

Division of Particles & Fields

University of

## Constraining Bosonic Asymmetric Dark Matter With Neutron Star Mass-Radius Measurements

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Authors: Nathan Rutherford, Geert Raaijmakers, Chanda Prescod-Weinstein, and Anna Watts Based on Rutherford et al. (2023) arXiv: 2208.03282

- <sup>1</sup> Uncertainties In The Neutron Star Equation Of State (EoS)
- Neutron stars may contain exotic states of matter, e.g., deconfined quarks or hyperons.
- The effects of the hypothetical components are captured by the equation of state.
- The EoS can be deduced from measurable properties, e.g., the mass and radius.



Figure made by Anna Watts

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Neutron stars are an excellent hunting ground for dark matter! For the new mass-radius measurement of PSR J0437-4517 check out Choudhury et al. 2024 and Rutherford et al. 2024 (both in review at APJ Letters!)

Figure made by Anna Watts

![](_page_3_Figure_7.jpeg)

- <sup>2</sup> Asymmetric Dark Matter (ADM) In Neutron Stars
- ADM can accumulate in two spatial regimes: the neutron star core and in the exterior spacetime.
- ADM cores reduce the gravitational mass, radius, and tidal deformability.
- ADM halos increase the gravitational mass and tidal deformability.

#### ADM core

**ADM halo** 

#### The Bosonic ADM Model

Modeled after the Nelson et al. model<sup>1</sup>.

 Describes MeV-GeV mass-scale bosonic ADM particles with repulsive self-interactions mediated by an eV-MeV mass-scale vector gauge boson.

The defining parameters of this model are: 1) The bosonic ADM particle mass  $(m_{\chi})$ 2) The effective self-repulsion strength  $(\frac{g_{\chi}}{m_{\phi}})$ 

3) The fraction of ADM mass inside the neutron star  $(F_{\chi})$ 

![](_page_5_Picture_6.jpeg)

<sup>4</sup> Using Bayesian Inference To Study Bosonic ADM In Neutron Stars

We perform a Bayesian analysis where we:

- Vary the baryonic matter and ADM EoS
- Vary the ADM EoS, but fix the baryonic matter EoS

For both cases, we consider synthetic mass and radius measurements and not allow for ADM halos since:

ADM halos modify the exterior spacetime

ADM could modify the universal relations that are used to model the oblateness

![](_page_6_Figure_7.jpeg)

#### Source Selection

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- Radius of the sources calculated using two ground-truth models:
- "ADM Core Model": Baryonic neutron star with ADM core defined by [ $m_{\chi} = 15 \text{ GeV}, \frac{g_{\chi}}{m_{\phi}/\text{MeV}} = 0.1, F_{\chi} = 7 \%$ ]
- "No ADM Model": Identical to "ADM Core Model", except we set  $F_{\gamma} = 0$  %.

![](_page_7_Figure_4.jpeg)

![](_page_7_Figure_5.jpeg)

#### The ADM Priors

![](_page_8_Figure_1.jpeg)

4:

We sample uniformly in all three intervals and then eliminate any halo configurations.

<u>Future-X</u>

![](_page_8_Figure_4.jpeg)

![](_page_8_Figure_5.jpeg)

#### The ADM Priors

3

2

 $\mathbf{\hat{\mathbf{y}}}$ 

0

 $\gamma$ 

r

~%

v

ծ

 $\log_{10}(\frac{g_{\chi}}{m_{\varphi}/MeV})$ 

 $\vec{\lambda}$ 

![](_page_9_Figure_1.jpeg)

: <u>Kourvaris & Tinyakov (2011),</u> 2: <u>Bramante et al. (2013),</u> 3: <u>Nelson et al</u> 2018), 4: <u>Karkevandi et al. (2022)</u> <sup>7</sup> Future / Future - X: Varying Baryonic EoS

ADM

Core

Model

- The 'Including ADM' band is noticeably wider than the 'Neglecting ADM' band.
- A stiffer baryonic EoS ⇒ posterior constraints from all NICER and STROBE-X sources can be relaxed if ADM is considered.

Future-X can more tightly constrain the neutron star EoS than Future.

![](_page_10_Figure_4.jpeg)

<sup>7</sup> Future / Future - X: Varying Baryonic EoS

> No ADM

Model

- The 'Including ADM' band is noticeably wider than the 'Neglecting ADM' band.
- A stiffer baryonic EoS ⇒ posterior constraints from all NICER and STROBE-X sources can be relaxed if ADM is considered.
  - Future-X can more tightly constrain the neutron star EoS than Future.

![](_page_11_Figure_4.jpeg)

8 Future / Future - X: Fixed ADM Baryonic EoS Core Model The ratio of  $\log_{10}\left(\frac{g_{\chi}}{m_{\phi}/\text{MeV}}\right)$  and  $\log_{10}(m_{\gamma}/\text{MeV})$  is constrained to a stripe. Future-X Future  $\log_{10}(m_{\chi}/MeV) = 5.58^{+1}_{-1}$  $\log_{10}(m_{\gamma}/MeV) = 5.59^{+1.59}_{-1.63}$ Posterior Prior The stripe widens for "No ADM" Ground Truth model.  $\log_{10}(\frac{g_{\chi}}{m_{\star}/MeV}) = 0.43^{+1.59}_{-1.66}$  $\log_{10}(\frac{g_{\chi}}{m_{\phi}/MeV}) = 0.31^{+1.60}_{-1.58}$  $g_{\chi}$  /  $m_{\phi}$  and  $m_{\chi}$  are individually  $og_{10}(\frac{g_{\chi}}{m_{\phi}/MeV})$  $\log_{10}(\frac{g_{\chi}}{m_{\phi}/MeV})$ r unconstrained.  $\sim$  $F_{\gamma} = 7.43^{+2.36}_{-2.36}$  $F_{\chi} = 7.77^{+3.65}_{-3.71}$ Gaussian-like shape of the 1-D  $F_{\gamma}$ 20 ~% posteriors  $\implies$  tight constraints on  $F_{\gamma}$  $\sqrt{r}$  $\vec{x}$  $_{\times}$ can be imposed. ծ D 2 20 20 60 3.0 N.S. 20 ୫ 20 6. 19  $\sim \lambda$ 0 2 20 20 N.S 20 3 ୫ Future-X will be able to provide  $\log_{10}(m_{\chi}/MeV)$ Fχ  $\log_{10}(m_{\chi}/MeV)$ Fχ  $\log_{10}(\frac{g_{\chi}}{m_{\star}/MeV})$  $\log_{10}(\frac{g_{\chi}}{m_{\star}/MeV})$ 

tighter constraints on  $F_{\gamma}$  than Future.

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8 Future / Future - X: Fixed No Baryonic EoS ADM Model The ratio of  $\log_{10}\left(\frac{g_{\chi}}{m_{\phi}/\text{MeV}}\right)$  and  $\log_{10}(m_{\gamma}/\text{MeV})$  is constrained to a stripe. Future Future-X  $\log_{10}(m_{\chi}/MeV) = 5.77^{+1.4}_{-1.5}$  $\log_{10}(m_{\chi}/MeV) = 5.70^{+1.5}_{-1.5}$ The stripe widens for "No ADM" model.  $\log_{10}(\frac{g_{\chi}}{m_{\star}/MeV}) = 0.12^{+1.53}_{-1.42}$  $\log_{10}(\frac{g_{\chi}}{m_{\bullet}/MeV}) = 0.22^{+1.59}_{-1.49}$  $g_{\chi}$  /  $m_{\phi}$  and  $m_{\chi}$  are individually  $\log_{10}(\frac{g_{\chi}}{m_{\phi}/MeV})$ log<sub>10</sub>(<u>g<sub>x</sub></u>) unconstrained. 0  $F_{\gamma} = 1.51^{+1.65}_{-1.05}$  $F_{\chi} = 3.26^{+3.17}_{-2.25}$ Gaussian-like shape of the 1-D  $F_{\gamma}$ 2º posteriors  $\Rightarrow$  tight constraints on  $F_{\gamma}$ Ŷ  $\vec{x}$ ч× can be imposed. ଚ 3.0 N.S. 6.0 19 0 60 1?  $\mathcal{V}$ \$ 20 *Future-X* will be able to provide  $\log_{10}(m_{\gamma}/MeV)$ Fχ  $\log_{10}(\frac{g_{\chi}}{m_{\star}/MeV})$  $\log_{10}(m_{\chi}/MeV)$ Fχ  $\log_{10}(\frac{g_{\chi}}{m_*/MeV})$ tighter constraints on  $F_{\gamma}$  than Future.

### Takeaways

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- The current uncertainties of the baryonic EoS are being underestimated because the possibility of ADM cores is not currently being accounted for.
- If the baryonic EoS is constrained independent of ADM, i.e., fixed, the ratio of  $\log_{10}\left(\frac{g_{\chi}}{m_{\phi}/\text{MeV}}\right)$  and  $\log_{10}(m_{\chi}/\text{MeV})$ , and  $F_{\chi}$  can be tightly constrained.
- We have shown the value in performing full inference runs on the ADM parameter space, rather than drawing conclusions only from the effects of ADM on the mass-radius relation.

![](_page_15_Figure_4.jpeg)

![](_page_16_Picture_0.jpeg)

# Thank you!

![](_page_16_Picture_2.jpeg)

![](_page_16_Picture_3.jpeg)

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# Extra Slides

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### <sup>2</sup>Accumulation methods of ADM in neutron stars

- One possibility: neutron bremsstrahlung of ADM<sup>1</sup> and neutron conversion to scalar ADM<sup>2</sup>
- Both processes combined can produce ADM masses of 0.07 M<sub>Ns</sub>.
- To achieve high ADM fractions for higher ADM particle masses other possibilities must be considered:
  - Accretion of baryonic matter onto a pre-existing ADM core<sup>3</sup>
    - A neutron star passed through a local ADM over-density<sup>3</sup>
      - Absorption of an ADM star by baryonic matter<sup>2</sup>

<u>Nelson et al. (2018) arXiv: 1803.03266</u> <u>Ellis et al. (2018) arXiv: 1804.01418</u> Karkevandi et al. (2022) arXiv: 2109.0380

- Baryonic EoS (extra slide) The PDF contours widen along the  $\log_{10}\left(\frac{g_{\chi}}{m_{\phi} \text{ MeV}^{-1}}/(m_{\chi}/\text{MeV})\right)$ 
  - axis for low  $F_{\gamma}$ .

If the actual  $F_{\chi}$  in neutron stars is sufficiently large, the ratio of  $g_{\chi}/m_{\phi}$  and  $m_{\chi}$  can be well constrained.

![](_page_19_Figure_3.jpeg)

- The PDF contours widen along the  $\log_{10}\left(\frac{g_{\chi}}{m_{\phi} \text{ MeV}^{-1}}/(m_{\chi}/\text{MeV})\right)$ 
  - axis for low  $F_{\gamma}$ .

If the actual  $F_{\chi}$  in neutron stars is sufficiently large, the ratio of  $g_{\chi}/m_{\phi}$  and  $m_{\chi}$  can be well constrained.

![](_page_20_Figure_3.jpeg)