

Kinetic Simulations of Collisionless Shock Formation in the Dark Sector

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Long Range Effects

- Self-interacting dark matter is not only $2 \rightarrow 2$ scattering
- >99.9% of visible matter in the universe is a plasma, governed by many → many scattering
- Long range collective effects can probe many orders of magnitude deeper into parameter space





Current Constraints

Some of the strongest 2→2 constraints come from dissociative cluster mergers such as the Bullet Cluster [1]
 σ/m ≤ 1 cm²g⁻¹

➢ Main Observables

- Evaporation of dark matter halo
- Offset of dark matter and standard model centers



Credit: European Space Agency



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Collisionless Regime

>Introduce model

$$\mathcal{L} = \mathcal{L}_{SM} - \frac{1}{4} F'_{\mu\nu} F'^{\mu\nu} + \bar{\chi} (\gamma^{\mu} (i\partial_{\mu} - qA') - m_{\chi}) \chi$$

Size of Bullet Cluster core ~100 kpc
Mean free path of dark matter

$$\lambda \sim 300 \operatorname{kpc} \left(\frac{v_{rel}}{0.01c}\right)^4 \left(\frac{q_{\chi}}{q_e}\right)^{-4} \left(\frac{m_{\chi}}{\operatorname{GeV}}\right)^3 \left(\frac{\rho_{\chi}}{0.01 \operatorname{GeV/cm^3}}\right)$$





Plasma Dynamics

≻Vlasov Equation

$$\left(\partial_t + \frac{q_s}{m_s} (\mathbf{E} + \mathbf{v} \times \mathbf{B}) \cdot \nabla_v + \mathbf{v} \cdot \nabla_x\right) f_s(\mathbf{x}, \mathbf{v}, t) = 0$$

≻Linear Regime

Analytical estimates predict growth rates and saturation times of instabilities

≻Nonlinear Regime

- Analytical estimates break down as perturbations grow
- In order to determine dynamics over long timescales, simulations are needed



Simulations

> Plasma frequency:
$$\omega_{\chi} = \sqrt{\frac{4\pi q_{\chi}^2 n_{0,\chi}}{m_{\chi}}} = \frac{q_{\chi}}{m_{\chi}} \sqrt{4\pi\rho_{\chi}}$$

"Smilei is a Particle-In-Cell code for plasma simulation. Open-source, collaborative, userfriendly and designed for high performances on super-computers, it is applied to a wide range of physics studies: from relativistic laser-plasma interaction to astrophysics."^[2]





Plasma Shocks



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Plasma Shocks





Plasma Shocks





Effective Collisions



- After the electrostatic shock saturates, the beam's bulk velocity drastically decreases.
- The transverse electromagnetic mode continues to grow on very large timescales creating long filaments of strong magnetic fields.
- These fields saturate when the exponential growth rate approaches the bounce frequency

$$\omega_B = \left(\frac{q_\chi k_{max}}{m_\chi} v_T B\right)^{1/2}$$

Regions of strong magnetic fields create Bohm-like diffusion

$$\gamma_{eff} \sim \frac{q_{\chi}}{m_{\chi}} B$$









Conclusions

- Collective effects have potential to constrain several orders of magnitude of parameter space.
- >Simulations are necessary to understand nonlinear behavior of plasmas.
- ➢ With more precise treatment and stronger computational power, further constraints can be placed.
- ≻Future studies
 - Instabilities triggered from background magnetic fields
 - Exploration of other models such as millicharged particles



References

- [1] A. Robertson, R. Massey, V. Eke, What does the Bullet Cluster tell us about self-interacting dark matter?, MNRAS, 465, 569-587 (2017)
- [2] J. Derouillat, A. Beck, F. Pérez, T. Vinci, M. Chiaramello, A. Grassi, M. Flé, G. Bouchard, I. Plotnikov, N. Aunai, J. Dargent, C. Riconda, M. Grech, *SMILEI: a collaborative, open-source, multi-purpose particle-in-cell code for plasma simulation*, Comput. Phys. Commun. 222, 351-373 (2018)
- [3] P. Agrawal, F-Y. Cyr-Racine, L. Randall, J. Scholtz, Make Dark Matter Charged Again, JCAP, 2017, 5 (2017)



Longitudinal Instabilities





Transverse Instabilities

- Small perturbations in the transverse magnetic field attract particles to nodes
- Current sheets form as particles collect near nodes
- Current sheets induce a magnetic field that strengthens the initial perturbation
- Expected growth rate: $\gamma_W \approx v_{rel} \omega_{\chi}$







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