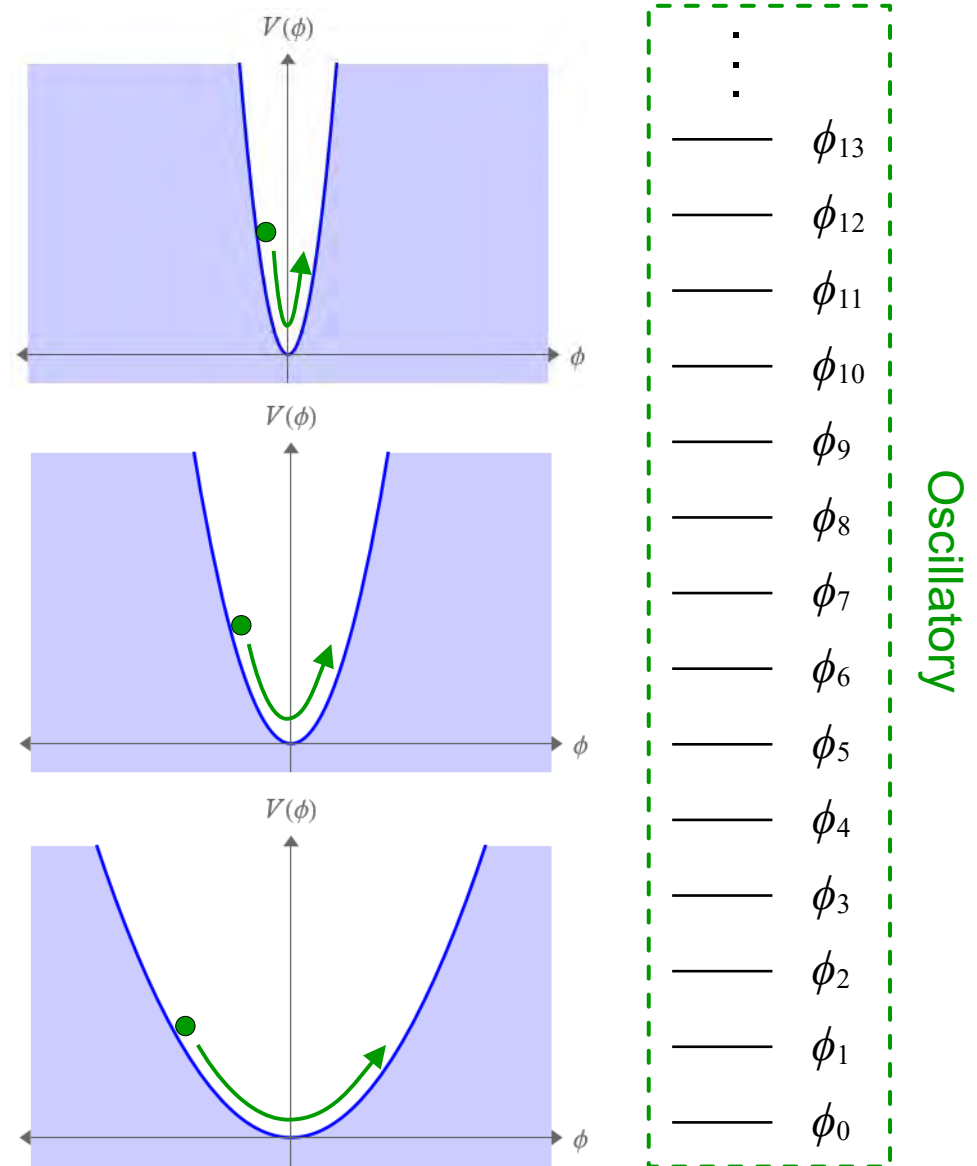
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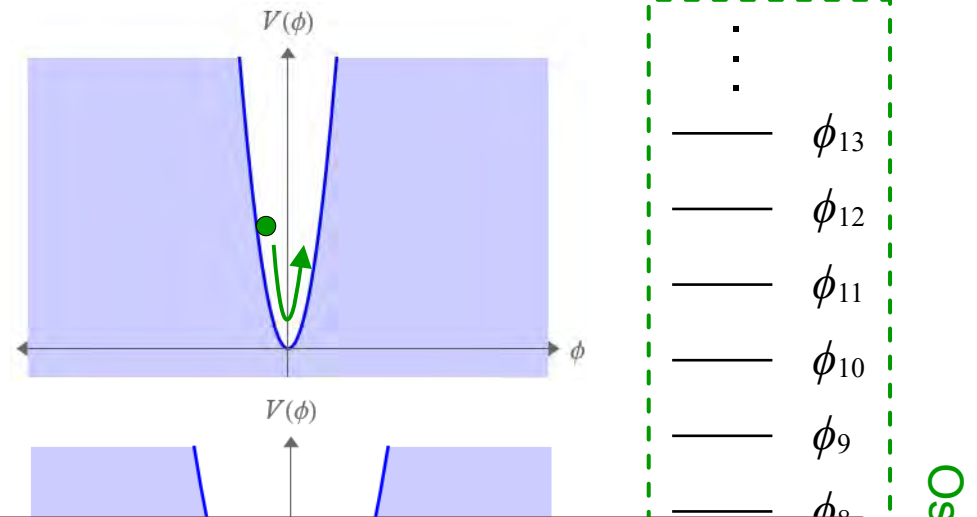
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The Question:

Can we achieve a stasis between these slow-roll and oscillatory cosmological energy components, which act like vacuum energy and matter, respectively?

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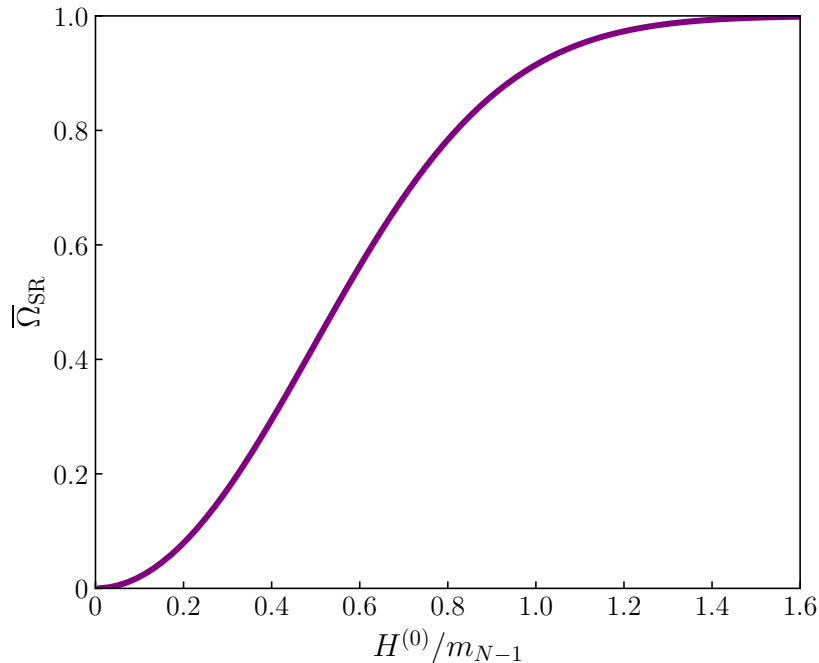
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Towers which satisfy this relation give rise to stasis. For $\delta = 1$, this corresponds to $\phi_\ell^{(0)} \sim \ell^{-1/2}$

Effect of Initial Conditions

- Unlike in previous realizations of stasis, the stasis abundances $\bar{\Omega}_{\text{SR}}$ and $\bar{\Omega}_{\text{SR}}$ depend on the **initial conditions** for the scalar tower.
- In particular, $\bar{\Omega}_{\text{SR}}$ and $\bar{\Omega}_{\text{SR}}$ are sensitive to the ratio $\phi_0^{(0)}/M_P$ which parametrizes the overall scale of the initial zero-mode displacements.

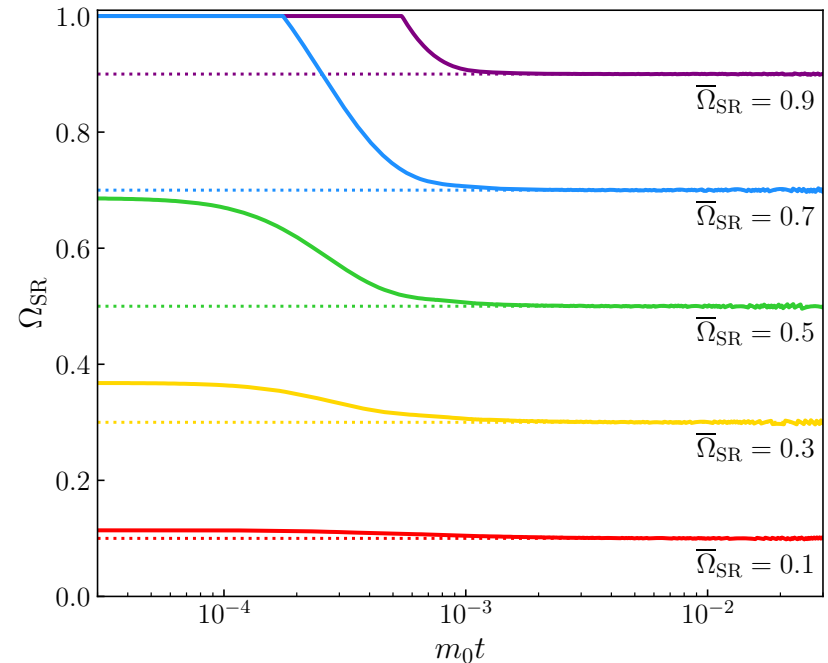
Impact of Initial Conditions



Duration of Stasis

$$\mathcal{N}_s \approx \frac{\bar{\kappa}}{3} \left[\delta \log N + \log \left(\frac{\Delta m}{m_0} \right) + \log \left(\frac{3H^{(0)}}{2m_{N-1}} \right) \right]$$

Evolution Toward Stasis



$$\alpha = 1$$

$$\delta = 1$$

$$N = 5000$$

Background Components and Tracking

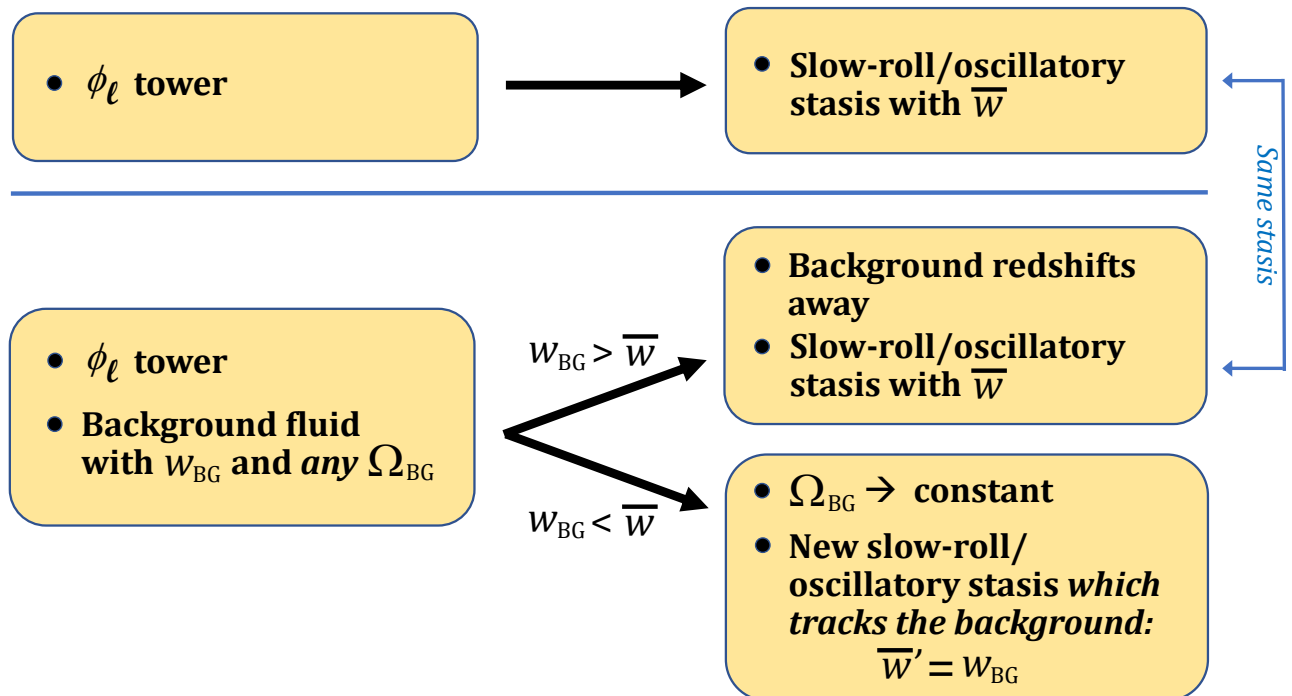
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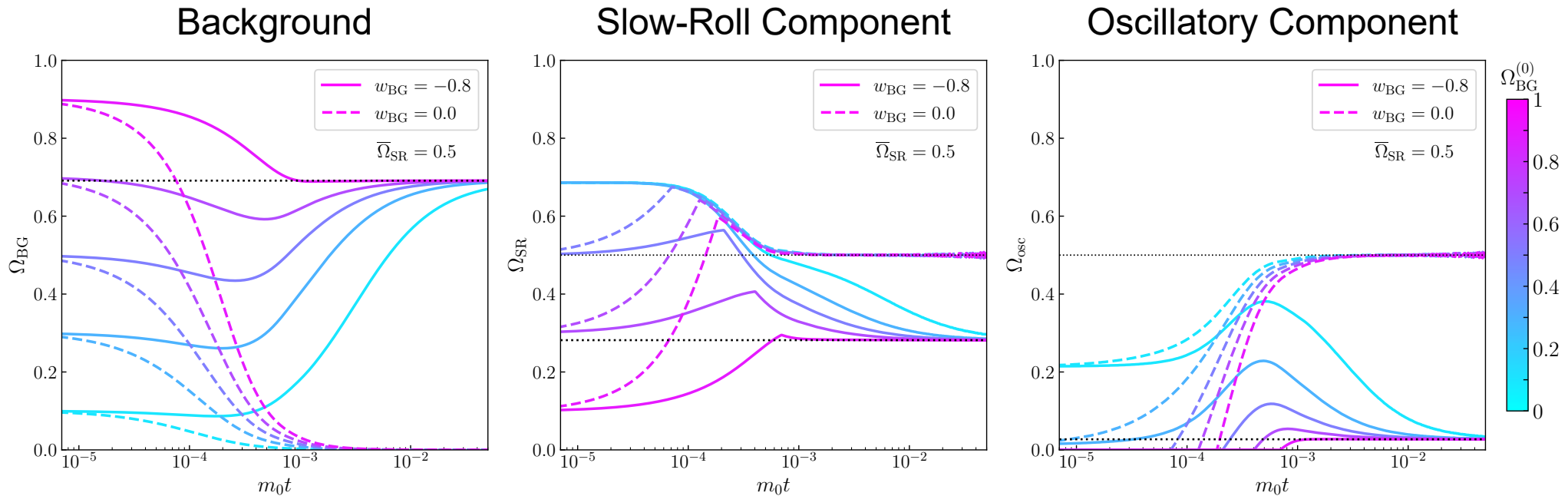
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- It turns out that in the presence of such an energy component, the universe still evolves toward stasis (or something like it).
- However, the outcome depends on the relationship between w_{BG} and the equation-of-state parameter \bar{w} the tower *would* have had during stasis if the background component weren't present.



Background Components and Tracking

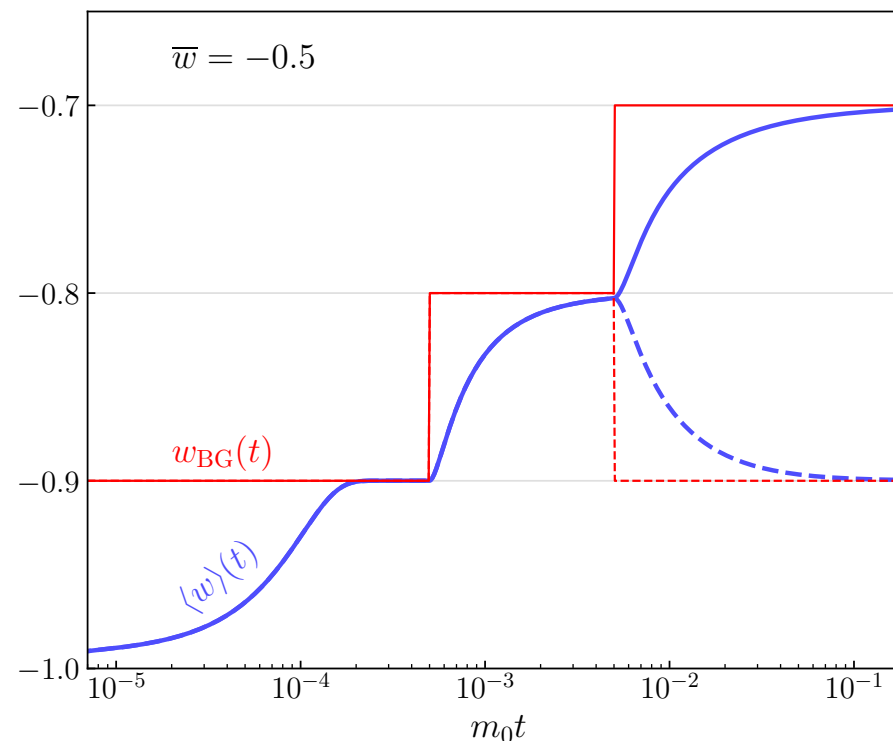
- The **tracking phenomenon** which arises in $w_{BG} < w$ has not been observed in other realizations of stasis.



- These results provide insight about how the universe might enter into – or exit from – an stasis epoch involving dynamical scalars.

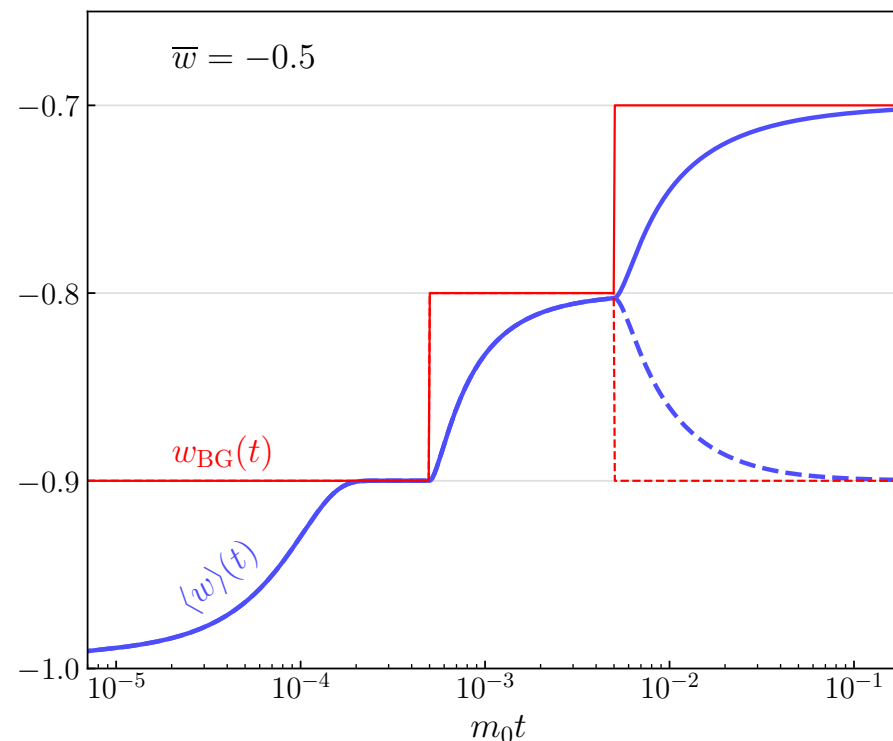
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- Indeed, as long as w_{BG} remains below \bar{w} , the tower's equation-of-state parameter $\langle w \rangle$ continues to evolve toward the new value of w_{BG} after the shift, regardless of whether this shift is positive or negative.



Summary

- **Stable, mixed-component cosmological eras** – i.e. **stasis eras** – are indeed a viable cosmological possibility – and one that can arise naturally in many extensions of the Standard Model.
- A **tower of scalar fields** which undergo a transition from overdamped to underdamped evolution can give rise to stasis.
- Stasis itself is an **attractor** in these systems, but several fundamental characteristics of the stasis epoch toward which the universe evolves depend on the initial conditions.
- In the presence of an additional background component with equation-of-state parameter w_{BG} , the tower exhibits a **tracking behavior** in which its own equation-of-state parameter evolves toward w_{BG} .

