Probing Fifth Force Interactions by Planetary Science Data

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Beyond the Standard Model and Fifth Force

• In BSM physics, new light, weakly-coupled degrees of freedom are introduced in the form of hypothetical particles and interactions.

 \bullet New fundamental interactions introduced by these particles \rightarrow "Fifth Force"

• These particles can arise from string theory.

Can also be dark matter or dark energy candidates.

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The Fifth Force and its Modeling

• Phenomenological model for fifth force is given by:

$$
V(r) = -\frac{GMm_*}{(1 + \alpha_1) r} \left(1 + \alpha_2 \exp\left(-\frac{r}{\lambda}\right)\right)
$$

\n
$$
\lambda \to \text{fifth force range } \left(\propto \frac{1}{\text{mediator mass}}\right)
$$

\n
$$
m_* \to \text{mass of celestial body orbiting object of mass } M
$$

\n
$$
\alpha_1, \alpha_2 \to \text{Yukawa parameters (model dependent)}
$$
\n(1)

• First term: deviation from Newtonian gravity term

$$
G \to \widetilde{G} \equiv \frac{G}{1 + \alpha_1} \tag{2}
$$

• Second term: Yukawa term due to the fifth force

$$
V(r) \propto -\alpha_2 \frac{\exp(-r/\lambda)}{r}
$$
 (3)

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Probing the Fifth Force

- Can be probed by tracking the deviations of Keplerian orbits of celestial objects.
- Many such probes are similar to the tests for General Relativity (GR).
- Deviations can be quantified in terms of orbital parameters:

 $(a, e, i, \Omega, \omega, \nu)$

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Past Literature

In the literature, constraints on the fifth force due to deviations from Newtonian trajectory are given by:

- Lunar Laser Ranging (Earth \leftrightarrow Moon) (Bergé et al., 2018)
- **■** Timing of pulsar around Sagittarius A^* (Dong et al., 2022)
- Orbital precession of:
	- Near-Earth Object (NEO) asteroids (Tsai et al., 2023)
	- Planets in the Solar system (Poddar et al., 2020)
	- S2 star around Sagittarius A[★] (Borka et al., 2016)

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Jupiter Near-Polar Orbiter (JUNO) Mission

The JUNO mission aims to perform a comprehensive study on the planet Jupiter by:

Mapping its gravitational and magnetic fields.

Exploring its structure of polar magnetosphere and auroras.

Mapping variations in atmospheric composition, temperature and cloud opacity.

Determining the abundance of water and place an upper limit on the mass of Jupiter's possible solid core.

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Jupiter Near-Polar Orbiter (JUNO) Mission

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Probing the Fifth Force with JUNO

- Jupiter's Gravity field data reconstructed in the form of:
	- **1** Central mass value \rightarrow $G \times M$
	- **2** Gravity anomaly coefficients \rightarrow J_i etc.
	- Tidal Love numbers $\rightarrow h$
		- perturbation in orbital parameters
- \bullet Fifth force perturbation \Longrightarrow $(\delta a, \delta e, \delta \Omega, \delta \omega, \delta i)$
- We constrain fifth force via the precession angle of JUNO's orbit (ω) :

$$
\Delta \omega_{\text{fifth-force}}^2 < \sum_{i=2}^{12} \left| \frac{\partial \omega}{\partial J_i} \right|^2 \sigma_{J_i}^2 \tag{4}
$$
\n
$$
\sigma_{J_i} \to \text{Uncertainty in } J_i
$$

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Unique from geocentric/ heliocentric celestial body probes.

Precession of Orbits

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Fifth Force as Perturbation Force

Perturbed force in Satellite-Normal co-ordinate system (RTN) is of the form:

$$
\vec{F} = R\hat{e}_R + N\hat{e}_N + T\hat{e}_T
$$

For Yukawa force term we have:

$$
R = -\alpha_2 \frac{GM}{(1+\alpha_1)r^2} \left(1+\frac{r}{\lambda}\right) \exp(-r/\lambda)
$$

\n
$$
N = T = 0
$$
 (5)

- For Satellite-Normal CS (R-T-N):
	- \hat{e}_R points along the Jupiter \rightarrow JUNO.
	- *e*ˆ*^N* along ~ *h* (angular momentum)
	- $-\hat{e}_T$ is defined by the right hand rule: $\hat{e}_{\tau} \cdot v > 0$

The Zonal Harmonics

- The non-spherical feature of the planet, can be accounted by expanding Newtonian potential in spherical harmonics basis.
- For our work, we only focus on zonal harmonics:

$$
U_{\text{zonal}}(r,\phi) = \frac{GM}{r} \sum_{n=2}^{\infty} J_n \left(\frac{R_J}{r}\right)^n P_n(\sin \phi) \tag{6}
$$

- **Coordinates given by:**
	- $\phi \rightarrow$ declination from equatorial plane.
	- $r \rightarrow$ radius
	- $\lambda \rightarrow$ right ascension (from a meridian)

Precession Angle

• Plugging Perturbation force in:

$$
\dot{\Omega} = \sqrt{\frac{a(1-e^2)}{\mu}} \frac{N \sin(\omega + \nu)}{\sin i(1 + e \cos \nu)}
$$

$$
\dot{\omega} = -\dot{\Omega} \cos i + \sqrt{\frac{a(1-e^2)}{e^2 \mu}} \left(-R \cos \nu + T \frac{(2 + e \cos \nu) \sin \nu}{1 + e \cos \nu} \right)
$$

Average change over an orbit:

$$
\frac{d\omega}{d\nu} = \frac{\dot{\omega}}{h/r^2}
$$
\n
$$
\Rightarrow \frac{d\omega}{d\nu}\bigg|_{AV} = \frac{1}{2\pi} \int_0^{2\pi} \frac{d\omega}{d\nu} d\nu \tag{7}
$$

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Constraints on Yukawa Parameters

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 $\mathcal{P}(\mathcal{A}) \subset \mathcal{P}(\mathcal{A})$

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- Fifth force arises due to new light, weakly-coupled degrees of freedom in BSM physics.
- They can be modeled in the form of Yukawa potential.
- Fifth force can be constrained via deviations from Newtonian orbits of various celestial objects.
- We constrain fifth-force via precession angle (ω) of JUNO's orbit around the Jupiter.
- Constraints obtained are shown to be most relevant in the force range region $\lambda \approx 10^{-4} - 10^{-2}$ AU

 $\left\{ \begin{array}{ccc} 1 & 0 & 0 \\ 0 & 1 & 0 \end{array} \right.$, $\left\{ \begin{array}{ccc} \frac{1}{2} & 0 & 0 \\ 0 & 0 & 0 \end{array} \right.$

Curious to know more? Stay tuned!

Supplementary Slides

Orbital Parameters

- \bullet $a \rightarrow$ semi-major axis (half the distance between the apojove and perijove)
- \bullet $e \rightarrow$ eccentricity of the ellipse
- $\bullet \nu \rightarrow$ defines the position of the orbiting body along the ellipse at a specific time

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Orbital Parameters

- \bullet *i* \rightarrow vertical tilt (inclination) of the ellipse with respect to the reference plane, measured at the ascending node
- $\bullet \omega \rightarrow$ angle measured from the ascending node to the perijove of the orbit
- $\bullet \Omega \rightarrow$ Longitude of the ascending node

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Orbital Parameter Perturbations

Taking Newtonian orbit to be ellipse (unperturbed): $r=\frac{1}{2}$ $\frac{a(1-e^2)}{b}$ $1\!+\!e\cos\nu$ Perturbation in orbital parameters:

$$
\dot{a} = 2\sqrt{\frac{a^3}{\mu(1 - e^2)}} [eR \sin \nu + T(1 + e \cos \nu)]
$$

\n
$$
\dot{e} = \sqrt{\frac{a(1 - e^2)}{\mu}} \left[R \sin \nu + T \left(\cos \nu + \cos \left(\sqrt{\frac{1 - e}{1 + e}} \tan \frac{\nu}{2} \right) \right) \right]
$$

\n
$$
\frac{d}{dt} i = \sqrt{\frac{a(1 - e^2)}{\mu}} \frac{N \cos(\omega + \nu)}{1 + e \cos \nu}
$$

\n
$$
\dot{\Omega} = \sqrt{\frac{a(1 - e^2)}{\mu}} \frac{N \sin(\omega + \nu)}{\sin i(1 + e \cos \nu)}
$$

\n
$$
\dot{\omega} = -\dot{\Omega} \cos i + \sqrt{\frac{a(1 - e^2)}{e^2 \mu}} \left(-R \cos \nu + T \frac{(2 + e \cos \nu) \sin \nu}{1 + e \cos \nu} \right)
$$

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- For Satellite-Normal CS (R-T-N):
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	- *e*ˆ*^N* along ~ *h* (angular momentum)
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The Zonal Harmonics

- Non-spherical feature of the planet, can be accounted by expanding \bullet Newtonian potential in sphere harmonics basis.
- For our work, we only focus on zonal harmonics:

$$
U_{\text{zonal}}(r,\phi) = \frac{GM}{r} \sum_{n=2}^{\infty} J_n \left(\frac{R_J}{r}\right)^n P_n(\sin \phi) \tag{9}
$$

- **Coordinates given by:**
	- $\phi \rightarrow$ declination from equatorial plane.
	- $r \rightarrow$ radius
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Precession due to *Jⁿ* Perturbation

$$
U_{J_n}(r,\phi) = \frac{GM}{r}J_n\left(\frac{R_J}{r}\right)^n P_n(\sin\phi)
$$
 (10)

- In Jupiter-centered inertial (JCI) frame: $\phi = \frac{z}{r}$ *r*
- Perturbation force due to *Jⁿ* term:

$$
\vec{F} = -\frac{\partial U_{J_n}}{\partial r} \hat{e}_R + \frac{\partial U_{J_n}}{\partial z} \hat{e}_z \tag{11}
$$

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 \bullet For JCI \longrightarrow RTN frame:

$$
\hat{e}_z = \sin i \sin(\omega + \nu) \hat{e}_R + \sin i \cos(\omega + \nu) \hat{e}_T + \cos i \hat{e}_N
$$

\n
$$
z = r \sin \phi = r \sin i \sin(\omega + \nu)
$$
 (12)

• For brevity: $\theta = \omega + \nu$

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$JCI \longrightarrow RTN$ Coordinates

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Precession Angle

• Plugging Perturbation force in:

$$
\dot{\Omega} = \sqrt{\frac{a(1-e^2)}{\mu}} \frac{N \sin(\omega + \nu)}{\sin i(1 + e \cos \nu)}
$$

$$
\dot{\omega} = -\dot{\Omega} \cos i + \sqrt{\frac{a(1-e^2)}{e^2 \mu}} \left(-R \cos \nu + T \frac{(2 + e \cos \nu) \sin \nu}{1 + e \cos \nu} \right)
$$

Average change over an orbit:

$$
\frac{d\omega}{d\nu} = \frac{\dot{\omega}}{h/r^2}
$$
\n
$$
\Rightarrow \frac{d\omega}{d\nu}\bigg|_{AV} = \frac{1}{2\pi} \int_0^{2\pi} \frac{d\omega}{d\nu} d\nu \tag{13}
$$

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The Harmonics

Spherical harmonic expansion of potential is given by:

$$
U(\phi, \lambda, r) = \frac{GM}{r} + U_{\text{zonal}}(r, \phi) + U_{\text{sectorial}}(r, \phi, \lambda)
$$

+ U_{\text{tesseral}}(r, \phi, \lambda) (14)

 \leftarrow \Box

- Coordinates given by:
	- $-\phi \rightarrow$ declination from equatorial plane.
	- $r \rightarrow$ radius
	- $\lambda \rightarrow$ right ascension (from a meridian)

$$
U_{\text{zonal}}(r,\phi) = \frac{GM}{r} \sum_{n=2}^{\infty} J_n \left(\frac{R_J}{r}\right)^n P_n(\sin \phi) \tag{15}
$$

- \bullet $P_n \rightarrow$ Legendre Polynomial
- Zonal harmonics vary only with latitude.

$$
U_{\text{sect}}(r, \phi, \lambda) = \frac{GM}{r} \sum_{n=2}^{\infty} (C_{n, \text{ sect}} \cos(n\lambda) + S_{n, \text{sect}} \sin(n\lambda)) \left(\frac{R_J}{r}\right)^n P_n(\sin \phi)
$$
(16)

- Divides globe into slices by longitude.
- Varies only with longitude.

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$$
U_{\text{tesseral}}(r,\phi,\lambda) = \frac{\mu}{r} \sum_{i,j=2}^{\infty} (C_{i,j} \cos(i\lambda) + S_{i,j} \sin(i\lambda)) \left(\frac{R_j}{r}\right)^i P_{i,j} \left(\sin \phi\right)
$$
\n(17)

Divides globe into slices by longitude and latitude.

28/28

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