## Prospects for constraining lightquark electroweak couplings at Higgs factories

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based on work in collaboration with D. Jeans, J. Reuter, J. Tian, A.F. Żarnecki

DPF-Pheno 2024, Pittsburgh, May 2024



Measuring precision observables allows to constrain the SM parameters and to search for New Physics.



#### Z decays to hadrons are constrained from LEP and SLC...

#### $R_b = \Gamma(b\overline{b})/\Gamma(\text{hadrons})$

 $\Gamma_{12}/\Gamma_8$ 

OUR FIT is obtained by a simultaneous fit to several c- and b-quark measurements as explained in the note "The Z boson" and ref. LEP-SLC 06.

The Standard Model predicts  $R_{b}=0.21581$  for  $m_{t}=174.3$  GeV and  $M_{H}=150$  GeV.

VALUE	DOCUMENT ID		TECN	COMMENT
0.21629±0.00066 OUR FIT				
$0.21594 \pm 0.00094 \pm 0.00075$	$^1$ ABE	05F	SLD	<i>E<sup>ee</sup></i> =91.28 GeV
$0.2174 \ \pm 0.0015 \ \pm 0.0028$	<sup>2</sup> ACCIARRI	00	L3	<i>E<sup>ee</sup></i> _cm= 89–93 GeV
$0.2178\ \pm 0.0011\ \pm 0.0013$	<sup>3</sup> ABBIENDI	<b>99</b> B	OPAL	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV
$0.21634 \pm 0.00067 \pm 0.00060$	<sup>4</sup> ABREU	<b>99</b> B	DLPH	<i>E<sup>ee</sup></i> <sub>cm</sub> = 88–94 GeV
$0.2159 \ \pm 0.0009 \ \pm 0.0011$	<sup>5</sup> BARATE	97F	ALEP	<i>E</i> <sup><i>ee</i></sup> <sub>CM</sub> = 88–94 GeV

#### Review of Particle Physics, PDG, 2022

$R_c =$	Γ(	(c <u>c</u> )	/Γ(	had	Iron	IS)
-						

 $\Gamma_{11}/\Gamma_8$ 

OUR FIT is obtained by a simultaneous fit to several c- and b-quark measurements as explained in the note "The Z boson" and ref. LEP-SLC 06.

The Standard Model predicts  $R_c = 0.1723$  for  $m_t = 174.3$  GeV and  $M_H = 150$  GeV.

VALUE	DOCUMENT ID		TECN	<u>COMMENT</u>
0.1721±0.0030 OUR FIT				
$0.1744 \pm 0.0031 \pm 0.0021$	$^1$ ABE	05F	SLD	<i>E<sup>ee</sup></i> =91.28 GeV
$0.1665 \pm 0.0051 \pm 0.0081$	<sup>2</sup> ABREU	00	DLPH	<i>E</i> <sup><i>ee</i></sup> <sub>CM</sub> = 88–94 GeV
$0.1698 \pm 0.0069$	<sup>3</sup> BARATE	<b>00</b> B	ALEP	<i>E<sup>ee</sup></i> <sub>CM</sub> = 88–94 GeV
$0.180\ \pm 0.011\ \pm 0.013$	<sup>4</sup> ACKERSTAFF	98E	OPAL	<i>E<sup>ee</sup></i> = 88–94 GeV
$0.167\ \pm 0.011\ \pm 0.012$	<sup>5</sup> ALEXANDER	<b>96</b> R	OPAL	<i>E<sup>ee</sup></i> <sub>CM</sub> = 88–94 GeV

#### $\Gamma((u\overline{u}+c\overline{c})/2)/\Gamma(hadrons)$

 $\Gamma_9/\Gamma_8$ This quantity is the branching ratio of  $Z \rightarrow$  "up-type" quarks to  $Z \rightarrow$  hadrons. Except ACKERSTAFF 97T the values of  $Z \rightarrow$  "up-type" and  $Z \rightarrow$  "down-type" branchings are extracted from measurements of  $\Gamma$ (hadrons), and  $\Gamma$ ( $Z \rightarrow \gamma + \text{jets}$ ) where  $\gamma$  is a highenergy (>5 or 7 GeV) isolated photon. As the experiments use different procedures and slightly different values of  $M_7$ ,  $\Gamma$  (hadrons) and  $\alpha_s$  in their extraction procedures, our average has to be taken with caution.

ALUE	<u>DOCUMENT ID</u>		TECN	<u>COMMENT</u>		
.166±0.009 OUR AVERAGE						
$.172 \substack{+0.011 \\ -0.010}$	<sup>1</sup> ABBIENDI	04E	OPAL	$E_{ m Cm}^{ee}=$ 91.2 GeV		
$.160 \pm 0.019 \pm 0.019$	$^2$ ACKERSTAFF	97⊤	OPAL	$E_{\rm CM}^{ee}$ = 88–94 GeV		
$.137 \substack{+ 0.038 \\ - 0.054}$	<sup>3</sup> ABREU	95X	DLPH	<i>E</i> <sup><i>ee</i></sup> <sub>CM</sub> = 88–94 GeV		
.137±0.033	<sup>4</sup> ADRIANI	93	L3	$E_{\rm CM}^{\it ee}=$ 91.2 GeV		

#### ...but not all of them!

Quantity	Value	Value (universal)	Standard Model		
$\overline{\Gamma_{e^+e^-}}$	$83.87 \pm 0.12$	$83.942 \pm 0.085$	$83.960 \pm 0.009$		
$\Gamma_{\mu^+\mu^-}$	$83.95\pm0.18$	$83.941 \pm 0.085$	$83.959 \pm 0.009$		
$\Gamma_{ au^+ au^-}$	$84.03\pm0.21$	$83.759 \pm 0.085$	$83.777 \pm 0.009$		
$\Gamma_{ m inv}$	$498.9 \hspace{0.2cm} \pm 2.5 \hspace{0.2cm}$	$500.5 \pm 1.5$	$501.445 \pm 0.047$		
$\Gamma_{u ar{u}}$			$299.89 \pm 0.20$		
$\Gamma_{c\bar{c}}$	$300.3 \hspace{0.2cm} \pm 5.3 \hspace{0.2cm}$	$300.0 \pm 5.2$	$299.81 \pm 0.20$		
$\Gamma_{d\bar{d}}, \Gamma_{s\bar{s}}$			$382.77 \pm 0.14$		
$\Gamma_{bar{b}}$	$377.4 \pm 1.3$	$377.0 \pm 1.2$	$375.73 \mp 0.18$		
$\Gamma_{ m had}$	$1744.8 \pm 2.6$	$1743.2  \pm 1.9 $	$1740.97 \pm 0.85$		
$\Gamma_Z$	$2495.5 \pm 2.3$	$2495.5 \pm 2.3$	$2494.11 \pm 0.86$		

Review of Particle Physics, PDG, 2022

## Future $e^+e^-$ colliders operating at the Z-pole would be a perfect place to study the couplings.



Source	$e^-e^+ \rightarrow c\bar{c}$				$e^-e^+ \to b\bar{b}$				
	$P_{e^{-}e^{+}}(-0.8,+0.3)$ $P_{e^{-}e^{+}}(-0.8,+0.3)$		$P_{e^{-}e^{+}}(+$	$P_{e^{-}e^{+}}(+0.8,-0.3)$		$P_{e^{-}e^{+}}(-0.8,+0.3)$		$P_{e^{-}e^{+}}(+0.8,-0.3)$	
	$R_c$	$A_{FB}^{car{c}}$	$R_c$	$A_{FB}^{car{c}}$	$R_b$	$A_{FB}^{bar{b}}$	$R_b$	$A_{FB}^{bar{b}}$	
Statistics	0.18%	0.38%	0.27%	0.52%	0.12%	0.24%	0.23%	0.70%	
Preselection eff.	<0.01%	0.12%	0.02%	0.16%	<0.01%	0.08%	0.06%	0.12%	
Background	0.01%	0.01%	0.02%	0.02%	0.01%	0.01%	0.06%	<0.01%	
heavy quark mistag	0.11%	<0.01%	0.06%	<0.01%	0.12%	<0.01%	0.22%	<0.01%	
uds mistag	0.03%	<0.01%	0.02%	<0.01%	0.08%	<0.01%	0.14%	<0.01%	
Angular correlations	0.10%	0.10%	0.10%	0.10%	0.10%	0.10%	0.10%	0.10%	
<b>Beam Polarisation</b>	<0.01%	<0.01%	0.02%	0.01%	<0.01%	0.01%	0.03%	0.15%	
Systematics	0.15%	0.16%	0.12%	0.19%	0.18%	0.13%	0.29%	0.22%	
Total	0.24%	0.41%	0.30%	0.55%	0.21%	0.27%	0.37%	0.73%	

A. Irles *et al.*, [2306.11413]

The branching ratios to heavy quarks could be well constrained e.g. at ILC thanks to excellent flavour-tagging.

### But how to take the measurement if...

tagging is imperfect (s quark)?
tagging is unavailable (u, d quarks)?

# How to measure Z couplings to light quarks?

#### General idea

We want to measure quark couplings:

$$c_f = v_f^2 + a_f^2$$

They are given in the SM by:

$$v_f = 2I_{3,f} - 4Q_f \sin^2 \theta_W \qquad a_f = 2I_{3,f}$$

 $\Gamma_{had}$  reads (exact at fixed order):

$$\Gamma_{had} = N_c \frac{G_\mu M_Z^3}{24\pi\sqrt{2}} \left(1 + \frac{\alpha_s}{\pi} + \dots\right) \left(3c_d + 2c_u\right)$$

and  $\Gamma_{had+\gamma}$ :

$$\Gamma_{had+\gamma} = N_c \frac{G_{\mu} M_Z^3}{24\pi\sqrt{2}} f(y_{cut}) \frac{\alpha}{2\pi} \left( 3q_d^2 c_d + 2q_u^2 c_u \right)$$

The correction factor  $f(y_{cut})$  to be determined for a given value of the resolution parameter  $y_{cut}$ .

#### Measurement at the Z-pole

We can measure radiative and non-radiative cross sections separately and disentangle the couplings  $c_d$  and  $c_u$ :

$$\sigma_{Z \to had} = \mathcal{C}_1 \cdot (3c_d + 2c_u)$$

$$\sigma_{Z \to had + \gamma} = \mathcal{C}_2 \cdot (3c_d + 8c_u)$$

Note: in this picture, the couplings are universal among quarks of the same type. This assumption can be lifted by employing heavy-flavour tagging.

# How to generate Monte Carlo events?

#### Analysis setup

We want to consider:

$$e^+e^- 
ightarrow qar{q}(\gamma)$$
 .

Experimentally measured photons can originate not only from the Final State Radiation but also from the Initial State Radiation, hadronisation and decays...

One may encounter the following issues:

- Matrix Element calculations may be either divergent or very slow for small photonemission angles,
- **ISR structure functions** can be used for small angles but a proper matching procedure is needed,
- FSR showers are important to account for QCD emissions but they may cause doublecounting,
- photons coming from hadronisation and hadron decays have to be included properly.

### Matching procedure – *Whizard* perspective

- <u>matching</u>: *soft* physics invisible in the detector, hard physics properly reconstructed
- soft ISR photons simulated using built-in structure functions
- soft FSR photons simulated using parton showers
- hard photons simulated using full ME calculations

→ momentum transfer and energy to define *soft* and *hard* regimes

### Efficiency of the matching procedure

• About 4% of *Whizard* events are rejected to avoid double-counting.



## How to select events?

work in progress

### Event reconstruction

Measurable photons can originate from:

- Initial State Radiation,
- Final State Radiation,
- hadronisation and decays.

The interesting information comes only from FSR so reconstruction criteria should reduce other contributions.

#### Isolation criterion y<sub>cut</sub>



#### Conclusions

- Couplings of the Z boson to light quarks are poorly constrained but an improvement could be achieved at future lepton colliders.
- Proper assessment of the uncertainties requires a deep understanding of theoretical calculations, event simulation and reconstruction.
- Work in progress...