



Ameen Ismail  
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16 May 2024



(towards QCD energy correlators from holography)

arXiv:2403.12123

*with* C. Csáki

DPF-Pheno 2024

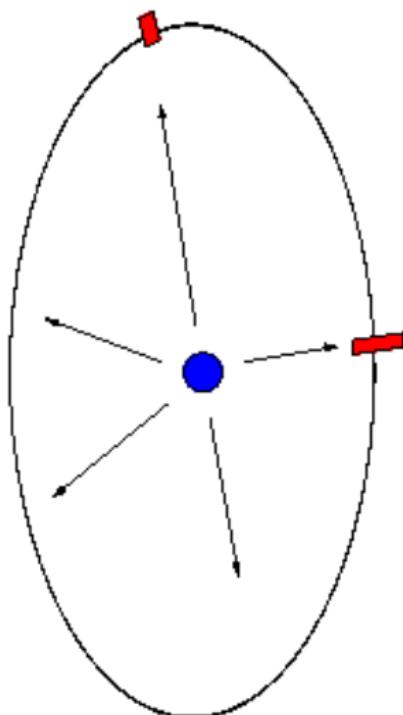
16 May 2024

# Energy correlators

Correlation functions of energy flow operators  $\langle \mathcal{E}(\vec{n}) \rangle, \langle \mathcal{E}(\vec{n}) \mathcal{E}(\vec{n}') \rangle$

$\mathcal{E}(\vec{n})$ : energy deposited in a calorimeter

CFT context by Hofman + Maldacena '08

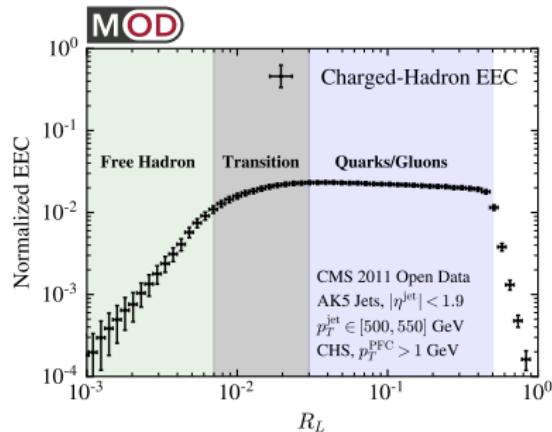


# Energy correlators and confinement

Two-point function visualizes  
confinement transition

The big idea:

Can we reproduce this in a  
holographic model?



Komiske, Moult, Thaler, Zhu 2201.07800

# Calculating energy correlators

Correlation functions of **energy flow operators**

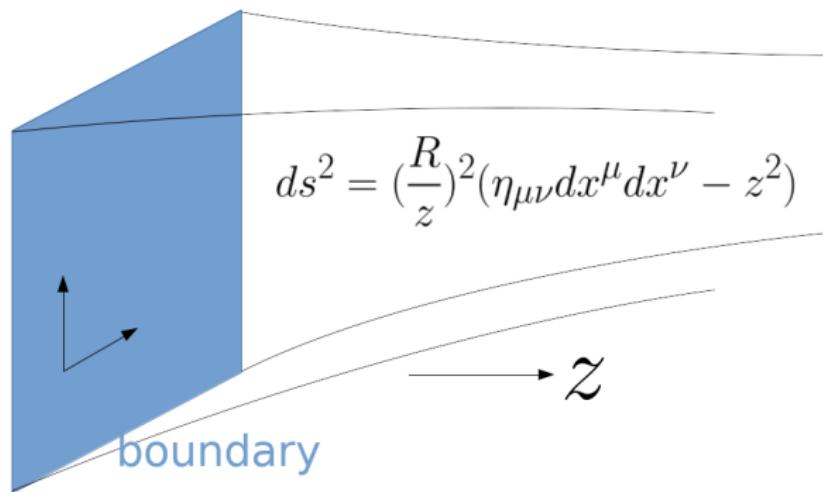
$$\mathcal{E}(\vec{n}) = \lim_{r \rightarrow \infty} r^2 \int_0^\infty dt T_{0i}(t, x^i = rn^i) n^i$$

Lightcone coordinates  $x^\pm = x^0 \pm x^3$  and inversion:

$$\mathcal{E}(\vec{n}) = \left(1 + |x^\perp|^2\right)^3 \int_{-\infty}^\infty dx^- T_{--}(x^+ = 0, x^-, x^\perp)$$

(mapping to celestial sphere:  $x^1 + ix^2 = e^{i\phi} \tan \theta/2$ )

# Holographic calculation setup



We want a **generating functional** for correlators

# Inserting an energy flow operator

Perturb

[Belin et al. 2011.13862](#)

$$\begin{aligned}\delta S_{\text{CFT}} &= \epsilon \int dx^- T_{--}(x^+ = 0, x^-, x^\perp = y^\perp) \\ &= \epsilon \int d^4x T_{--} \delta(x^+) \delta^2(x^\perp - y^\perp)\end{aligned}$$

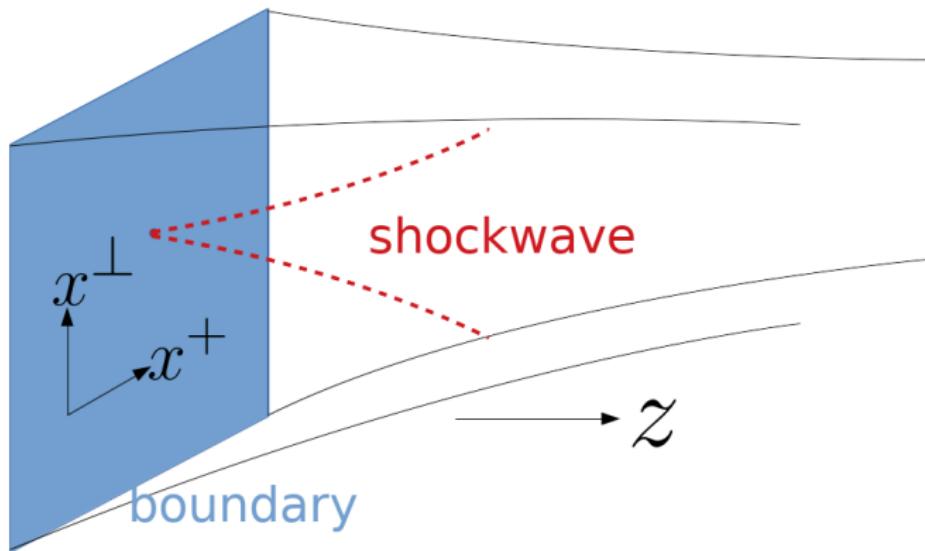
Inserts  $e^{\epsilon \mathcal{E}}$  in path integral

Dual picture: **shockwave**,

$$\delta ds^2 = \frac{\epsilon}{z^2} \delta(x^+) f(x^\perp - y^\perp, z) (dx^+)^2$$

(boundary condition:  $f(x^\perp, z = 0) = \delta^2(x^\perp)$ )

# Shockwave picture



# Shockwaves in AdS

Field equations are **linear**:  $(3/z\partial_z - \partial_z^2 - \partial_{\perp}^2) f = 0$

## Shockwave

$$f(x^\perp, z) = \frac{z^4}{(z^2 + |x^\perp|^2)^3}$$

Shockwaves are linear  $\Rightarrow$  **superpose** them:  $ds^2 = ds_{\text{AdS}}^2 + \frac{\delta(x^+)}{z^2} [\epsilon_1 f(x^\perp - y_1^\perp, z) + \epsilon_2 f(x^\perp - y_2^\perp, z)] (dx^+)^2$

**Exact solution** of field equations

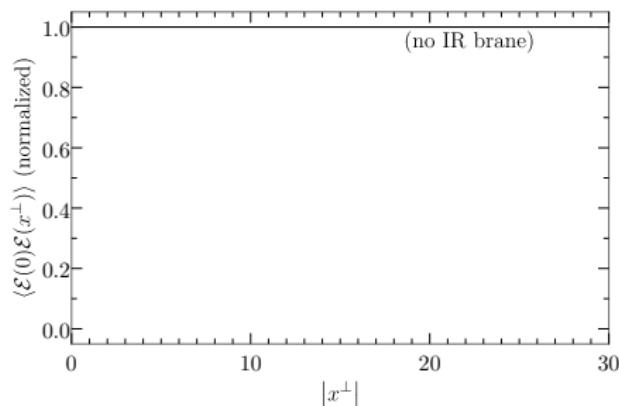
# Holographic correlators

Scalar source  $\Leftrightarrow$  bulk scalar  $\phi$

momentum  $q^\mu = (q, \vec{0}) \Leftrightarrow$  BC  $\phi(z=0) = e^{iqt}$

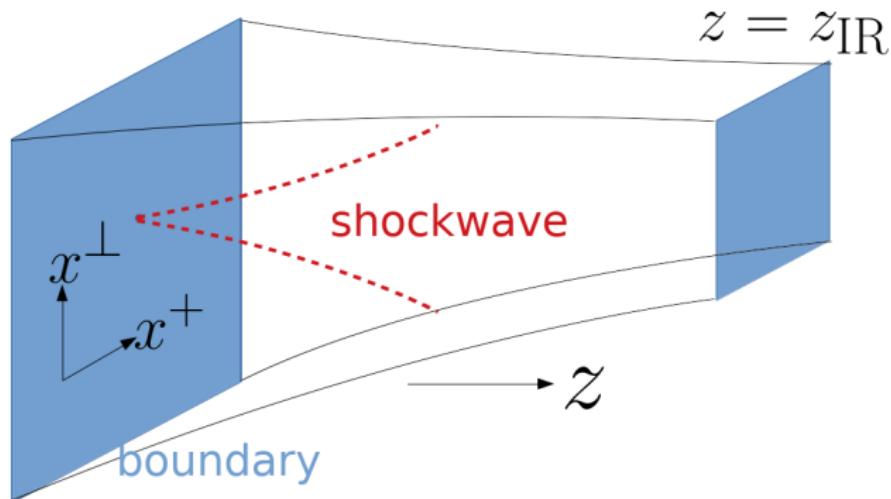
$\phi$  gets a “kick” at the  
shockwave(s):

$\Rightarrow \langle \mathcal{E} \rangle \sim 1$  (constant)  
 $\Rightarrow \langle \mathcal{E} \mathcal{E} \rangle \sim 1$



# Cutting off AdS

Simplest “hard-wall” confinement: IR brane at  $z = z_{\text{IR}}$   
(essentially RSI model)



# Shockwaves with a brane

Shockwave EOM unchanged,  $(3/z\partial_z - \partial_z^2 - \partial_1^2 - \partial_2^2)f = 0$

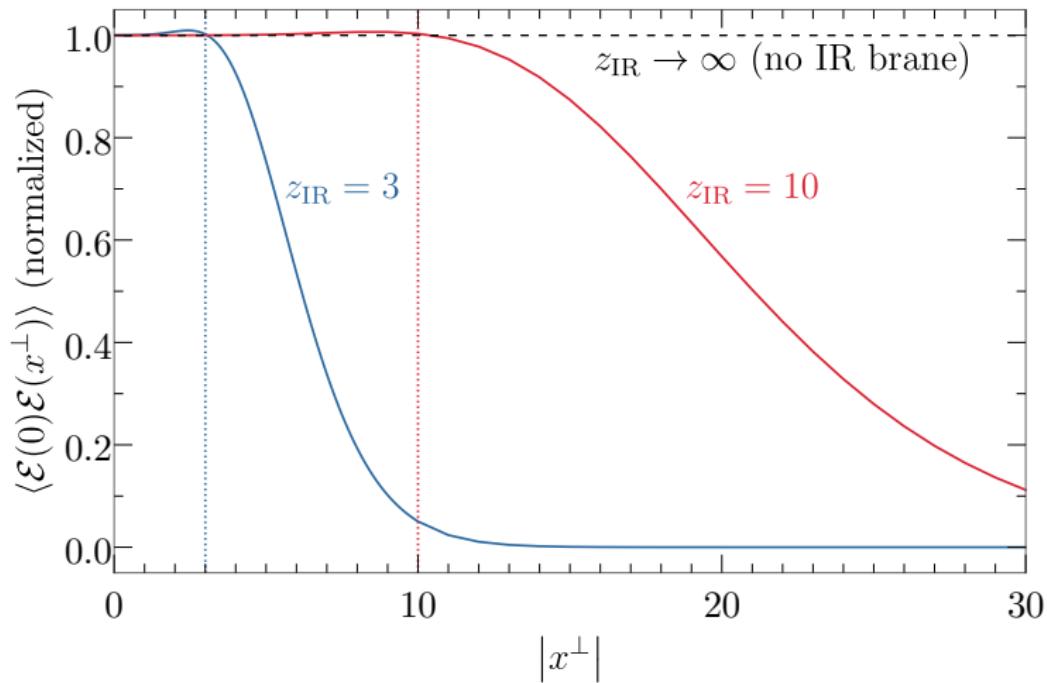
Brane modifies BC:  $\partial_z f(x^\perp, z) \Big|_{z=z_{\text{IR}}} = 0$

## Solution

$$f(x^\perp, z) = \frac{1}{8} \int_0^\infty dk J_0(kr) k^3 z^2 \left[ K_2(kz) + \frac{K_1(kz_{\text{IR}})}{I_1(kz_{\text{IR}})} I_2(kz) \right]$$

(where  $r = |x^\perp|$ )

# Results



# Outlook

Manifest transition between confined and deconfined regimes

Comparison with QCD:

- ▶ constant correlator vs. asymptotic freedom
- ▶ exponential vs. power-law decay
- ▶ jetty behaviour

Future: generalize to more realistic AdS/QCD backgrounds

# Thank you!



Photos: Jefferson Lab / Wikimedia Commons / CERN

more info:

[arxiv.org/abs/2403.12123](https://arxiv.org/abs/2403.12123)

[ai279@cornell.edu](mailto:ai279@cornell.edu)

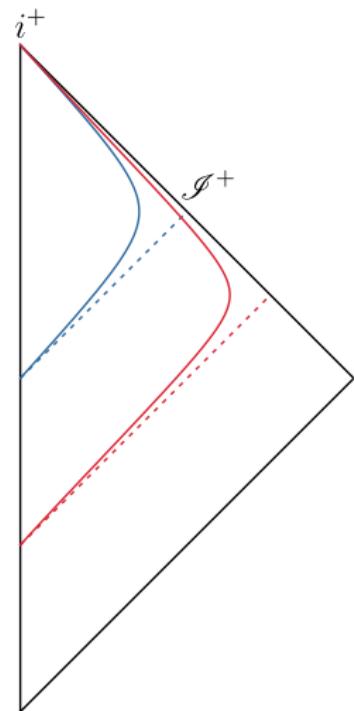
[ameenismail.github.io](https://ameenismail.github.io)

# A caveat about limits

Be **careful** about the order of limits!

- ▶  $x^- \rightarrow \pm\infty$  (integrate over time)
- ▶  $x^+ \rightarrow \infty$  (calorimeter to boundary)

Only important for gapped theories



# The inversion

Avoid large- $r$  limit with a conformal transformation:

$$x^+ \rightarrow -1/x^+, x^- \rightarrow x^- - |x^\perp|^2/x^+, x^\perp \rightarrow x^\perp/x^+$$

mapping  $x^+ = \infty$  to  $x^+ = 0$ ,  $ds^2 \rightarrow ds^2/(x^+)^2$

Energy flow calculated as

$$\mathcal{E}(\vec{n}) = \left(1 + |x^\perp|^2\right)^3 \int_{-\infty}^{\infty} dx^- T_{--}(x^+ = 0, x^-, x^\perp)$$

with mapping to celestial sphere:  $x^1 + ix^2 = e^{i\phi} \tan \theta/2$

# Scalar source

Scalar source, momentum  $q^\mu = (q, \vec{0})$

$\Rightarrow$  bulk scalar  $\phi$  with BC  $\phi(z=0) = e^{iqt}$

$\phi$  gets a “kick” at the shockwave:

$$\lim_{\delta \rightarrow 0} \partial_- \phi(x^+ = \delta, x^-, x^\perp, z) = e^{-\epsilon f(x^\perp, z) \partial_-} \partial_- \phi(x^+ = -\delta, x^-, x^\perp, z)$$

(away from shockwave, usual AdS evolution)

# Holographic correlators

One point:  $\langle e^{\epsilon \mathcal{E}(y^\perp)} \rangle \sim$

$$\int \frac{dz}{z^3} d^2x^\perp dx^- \phi^* \exp \left[ -\epsilon (1 + (y^\perp)^2)^3 f(x^\perp - y^\perp, z) \partial_- \right] \partial_- \phi \Big|_{x^+ = 0}$$

Expanding in  $\epsilon$ , find  $\langle \mathcal{E} \rangle \sim 1$  (constant)

Higher-point functions from inserting more shockwaves, e.g.  
 $\langle \mathcal{E} \mathcal{E} \rangle \sim 1$

# Wavefunction at the shock

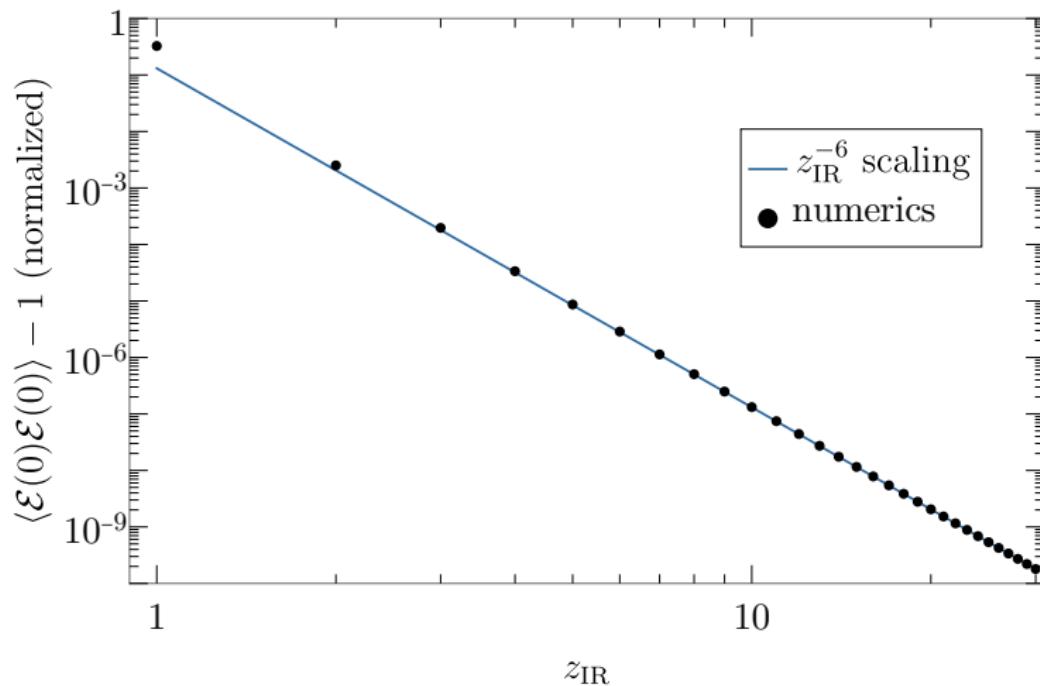
EOM:

$$\partial_- \partial_+ \phi + \epsilon \delta(x^+) f(x^\perp, z) \partial_-^2 \phi + \text{terms regular at shockwave} = 0,$$

Discontinuity:

$$\lim_{\delta \rightarrow 0} \partial_- \phi(x^+ = \delta, x^-, x^\perp, z) = e^{-\epsilon f(x^\perp, z) \partial_-} \partial_- \phi(x^+ = -\delta, x^-, x^\perp, z)$$

# More figures 1



## More figures 2

