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Decaying Sterile Neutrinos at MicroBooNE

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Outlines

- Why sterile neutrino?
- Decaying sterile neutrino
- Formalism
- Event rates
- Oscillation fits
- Decay fits
- Conclusion





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As an explanation for SBL anomalies

 $P_{\nu_{\mu} \to \nu_{e}} = 4 \left| U_{e4} \right|^{2} \left| U_{\mu 4} \right|^{2} \sin^{2} \left(\frac{\Delta m_{41}^{2} L}{4E} \right)$

A sterile neutrino only participates in oscillation.

However, significant tension remains

A simple 3+1 model is not enough

This motivates us to go beyond oscillations





MicroBooNE



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Decaying sterile neutrinos (3+1+decay)



MiniBooNE excess is interpreted as the active neutrinos from sterile decay product.

 ν_e appearance signal is only suppressed by the square of the mixing, not 4th power in the oscillation case ($4|U_{\mu4}|^2|U_{e4}|^2$). Therefore, smaller $U_{\mu4}$ is allowed, evading limits from ν_{μ} disappearance.



Decaying sterile neutrinos (3+1+decay)

$$-\mathcal{L} \supset g_{\phi}\overline{\nu_{s}}\nu_{s}\phi + \sum_{\alpha,\beta} m_{\alpha\beta}\overline{\nu}_{\alpha}\nu_{\beta}$$

Decay width:

$$\Gamma_{\nu_4}^{(\mathrm{I})} = \Gamma_{\nu_4 \to \hat{\nu}_s \phi} = |U_{s4}|^2 (1 - |U_{s4}|^2) \frac{g_{\phi}^2}{16\pi} \frac{m_4^2}{E_4}$$

Normalized active states:

$$|\hat{\nu}_{s}\rangle = \frac{\sum_{i=1}^{3} U_{fi}^{*} |\nu_{i}\rangle}{\left(\sum_{k=1}^{3} |U_{fk}|^{2}\right)^{1/2}}$$



 $|U_{\mu4}|^2 |U_{s4}|^2 g_{\varphi}$

 ν_s

 $\nu_{\mu} \quad \nu_{4}$

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 $(1 - |U_{s4}|^2)$

 $\hat{\nu}_{\rm s}$

Decaying sterile neutrinos (3+1+decay)

$$\nu_{\beta} \text{ flux from a } \nu_{\alpha} \text{ source:}$$

$$\Phi_{\nu_{\beta}}(L, E_{\nu}) = \int_{E_{4}^{\min}}^{\infty} dE_{4} \Phi_{\nu_{\alpha}}(L = 0, E_{4}) P_{\alpha\beta}(L, E_{4}, E_{\nu})$$

$$P_{\alpha\beta} = P_{\alpha\beta}^{\text{dec}} S_{\alpha\beta}^{\text{dec}} + P_{\alpha\beta}^{\text{osc}}$$

 $P_{\alpha\beta}^{\text{osc}}$: ν_4 that is yet to decay at baseline L produces ν_e through oscillation.

 $P_{\alpha\beta}^{\text{dec}}S^{\text{dec}}$: ν_4 that decays into active states.

For helicity-conserving decays, $S_{\alpha\beta}^{\text{dec}}(E_4, E_\nu) = \frac{1}{\Gamma_{\nu_4}} \frac{d\Gamma_{\nu_4 \to \nu\phi}}{dE_\nu} = \frac{E_\nu}{E_4}$



Disappearance of the intrinsic ν_e







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$| u_{\mu} ightarrow u_{e}$ Appearance signal



MiniBooNE favors decay because of the low energy events from decay. Although there is some penalty from detector efficiency, cross section and the helicity-conserving factor.



Oscillation Fits – varying $|U_{e4}|^2$





Decay Fits – varying coupling





Decay Fits – varying $|U_{e4}|^2$





Conclusion

- The minimal 3+1 model is not enough to reconcile all the anomalies.
- Decaying sterile neutrinos predicts LEE at MiniBooNE and fits better.
- We present the first comprehensive fit to MicroBooNE, accounting for disappearance, energy loss, etc.
- Decaying sterile solution to MiniBooNE is ruled out by MicroBooNE at more than 95% CL.
- In principle it can also explain the BEST anomaly.



Backup – 3+1 slices



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Backup – 3+1+decay slices



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Oscillation probability, best-fits

$$P_{\alpha\beta} = P_{\alpha\beta}^{\text{dec}} S_{\alpha\beta}^{\text{dec}} + P_{\alpha\beta}^{\text{osc}}$$

$$P_{\alpha\beta}^{\text{osc}}(L, E_{\nu}) = \delta_{\alpha\beta} - 2\delta_{\alpha\beta} |U_{\alpha4}U_{\beta4}| \left[1 - e^{-\frac{L}{2L_{\text{dec}}}} \cos\left(\pi \frac{L}{L_{\text{osc}}}\right) \right] + |U_{\alpha4}U_{\beta4}|^2 \left[1 - 2e^{-\frac{L}{2L_{\text{dec}}}} \cos\left(\pi \frac{L}{L_{\text{osc}}}\right) + e^{-\frac{L}{L_{\text{dec}}}} \right]$$

$$P_{\alpha\beta}^{\text{dec}}(L, E_4, E_{\nu}) = |U_{\alpha4}|^2 \frac{|\langle \hat{\nu}_s \mid \nu_\beta \rangle|^2}{|\langle \hat{\nu}_s \mid \hat{\nu}_s \rangle|^2} \left(1 - e^{-\frac{L}{L_{\text{dec}}}} \right)$$

New Physics Model	$ U_{e4} ^2$	$ U_{\mu 4} ^2$	g_arphi or g_e	Δm^2_{41}	$p_{ m MB}^{ m val}$	$\chi^2_{\mu m B}-\chi^2_{\mu m B,Null}$
Decay model (I)	0.21	0.15	3	0.1 eV^2	36%	35
	0.19	0.17	2.5 (fixed)	0.1 eV^2	39%	22
	0.013	0.012	1.0 (fixed)	$2.4 \times 10^2 \text{ eV}^2$	38%	25

