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Decaying Sterile Neutrinos at MicroBooNE

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Outlines

- Why sterile neutrino?
- Decaying sterile neutrino
- Formalism
- Event rates
- Oscillation fits
- Decay fits
- Conclusion

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As an explanation for SBL anomalies

$$
P_{\nu_{\mu}\to\nu_{e}} = 4|U_{e4}|^{2}|U_{\mu4}|^{2}\sin^{2}\left(\frac{\Delta m_{41}^{2}L}{4E}\right)
$$

A sterile neutrino only participates in oscillation.

However, significant tension remains

A simple 3+1 model is not enough

This motivates us to go beyond oscillations

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MicroBooNE

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Decaying sterile neutrinos (3+1+decay)

MiniBooNE excess is interpreted as the active neutrinos from sterile decay product.

 ν_e appearance signal is only suppressed by the square of the mixing, not 4th power in the oscillation case ($4|U_{\mu 4}|^2|U_{e 4}|^2$). Therefore, smaller $U_{\mu 4}$ is allowed, evading limits from ν_{μ} disappearance.

Decaying sterile neutrinos (3+1+decay)

$$
-\mathcal{L} \supset g_\phi \overline{\nu_s} \nu_s \phi + \sum_{\alpha,\beta} m_{\alpha\beta} \overline{\nu}_\alpha \nu_\beta
$$

Decay width:

$$
\Gamma^{(I)}_{\nu_4} = \Gamma_{\nu_4 \to \hat{\nu}_s \phi} = |U_{s4}|^2 (1 - |U_{s4}|^2) \frac{g_{\phi}^2}{16\pi} \frac{m_4^2}{E_4}
$$

Normalized active states:

$$
|\hat{\nu}_s\rangle = \frac{\sum_{i=1}^3 U_{fi}^* |\nu_i\rangle}{\left(\sum_{k=1}^3 |U_{fk}|^2\right)^{1/2}}
$$

 $|U_{\mu 4}|^2 |U_{s4}|^2 g_{\varphi}$

 ν_{s}

 ν_μ ν_4

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 $(1-|U_{s4}|^2)$

 $\boldsymbol{\hat{\nu}}_{\rm c}$

Decaying sterile neutrinos (3+1+decay)

$$
\nu_{\beta} \text{ flux from a } \nu_{\alpha} \text{ source:}
$$
\n
$$
\Phi_{\nu_{\beta}}(L, E_{\nu}) = \int_{E_4^{\text{min}}}^{\infty} dE_4 \Phi_{\nu_{\alpha}}(L = 0, E_4) P_{\alpha \beta}(L, E_4, E_{\nu})
$$
\n
$$
P_{\alpha \beta} = \boxed{P_{\alpha \beta}^{\text{dec}} S_{\alpha \beta}^{\text{dec}}} + \boxed{P_{\alpha \beta}^{\text{osc}}}
$$

 $\left|P_{\alpha\beta}^{\mathrm{osc}}\right|$: ν_4 that is yet to decay at baseline L produces ν_e through oscillation.

 $P^{\text{dec}}_{\alpha\beta}S^{\text{dec}}$: ν_4 that decays into active states.

For helicity-conserving decays, $S_{\alpha\beta}^{\text{dec}}(E_4, E_\nu) = \frac{1}{\Gamma_{\nu} \mu} \frac{d\Gamma_{\nu_4 \to \nu \phi}}{dE_{\nu_4}} = \frac{E_\nu}{E_A}$

Disappearance of the intrinsic

$\overline{\nu_{\mu}} \rightarrow \overline{\nu_{e}}$ Appearance signal

MiniBooNE favors decay because of the low energy events from decay. Although there is some penalty from detector efficiency, cross section and the helicity-conserving factor.

Oscillation Fits - varying $|\overline{U}_{e4}|^2$

Decay Fits - varying coupling

Decay Fits - varying $|U_{e4}|^2$

Conclusion

- The minimal $3+1$ model is not enough to reconcile all the anomalies.
- Decaying sterile neutrinos predicts LEE at MiniBooNE and fits better.
- We present the first comprehensive fit to MicroBooNE, accounting for disappearance, energy loss, etc.
- Decaying sterile solution to MiniBooNE is ruled out by MicroBooNE at more than 95% CL.
- In principle it can also explain the BEST anomaly.

$Backup - 3+1$ slices

$Backup - 3+1+decay$ slices

Oscillation probability, best-fits

 R^{D}

 Γ dec α dec

$$
P_{\alpha\beta} = P_{\alpha\beta}^{\text{acc}} S_{\alpha\beta}^{\text{acc}} + P_{\alpha\beta}^{\text{osc}}
$$

\n
$$
P_{\alpha\beta}^{\text{osc}}(L, E_{\nu}) = \delta_{\alpha\beta} - 2\delta_{\alpha\beta} |U_{\alpha4}U_{\beta4}| \left[1 - e^{-\frac{L}{2L_{\text{dec}}}} \cos\left(\pi \frac{L}{L_{\text{osc}}}\right) \right] + |U_{\alpha4}U_{\beta4}|^2 \left[1 - 2e^{-\frac{L}{2L_{\text{dec}}}} \cos\left(\pi \frac{L}{L_{\text{osc}}}\right) + e^{-\frac{L}{L_{\text{dec}}}} \right]
$$

\n
$$
P_{\alpha\beta}^{\text{dec}}(L, E_4, E_{\nu}) = |U_{\alpha4}|^2 \frac{|\langle \hat{\nu}_s | \nu_\beta \rangle|^2}{|\langle \hat{\nu}_s | \hat{\nu}_s \rangle|^2} \left(1 - e^{-\frac{L}{L_{\text{dec}}}} \right)
$$

