

# Energy Correlators, Heavy Flavor, & Precision QCD

**Evan Craft — Yale University DPF-Pheno 2024** 









Based on work with K. Lee, B. Mecaj, I. Moult, & M. Gonzalez

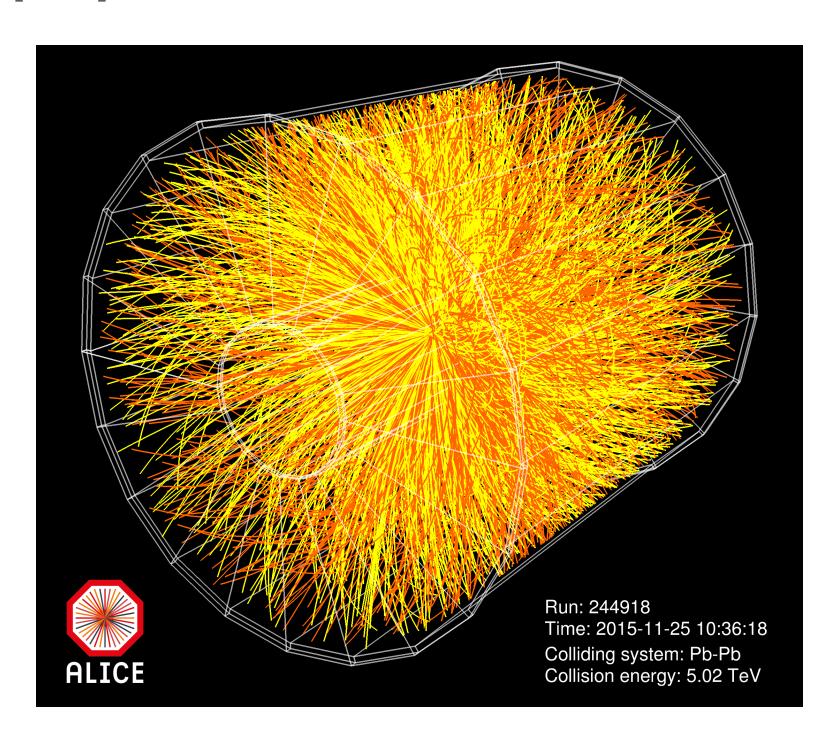
Yale University

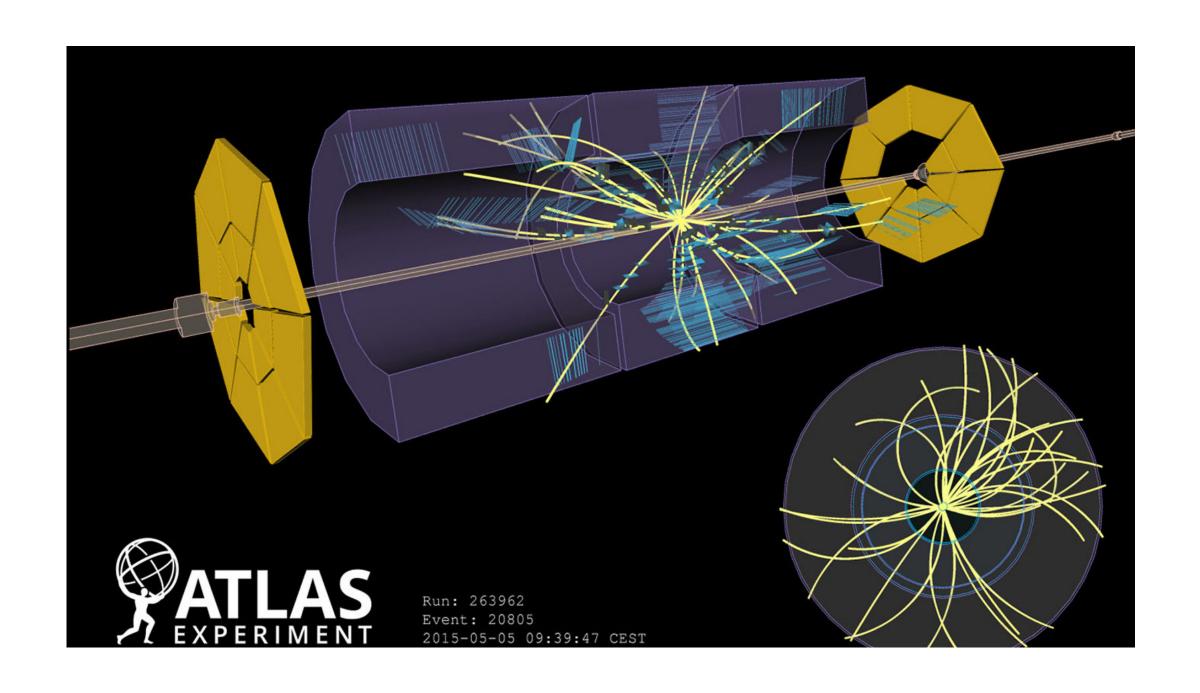


## Collider Experiments

Many important questions have been addressed at collider experiments

→ Great historical success in verifying properties of the standard model





- → But the detailed structure of QCD produces immensely complicated datasets.
- → Need new tools for future success

A unique frontier for novel collaborations between both theory and experiment

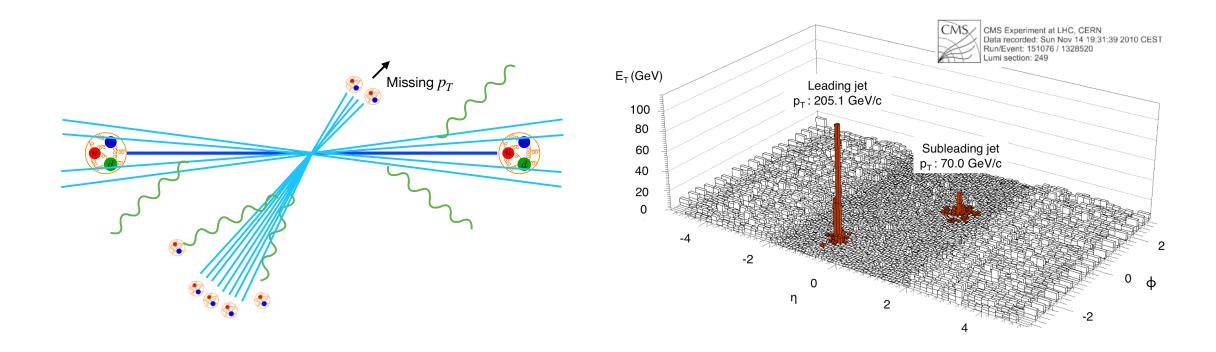
## How do we study collisions?

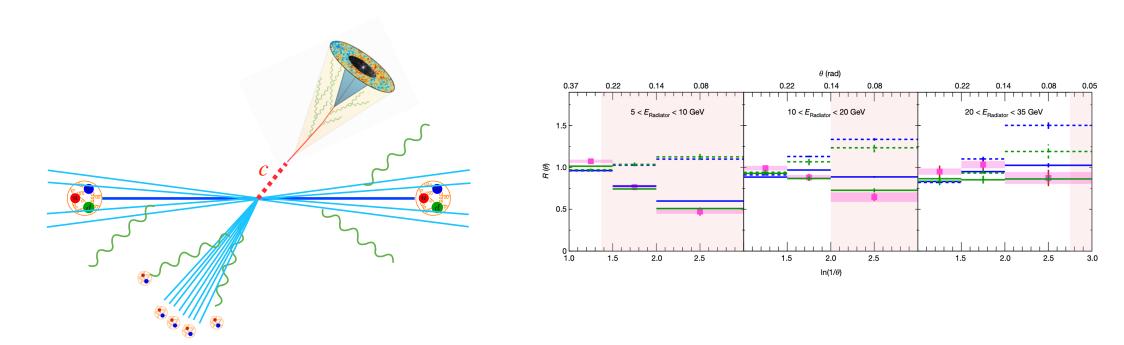
## Along with many more exciting observations!

Several remarkable observations made by studying jets

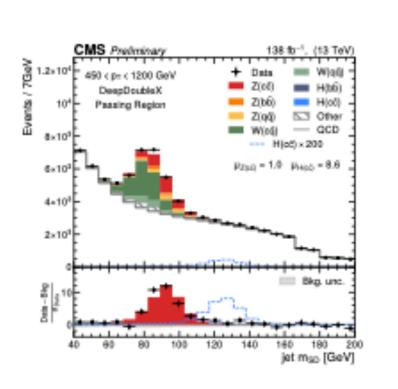
→ provided initial evidence for the existence of the Quark Gluon Plasma

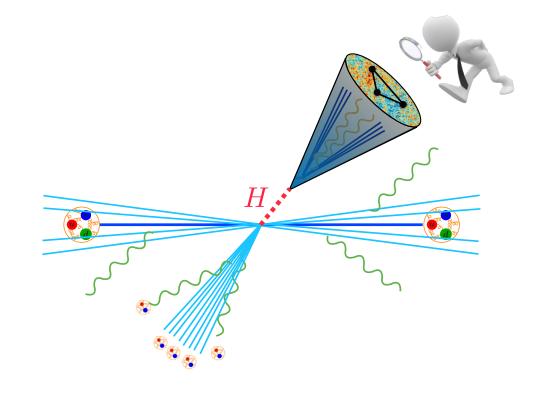
→ allowed first ever observation of the "dead cone" effect of QCD



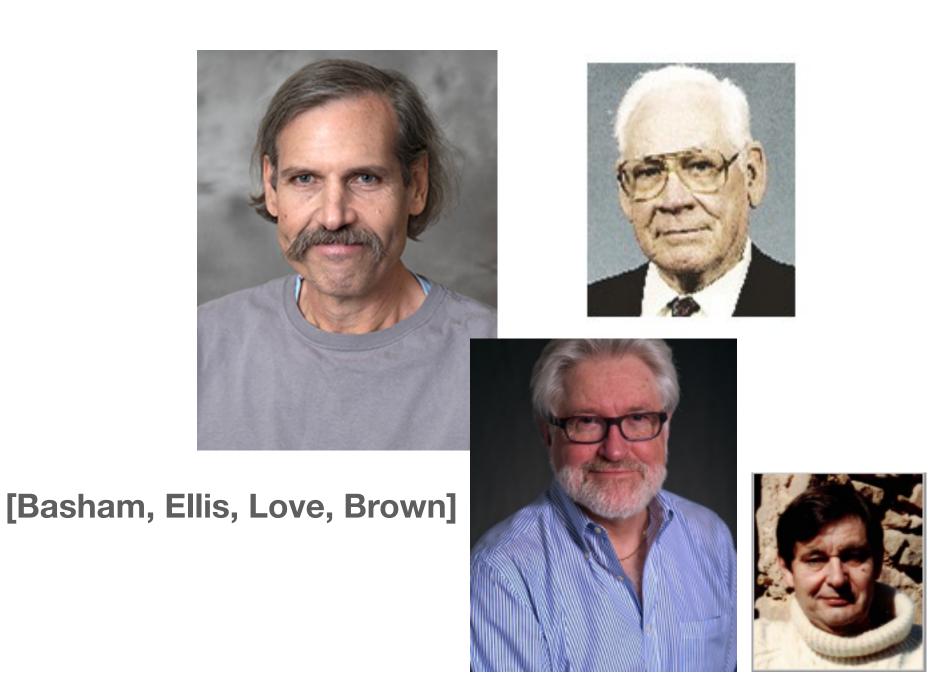


→ provides the most stringent bounds on charm Yukawa couplings



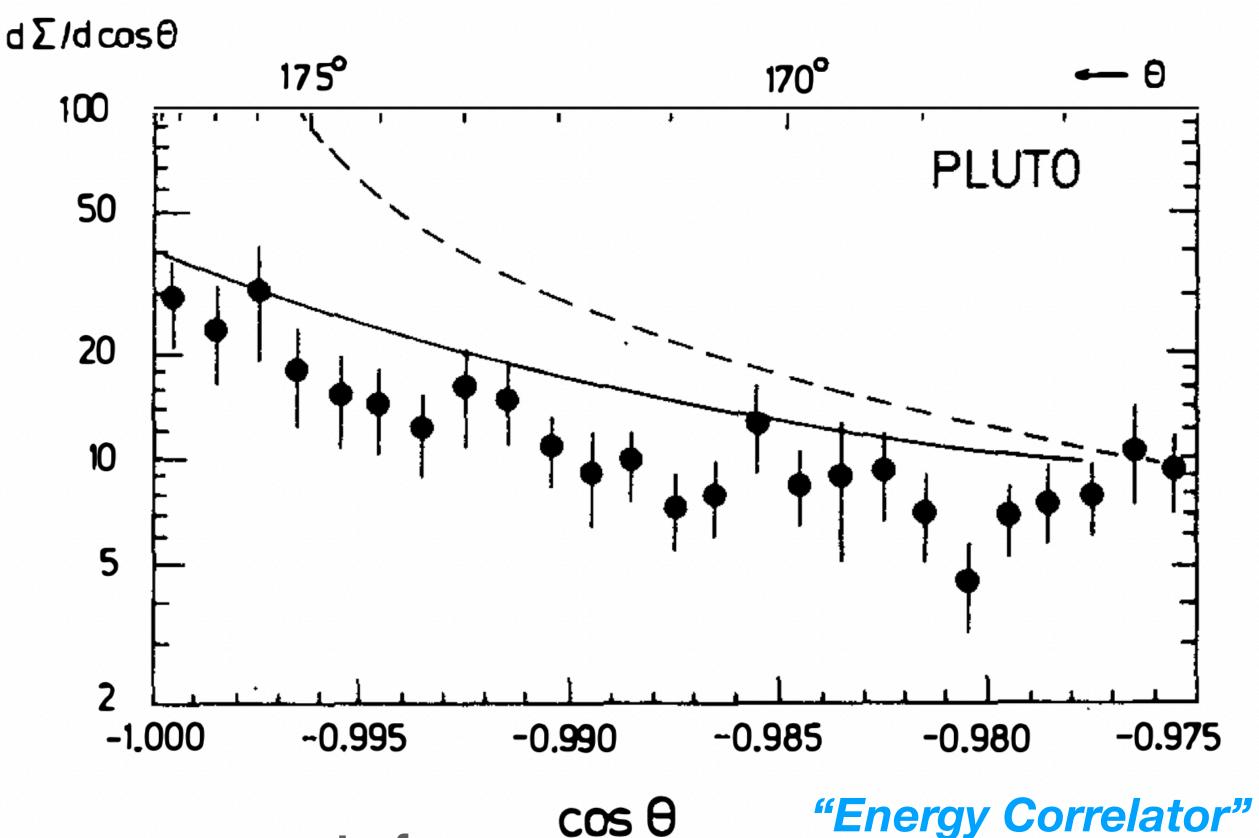


## How do we study collisions?



- → First introduced EEC as experimental observable sensitive to  $\alpha_s$  (1978)
- $\rightarrow$  Measured in  $e^+e^-$  by PLUTO (1978)

$$EEC(\chi) = \frac{1}{\Delta \chi N} \int_{\chi - \Delta \chi/2}^{\chi + \Delta \chi/2} \sum_{\text{events}}^{N} \sum_{i,j} \frac{E_i E_j}{E_{\text{vis}}^2} \delta(\chi' - \chi_{ij}) \, d\chi'$$



## How do we study collisions?





$$EEC(\chi) = \frac{\chi + \Delta \chi/2}{N} = \frac{K + \Delta \chi/2}{N} = \frac$$

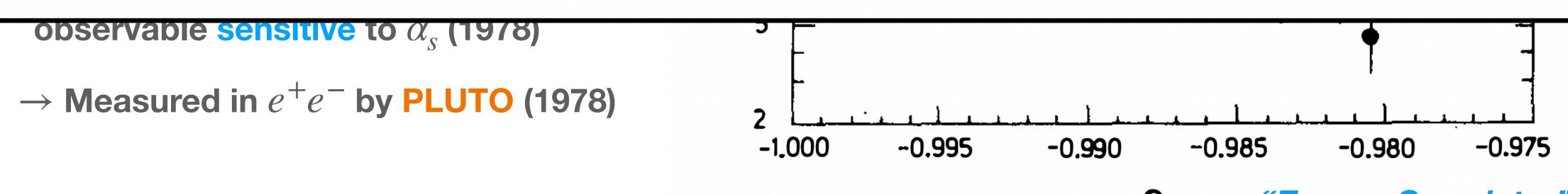
#### ADVANTAGES OF THE COLOR OCTET GLUON PICTURE<sup>☆</sup>

H. FRITZSCH\*, M. GELL-MANN and H. LEUTWYLER\*\*

California Institute of Technology, Pasadena, Calif. 91109, USA

Received 1 October 1973

It is pointed out that there are several advantages in abstracting properties of hadrons and their currents from a Yang-Mills gauge model based on colored quarks and color octet gluons.

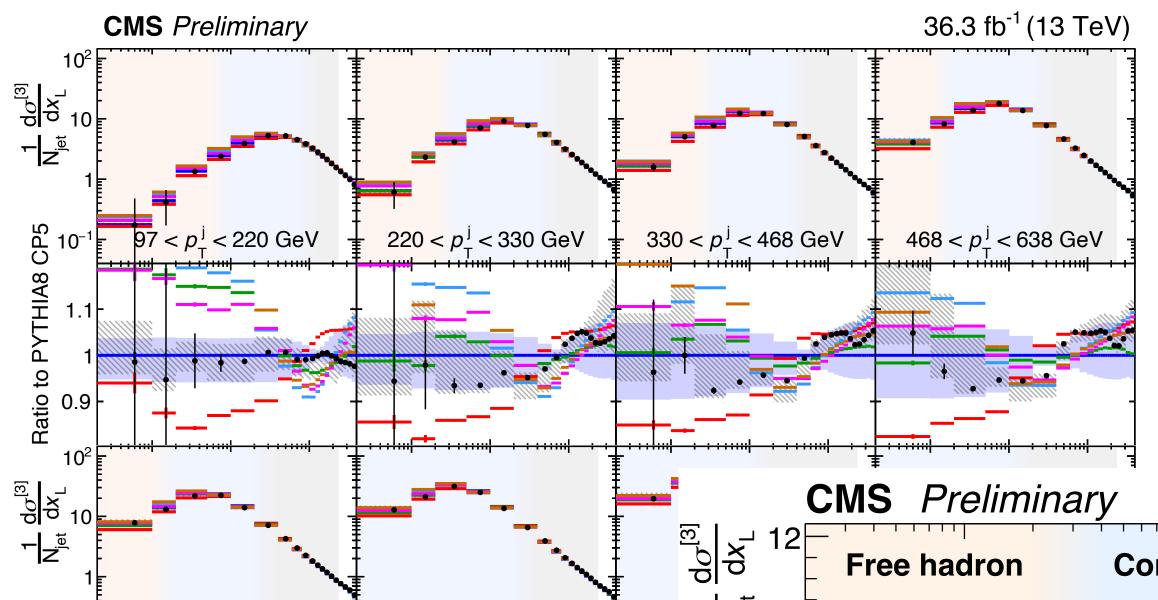


 $\cos \theta$ 

"Energy Correlator" (EEC)

### The EEC in 2024

#### "Energy Correlator" (EEC)



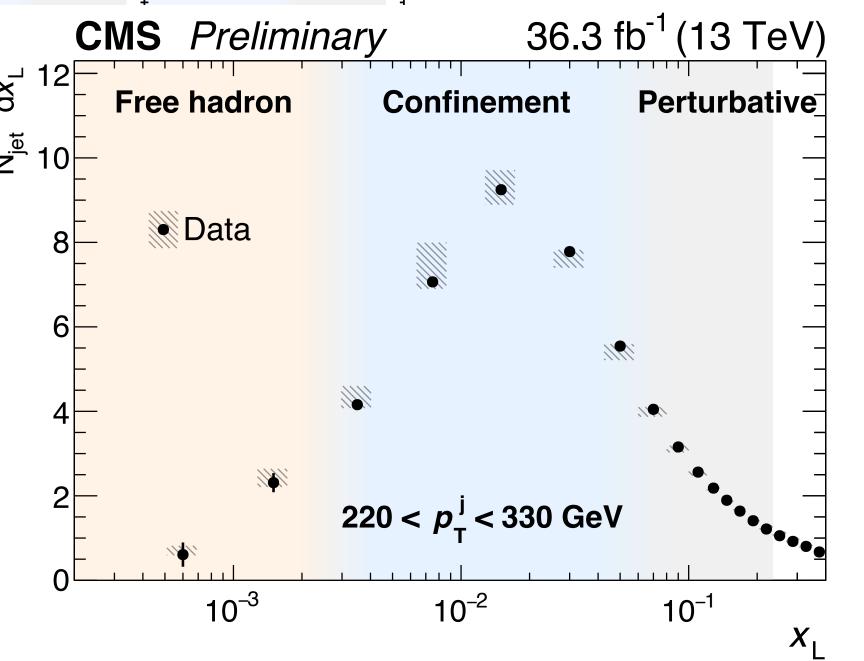
 $846 < p_{_{\rm T}}^{\rm j} < 1101 \, \text{GeV} \stackrel{\text{I}}{=} 1101$ 

PYTHIA8 CP5(simple sl

"This is the most precise measurement of  $\alpha_s(M_Z)$  by a jet substructure observable to date"

#### Quote taken directly from:

**CMS** Collaboration, 2024





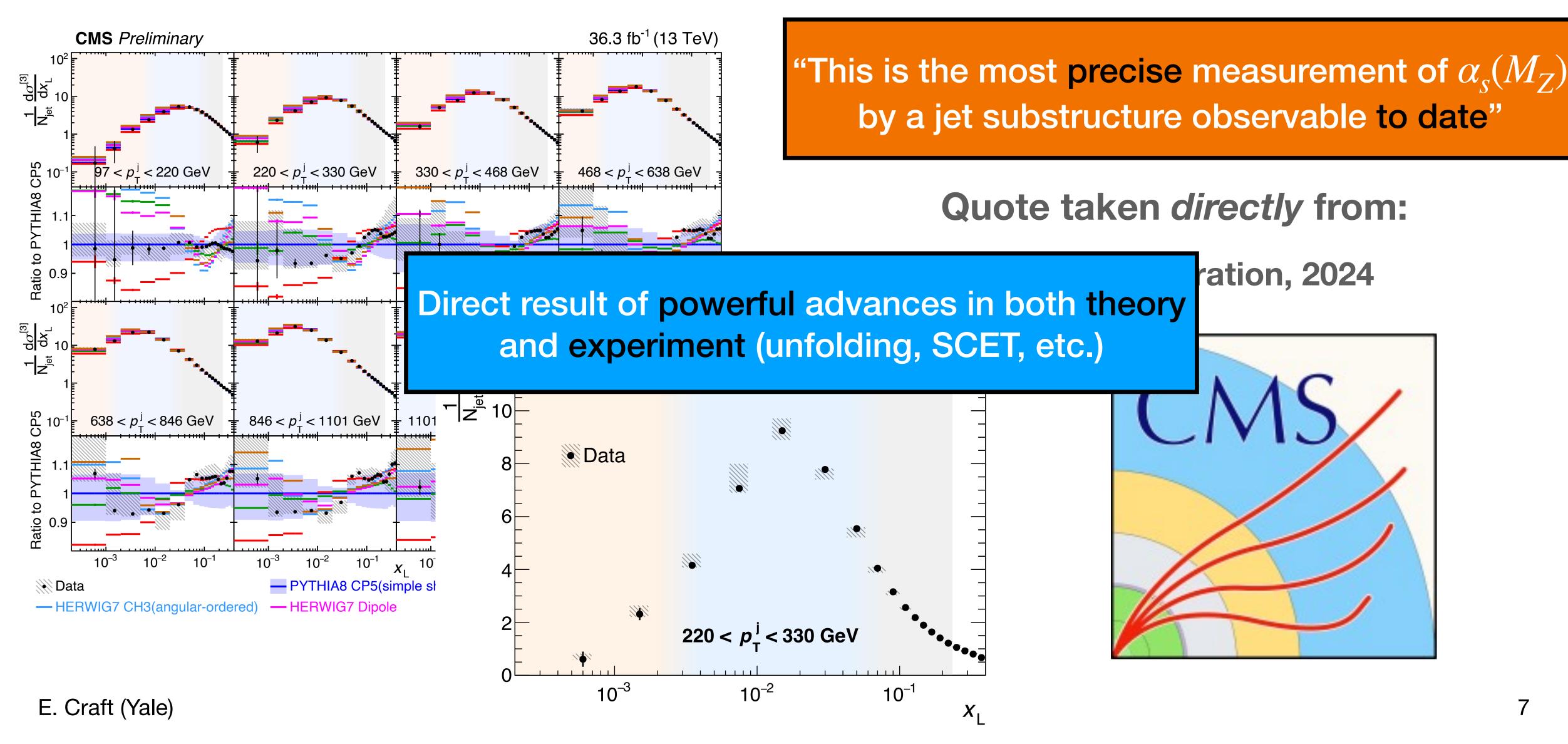
**N** Data

 $638 < p_{\tau}^{j} < 846 \text{ GeV}$ 

— HERWIG7 CH3(angular-ordered) — HERWIG7 Dipole

### The EEC in 2024

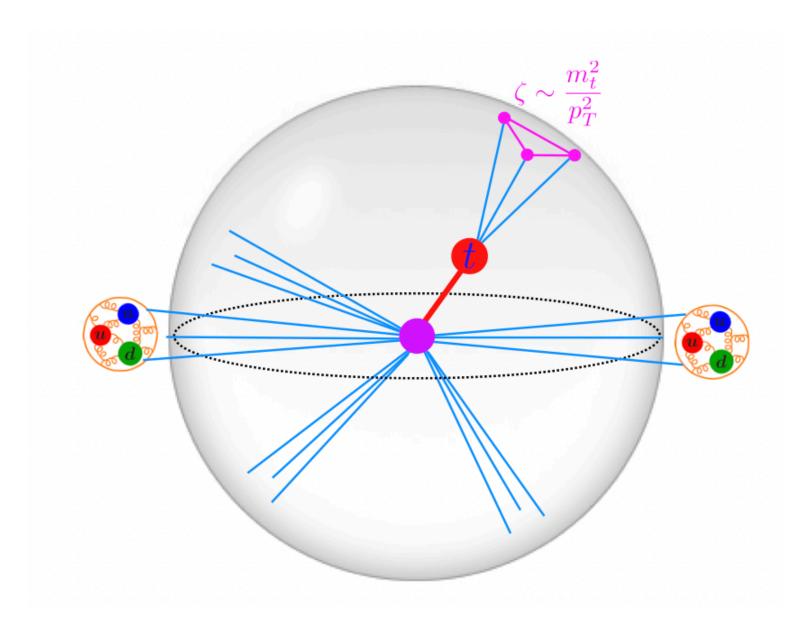
#### "Energy Correlator" (EEC)

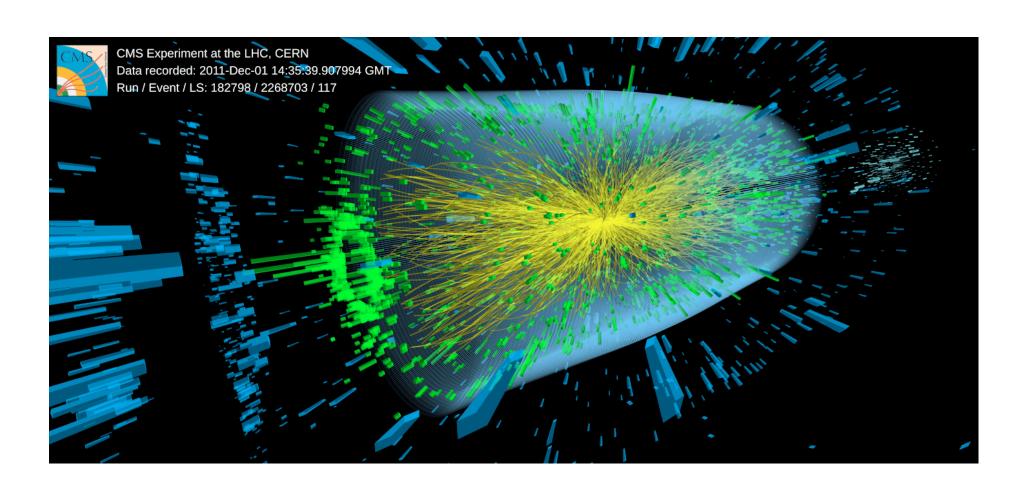


### From Searches to Measurements

To fully take advantage of the LHC, it is necessary to bolster our current physics searches with first principles theory calculations

 $\rightarrow$  Many interesting opportunities to study QCD at high energies: understanding confinement, precision measurements,  $\alpha_{\rm S}, m_t \dots$ 





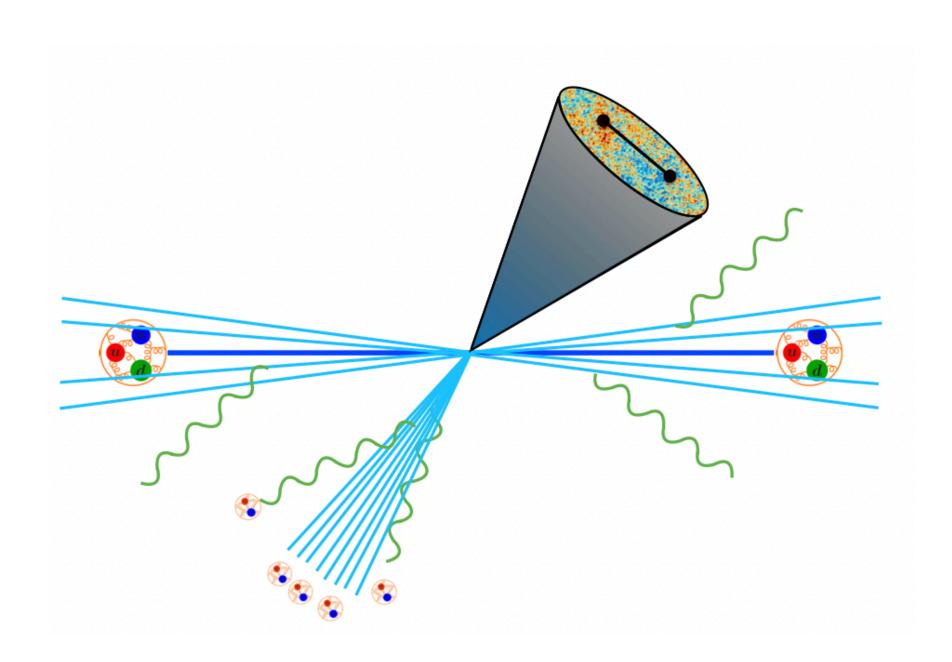
Requires the development of a new set of theoretical tools

## Reformulating Jet Substructure

Field Theoretic Foundations

## **Energy Flow Operators**

From the perspective of QFT, jet substructure is the study of correlation functions of energy flow operators



$$\mathcal{E}(\overrightarrow{n}) = \lim_{t \to \infty} t^3 \int_0^1 dv \, v^2 \left[ n^i T_i^0(t, tv \overrightarrow{n}) \right]$$
Sveshnikov, Tkachov (1995)

→ "Energy Flow Operator"

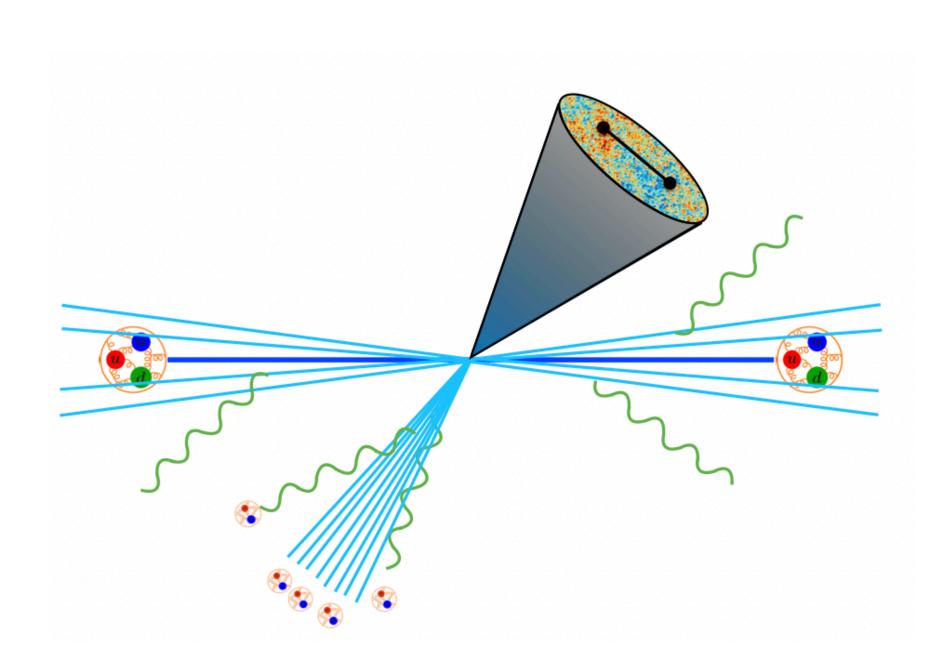
$$\langle \Psi | \mathcal{E}(\hat{n}_1) \dots \mathcal{E}(\hat{n}_k) | \Psi \rangle$$

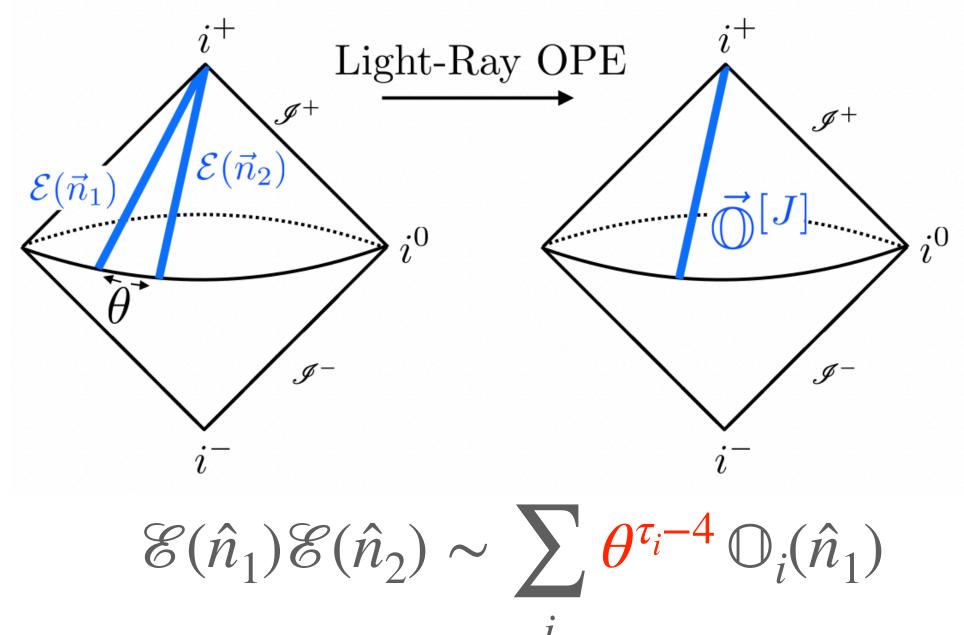
→ "Statistical Correlations"

These correlation functions measure the flow of energy at infinity.

## **Energy Flow Operators**

Situations of interest at the LHC involve non-generic configurations of lightray operators: interested in the small angle (OPE) limit.

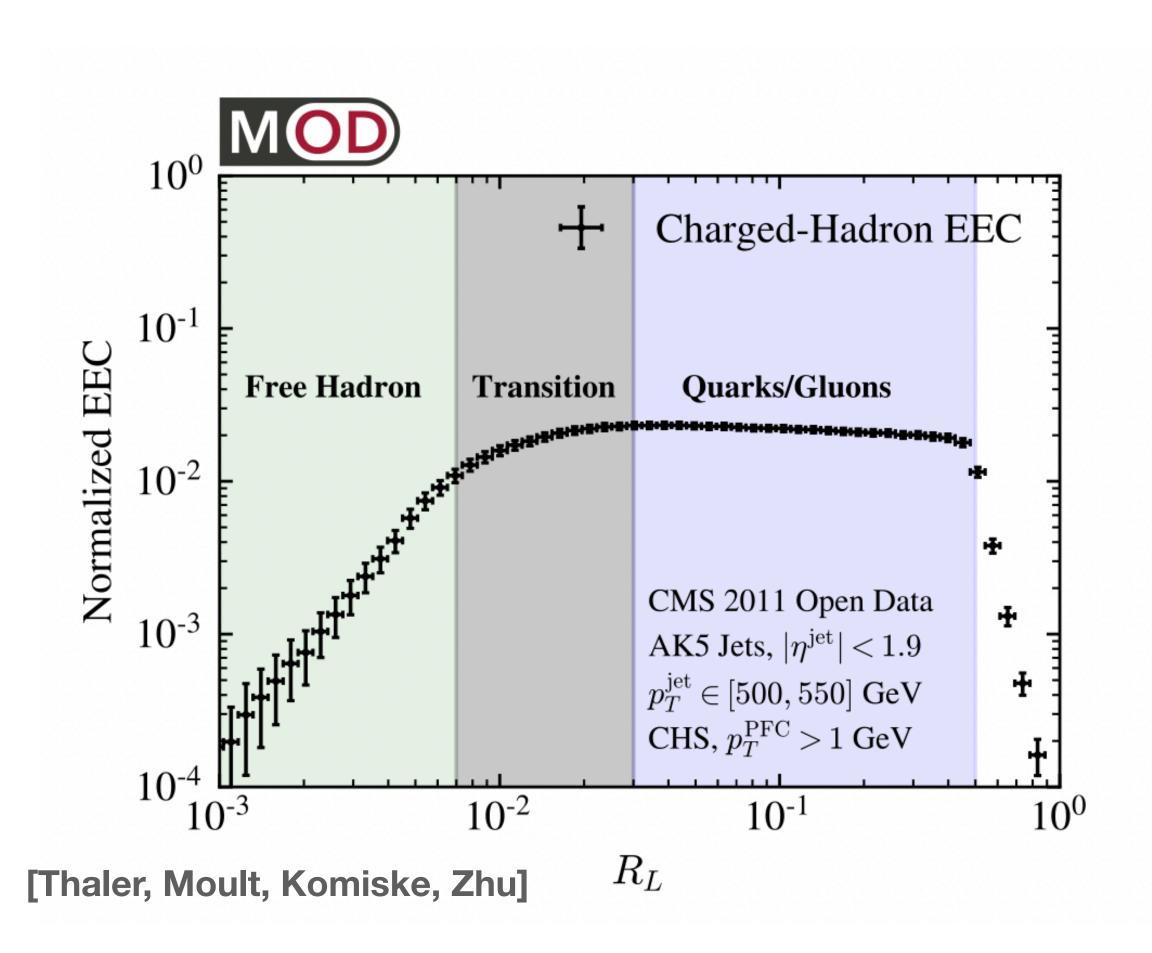


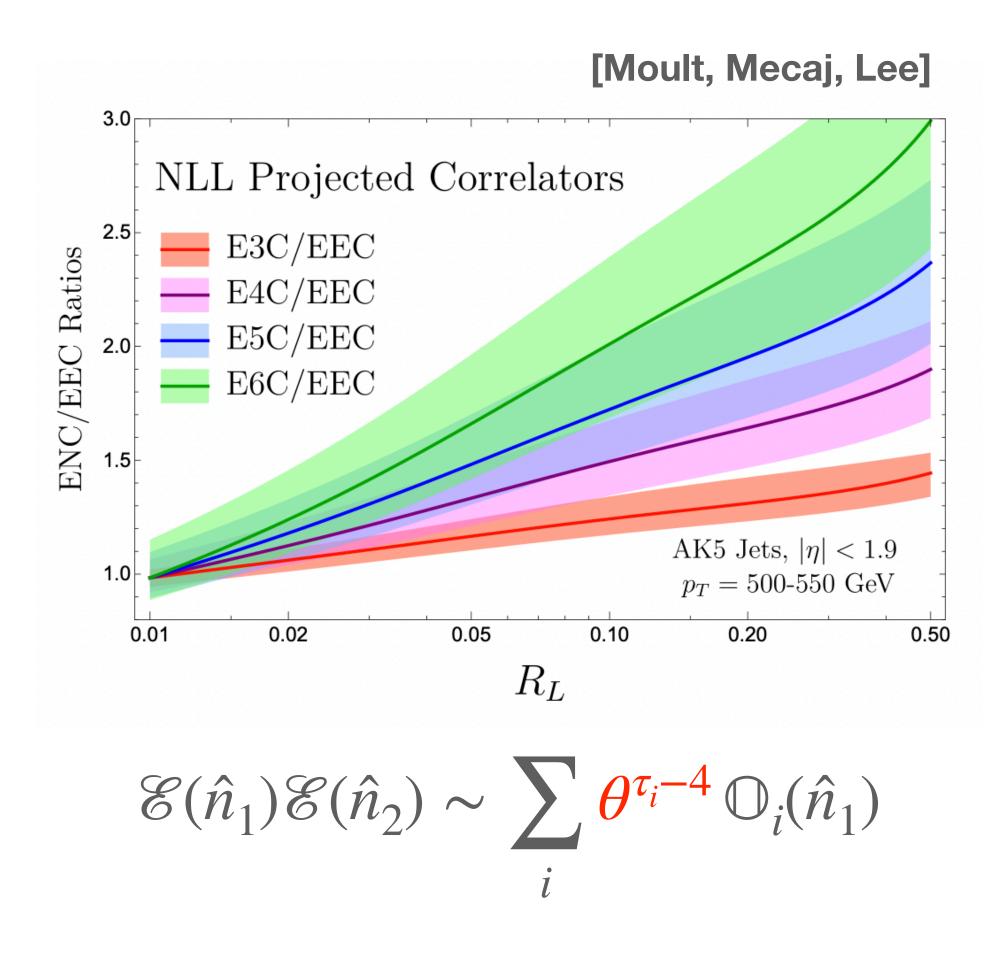


$$\mathcal{E}(\hat{n}_1)\mathcal{E}(\hat{n}_2) \sim \sum_i \theta^{\tau_i-4} \mathbb{O}_i(\hat{n}_1)$$
[Hofman, Maldacena]

In the small angle limit, these lightray operators should exhibit the universal behavior of QCD

## **Energy Flow Operators**



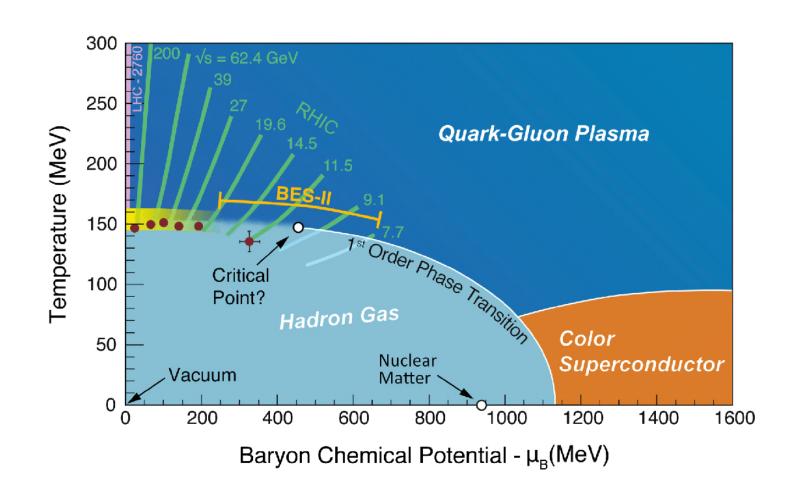


It is precisely this scaling behaviour which allows for such precise measurements of  $\alpha_{\rm s}(M_7)$ 

## Going Beyond $\alpha_s(M_Z)$

Several open questions remain across both Particle and Nuclear Physics

→ Many of these open problems are deeply connected to Quantum Chromodynamics



→ Why is color charge so complicated?

#### Hot QCD Medium QCD Cold QCD

→ Quark Gluon Plasma
→ Strong CP

→ Rare Higgs Decays

→ Confinement

→ Gluon Saturation

→ Proton Spin and Radius Puzzle

→ 3D Structure of protons and nuclei

Numerous collider experiments spanning several continents working to resolve these fundamental questions

E. Craft (Yale)

→ Hadronization

→ Quarkonia



## **Beautiful and Charming Energy Correlators**

Evan Craft — Yale University arXiv: 2210.09311



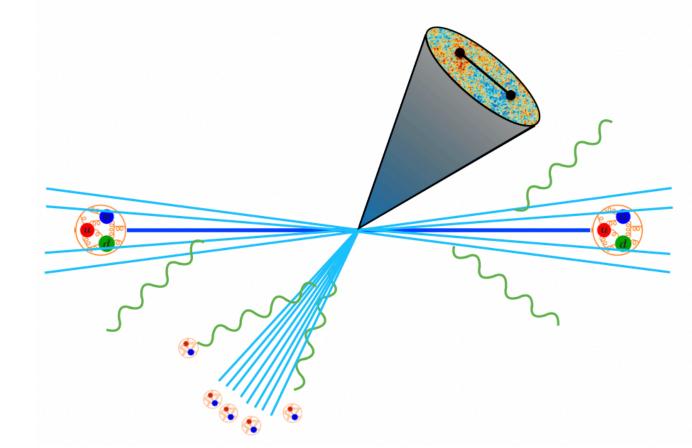




Based on work with K. Lee, B. Mecaj, I. Moult

**Yale University** 

MIT

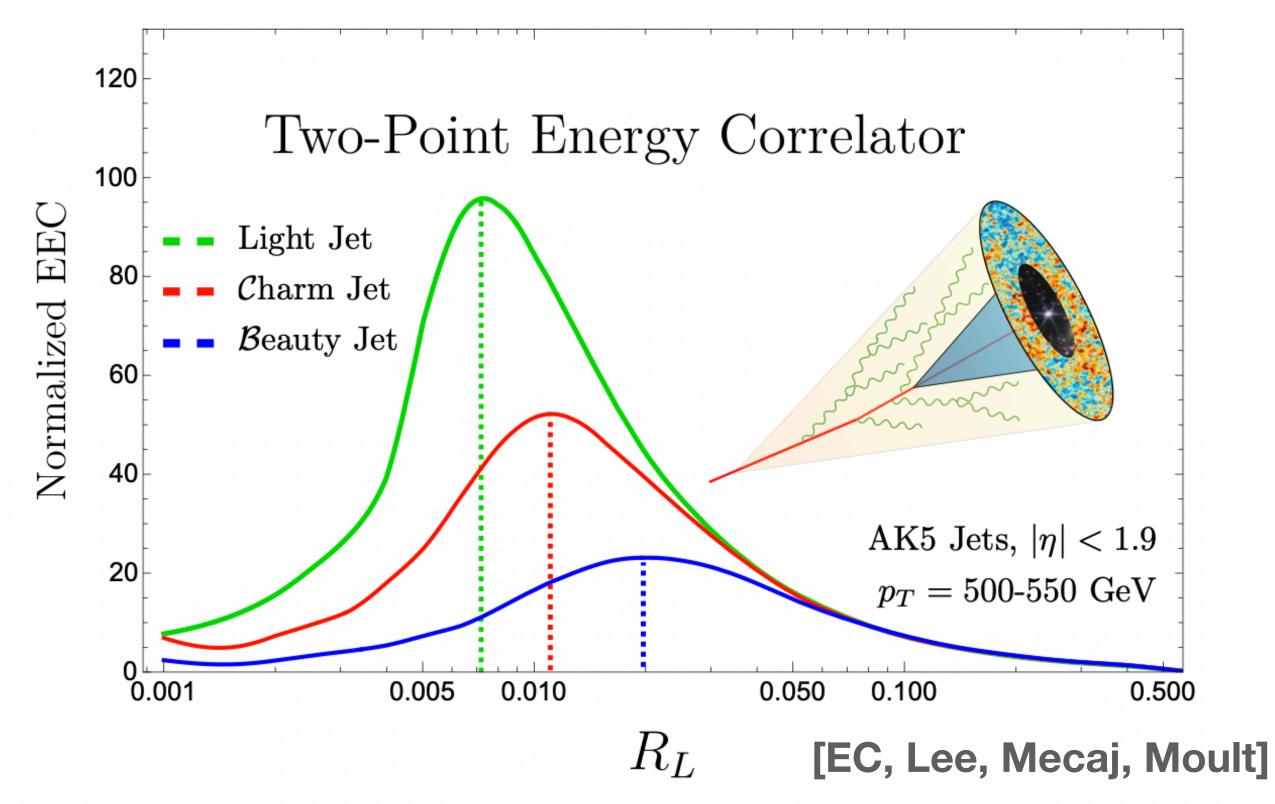


#### Intrinsic masses of QCD imprinted onto energy correlators

- → allows for an unprecedented window into hadronization effects
- → provides a powerful perspective for probing jet substructure
- → provides a new, unifying technique for understanding intrinsic mass

$$\langle \Psi | \mathcal{E}(\hat{n}_1) ... \mathcal{E}(\hat{n}_k) | \Psi \rangle$$

the "perfect" observable



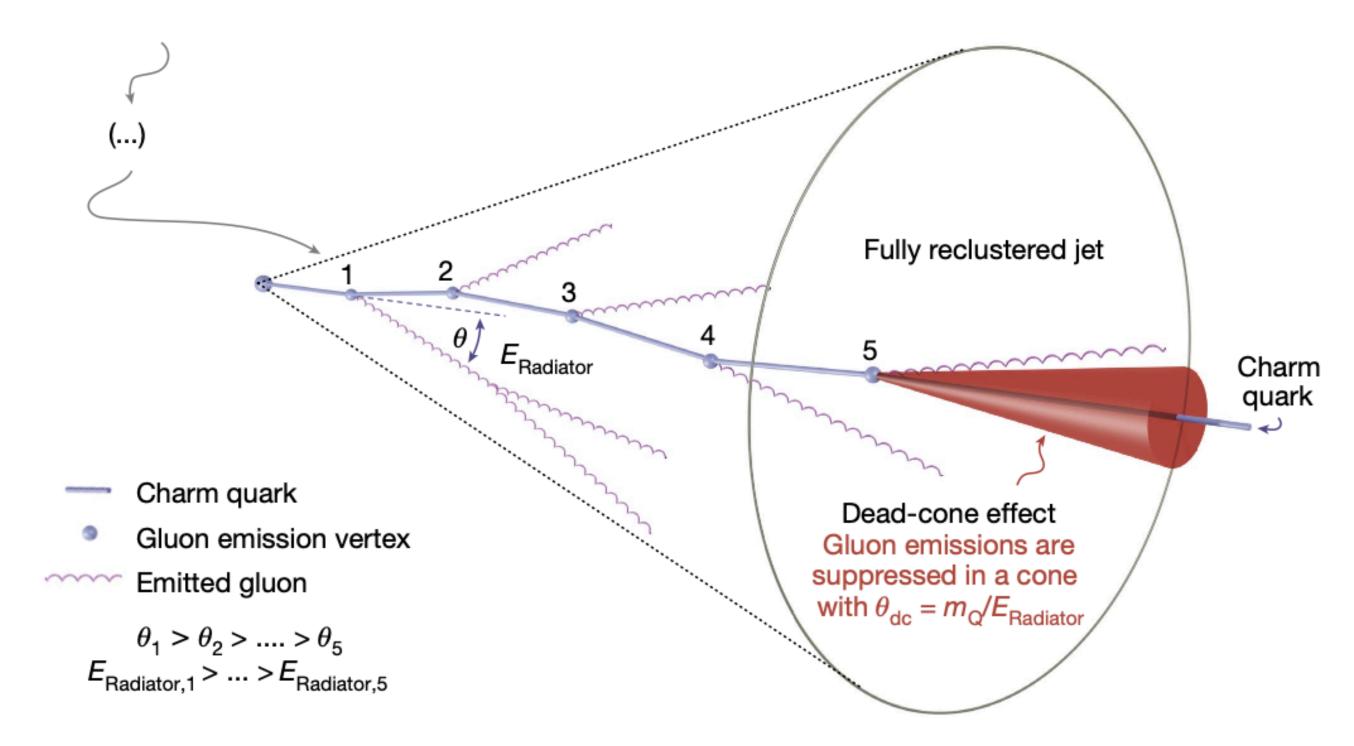
[ALICE Collaboration, Nature Physics]

Dokshitzer, Khoze, Troyan (1991)

Heavy quark radiation of gluons is suppressed within a cone of radius  $m_q/E_q$  around its center.

- → Fundamental property of all gauge field theories
- → Direct signature of intrinsic mass before confinement

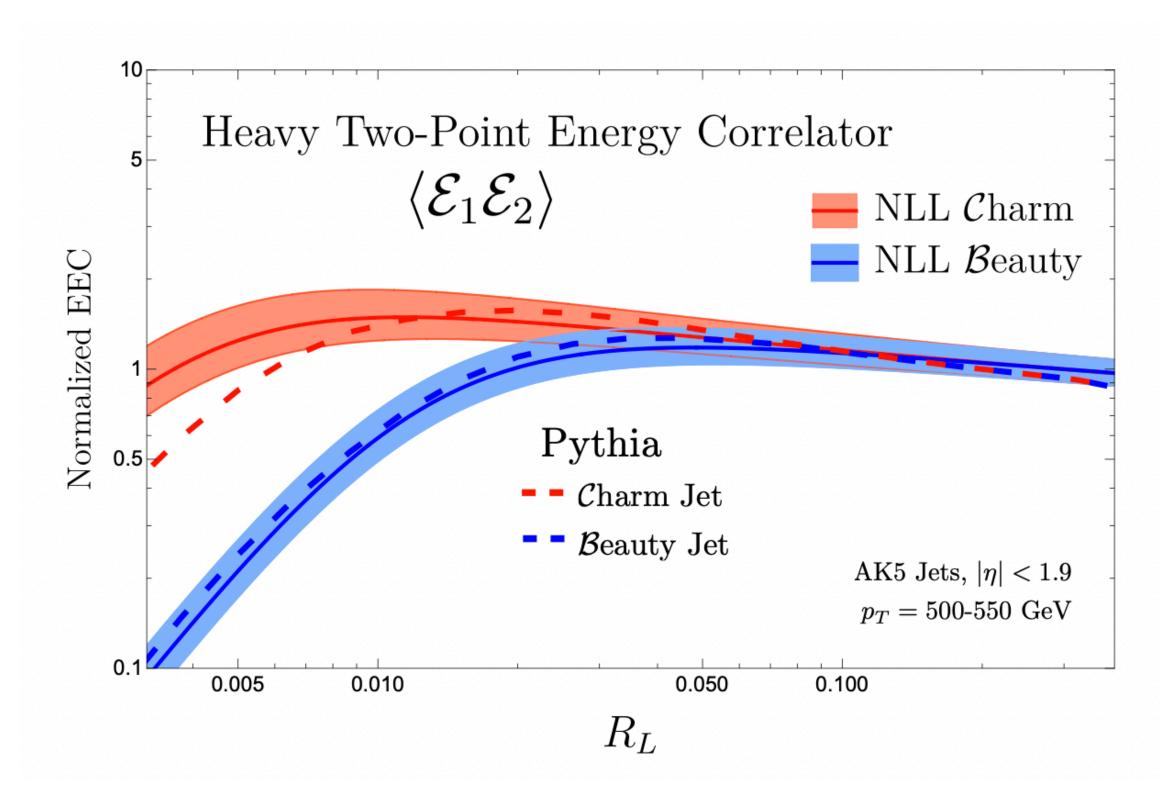
We can access this effect simply with statistical correlations (light-ray operators) — providing a precise, field theoretic description of the dead cone.



Measured this year by ALICE using a complex iterative declustering technique

- → Inferred all gluon emissions *directly*
- → State of the art analysis techniques

Heavy quark radiation of gluons is *suppressed* within a cone  $\theta_q \sim m_q/E_q$  and this suppression is visibly imprinted on energy correlators



[EC, Lee, Mecaj, Moult]

Exposes the "dead-cone" effect of fundamental QCD, using correlations of light-ray operators

→ first collinear NLL calculation of a heavy quark jet substructure observable at the LHC

$$EEC = H \times J \times S$$
 fundamental test of SCET factorization at the LHC

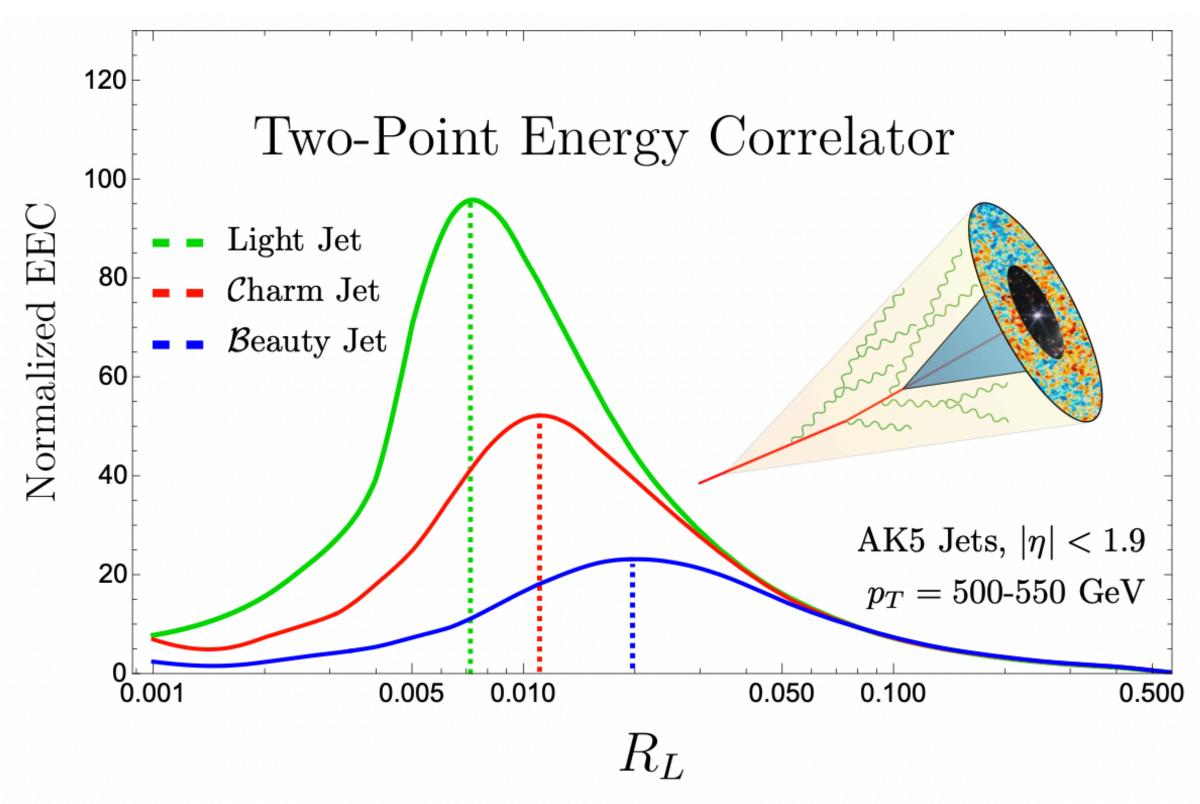
In the UV regime, scaling should be independent of mass

$$\mathcal{E}(\hat{n}_1)\mathcal{E}(\hat{n}_2) \sim \sum_{i} \theta^{\tau_i - 4} \mathbb{O}_i(\hat{n}_1)$$

In the IR regime, mass is an intrinsic scale, and should be imprinted on the correlator

$$\langle \Psi | \mathscr{E}(\hat{n}_1) ... \mathscr{E}(\hat{n}_k) | \Psi \rangle$$

[EC, Lee, Mecaj, Moult]



EECs provide a precise, field-theoretic description of the dead-cone effect

Transition Scale 
$$\sim \frac{m_q}{p_{T,jet}}$$



# Pushing the Boundaries of Jet Substructure

Evan Craft — Yale University









Work in prep. with K. Lee, B. Mecaj, I. Moult, & M. Gonzalez

Yale University

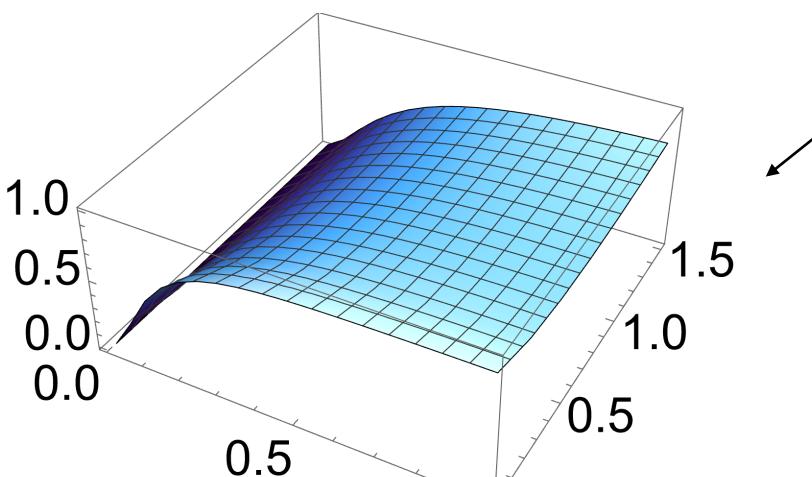


## Extension: Higher Points

Natural to also consider higher point correlators



shape of the dead-cone!







transverse spin 0

$$\mathcal{O}_{q}^{[J]} = \frac{1}{2^{J}} \bar{\psi} \gamma^{+} (iD^{+})^{J-1} \psi$$

$$\mathcal{O}_{g}^{[J]} = -\frac{1}{2^{J}} F_{a}^{\mu+} \gamma^{+} (iD^{+})^{J-2} F_{a}^{\mu-1}$$

excited by 2-point

transverse spin 2

Collinear

Limit

$$\mathcal{O}_{q}^{[J]} = \frac{1}{2^{J}} \bar{\psi} \gamma^{+} (iD^{+})^{J-1} \psi \qquad \qquad \mathcal{O}_{\tilde{g}\lambda}^{[J]} = -\frac{1}{2^{J}} F_{a}^{\mu +} \gamma^{+} (iD^{+})^{J-2} F_{a}^{\nu +} \epsilon_{\lambda \mu} \epsilon_{\lambda \nu}$$

$$\mathcal{O}_{g}^{[J]} = -\frac{1}{2^{J}} F_{a}^{\mu +} \gamma^{+} (iD^{+})^{J-2} F_{a}^{\mu +} \qquad \qquad \uparrow$$
helicity  $\pm 1$ 

excited by 3-point

$$\mathscr{E}(\hat{n}_1) \dots \mathscr{E}(\hat{n}_k) \sim \sum_{i} \theta^{\tau_i - 4} \mathbb{O}_i(\hat{n}_1)$$

→ Probe fundamental operators of QCD

→ Access to non-Gaussianities

0.0

→ Full Shape Dependence

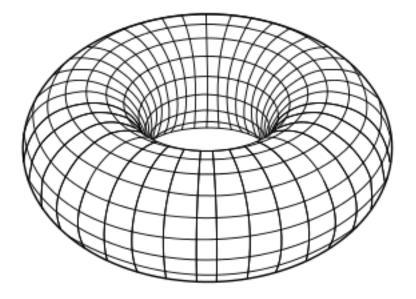
 $\mathcal{E}(\vec{n}_1)$ 

 $\mathcal{E}(\vec{n}_2)$ 

## Topological Aspects

There is a direct mapping from the kinematic configuration of the EEC, to the torus

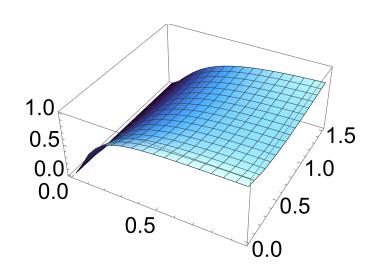
$$y^2 = 4x^3 - g_1x - g_3 \longrightarrow$$



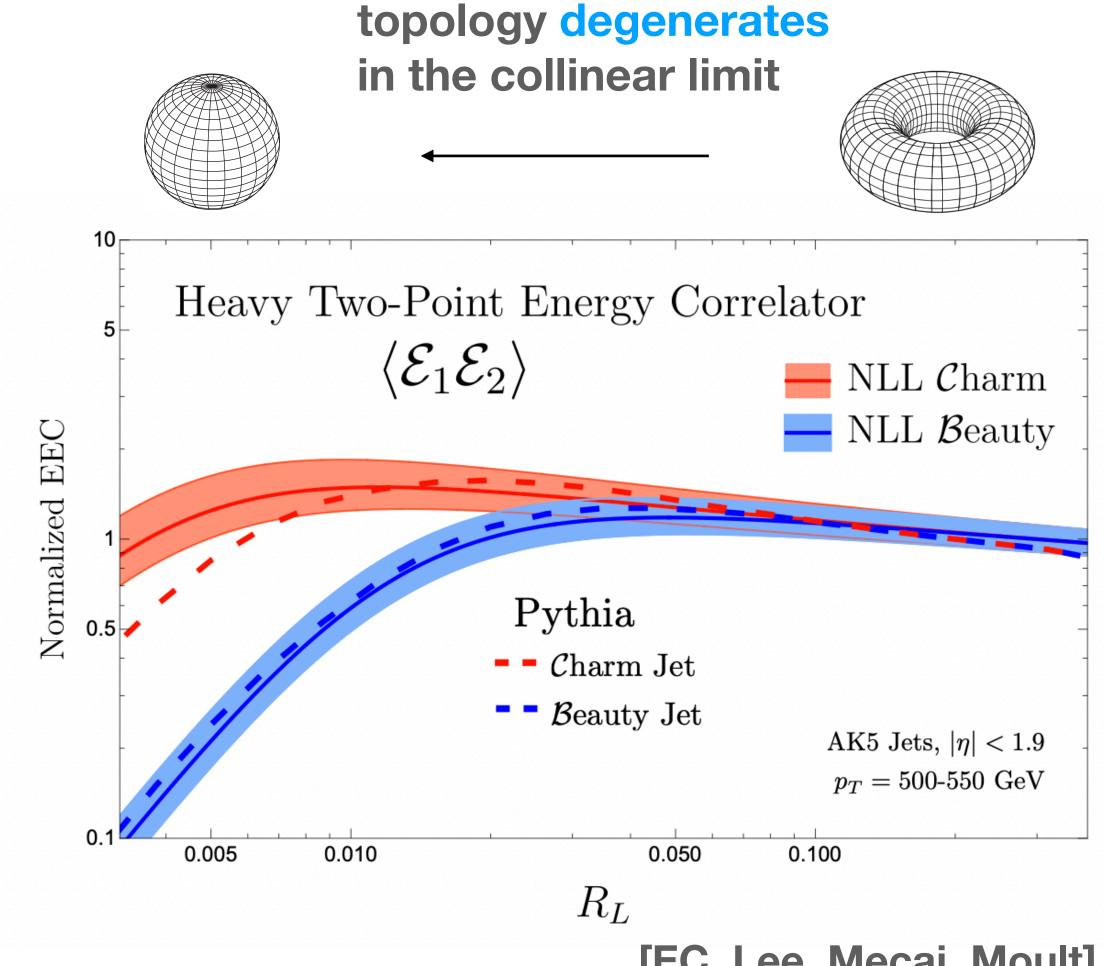
$$\omega_1 \sim {}_2F_1(1/2, 1/2, 1; \lambda)$$

$$\omega_2 \sim {}_2F_1(1/2, 1/2, 1; 1 - \lambda)$$

periods deformed by kinematics



Similar degeneration for the three point!



[EC, Lee, Mecaj, Moult]

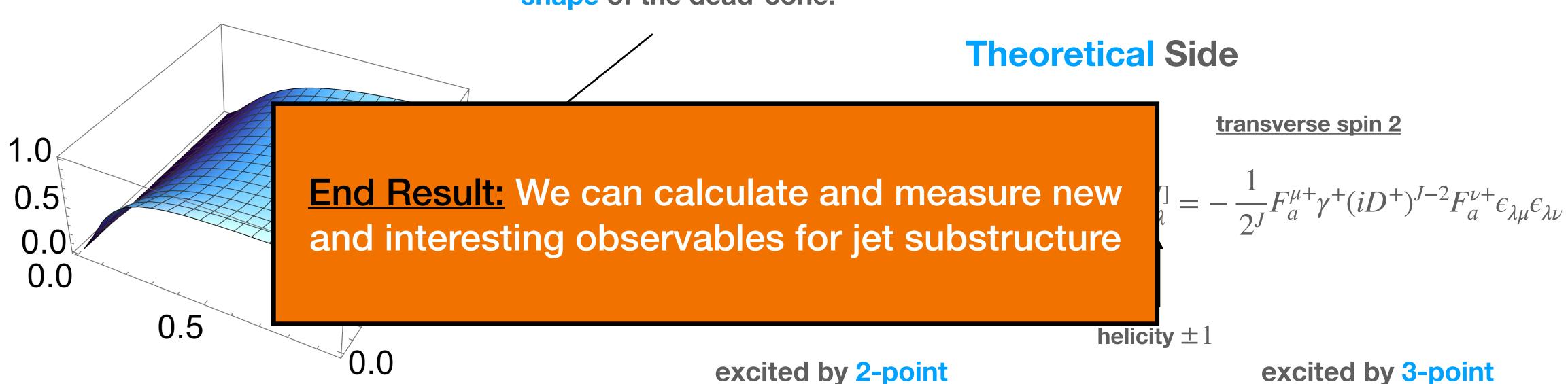
## Extension: Higher Points

Natural to also consider higher point correlators

 $\mathcal{E}(\vec{n}_1)$   $\mathcal{E}(\vec{n}_2)$   $\mathcal{E}(\vec{n}_3)$   $\mathcal{E}(\vec{n}_3)$   $\mathcal{E}(\vec{n}_3)$   $\mathcal{E}(\vec{n}_3)$   $\mathcal{E}(\vec{n}_3)$ 

**Experimental Side** 

3-point EEC allows access to the shape of the dead-cone!



→ Access to non-Gaussianities

→ Full Shape Dependence

$$\mathcal{E}(\hat{n}_1) \dots \mathcal{E}(\hat{n}_k) \sim \sum_{i} \theta^{\tau_i - 4} \mathbb{O}_i(\hat{n}_1)$$

→ Probe fundamental operators of QCD

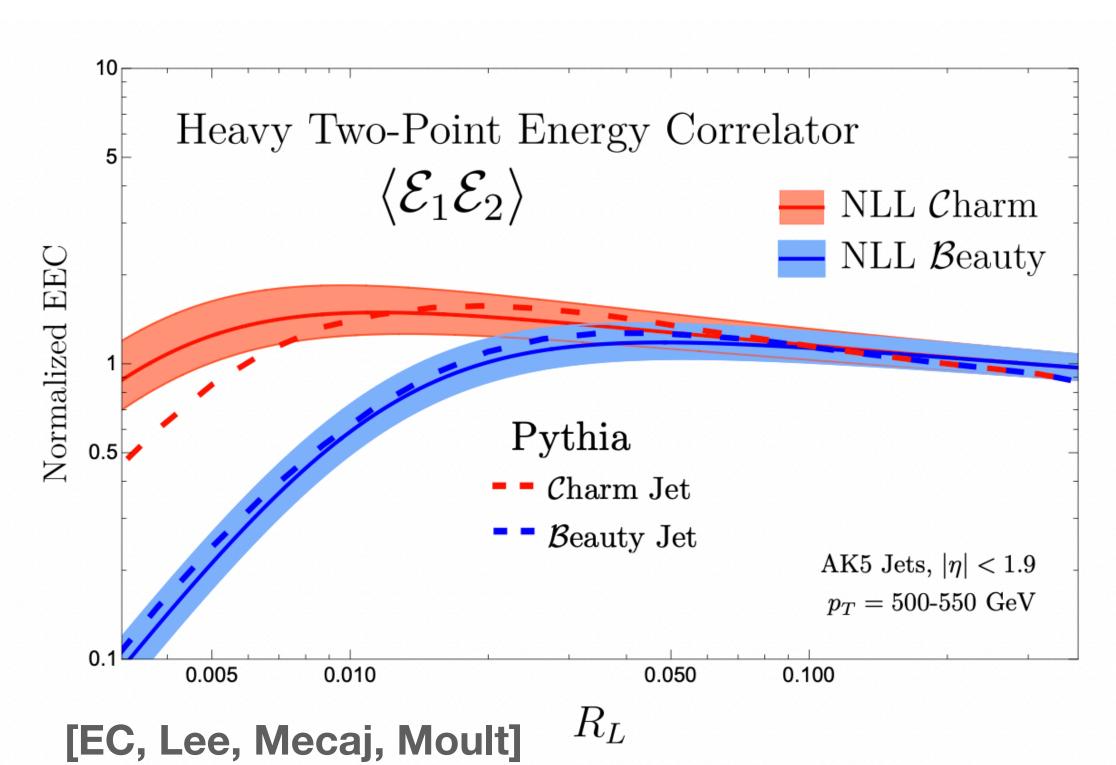
 $\mathcal{E}(\vec{n}_1)$ 

## Concluding Remarks

**Unifying Theory and Experiment** 

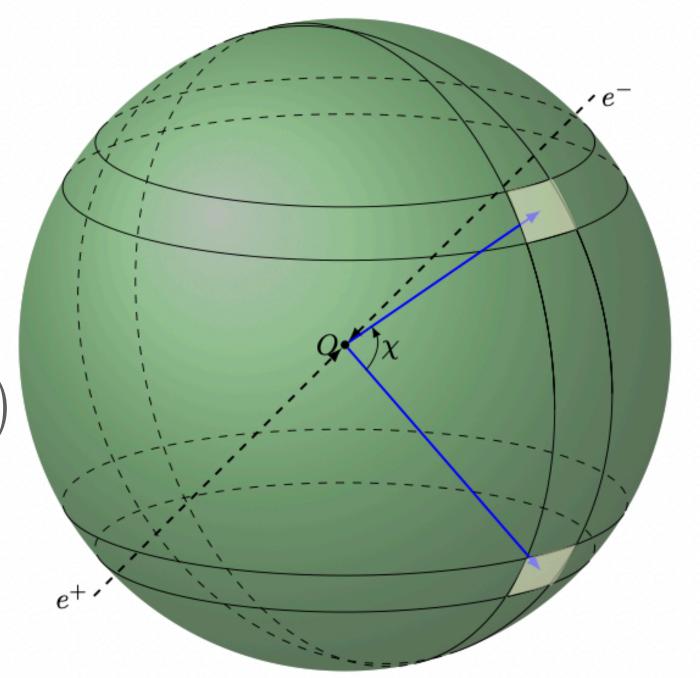
## Two Symbiotic Perspectives

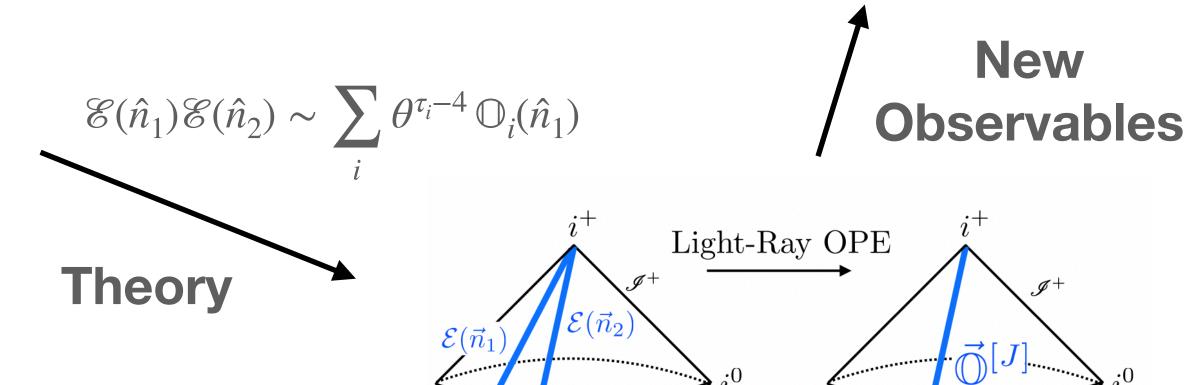
#### **Beautiful and Charming Interplay!**



$$\frac{d\sigma}{d\cos\chi} = \sum_{i < j} \int d\sigma \frac{E_i E_j}{Q^2} \delta\left(\overrightarrow{n_i} \cdot \overrightarrow{n_j} - \cos\chi\right)$$







This sort of collaboration is crucial for the success of future collider studies

## Summary

Jet substructure provides a physical realization of the OPE limit of light-ray operators

→ Direct bridge between recent theoretical advancements and QCD Phenomenology

Creates an unprecedented symbiosis between theory and experiment

→ Allowing for sharp probes of interesting physics, new and old

