

Exploring Dark Forces with Multimessenger Studies of Extreme Mass Ratio Inspirals

[arXiv:2405.06011](https://arxiv.org/abs/2405.06011)

Badal Bhalla

With Kuver Sinha and Tao Xu



The UNIVERSITY of OKLAHOMA

HOMER L. DODGE DEPARTMENT OF PHYSICS AND ASTRONOMY



Binary inspirals can serve as probes of dark mediators,
under the assumption that the spiraling objects
accumulate dark charge

NS, BHs, or other compact objects.



Binary inspirals can serve as probes of dark mediators,
under the assumption that the spiraling objects
accumulate dark charge

NS, BHs, or other compact objects.

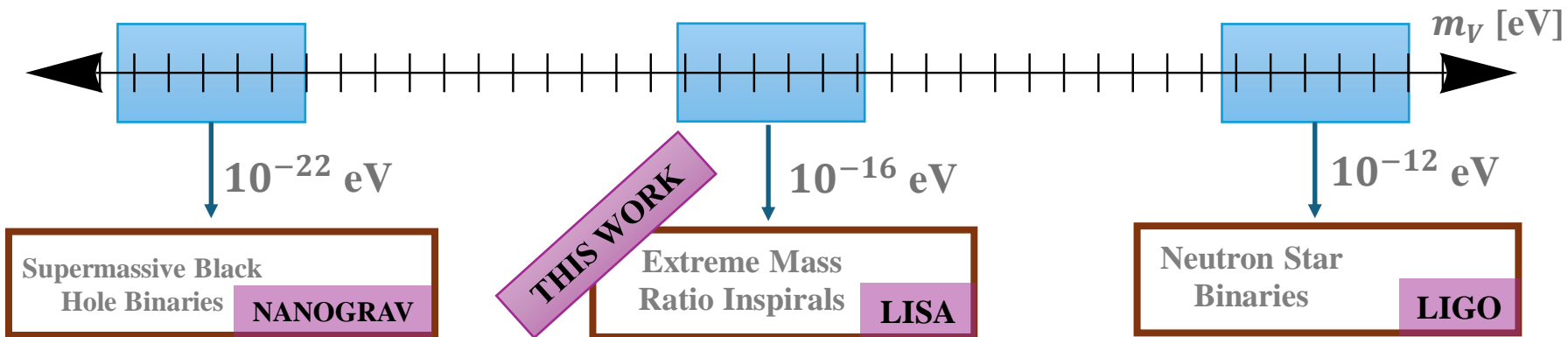
The Compton wavelength of the mediator (and hence the mass) that
can be probed is primarily determined by the characteristic length
scale of the binary system.



Binary inspirals can serve as probes of dark mediators, under the assumption that the spiraling objects accumulate dark charge

NS, BHs, or other compact objects.

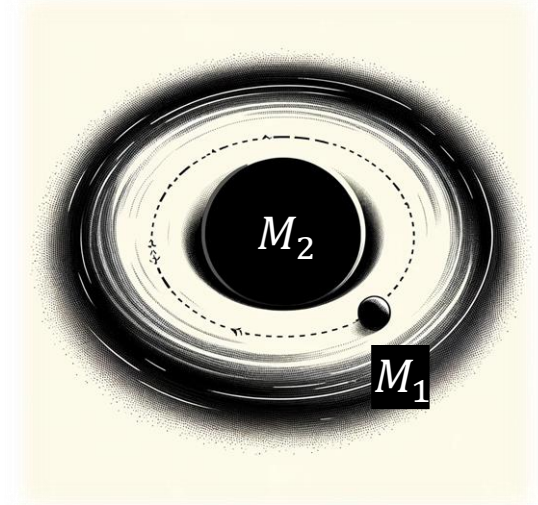
The Compton wavelength of the mediator (and hence the mass) that can be probed is primarily determined by the characteristic length scale of the binary system.



mHz Gravitational Waves.



In-spiral of a compact, stellar-mass object into a massive black hole.



$$\frac{M_1}{M_2} = 10^{-5}$$

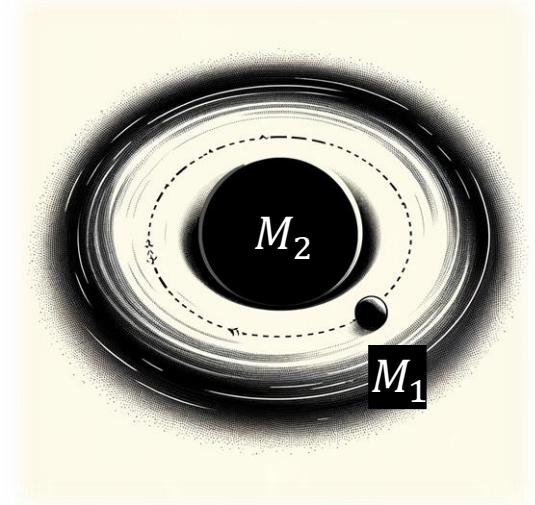


mHz Gravitational Waves.



In-spiral of a compact, stellar-mass object into a massive black hole.

Few EMRIs should be detectable by LISA per year.



$$\frac{M_1}{M_2} = 10^{-5}$$



mHz Gravitational Waves.



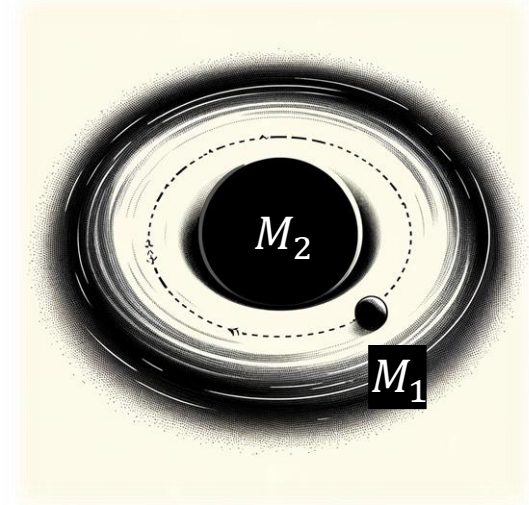
In-spiral of a compact, stellar-mass object into a massive black hole.

Few EMRIs should be detectable by LISA per year.

What is interesting about EMRI?

Precise Measurement

Long Observation Time



$$\frac{M_1}{M_2} = 10^{-5}$$



We consider a Yukawa force mediated by a particle of mass m_ν . Furthermore, we assume that the dark force only acts between two black holes. All other compact objects do not experience it.

$$F_{DARK} = \frac{\alpha' Q_1 Q_2}{r^2} e^{-m_\nu r} (1 + m_\nu r)$$



We consider a Yukawa force mediated by a particle of mass m_ν . Furthermore, we assume that the dark force only acts between two black holes. All other compact objects do not experience it.

$$F_{DARK} = \frac{\alpha' Q_1 Q_2}{r^2} e^{-m_\nu r} (1 + m_\nu r)$$

α' = Coupling Constant of the Dark Force
 m_ν = mass of the dark force mediator
 r = separation between the two objects
 Q_i = dark charge on i^{th} object.



We consider a Yukawa force mediated by a particle of mass m_ν . Furthermore, we assume that the dark force only acts between two black holes. All other compact objects do not experience it.

$$F_{DARK} = \frac{\alpha' Q_1 Q_2}{r^2} e^{-m_\nu r} (1 + m_\nu r)$$

α' = Coupling Constant of the Dark Force
 m_ν = mass of the dark force mediator
 r = separation between the two objects
 Q_i = dark charge on i^{th} object.

In the presence of the dark force, the evolution of orbital frequency is given by

$$\omega^2 = \frac{G(M_1 + M_2)}{r^3} [1 + \tilde{\alpha}' e^{-m_\nu r} (1 + m_\nu r)]$$

$$\tilde{\alpha}' = \frac{\alpha' Q_1 Q_2}{GM_1 M_2}$$



Why is it necessary to conduct a Multimessenger study?



The UNIVERSITY *of* OKLAHOMA

HOMER L. DODGE DEPARTMENT OF PHYSICS AND ASTRONOMY



Why is it necessary to conduct a Multimessenger study?

Because

$$F_{TOTAL} = \frac{GM_1M_2}{r^2} [1 + \tilde{\alpha}'e^{-m_v r}(1 + m_v r)]$$

is same as

$$F_{TOTAL} = \frac{GM_1'M_2'}{r^2}$$



Why is it necessary to conduct a Multimessenger study?

Because

$$F_{TOTAL} = \frac{GM_1M_2}{r^2} [1 + \tilde{\alpha}'e^{-m_\nu r}(1 + m_\nu r)]$$

is same as

$$F_{TOTAL} = \frac{GM_1'M_2'}{r^2}$$

Rescaling the mass of binary components can imitate the effects of Dark Force. Therefore, it is crucial to measure the mass of binary components without any dark force acting between them.



Other methods, free from dark forces, can determine the mass of SMBH.

1. Multimessenger Studies with XMRI-EMRI Combinations :

$$\frac{\Delta M_{SMBH}}{M_{SMBH}} = 10^{-3}\%$$



Other methods, free from dark forces, can determine the mass of SMBH.

1. Multimessenger Studies with XMRI-EMRI Combinations :

$$\frac{\Delta M_{SMBH}}{M_{SMBH}} = 10^{-3}\%$$

2. Multimessenger Studies with Stellar Kinematics and EMRI :

$$\frac{\Delta M_{SMBH}}{M_{SMBH}} = 10 - 30\%$$



Other methods, free from dark forces, can determine the mass of SMBH.

1. Multimessenger Studies with XMRI-EMRI Combinations :

$$\frac{\Delta M_{SMBH}}{M_{SMBH}} = 10^{-3}\%$$

2. Multimessenger Studies with Stellar Kinematics and EMRI :

$$\frac{\Delta M_{SMBH}}{M_{SMBH}} = 10 - 30\%$$

3. Studies of Spectroscopic Reverberation Mapping and EMRIs:

$$\frac{\Delta M_{SMBH}}{M_{SMBH}} = 10\%$$



The time evolution of orbital frequency is

$$\frac{d\omega}{dt} = -\frac{32}{5} G\mu r^2 \omega^5 g(\alpha', m_V, r) N^{-1}$$

Relativistic Correction

$$g = -\frac{3 + \alpha' e^{-m_V r} (3 + m_V r (3 + m_V r))}{1 + \alpha' e^{-m_V r} (1 + m_V r (1 - m_V r))}$$

We assume a 10% uncertainty in the mass measurement of the compact object in this study.



The time evolution of orbital frequency is

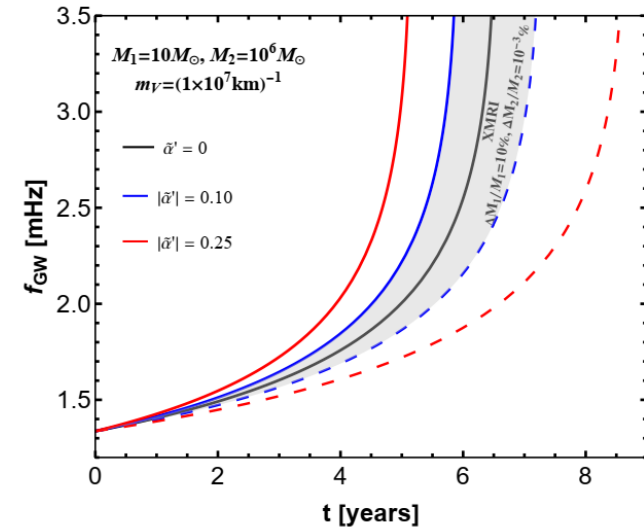
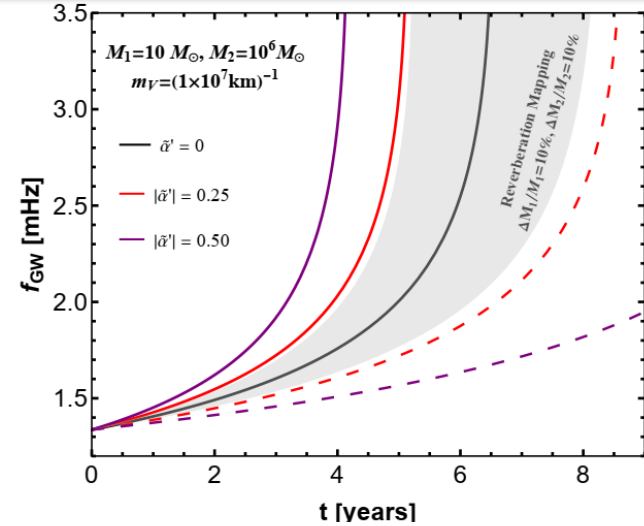
$$\frac{d\omega}{dt} = -\frac{32}{5} G\mu r^2 \omega^5 g(\alpha', m_V, r) N^{-1}$$

Relativistic Correction

$$g = -\frac{3 + \alpha' e^{-m_V r} (3 + m_V r (3 + m_V r))}{1 + \alpha' e^{-m_V r} (1 + m_V r (1 - m_V r))}$$

We assume a 10% uncertainty in the mass measurement of the compact object in this study.

- Attractive Dark Force
- - - Repulsive Dark Force



The rms amplitude of the gravitational wave emitted for the dominant, $m = 2$, radiation is

$$h_{o,2} = \frac{3.6 \times 10^{-23}}{r_0/1\text{Gpc}} \left(\frac{\mu}{M_\odot} \right) \left(\frac{M}{100M_\odot} \right)^{2/3} \left(\frac{f_2}{100\text{Hz}} \right)^{2/3} \mathcal{H}_{o,2}$$

Here

r_0 is the distance from EMRI.

$\mathcal{H}_{o,2}$ is the relativistic correction to the amplitude

$$h_{c,2} \stackrel{\text{def}}{=} h_{o,2} \sqrt{\frac{2f_2^2}{\dot{f}_2}}$$

Is the characteristic amplitude for the waves



The rms amplitude of the gravitational wave emitted for the dominant, $m = 2$, radiation is

$$h_{o,2} = \frac{3.6 \times 10^{-23}}{r_0/1\text{Gpc}} \left(\frac{\mu}{M_\odot}\right) \left(\frac{M}{100M_\odot}\right)^{2/3} \left(\frac{f_2}{100\text{Hz}}\right)^{2/3} \mathcal{H}_{o,2}$$

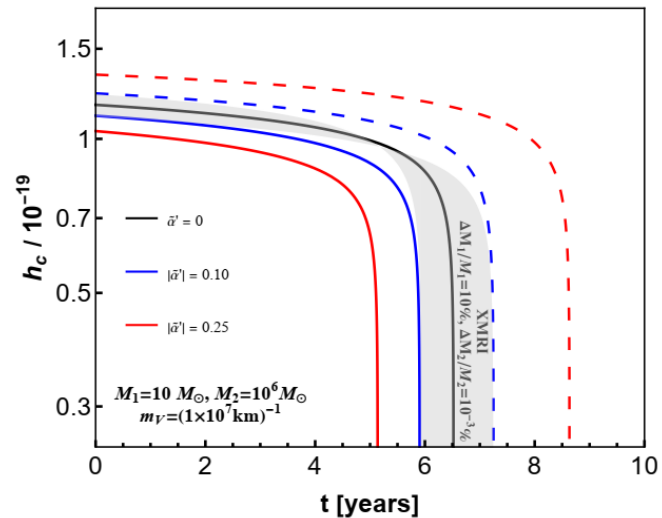
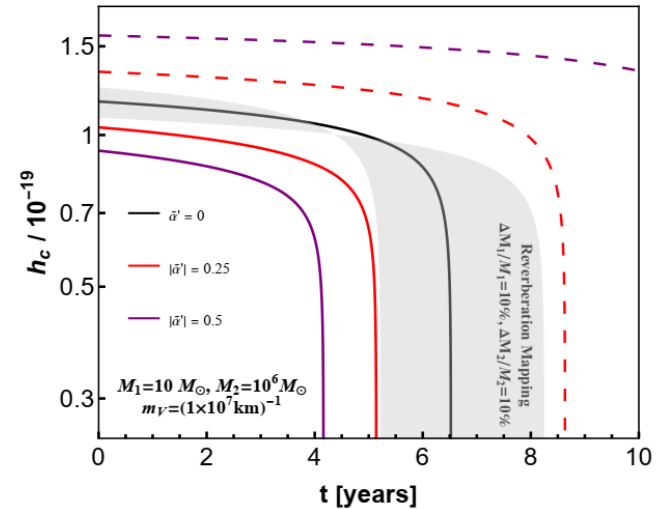
Here

r_0 is the distance from EMRI.

$\mathcal{H}_{o,2}$ is the relativistic correction to the amplitude

$$h_{c,2} \stackrel{\text{def}}{=} h_{o,2} \sqrt{\frac{2f_2^2}{\dot{f}_2}}$$

Is the characteristic amplitude for the waves

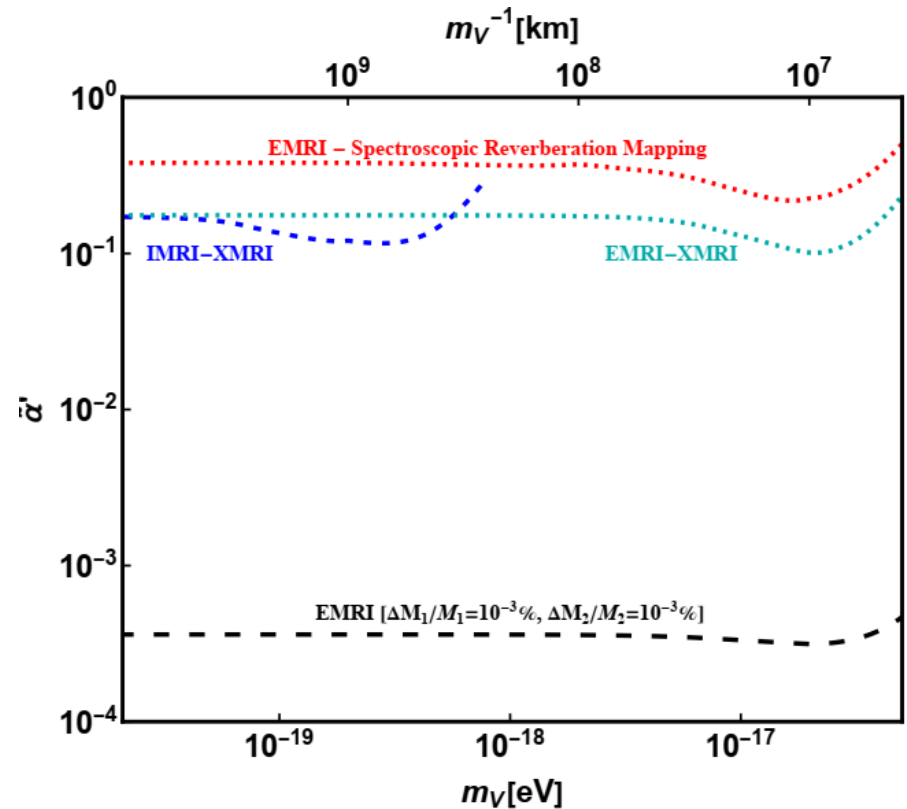


LIKELIHOOD ANALYSIS

$$\Lambda(s|\tilde{\alpha}') = \mathcal{K} \exp[(h_{\tilde{\alpha}'}|s) - \frac{1}{2}(h_{\tilde{\alpha}'}|h_{\tilde{\alpha}'}) - \frac{1}{2}(s|s)]$$

$$(h_{\tilde{\alpha}'}|s) = \text{Re} \int_{f_{\min}}^{f_{\max}} df_{\text{GW}} \frac{h_{c,\tilde{\alpha}'}(f_{\text{GW}}) s_c(f_{\text{GW}})}{f_{\text{GW}}^2 S_n(f_{\text{GW}})}$$

	EMRI-SRM	EMRI-XMRI	IMRI-XMRI
M_1	$10 M_{\odot}$	$10 M_{\odot}$	$10^4 M_{\odot}$
M_2	$10^6 M_{\odot}$	$10^6 M_{\odot}$	$10^8 M_{\odot}$
$\frac{\Delta M_1}{M_1}$	10 %	10 %	10 %
$\frac{\Delta M_2}{M_2}$	10 %	10^{-3} %	10^{-3} %



EMRIs can be used to probe Dark Forces.



EMRIs can be used to probe Dark Forces.

EMRI is sensitive to a range of mediator masses. The best sensitivity arises for mediator mass which is comparable to the separation of the EMRI.



EMRIs can be used to probe Dark Forces.

EMRI is sensitive to a range of mediator masses. The best sensitivity arises for mediator mass which is comparable to the separation of the EMRI.

An independent mass measurement of the central SMBH breaks the degeneracy of the effect of the dark force with a simple rescaling of the binary component masses.



THANK YOU



The UNIVERSITY *of* OKLAHOMA

HOMER L. DODGE DEPARTMENT OF PHYSICS AND ASTRONOMY



BACKUP SLIDE: DARK-CHARGED BLACK HOLES

Consider a millicharged DM fermion interacting by the exchange of a vector mediator

$$\mathcal{L} = -\frac{1}{4}V_{\mu\nu}V^{\mu\nu} + \frac{1}{2}m_V V_\mu V^\mu + \bar{\chi}(i\gamma_\mu D^\mu - m_\chi)\chi .$$

$$Q \sim M \longrightarrow$$

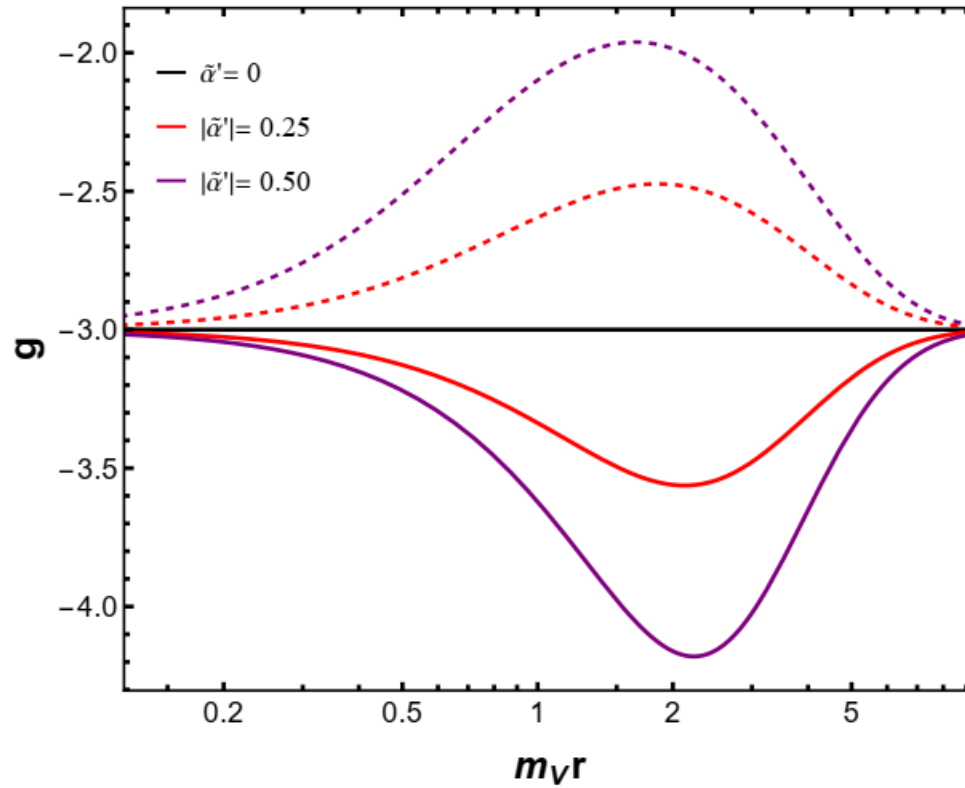
$$Q_2 = N_2 q_\chi$$
$$m_\chi \sim m_p, \quad \text{and} \quad q_\chi \sim 10^{-18} e$$

Assuming that the SMBH is in an environment with density ρ of oppositely charged dark fermions moving with velocity v

$$\tau_{\text{discharge}} \sim \left(\frac{v}{220 \text{ km/s}}\right)^3 \left(\frac{0.4 \text{ GeV/cm}^3}{\rho}\right) \left(\frac{10M_\odot}{M_2}\right) \left(\frac{m_\chi}{m_p}\right) \left(\frac{e}{q_\chi}\right) \text{ yr}$$



BACKUP SLIDE: MORE RESULTS



The UNIVERSITY of OKLAHOMA

HOMER L. DODGE DEPARTMENT OF PHYSICS AND ASTRONOMY



BACKUP SLIDE: MORE RESULTS

