

# Amplifying CMB Phase Shift with Dark Matter-Radiation Interactions

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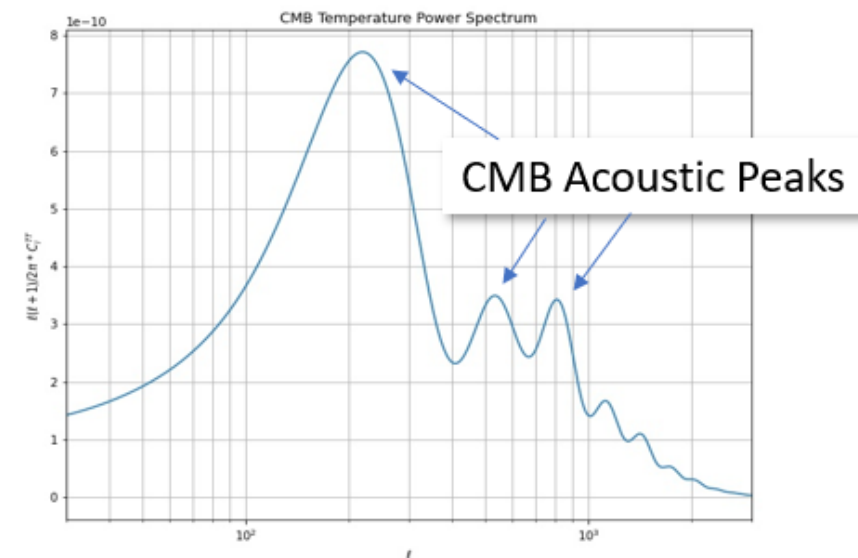
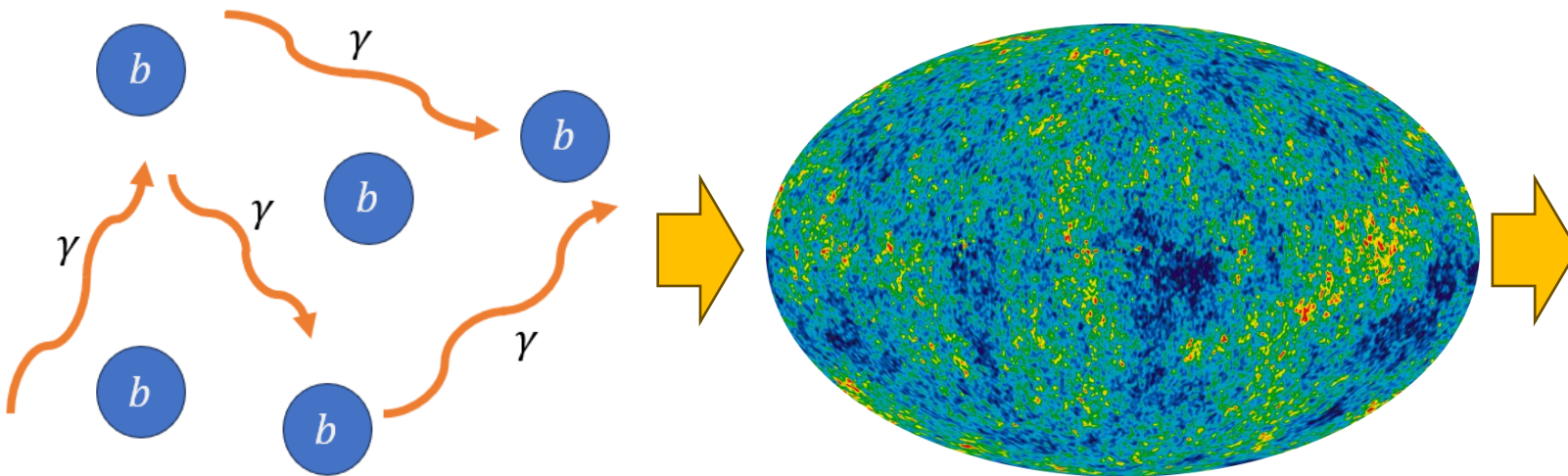


In collaboration with Subhajit Ghosh and Yuhsin Tsai

DPF-PHENO 2024, University of Pittsburgh

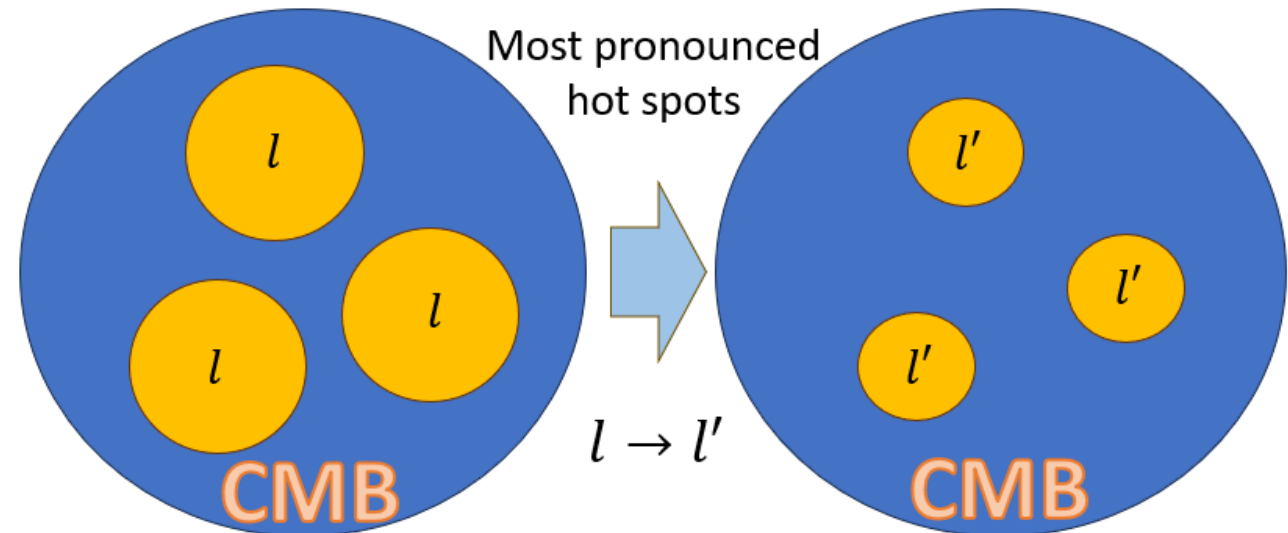
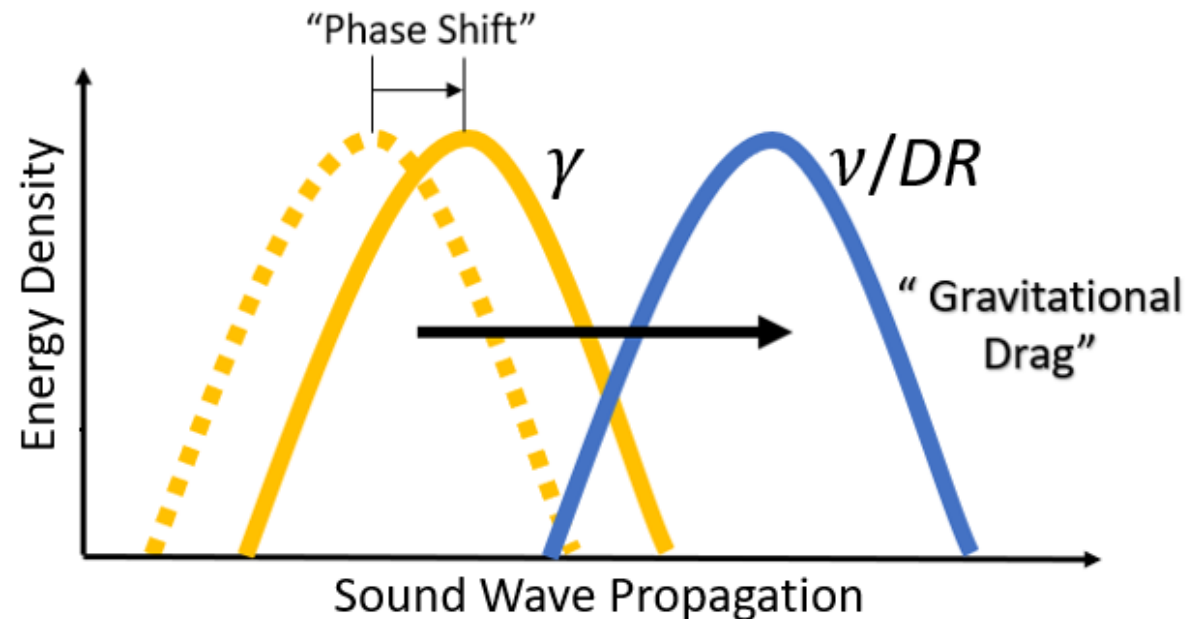
# CMB Acoustic Peaks

- Sound waves propagating in photon-baryon plasma before recombination leaves peak structure in CMB power spectrum.
- Consider the phase of acoustic oscillations: shift in phase manifests as “shift” in CMB peaks (to be clarified later)
- **Phenomenology**: what kind of physics can produce phase shift?



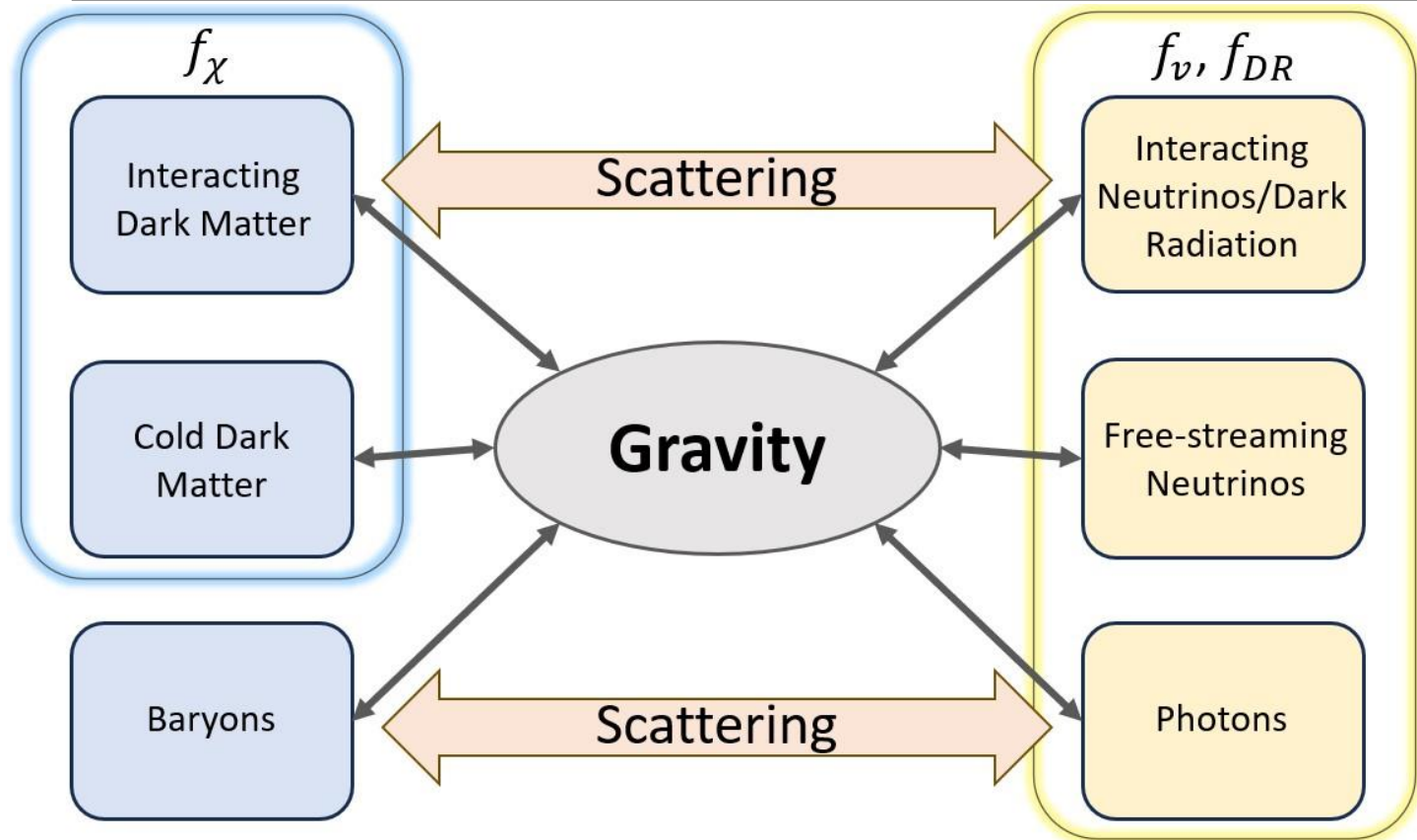
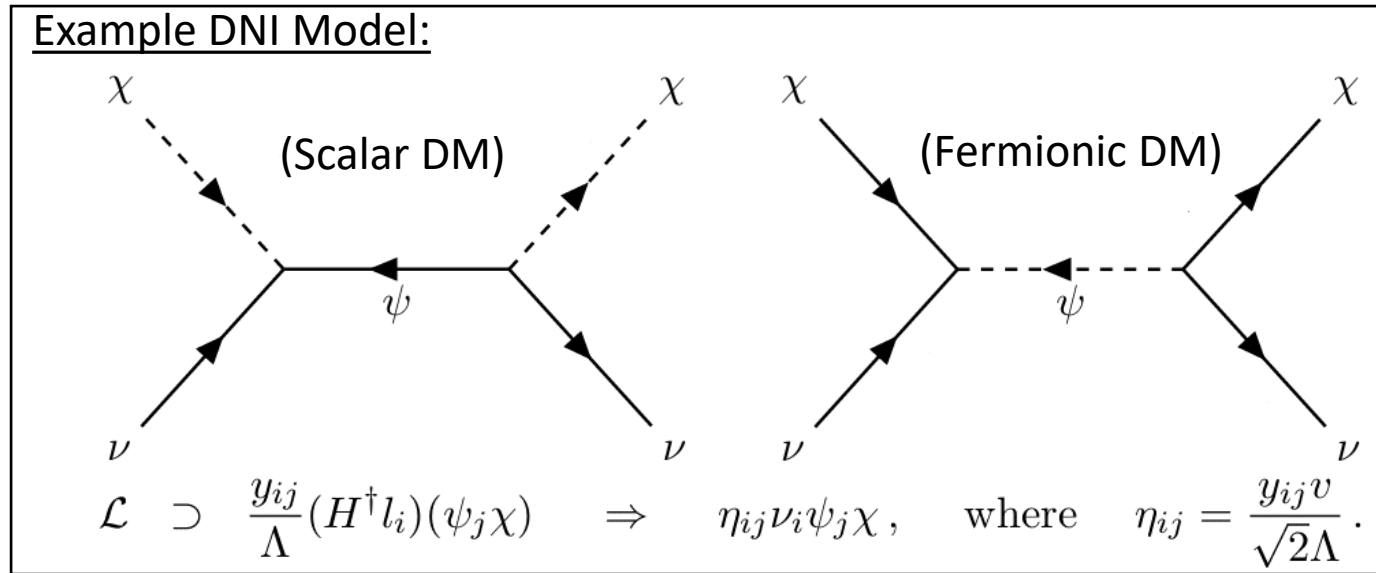
# Phase Shift in the CMB

- CMB phase shift is sensitive to **propagation behaviour of non-photon radiation** (e.g. SM neutrinos, light dark photon...) before recombination
- Non-photon radiation exerts “gravitational drag” on photon-baryon waves: sensitive to physics that interact **only gravitationally with us** (i.e. no other interaction with SM)
- Studied before in the context of free-streaming vs self-interacting radiation  
(Bashinsky & Seljak [arXiv:astro-ph/0310198](https://arxiv.org/abs/astro-ph/0310198), Baumann et. al. [arXiv:1508.06342v3](https://arxiv.org/abs/1508.06342v3))



# Dark Matter-Radiation Interactions

- Phase shift effect can be amplified compared to self-interacting radiation scenario
- Consider multi-component dark matter and radiation sectors. Scattering of radiation comes from interaction with DM
- **Demonstrative example:** let (massless) neutrinos play the role of interacting radiation first



# Dark Matter Loading

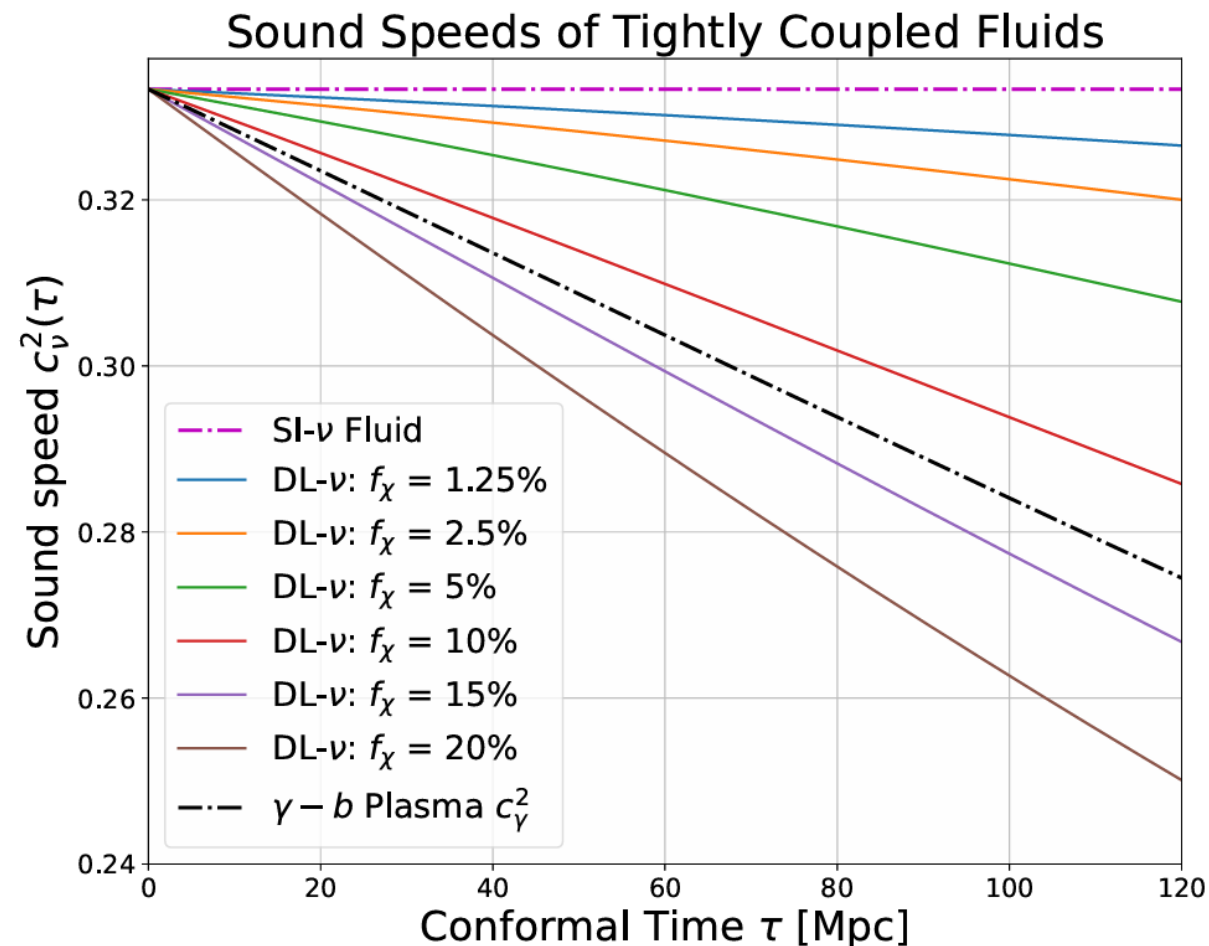
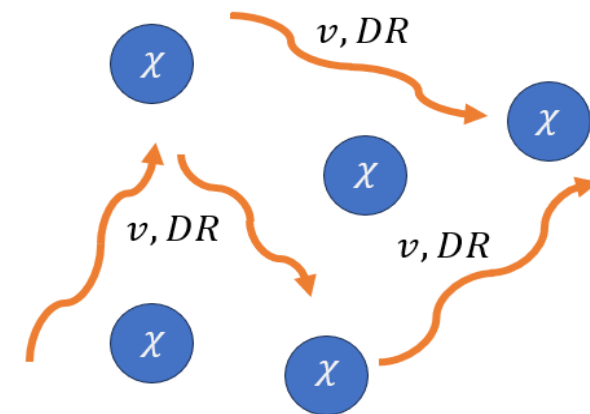
- **Efficient scattering:** scattering rate large compared to Hubble rate
- Interacting radiation (r) and matter (m) forms *tightly-coupled* fluid, with sound speed

$$c_r^2 = \frac{1}{3(1+R_r)}, \quad \text{where} \quad R_r = \frac{3 \rho_m}{4 \rho_r}$$

- Matter-loading effect suppresses sound speed over time; larger suppression for larger  $f_\chi$

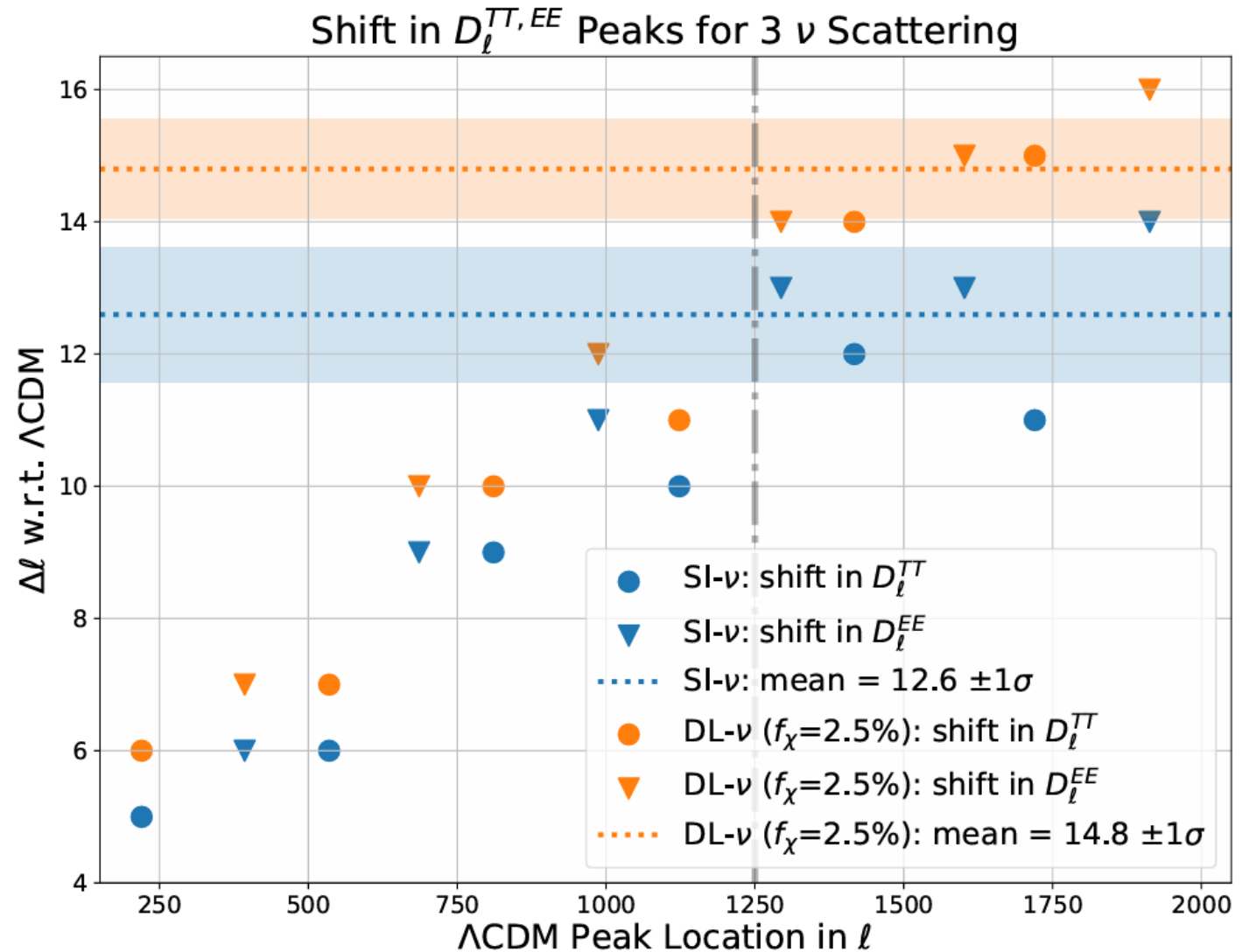
## Radiation propagation behaviour:

1. Free-Streaming (FS)
2. Self-Interacting (SI)
3. Dark-matter Loading (DL)



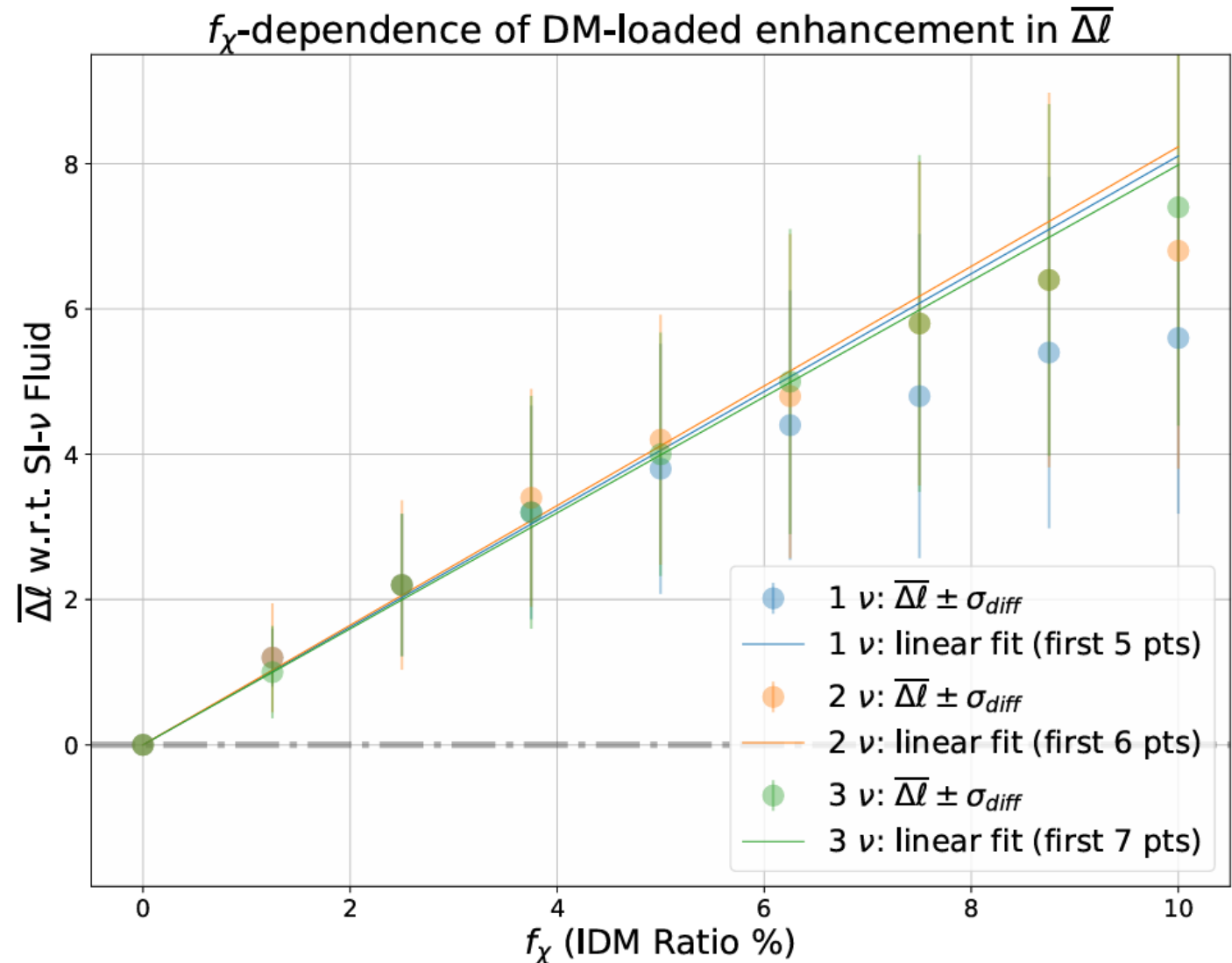
# Amplifying CMB Phase Shift

- Calculate shift in CMB peaks w.r.t.  $\Lambda$ CDM model using CLASS (all neutrinos FS)
- Peaks shift to positive  $l$  for SI neutrino, shift enhanced further for DL neutrino ( $f_\chi = 2.5\%$ )



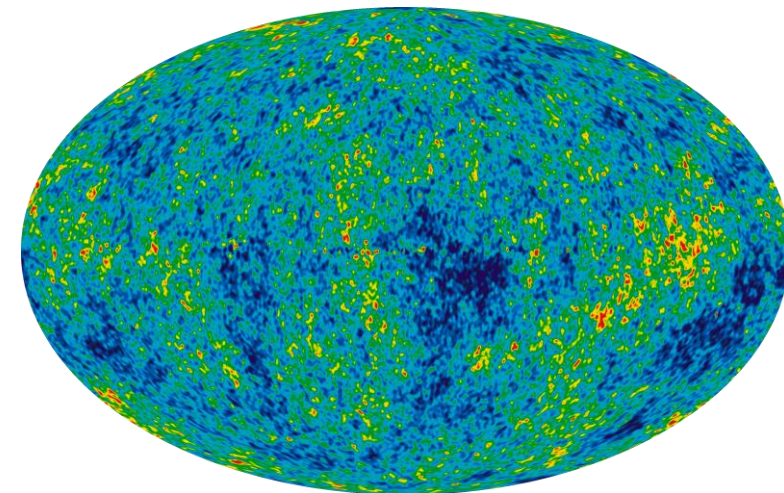
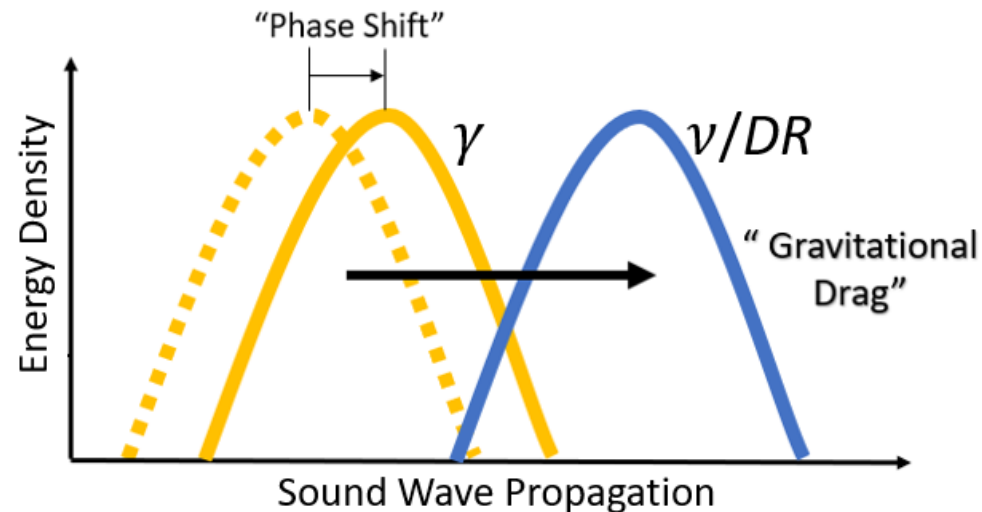
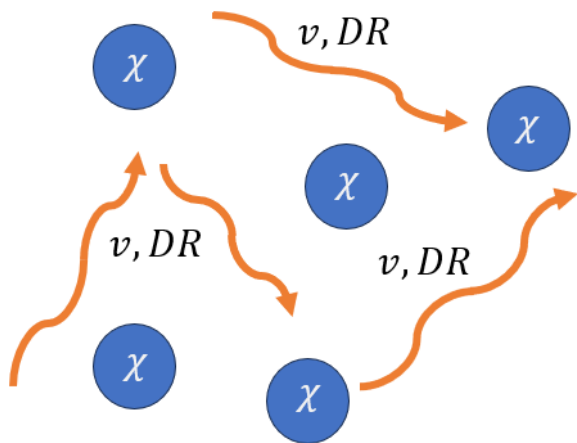
# Amplifying CMB Phase Shift

- Vary number of interacting neutrinos: overall shift w.r.t.  $\Lambda$ CDM, but similar DM-loading enhancement
- DL vs SI enhancement has linear dependence on  $f_\chi$  and independent of  $f_\nu$  (for small  $f_\chi$  vs  $f_\nu$ )



# Brief Outline

- Numerical calculations (CLASS) show CMB phase shift amplified by dark matter loading
- Further questions for this talk:
  1. What's going on? Understand mechanism with simplistic toy model
  2. What to look for? Study observability with more realistic model



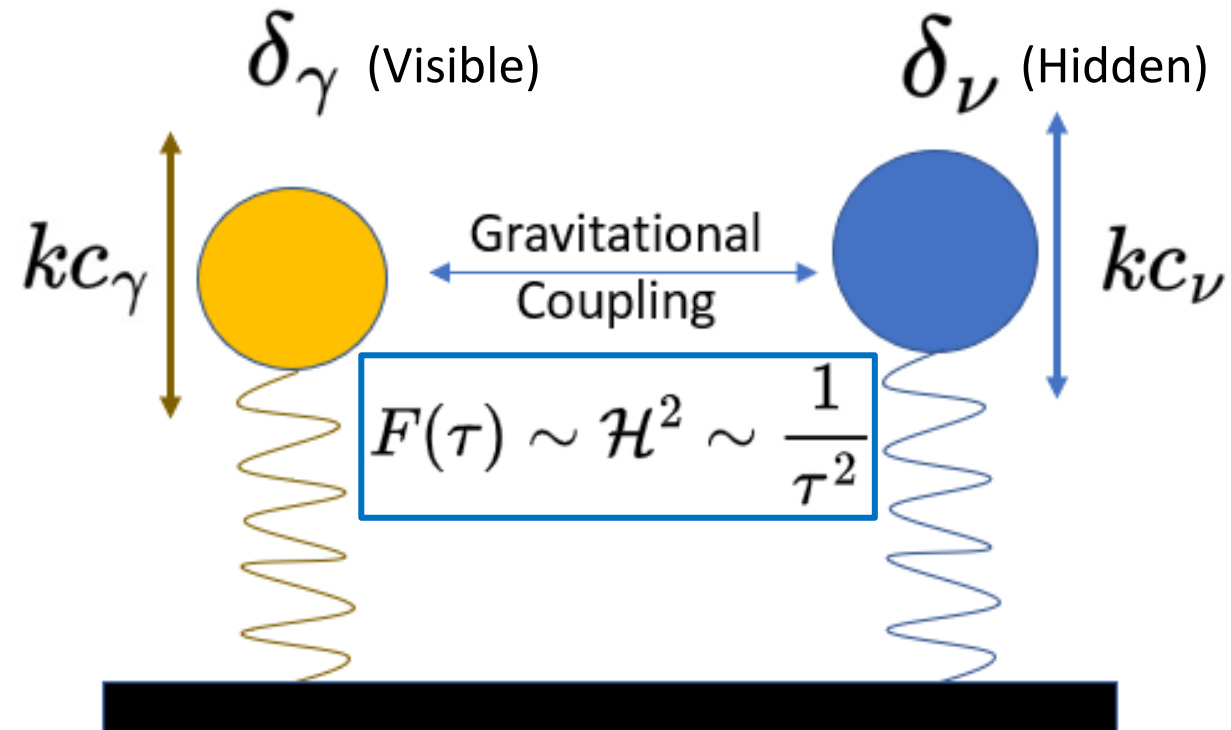


# Toy Model: Coupled Oscillators

1. Two tightly-coupled fluids: photon-baryon and neutrino-DM, described by radiation energy density contrast  $\delta_r$
2. Fluids carry acoustic oscillations suppressed by **matter-loading**, interact with each other **only gravitationally**
3. Phase shift imprinted in photons by hidden oscillator: size and direction of shift depends on relative sound speed
4. Gravitational interaction strength decreases over time with Hubble: phase shift gets “fixed”

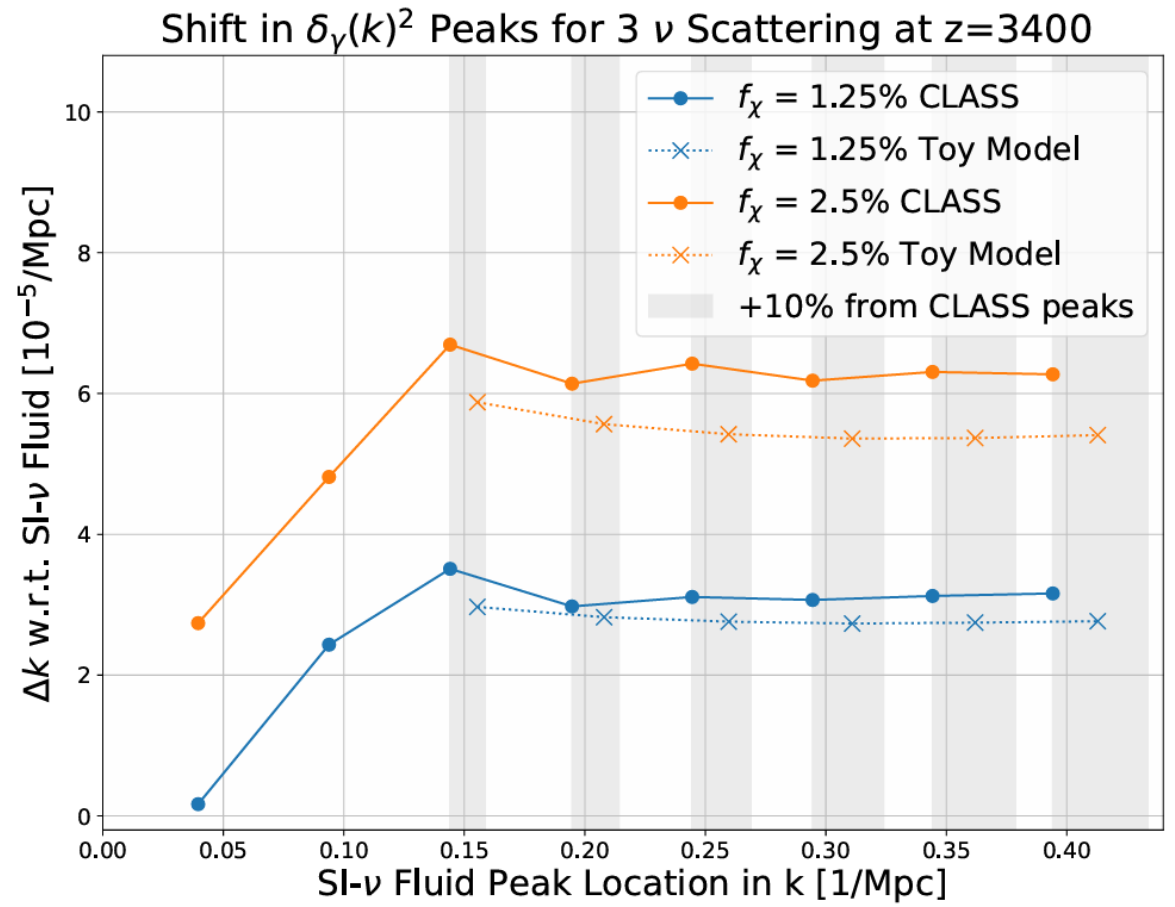
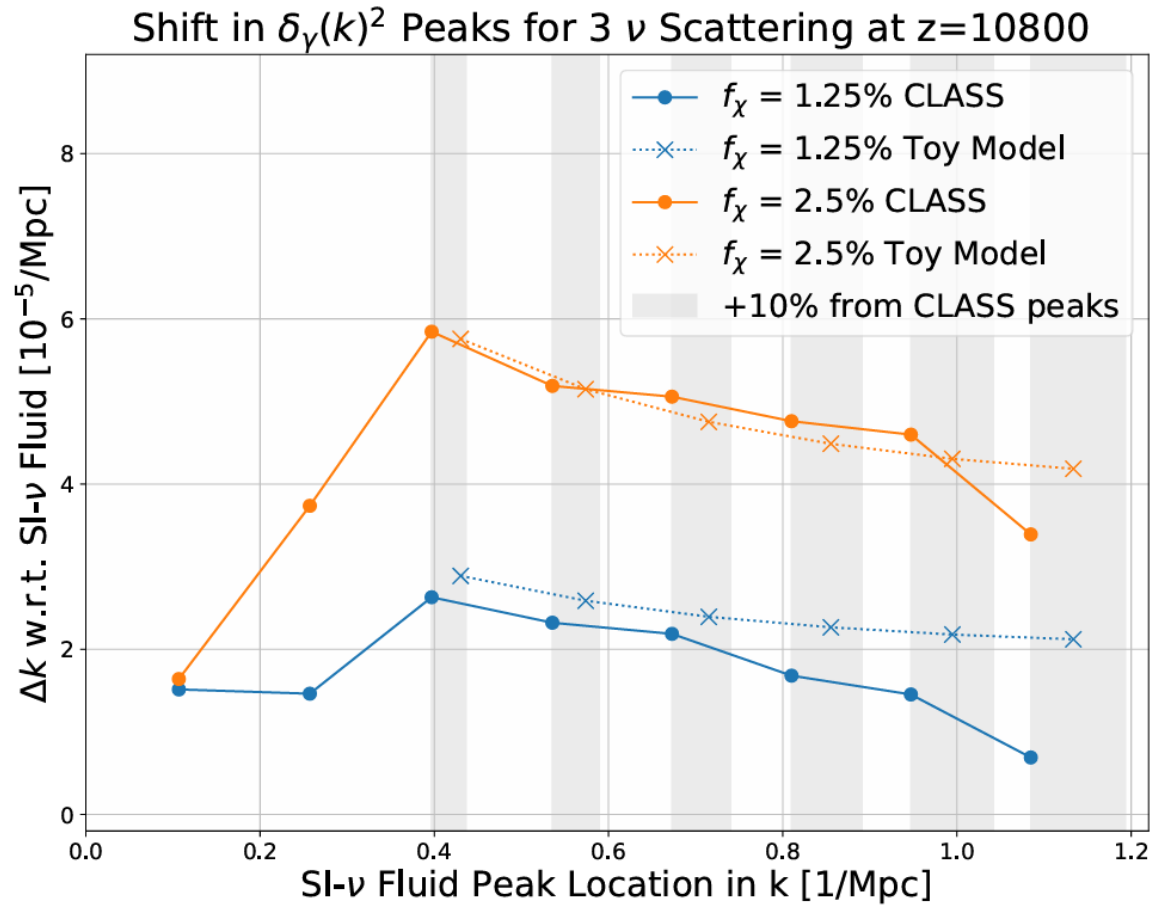
Gravitational Coupling

$$\ddot{\delta}_\gamma + k^2 c_\gamma^2 \delta_\gamma = F(\tau)(f_\gamma \delta_\gamma + f_\nu \delta_\nu)$$
$$\ddot{\delta}_\nu + k^2 c_\nu^2 \delta_\nu = F(\tau)(f_\gamma \delta_\gamma + f_\nu \delta_\nu)$$



# Toy Model Approximates CLASS Well

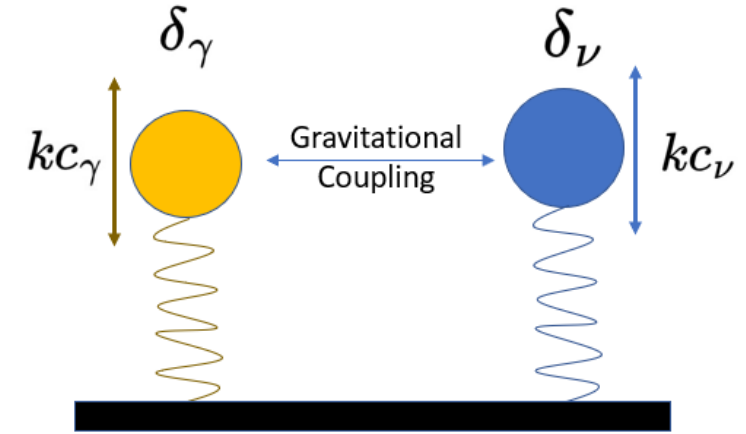
- Numerical check: toy model captures shift in photon transfer function from CLASS  
(*details omitted*)



# Simple Parametric Dependence

$$\ddot{\delta}_\gamma + k^2 c_\gamma^2 \delta_\gamma = F(\tau)(f_\gamma \delta_\gamma + f_\nu \delta_\nu)$$

$$\ddot{\delta}_\nu + k^2 c_\nu^2 \delta_\nu = F(\tau)(f_\gamma \delta_\gamma + f_\nu \delta_\nu)$$



(i) Small coupling: homogeneous solutions

$$\delta_\gamma \sim \cos(c_\gamma k \tau)$$

$$\delta_\nu \sim \cos(c_\nu k \tau)$$

(ii) Small matter-loading: small deviation

$$c_\gamma - c_\nu = \delta c \ll 1$$

$$\delta c \sim R_\nu \sim f_\chi / f_\nu$$

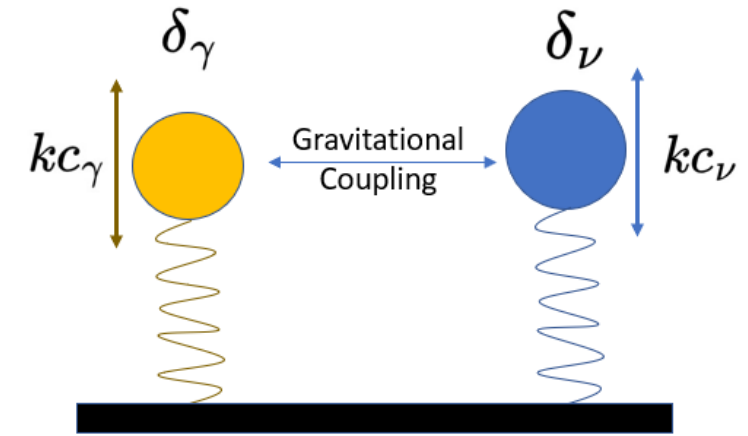
Phase shift  $\Delta\phi$  gets imprinted in homogeneous photon oscillations by (perturbative) gravitational influence of neutrinos

$$\delta_\gamma \sim \cos(c_\gamma k \tau - \Delta\phi)$$

# Simple Parametric Dependence

$$\ddot{\delta}_\gamma + k^2 c_\gamma^2 \delta_\gamma = F(\tau) (f_\gamma \delta_\gamma + f_\nu \delta_\nu)$$

$$\ddot{\delta}_\nu + k^2 c_\nu^2 \delta_\nu = F(\tau) (f_\gamma \delta_\gamma + f_\nu \delta_\nu)$$



Gravitational  
"driving force"

$$f_\gamma \cos(c_\gamma k \tau) + f_\nu \cos((c_\gamma - \delta c) k \tau) \approx \cos(c_\gamma k \tau) + f_\nu \delta c \sin(c_\gamma k \tau)$$

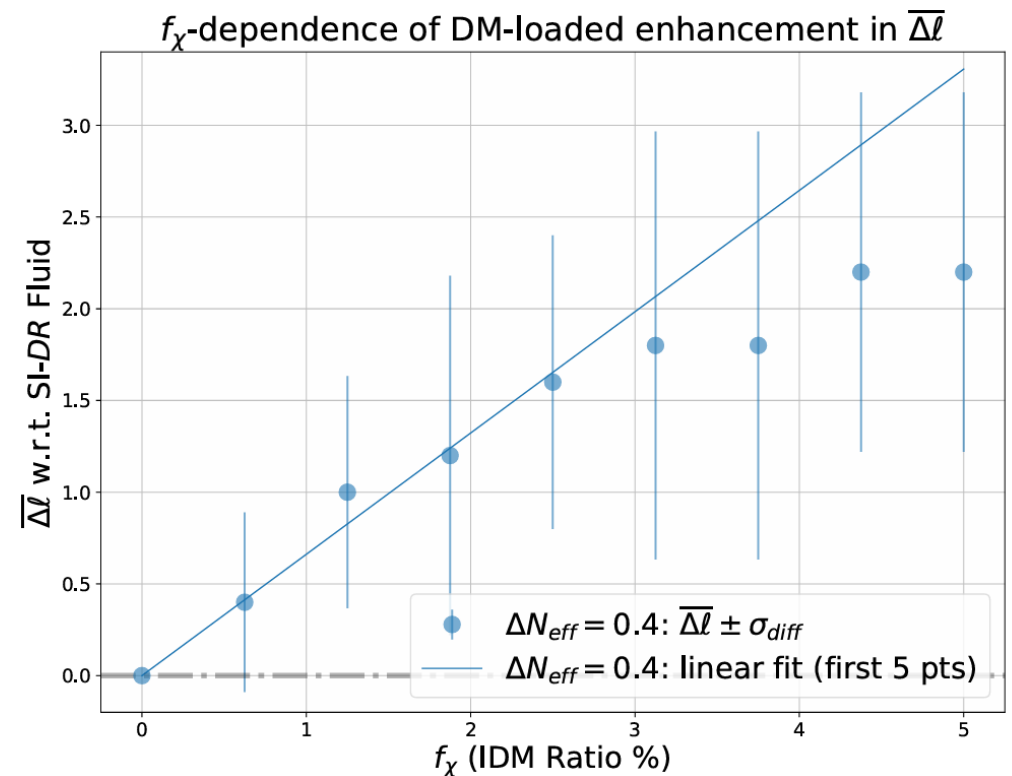
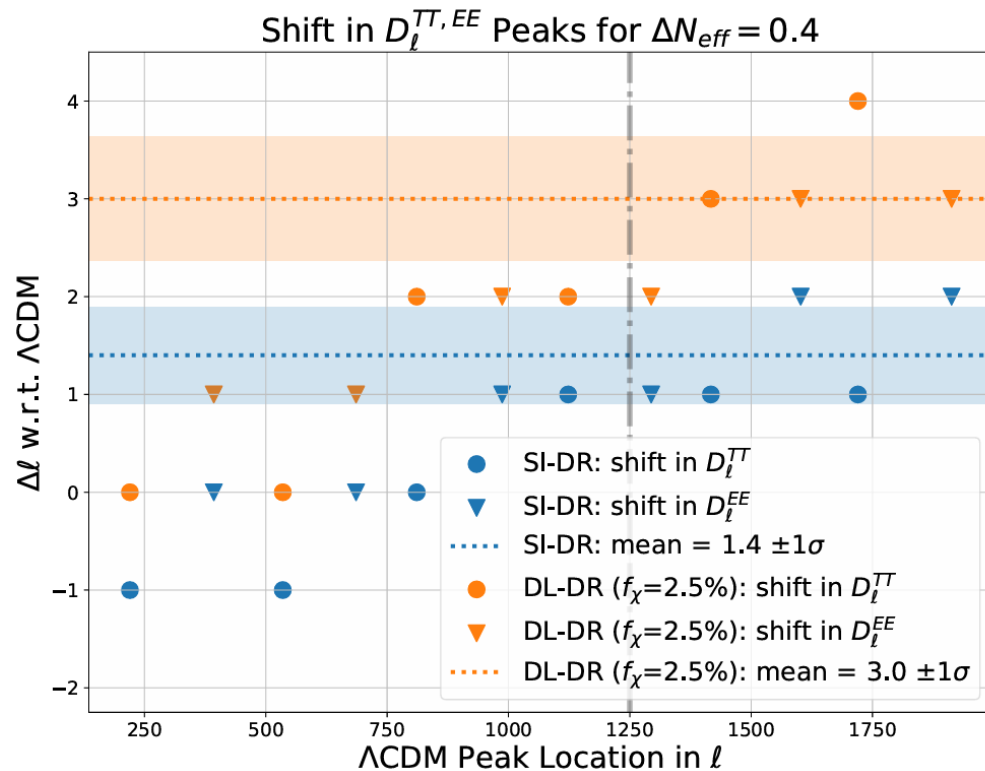
Phase shift to  $\delta_\gamma \sim \cos(kc_\gamma \tau)$  comes from sine part, which is linear in  $f_\chi$  and independent of  $f_\nu$

$$\Delta\phi \sim f_\nu \delta c \sim f_\nu \left( \frac{f_\chi}{f_\nu} \right) \sim f_\chi$$

$\Delta\phi$ : *relative shift w.r.t. self-interacting case*

# Interacting Dark Radiation Model

- Additional dark radiation component ( $\Delta N_{eff}$ ) scatters efficiently with DM
- All neutrinos free-streaming, new physics only in dark sector: interacts with Standard Model only gravitationally
- Similar phase shift amplification and parametric dependence (*details omitted*)



# Signal: Angular Sound Horizon $\theta_s$

- CMB peak positions well measured (e.g. Planck2018)
- Degeneracy in parameters: how far away is the CMB (2D surface of last scattering)?
- Previous CLASS analysis: fixed  $\theta_s \Rightarrow$  peaks shifted
- But peak positions fixed by data: determine phase shift  $\Delta\phi$  from fitting  $\theta_s$

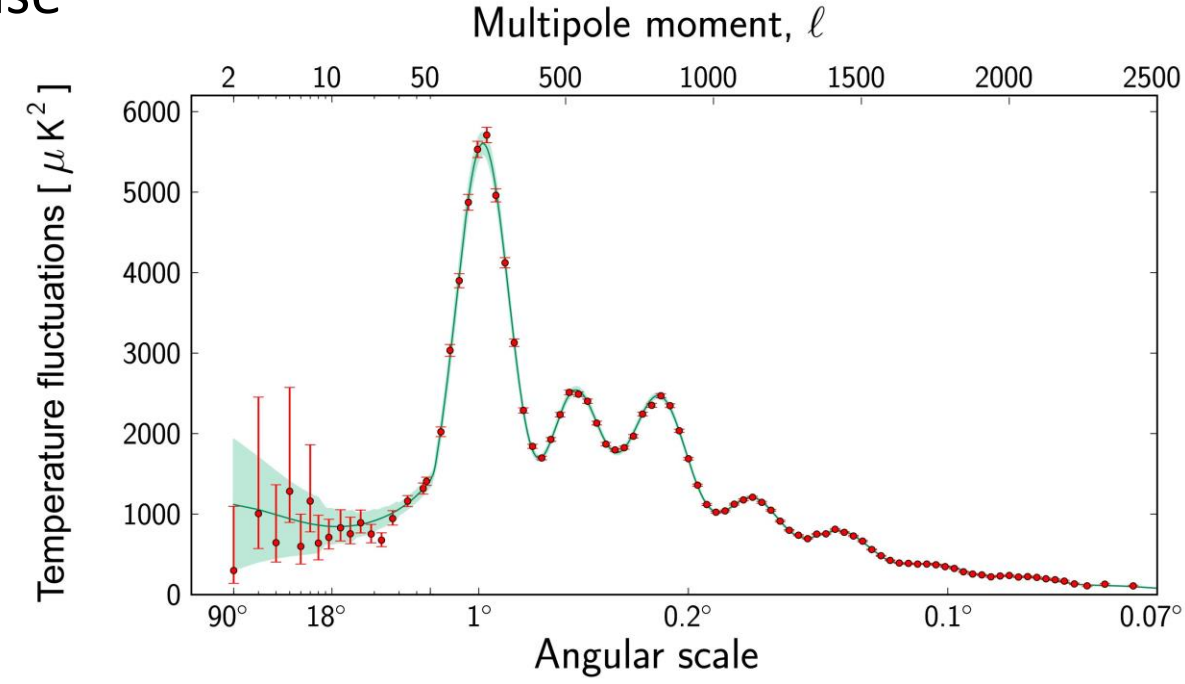
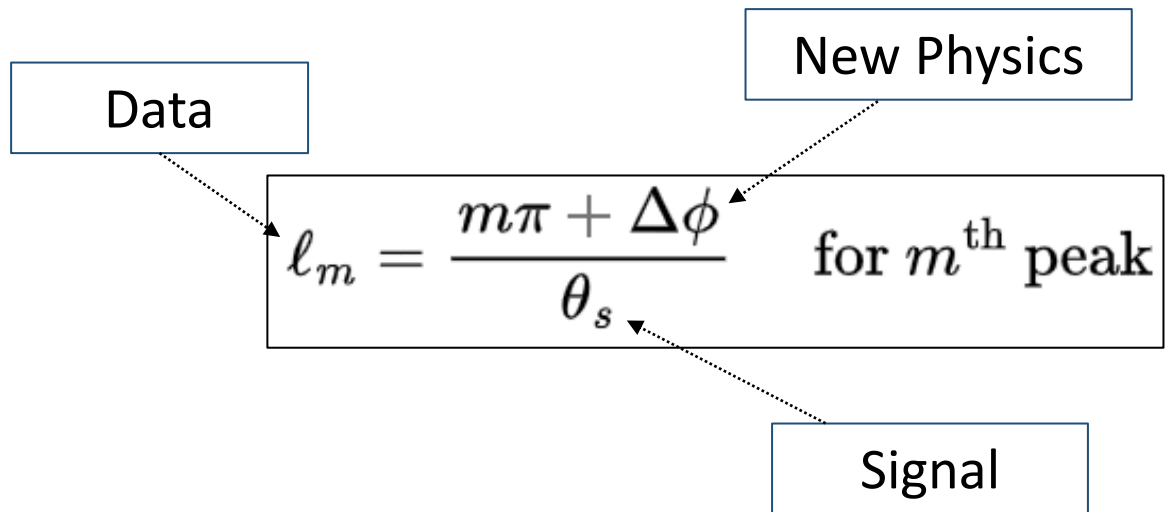
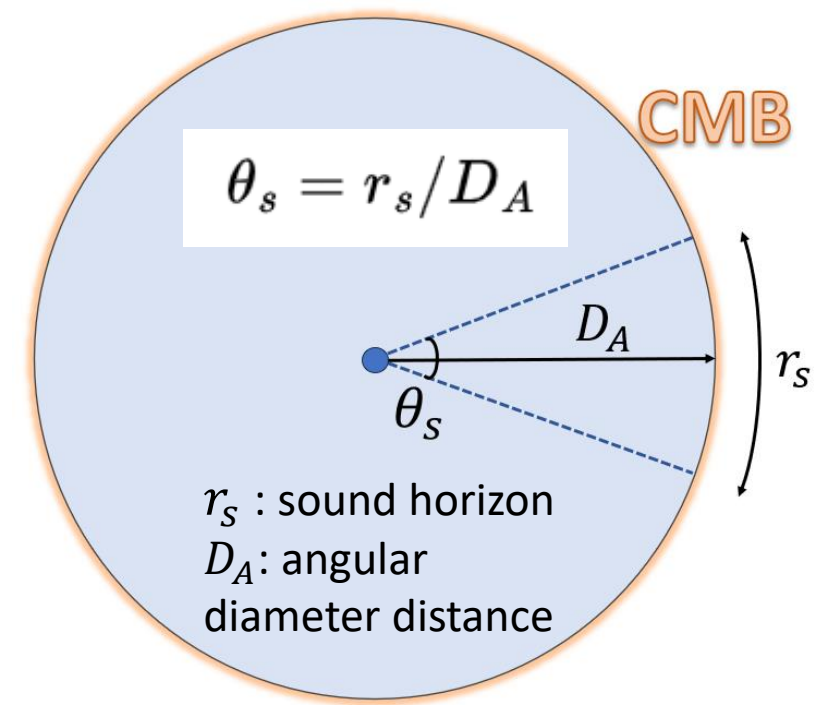
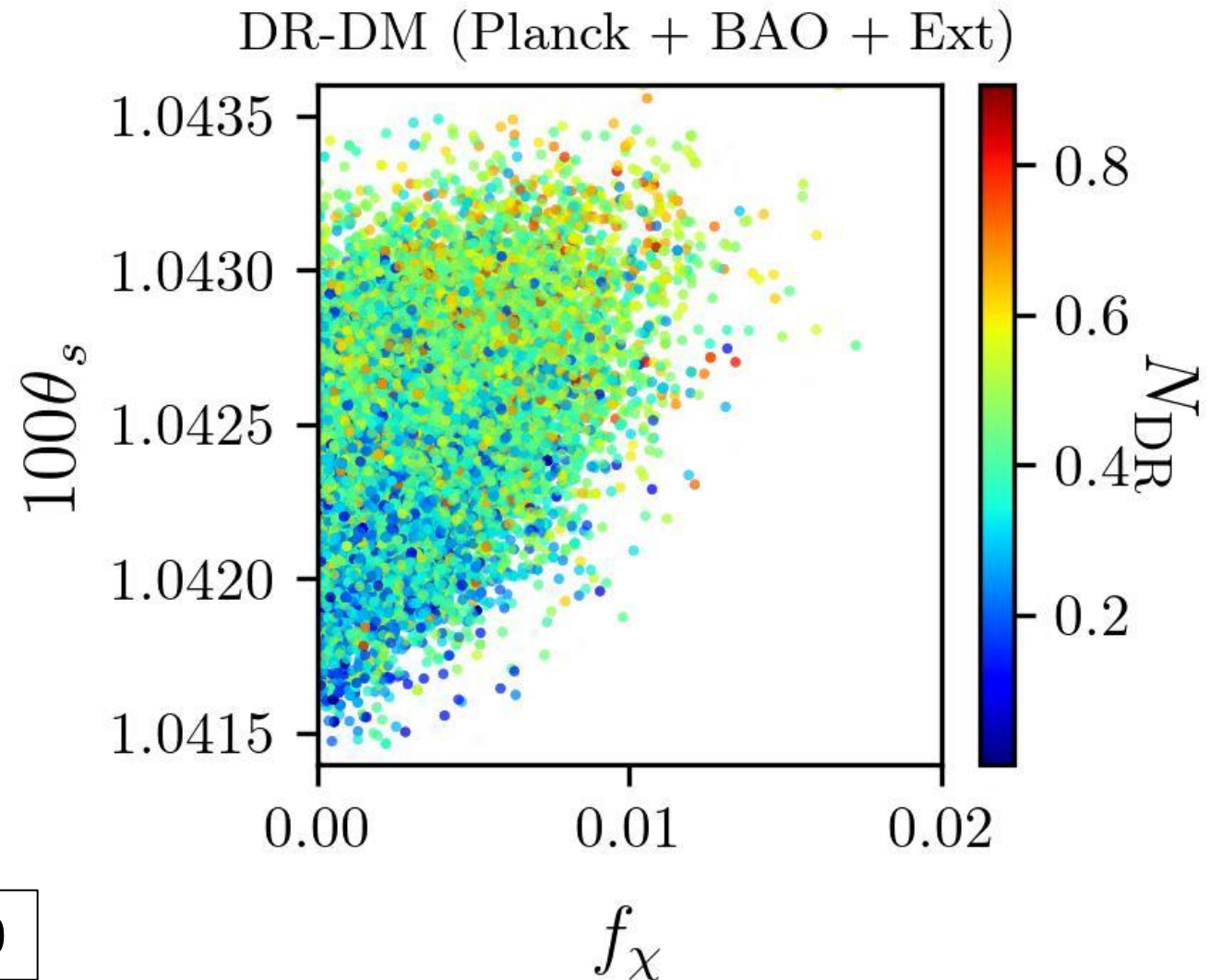


Image: ESA/Planck Collaboration – Planck Power Spectrum

# MCMC Analysis: Signature in $\theta_s$

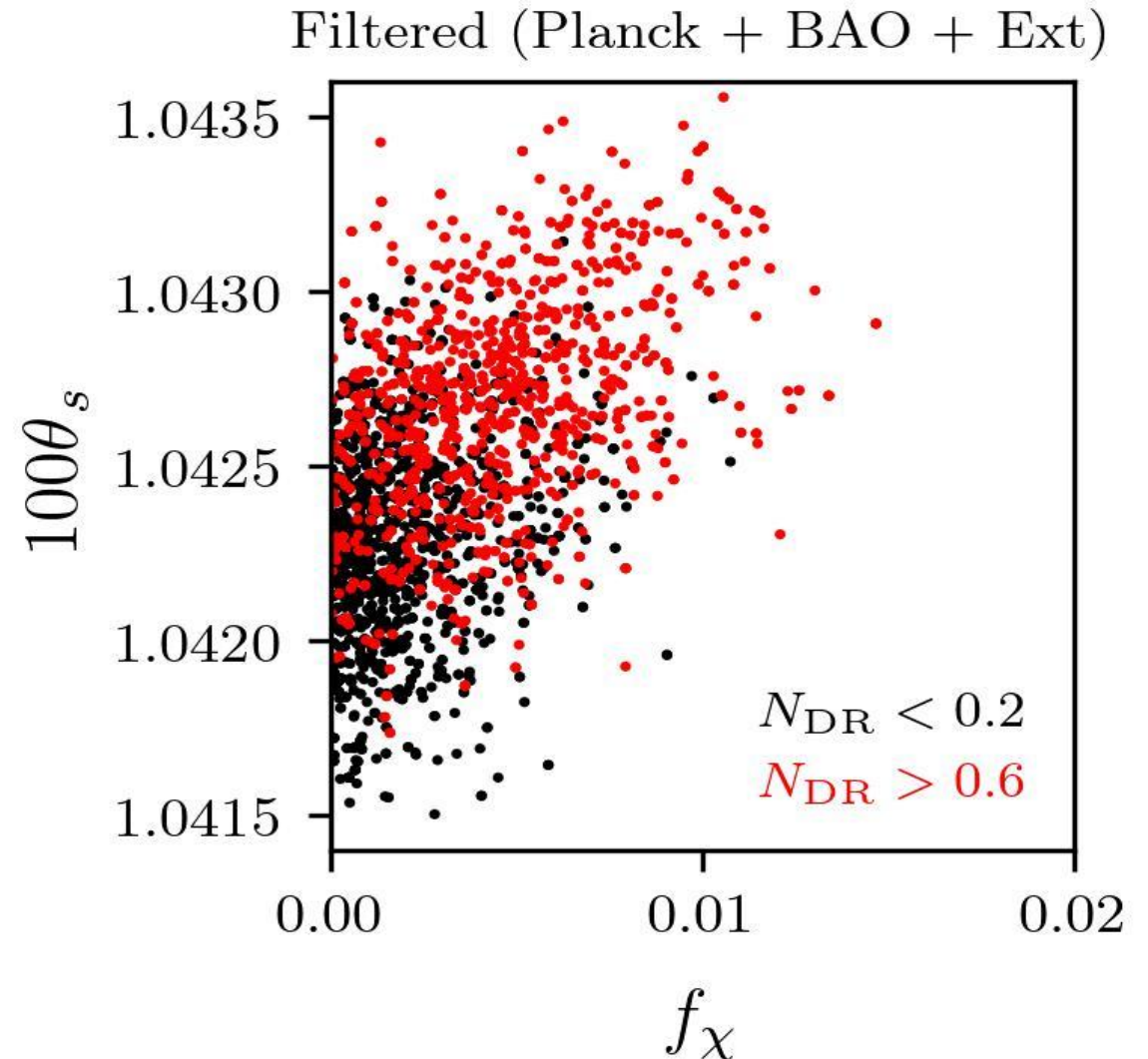
- Use Montepython to fit model. Allow amount of interacting DR ( $N_{DR}$ ) to vary
- DM-loading signature: angular sound horizon  $\theta_s$  positively correlated with interacting DM fraction  $f_\chi$
- For comparison:  $N_{DR} = 0$  corresponds to  $\Lambda$ CDM and the  $f_\chi \rightarrow 0$  limit for each  $N_{DR} > 0$  approaches the self-interacting DR scenario



Datasets: Planck2018 + BAO + SH0ES + kv450

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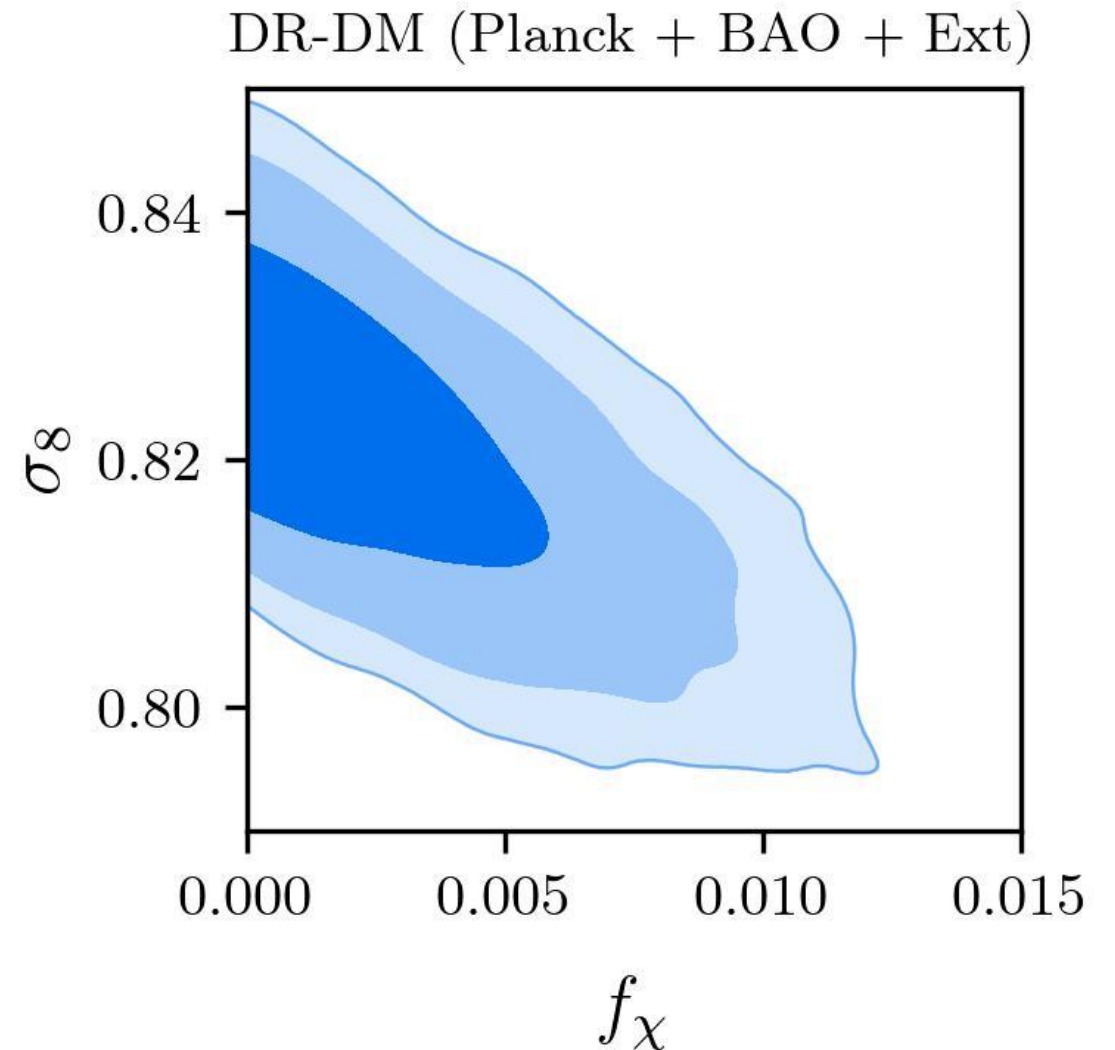
Datasets: Planck2018 + BAO + SH0ES + kv450



# MCMC Analysis: Dual signature in $\sigma_8$

- $\sigma_8$  parameter measures amplitude of matter density fluctuations on scales  $k \sim 8 h/\text{Mpc}$
- From matter POV, scattering with radiation interferes with clumping/structure formation
- **Dual signature**:  $\sigma_8$  suppression appears alongside  $\theta_s$  enhancement with increasing  $f_\chi$

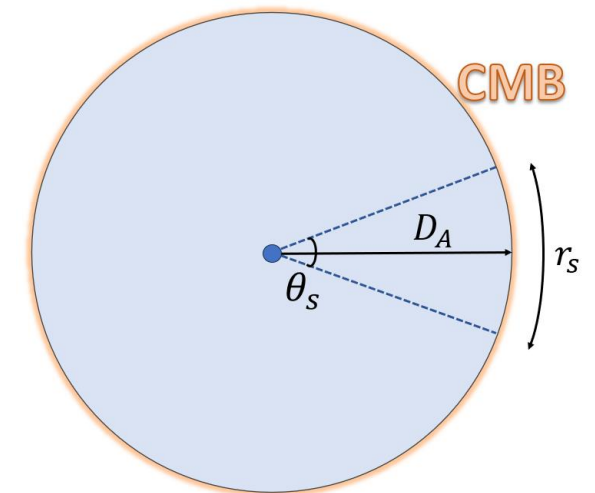
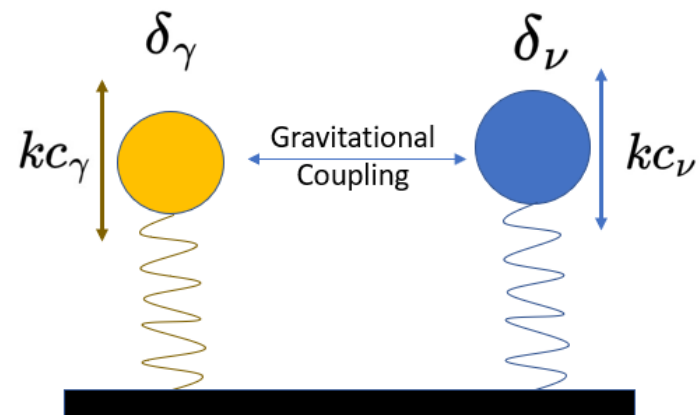
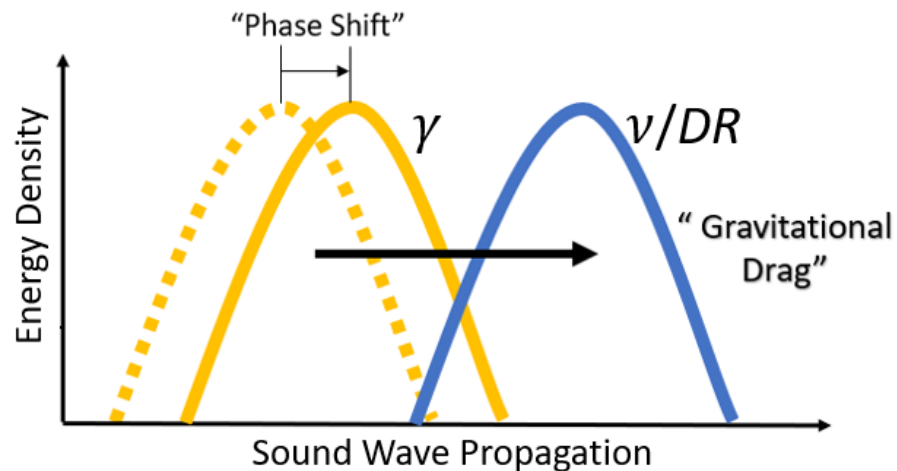
Datasets: Planck2018 + BAO + SH0ES + kv450



# Conclusion

arXiv:2405.08064

1. CMB phase shift provides sensitive **gravitational** probe of propagation behaviour of non-photon radiation before recombination. (Useful for probing physics that only interacts with SM gravitationally).
2. Radiation propagation slowed further (compared to self-interacting case) by scattering with dark matter. Generates **amplified phase shift** in CMB.
  - i. Effect can be understood using *simple coupled oscillator* picture
  - ii. Effect is observable by looking for  $\theta_s$  *enhancement* (dual signal in  $\sigma_8$  *suppression*)



# Backup Slides

# More on the DNI Model

$$\mathcal{L} \supset \frac{y_{ij}}{\Lambda} (H^\dagger l_i) (\psi_j \chi) \quad \Rightarrow \quad \eta_{ij} \nu_i \psi_j \chi, \quad \text{where} \quad \eta_{ij} = \frac{y_{ij} v}{\sqrt{2} \Lambda}.$$

Temperature independent cross-section when DM and mediator mass difference much smaller than neutrino temperature

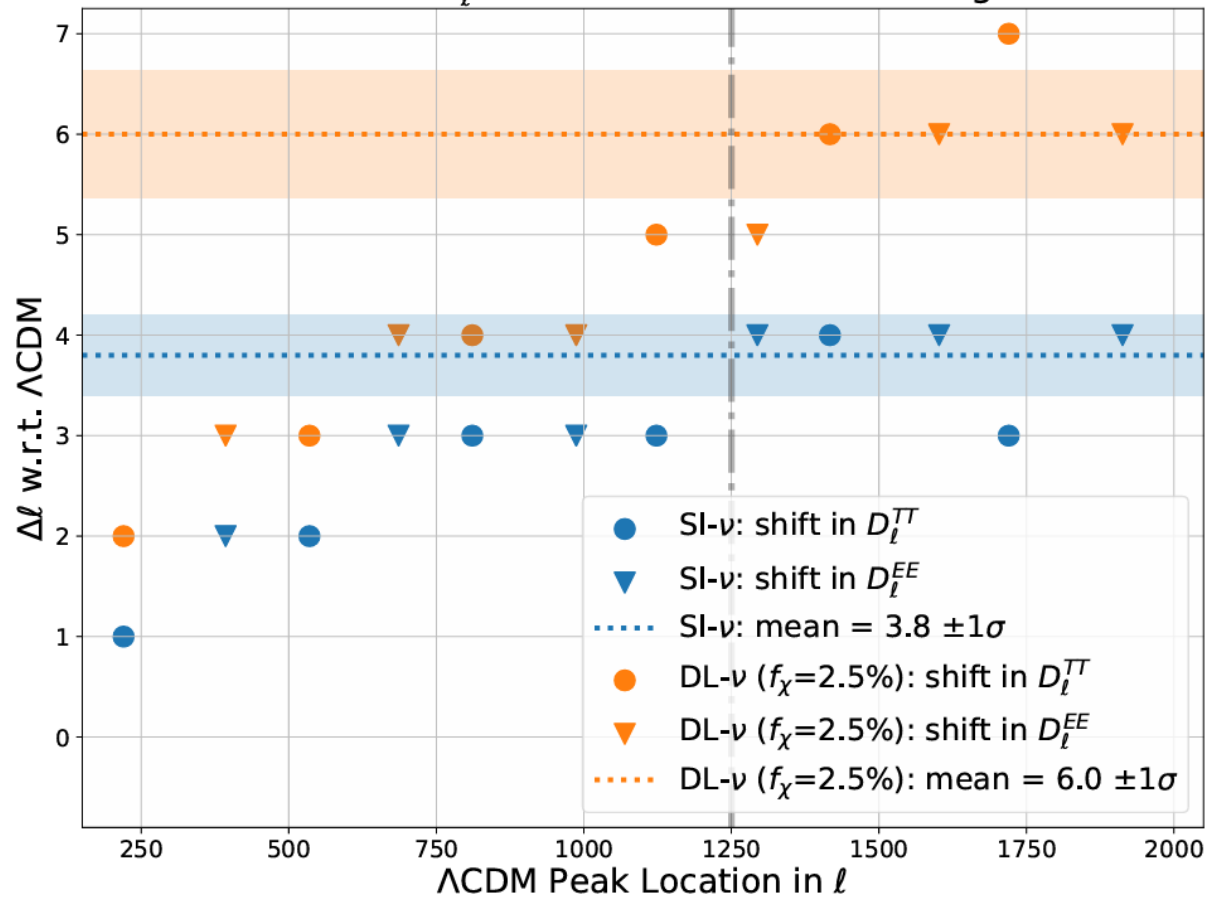
$$\sigma = 1.7 \times 10^{-6} \left( \frac{\eta}{0.1} \right)^4 \left( \frac{\text{GeV}}{m_\chi} \right)^2 \text{GeV}^{-2}$$

Possible UV completion with massive vector-like fermion  $N$

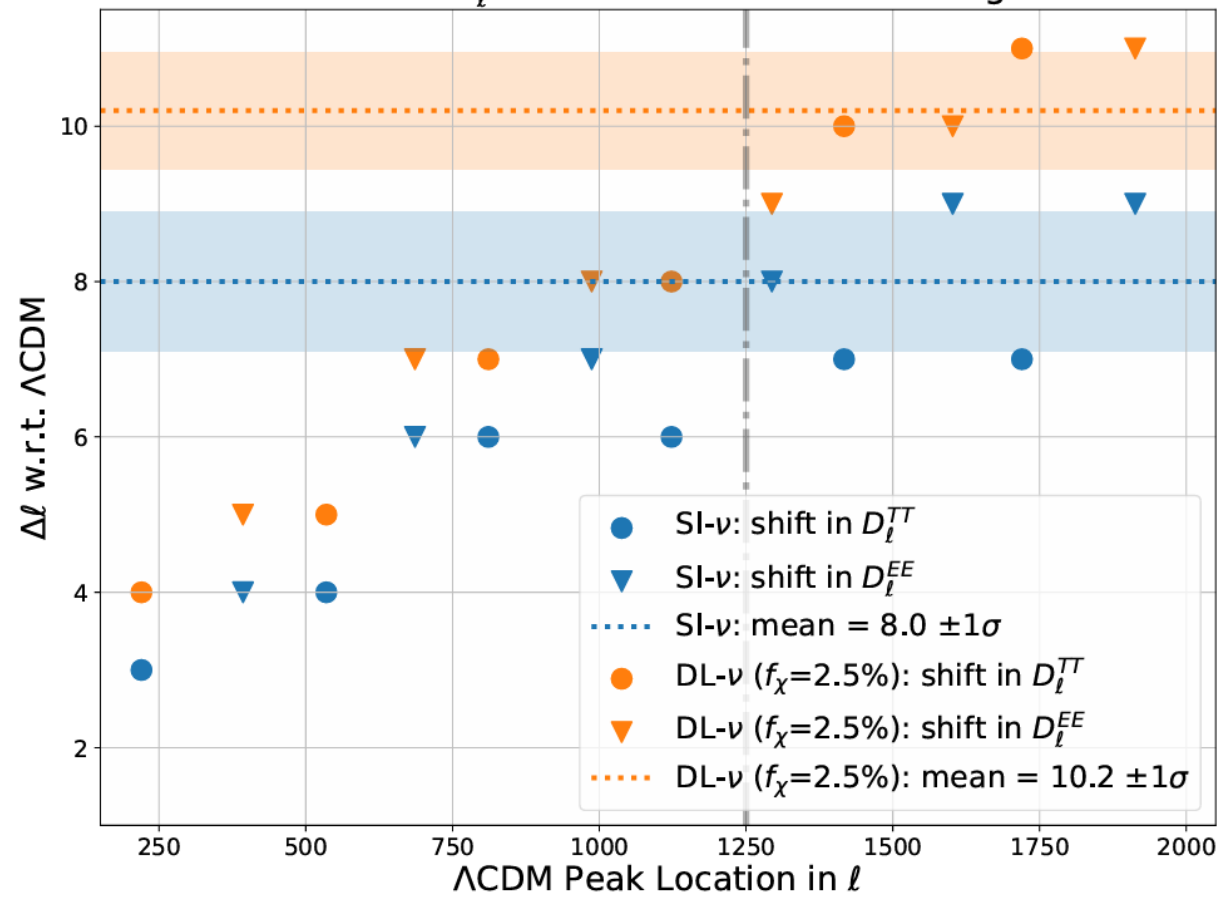
$$\mathcal{L} \supset Y_{N,ij} N_i (H^\dagger l_j) + Y_{\bar{N},ij} N_i^c (\psi_j \chi) + M_{N,ij} N_i N_j^c, \quad \text{where} \quad \frac{y_{ij}}{\Lambda} \sim 2 \frac{Y_{N,ik} Y_{\bar{N},kj}}{M_N}.$$

# CLASS 1 and 2 $\nu$

Shift in  $D_\ell^{TT, EE}$  Peaks for 1  $\nu$  Scattering



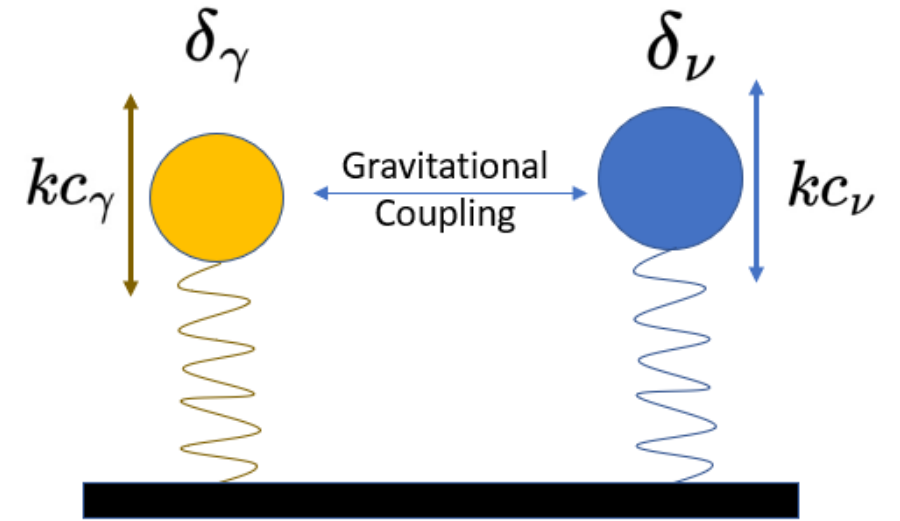
Shift in  $D_\ell^{TT, EE}$  Peaks for 2  $\nu$  Scattering



# Toy Model: Coupled Oscillators

$$\ddot{\delta}_\gamma(\tau) + k^2 c_\gamma^2(\tau) \delta_\gamma(\tau) = \frac{4\mathcal{H}^2(\tau)}{1 + \frac{a(\tau)}{a_{\text{eq}}}} (f_\gamma \delta_\gamma(\tau) + f_\nu \delta_\nu(\tau)),$$
$$\ddot{\delta}_\nu(\tau) + k^2 c_\nu^2(\tau) \delta_\nu(\tau) = \frac{4\mathcal{H}^2(\tau)}{1 + \frac{a(\tau)}{a_{\text{eq}}}} (f_\gamma \delta_\gamma(\tau) + f_\nu \delta_\nu(\tau)).$$

- Two tightly-coupled fluids interacting **only gravitationally**; gravitational interaction weakens over time with Hubble expansion
- Each fluid has natural frequency set by sound speed; matter-loading effect drives sound speeds apart
- Phase shift in photon oscillator due to gravitational influence of hidden oscillator. Direction of shift depends on relative sound speed



$$c_r^2 = \frac{1}{3(1+R_r)}, \quad R_r = \frac{3}{4} \frac{\rho_m}{\rho_r}$$

(for  $r = \gamma$  or  $\nu$ )

## Toy Model Assumptions

1. No free-streaming radiation
2. Small matter-loading
3. Sub-horizon (simplified horizon-entry)
4. Radiation dominated perturbations

# Toy Model Analysis

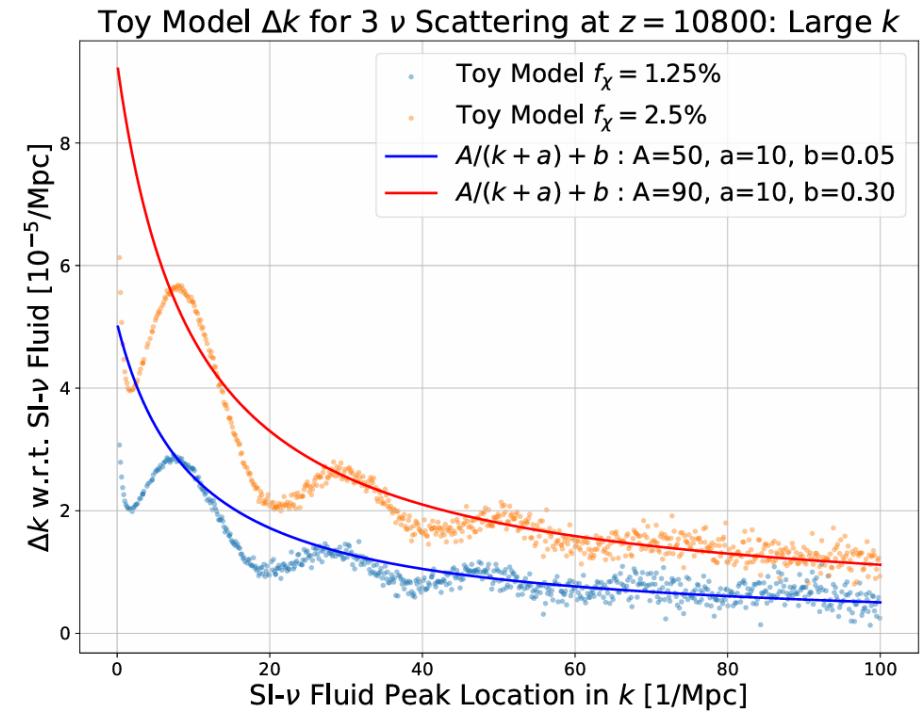
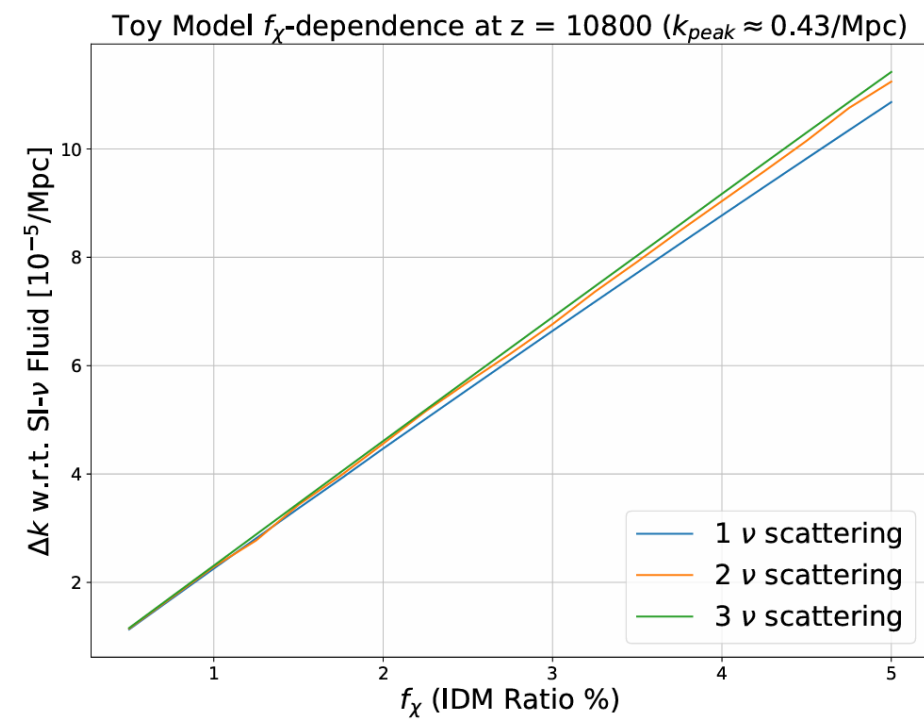
- Analyse toy model for parametric dependences in radiation era, assuming small difference in photon and neutrino sound speeds (and other simplifying assumptions)
- Consider phase shift induced in photon oscillations due to gravitational driving from neutrinos

$$\cos(\omega\tau + \Delta\phi_{\text{load}}) \quad \omega^2 = k^2 c_\gamma^2$$

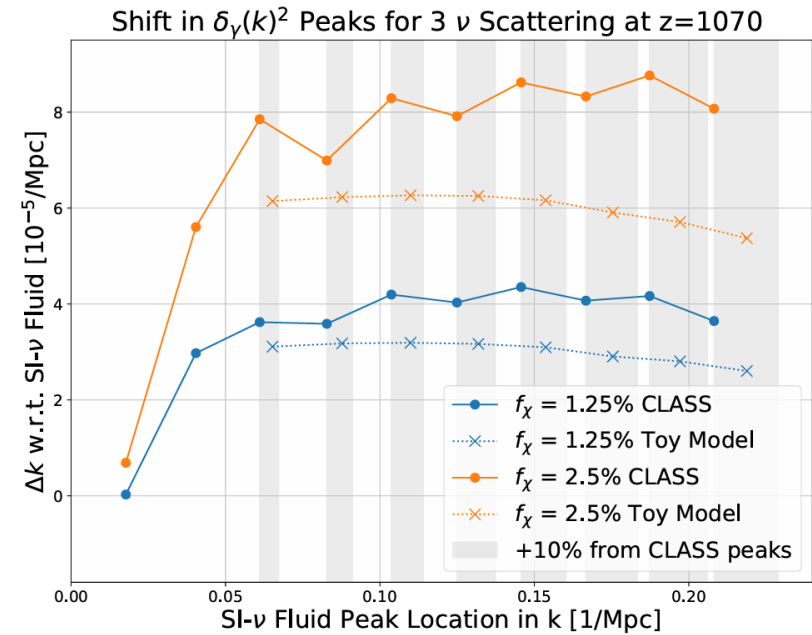
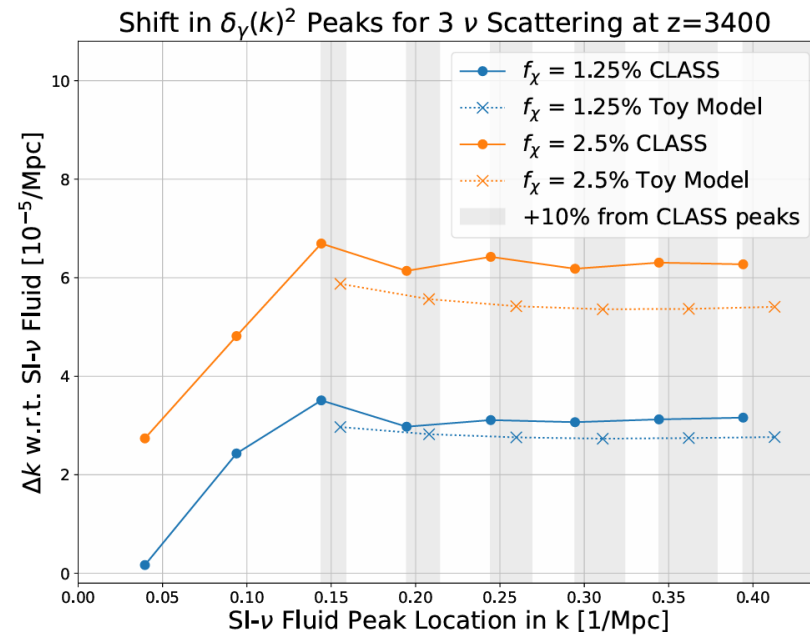
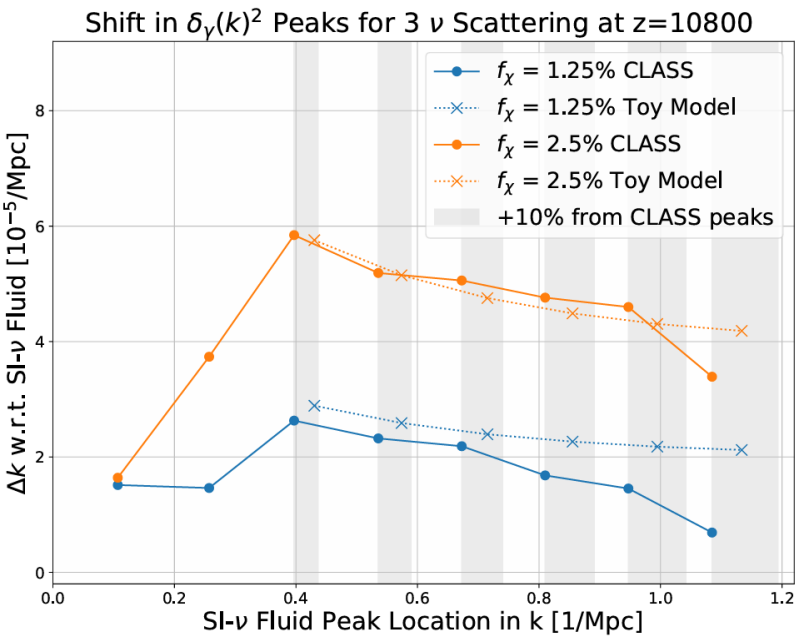
- Analytic approximations

$$\Delta\phi_{\text{load}} \approx -\frac{3\alpha^2 f_{\text{DM}}}{2c_\gamma\tau_{\text{eq}}} \frac{f_\chi}{k+a}, \quad a = \frac{1}{c_\gamma\tau_{\text{eq}}} \left( 2 + \frac{3\alpha f_{\text{DM}}}{4} \frac{f_\chi}{f_\nu} \right)$$

$$\delta k \approx \frac{-\Delta\phi_{\text{load}}}{c_\gamma\tau} \approx \frac{3\alpha^2 f_{\text{DM}}}{2c_\gamma^2\tau_{\text{eq}}} \frac{f_\chi}{(k+a)\tau} \approx 0.07 f_\chi (k\tau)^{-1} \text{Mpc}^{-1}$$



# Numerical Check Toy Model vs CLASS (3 neutrinos)

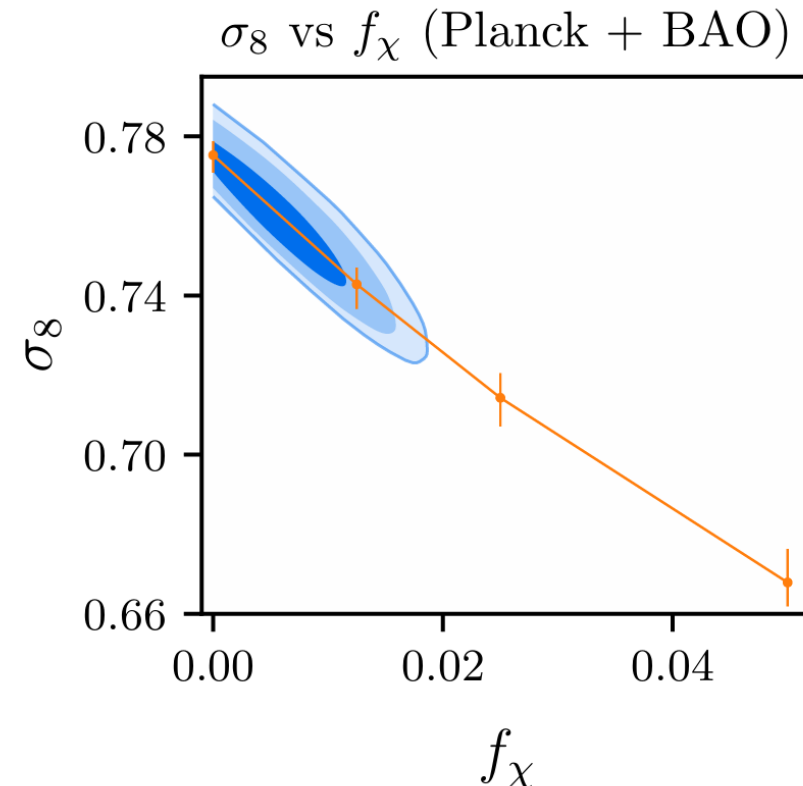
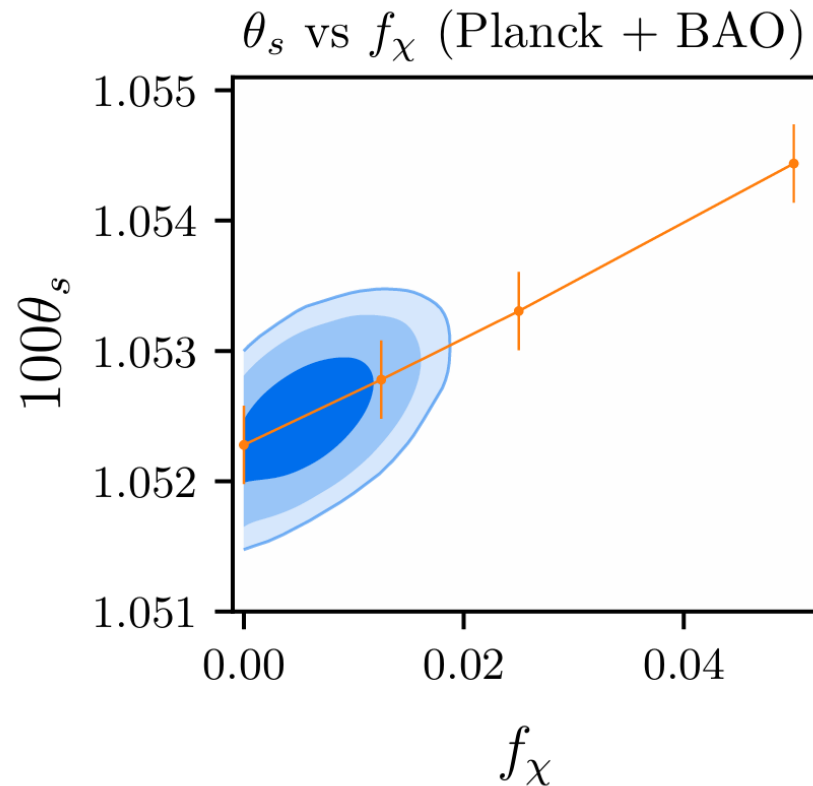


See paper for more details [arXiv:2405.08064](https://arxiv.org/abs/2405.08064)



# MCMC: Proof-of-Principle

- Consider case where all neutrinos scatter efficiently first to isolate  $f_\chi$ -dependence of observables due to DM-loading
- Look at correlations of  $\theta_s$  (CMB phase shift) and  $\sigma_8$  (matter power spectrum) parameters with DL parameter  $f_\chi$



# MCMC Analysis: Signature in $\theta_s$ (Planck+BAO)

- Planck and BAO datasets. DM-loading apparent only when amount of DR is significant
- When DR negligible,  $f_\chi$  becomes unconstrained (does nothing)

