

Opening up baryon number violating operators

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Baryon number violation

Standard Model of Elementary Particles

	three generations of matter (fermions)			interactions / force carriers (bosons)	
	I	II	III		
mass	$\approx 2.2 \text{ MeV}/c^2$	$\approx 1.28 \text{ GeV}/c^2$	$\approx 173.1 \text{ GeV}/c^2$	0	$\approx 125.11 \text{ GeV}/c^2$
charge	$\frac{2}{3}$	$\frac{2}{3}$	$\frac{2}{3}$	0	0
spin	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1	0
	u up	c charm	t top	g gluon	H higgs
	d down	s strange	b bottom	γ photon	
	e electron	μ muon	τ tau	Z Z boson	
	ν_e electron neutrino	ν_μ muon neutrino	ν_τ tau neutrino	W W boson	

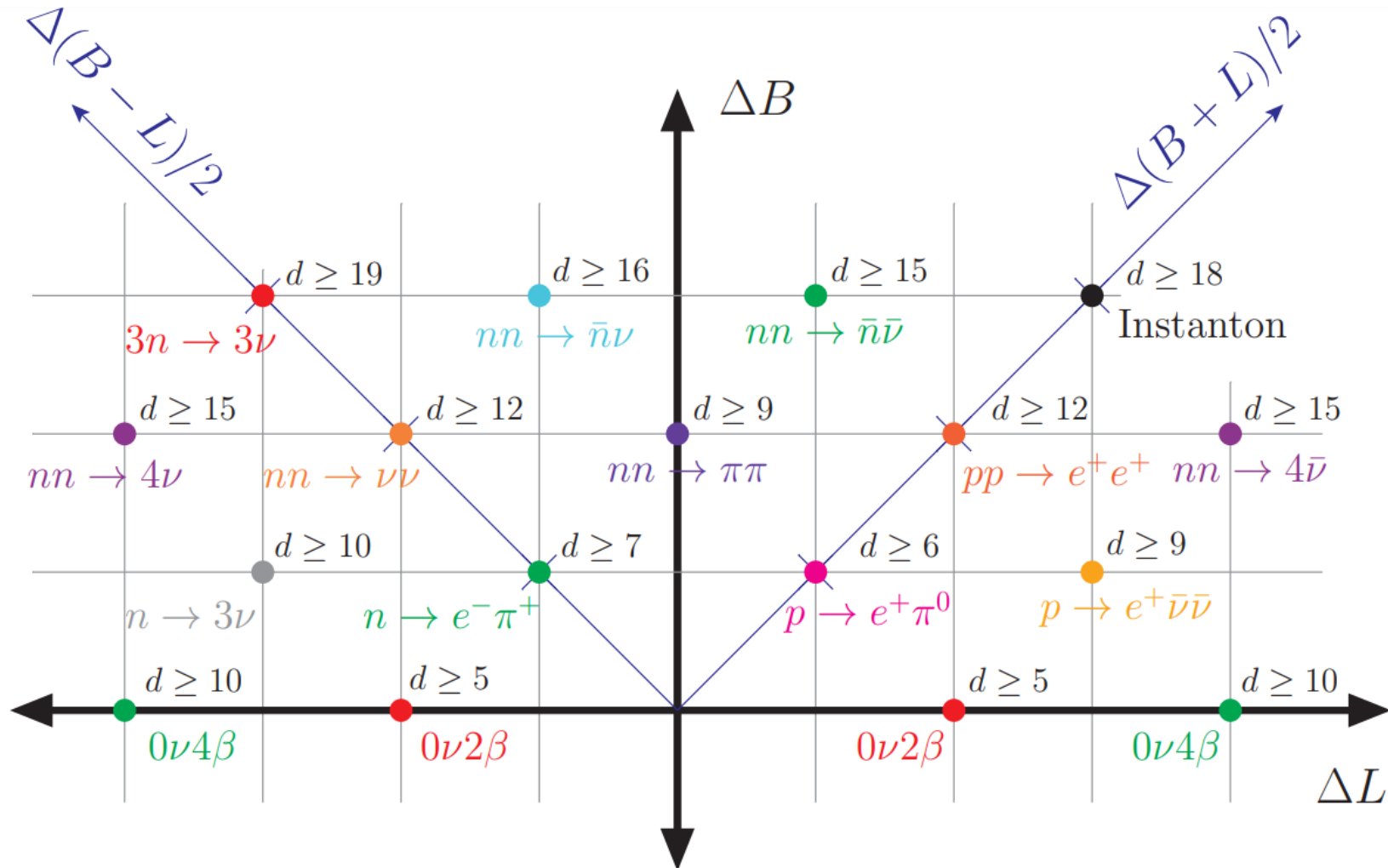
QUARKS (left side of quark boxes), **LEPTONS** (left side of lepton boxes), **GAUGE BOSONS VECTOR BOSONS** (right side of gauge boson boxes), **SCALAR BOSONS** (right side of Higgs box).

- Baryon and lepton number violation is sensitive to $d \gg 6$, unlike any other experiment.
- Use SMEFT to describe:

$$\mathcal{L} = \mathcal{L}_{SM} + \sum_{d=5} \frac{\sum_i C_i \mathcal{O}_i}{\Lambda^{d-4}}$$

We can assign $U(1)_B$ charge 1/3 to baryons

BNV landscape

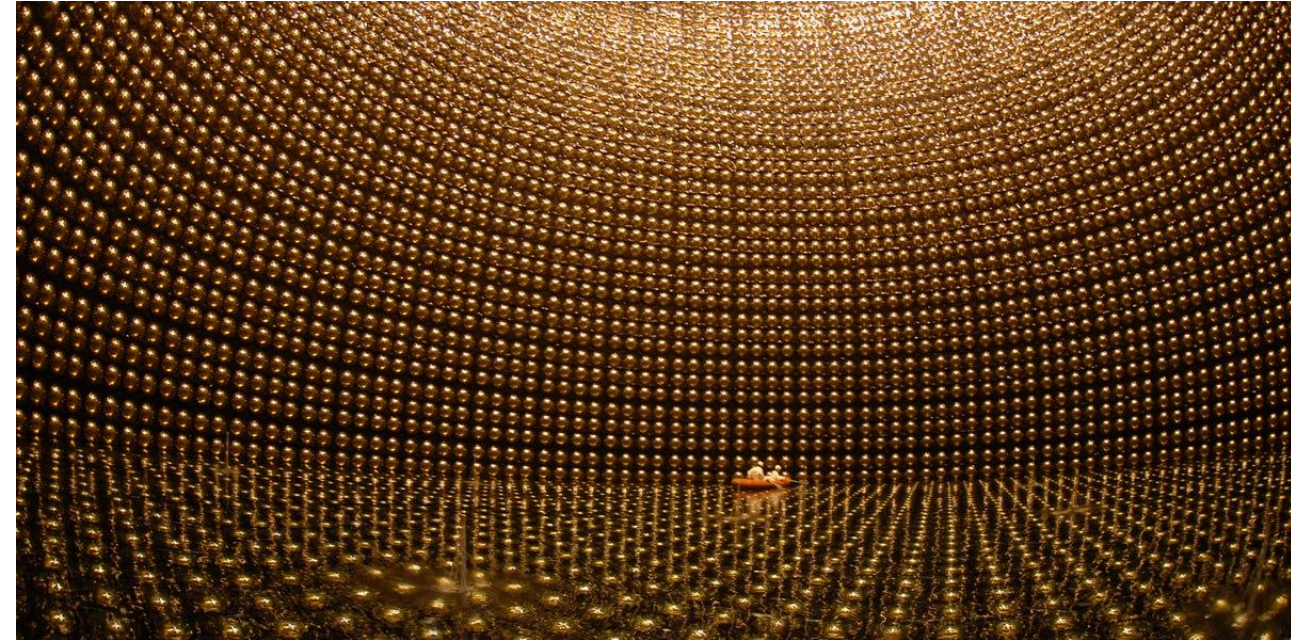
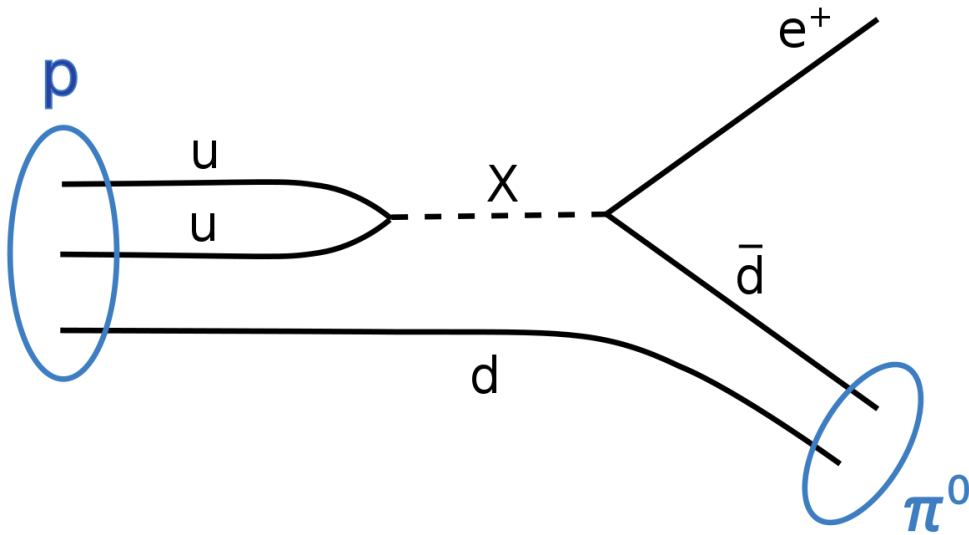


[J. Heeck, V. Takhistov ([arXiv:1910.07647](https://arxiv.org/abs/1910.07647))]

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Example: Proton decay

Super-Kamiokande.



$$\frac{uude}{\Lambda^2} \rightarrow \Gamma(p \rightarrow e^+ \pi^0) \sim \frac{m_p^5}{\Lambda^4} \sim \frac{1}{10^{34} \text{ years}} \left(\frac{10^{15} \text{ GeV}}{\Lambda} \right)^4$$

$$\tau(p \rightarrow e^+ \pi^0) > 10^{34} \text{ yr.}$$

Generate operators

$$O_{abcd}^{(1)} = (\bar{d}_{\alpha a R}^C u_{\beta b R}) (\bar{q}_{i \gamma c L}^c l_{j d L}) \epsilon_{\alpha \beta \gamma} \epsilon_{ij}, \quad (1)$$

$$O_{abcd}^{(2)} = (\bar{q}_{i \alpha a L}^c q_{j \beta b L}) (\bar{u}_{\gamma c R}^c l_{d R}) \epsilon_{\alpha \beta \gamma} \epsilon_{ij}, \quad (2)$$

$$O_{abcd}^{(3)} = (\bar{q}_{i \alpha a L}^c q_{j \beta b L}) (\bar{q}_{k \gamma c L}^c l_{l d L}) \epsilon_{\alpha \beta \gamma} \epsilon_{ij} \epsilon_{kl}, \quad (3)$$

$$O_{abcd}^{(4)} = (\bar{q}_{i \alpha a L}^c q_{j \beta b L}) (\bar{q}_{k \gamma c L}^c l_{l d L}) \epsilon_{\alpha \beta \gamma} \times (\hat{\tau} \epsilon)_{ij} \cdot (\hat{\tau} \epsilon)_{kl}, \quad (4)$$

$$O_{abcd}^{(5)} = (\bar{d}_{\alpha a R}^C u_{\beta b R}) (\bar{u}_{\gamma c R}^c l_{d R}) \epsilon_{\alpha \beta \gamma}, \quad (5)$$

$$O_{abcd}^{(6)} = (\bar{u}_{\alpha a R}^c u_{\beta b R}) (\bar{d}_{\gamma c R}^c l_{d R}) \epsilon_{\alpha \beta \gamma}. \quad (6)$$

[Weinberg, 1979]

Linearly independent operators
(including right-handed neutrinos).

$$O_{6abcd}^1 \equiv \epsilon^{\alpha \beta \gamma} \epsilon_{ij} (\bar{d}_{a, \alpha}^C u_{b, \beta}) (\bar{Q}_{i, c, \gamma}^C L_{j, d}),$$

$$O_{6abcd}^2 \equiv \epsilon^{\alpha \beta \gamma} \epsilon_{ij} (\bar{Q}_{i, a, \alpha}^C Q_{j, b, \beta}) (\bar{u}_{c, \gamma}^C l_d),$$

$$O_{6abcd}^3 \equiv \epsilon^{\alpha \beta \gamma} \epsilon_{il} \epsilon_{jk} (\bar{Q}_{i, a, \alpha}^C Q_{j, b, \beta}) (\bar{Q}_{k, c, \gamma}^C L_{l, d}),$$

$$O_{6abcd}^4 \equiv \epsilon^{\alpha \beta \gamma} (\bar{d}_{a, \alpha}^C u_{b, \beta}) (\bar{u}_{c, \gamma}^C l_d),$$

$$O_{6abcd}^5 \equiv \epsilon^{\alpha \beta \gamma} \epsilon_{ij} (\bar{Q}_{i, a, \alpha}^C Q_{j, b, \beta}) (\bar{d}_{c, \gamma}^C \nu_d),$$

$$O_{6abcd}^6 \equiv \epsilon^{\alpha \beta \gamma} (\bar{d}_{a, \alpha}^C u_{b, \beta}) (\bar{d}_{c, \gamma}^C \nu_d),$$

[Weinberg, 1979]

[Beltran, Cepedello, Hirsch ([arXiv:2306.12578](https://arxiv.org/abs/2306.12578))]

Generate operators

field	chirality	generations	$SU(3)_C \times SU(2)_L \times U(1)_Y$ representation
Q	left	3	$(\mathbf{3}, \mathbf{2}, \frac{1}{6})$
u	right	3	$(\mathbf{3}, \mathbf{1}, \frac{2}{3})$
d	right	3	$(\mathbf{3}, \mathbf{1}, -\frac{1}{3})$
L	left	3	$(\mathbf{1}, \mathbf{2}, -\frac{1}{2})$
ℓ	right	3	$(\mathbf{1}, \mathbf{1}, -1)$
ν	right	3	$(\mathbf{1}, \mathbf{1}, 0)$
H	scalar	1	$(\mathbf{1}, \mathbf{2}, \frac{1}{2})$



SYM2INT

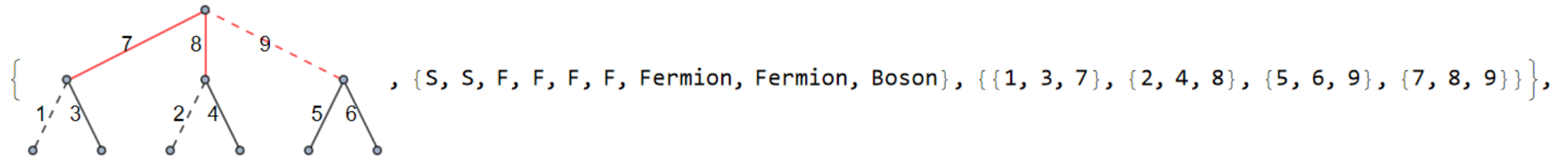
<https://renatofonseca.net/sym2int>

- Dim 6** $\mathcal{O}_6^1 \equiv duQL, \mathcal{O}_6^2 \equiv QQul, \mathcal{O}_6^3 \equiv QQQQL, \mathcal{O}_6^4 \equiv duul, \mathcal{O}_6^5 \equiv QQd\nu, \mathcal{O}_6^6 \equiv dud\nu.$
- Dim 7** $\mathcal{O}_7^1 \equiv QQd\bar{L}\bar{H}, \mathcal{O}_7^2 \equiv ddQ\bar{e}\bar{H}, \mathcal{O}_7^3 \equiv ddu\bar{L}\bar{H}, \mathcal{O}_7^4 \equiv ddd\bar{L}\bar{H},$
 $\mathcal{O}_7^5 \equiv duQ\bar{\nu}\bar{H}, \mathcal{O}_7^6 \equiv QQQ\bar{\nu}\bar{H}, \mathcal{O}_7^7 \equiv ddQ\bar{\nu}\bar{H}.$
- Dim 8** $\mathcal{O}_8^1 \equiv uuQL\bar{H}\bar{H}, \mathcal{O}_8^2 \equiv udu\bar{e}\bar{H}\bar{H}, \mathcal{O}_8^3 \equiv udQL\bar{H}\bar{H}, \mathcal{O}_8^4 \equiv QQ\bar{e}\bar{H}\bar{H}, \mathcal{O}_8^5 \equiv QQQ\bar{L}\bar{H}\bar{H},$
 $\mathcal{O}_8^6 \equiv ddQL\bar{H}\bar{H}, \mathcal{O}_8^7 \equiv QQd\bar{e}\bar{H}\bar{H}, \mathcal{O}_8^8 \equiv QQ\bar{u}\bar{\nu}\bar{H}\bar{H}, \mathcal{O}_8^9 \equiv udd\nu\bar{H}\bar{H}, \mathcal{O}_8^{10} \equiv QQd\nu\bar{H}\bar{H}.$
- Dim 9** $\mathcal{O}_{9,(1,3)}^1 \equiv uuQLLL$ + 61 more operators
- Dim 10** $\mathcal{O}_{10,(1,1)}^1 \equiv \bar{d}ddQQLH$ + 111 more operators
- Dim 11** $\mathcal{O}_{11,(1,-1)}^1 \equiv \bar{H}\bar{e}\bar{d}dddH$ + 130 more operators
- Dim 12** $\mathcal{O}_{12,(1,-3)}^1 \equiv \bar{L}\bar{L}\bar{e}\bar{u}dddd$ + 479 more operators
- Dim 13** $\mathcal{O}_{13,(1,3)}^1 \equiv \bar{d}QQQQQLLLH$ + 781 more operators
- Dim 14** $\mathcal{O}_{14,(1,5)}^1 \equiv uuQLLLLLHH$ + 983 more operators
- Dim 15** $\mathcal{O}_{15,(3,1)}^1 \equiv ddQQQQQQQL$ + 2626 more operators

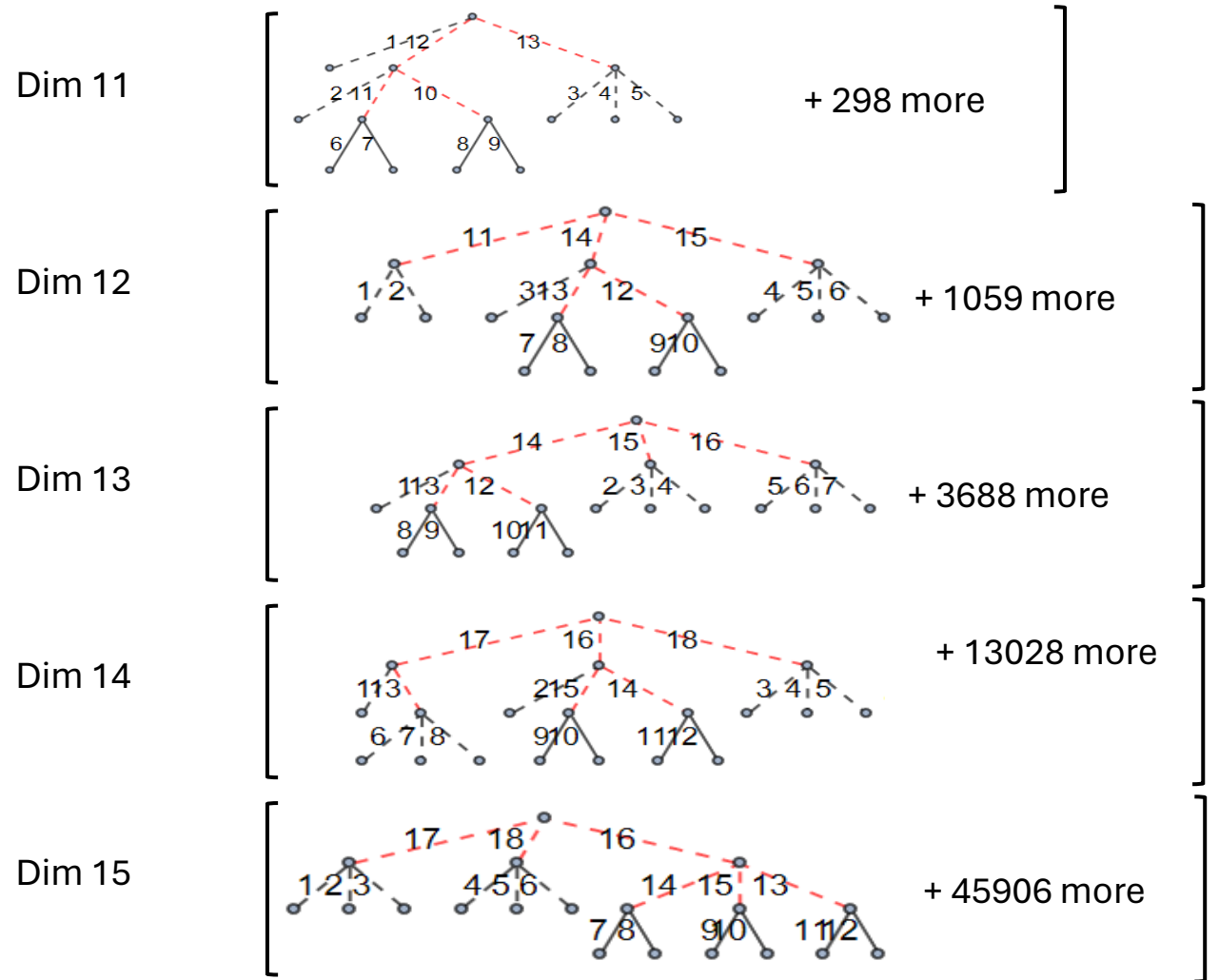
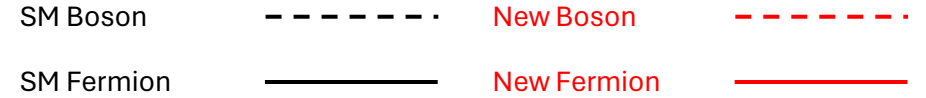
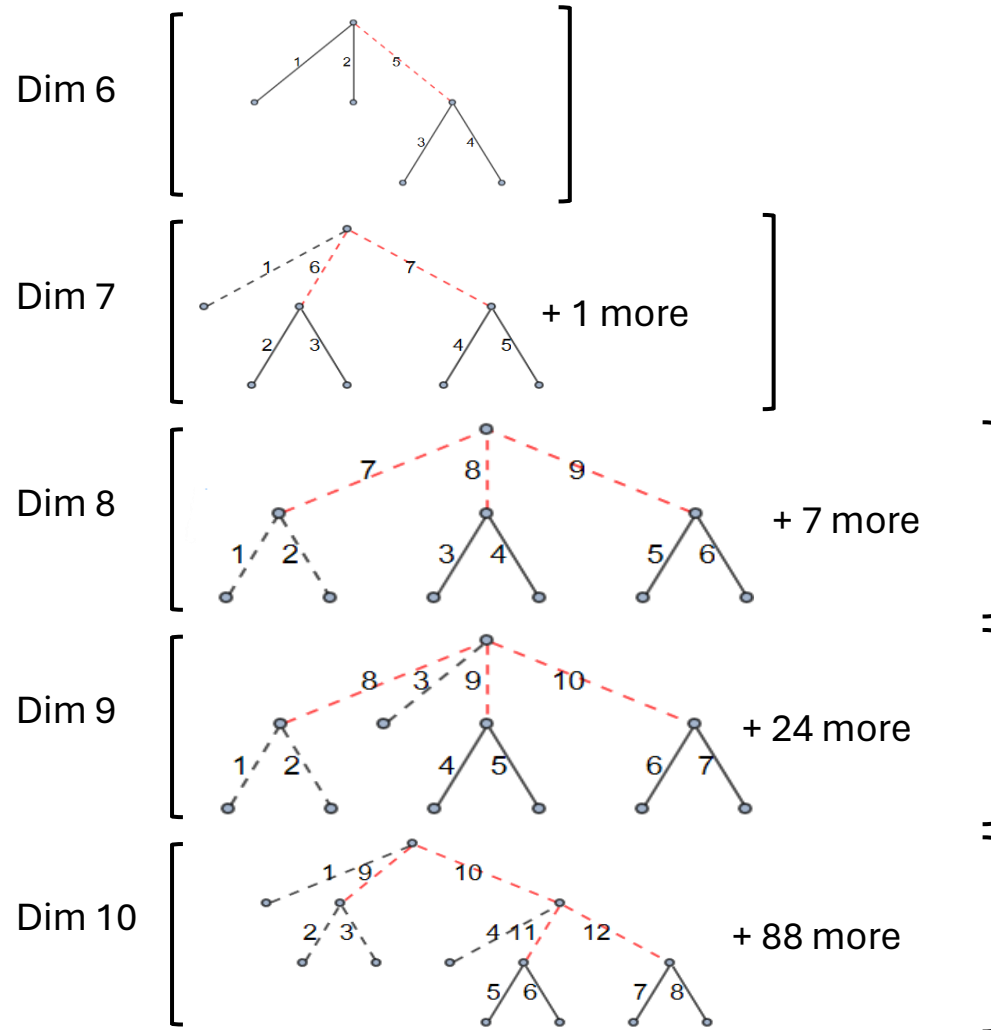
Generate topologies

- To obtain topologies, we use FeynArts Mathematica package to generate all possible diagrams and removed topologically same ones.

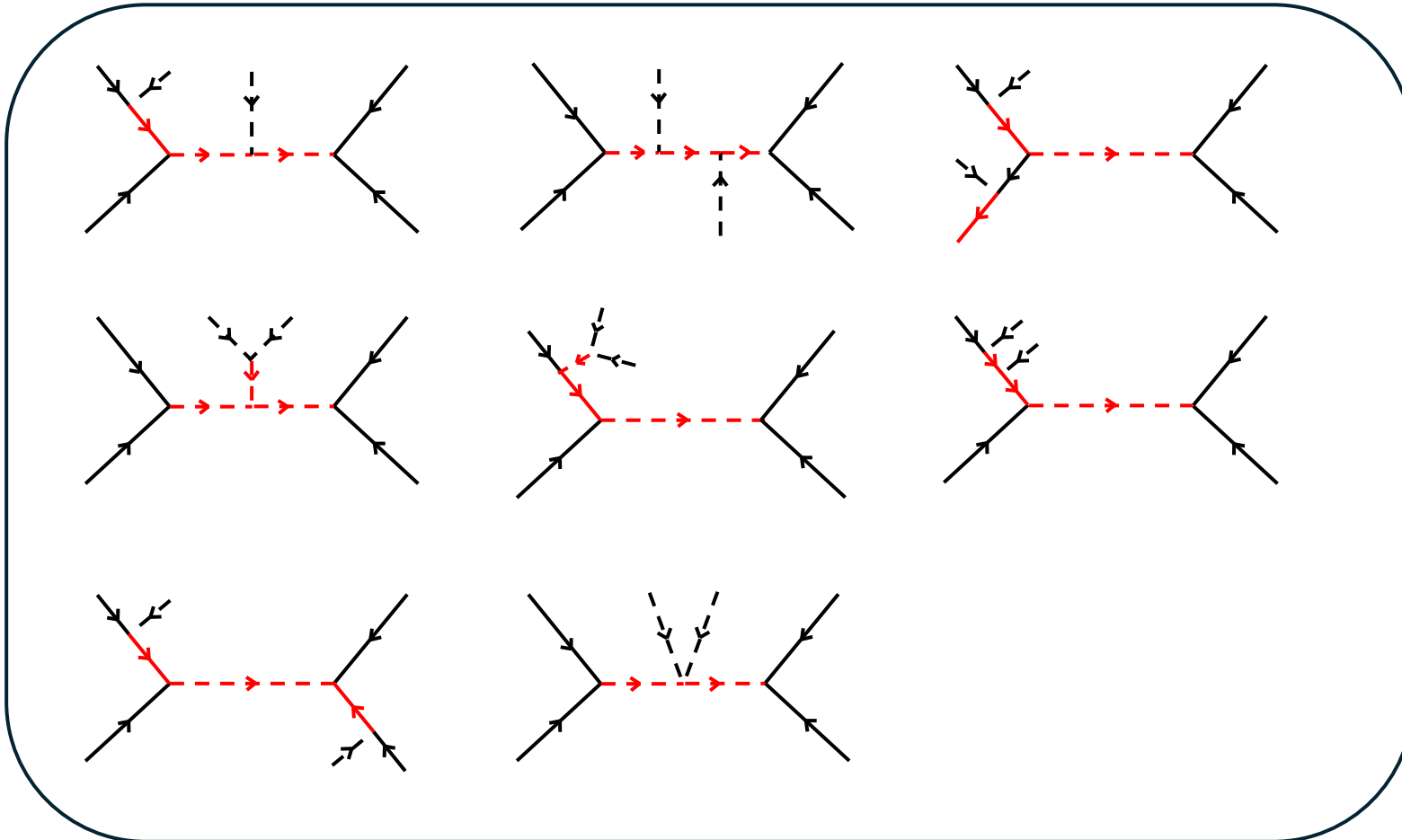
(<https://feynarts.de>)



Topologies

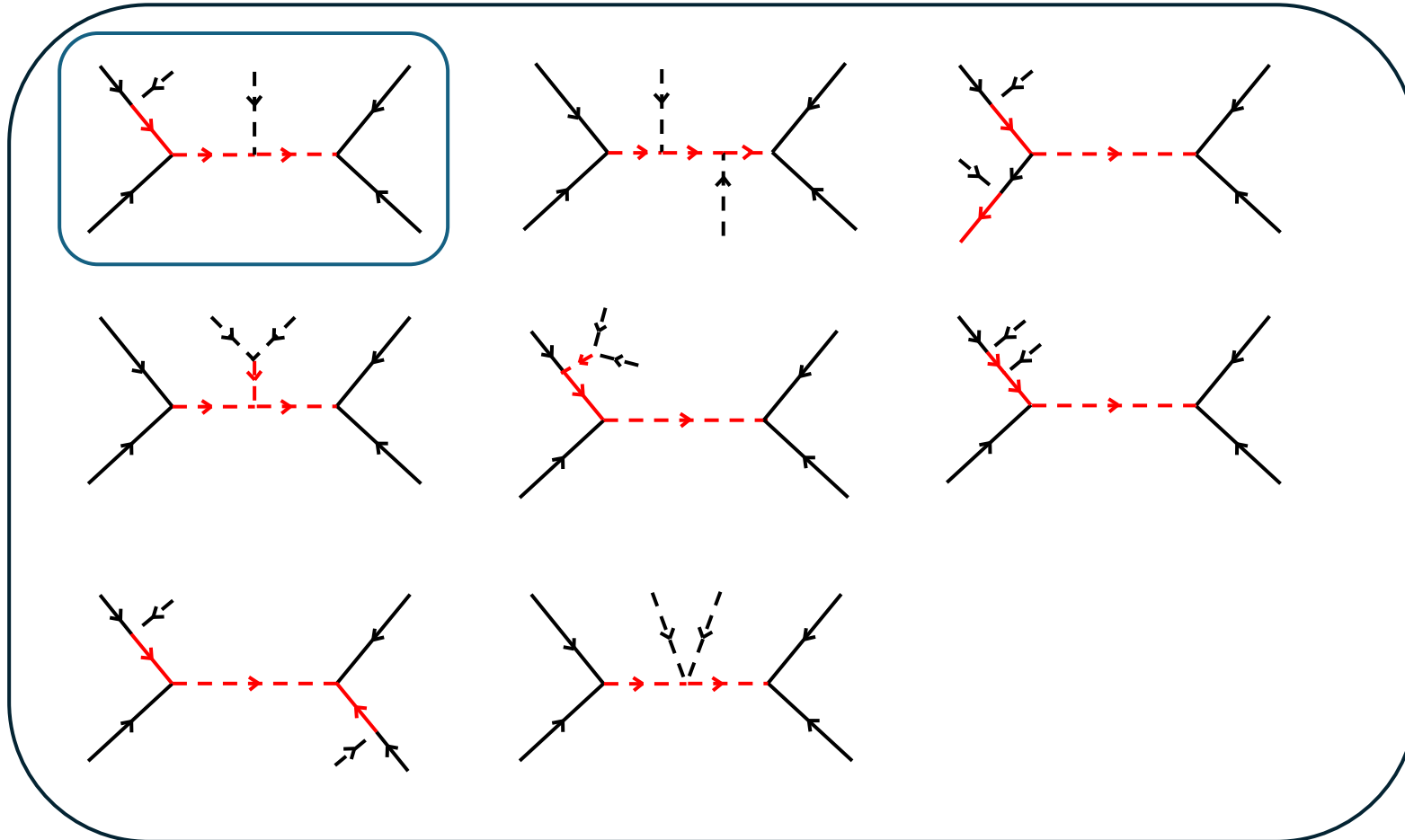


Example (Dimension 8)



$$\begin{aligned}
 \mathcal{O}_8^1 &\equiv uuQL\bar{H}\bar{H}, \\
 \mathcal{O}_8^2 &\equiv udu\bar{e}\bar{H}H, \\
 \mathcal{O}_8^3 &\equiv udQL\bar{H}H, \\
 \mathcal{O}_8^4 &\equiv QQe\bar{H}H, \\
 \mathcal{O}_8^5 &\equiv QQQ\bar{L}\bar{H}H, \\
 \mathcal{O}_8^6 &\equiv ddQLHH, \\
 \mathcal{O}_8^7 &\equiv QQdeHH, \\
 \mathcal{O}_8^8 &\equiv QQ\nu\nu\bar{H}\bar{H}, \\
 \mathcal{O}_8^9 &\equiv udd\nu\bar{H}H, \\
 \mathcal{O}_8^{10} &\equiv QQd\nu\bar{H}H.
 \end{aligned}$$

Example (Dimension 8)

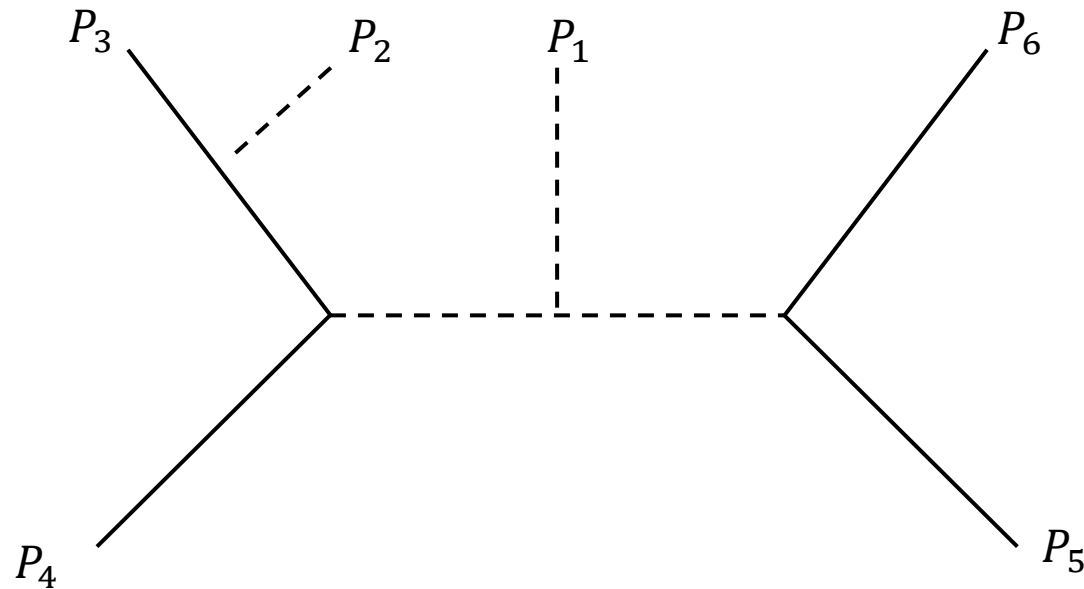


- $\mathcal{O}_8^1 \equiv uuQL\bar{H}\bar{H},$
- $\mathcal{O}_8^2 \equiv udu\bar{e}\bar{H}H,$
- $\mathcal{O}_8^3 \equiv udQL\bar{H}H,$
- $\mathcal{O}_8^4 \equiv QQe\bar{H}H,$
- $\mathcal{O}_8^5 \equiv QQQ\bar{L}\bar{H}H,$
- $\mathcal{O}_8^6 \equiv ddQLHH,$
- $\mathcal{O}_8^7 \equiv QQdeHH,$
- $\mathcal{O}_8^8 \equiv QQ\nu\nu\bar{H}\bar{H},$
- $\mathcal{O}_8^9 \equiv udd\nu\bar{H}H,$
- $\mathcal{O}_8^{10} \equiv QQd\nu\bar{H}H.$

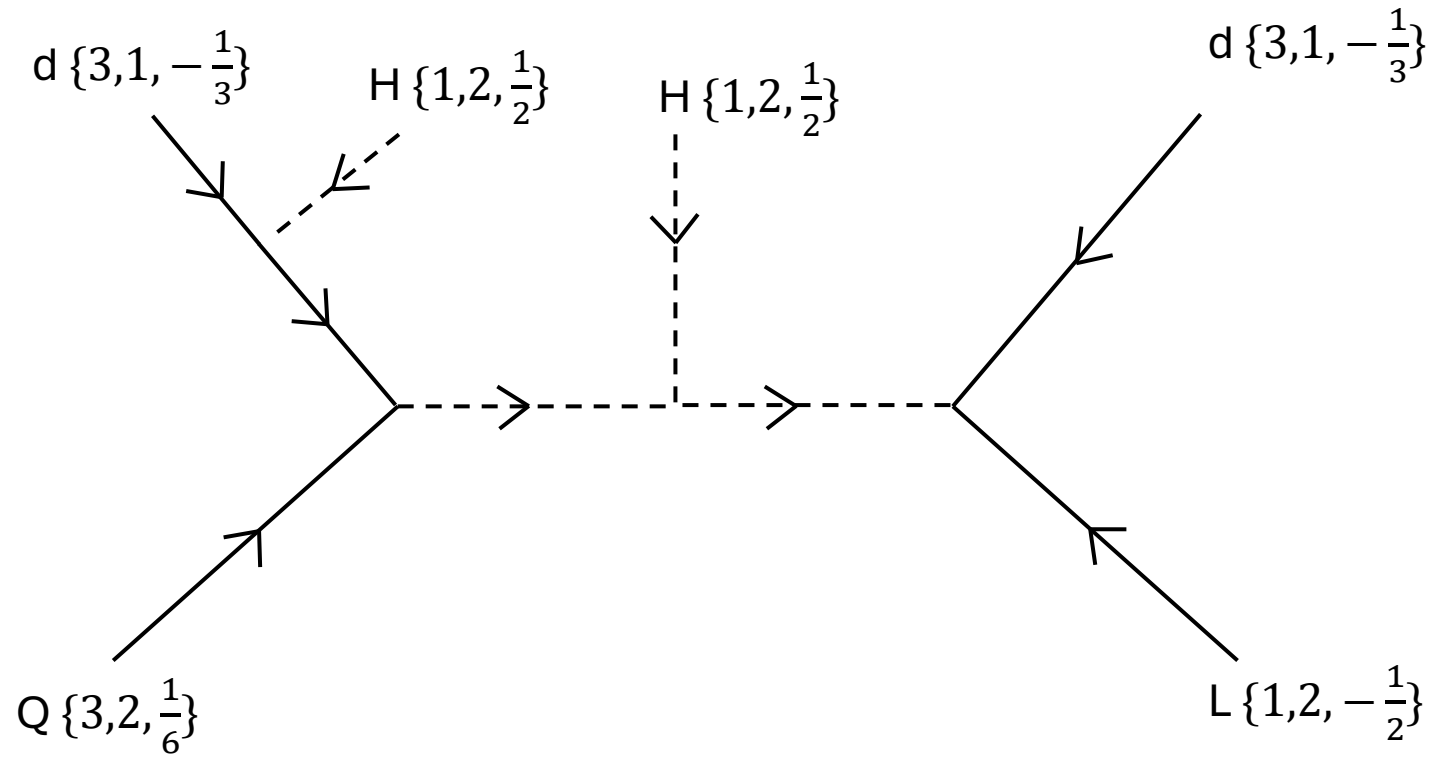
Possible permutations

Operator: ddLQHH

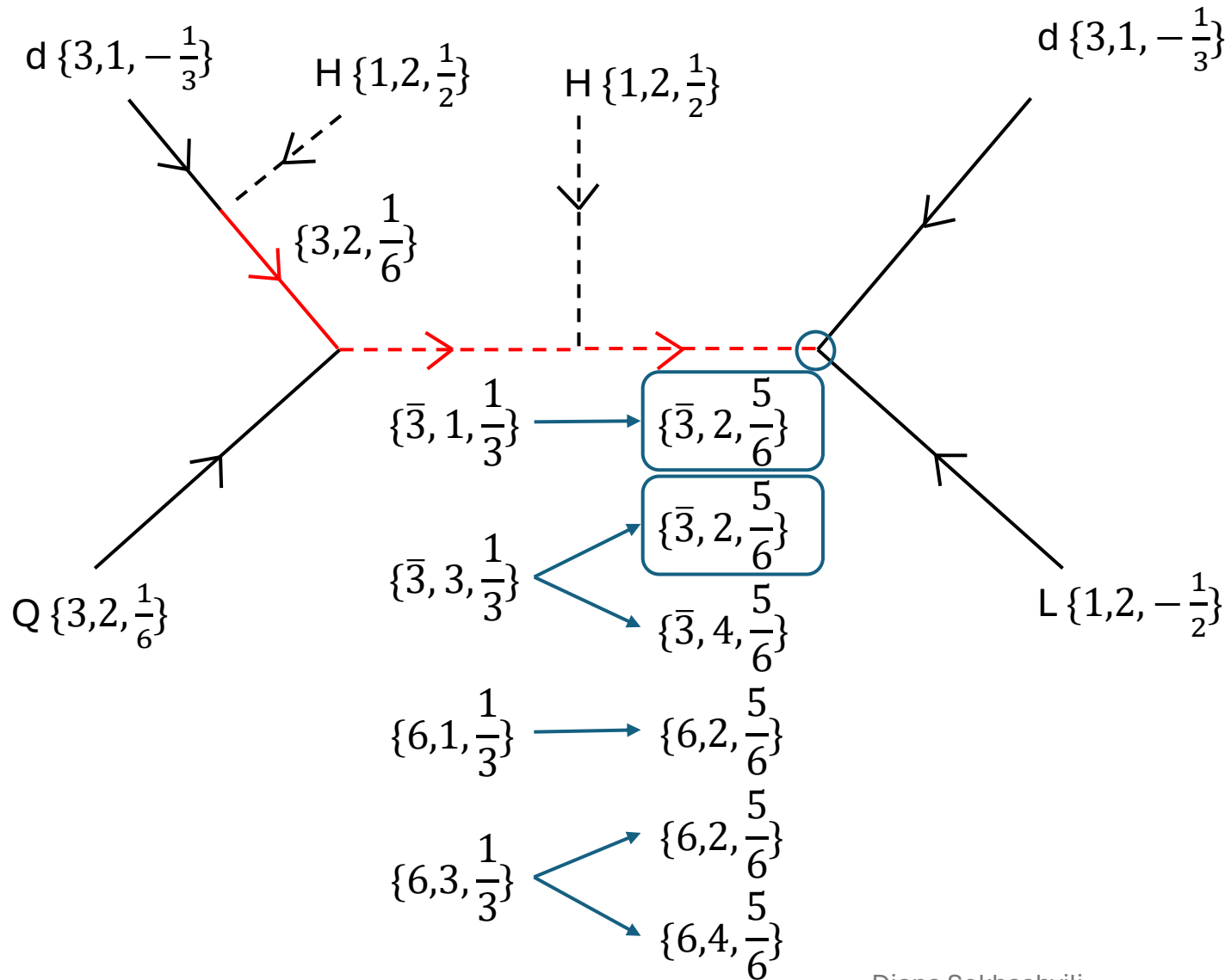
$\{\{H H d d L Q\}, \{H H d d Q L\}, \{H H d L d Q\}, \{H H d L Q d\}, \{H H d Q d L\}, \{H H d Q L d\},$
 $\{H H L d d Q\}, \{H H L d Q d\}, \{H H L Q d d\}, \{H H Q d d L\}, \{H H Q d L d\}, \{H H Q L d d\}\}$



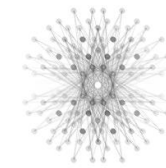
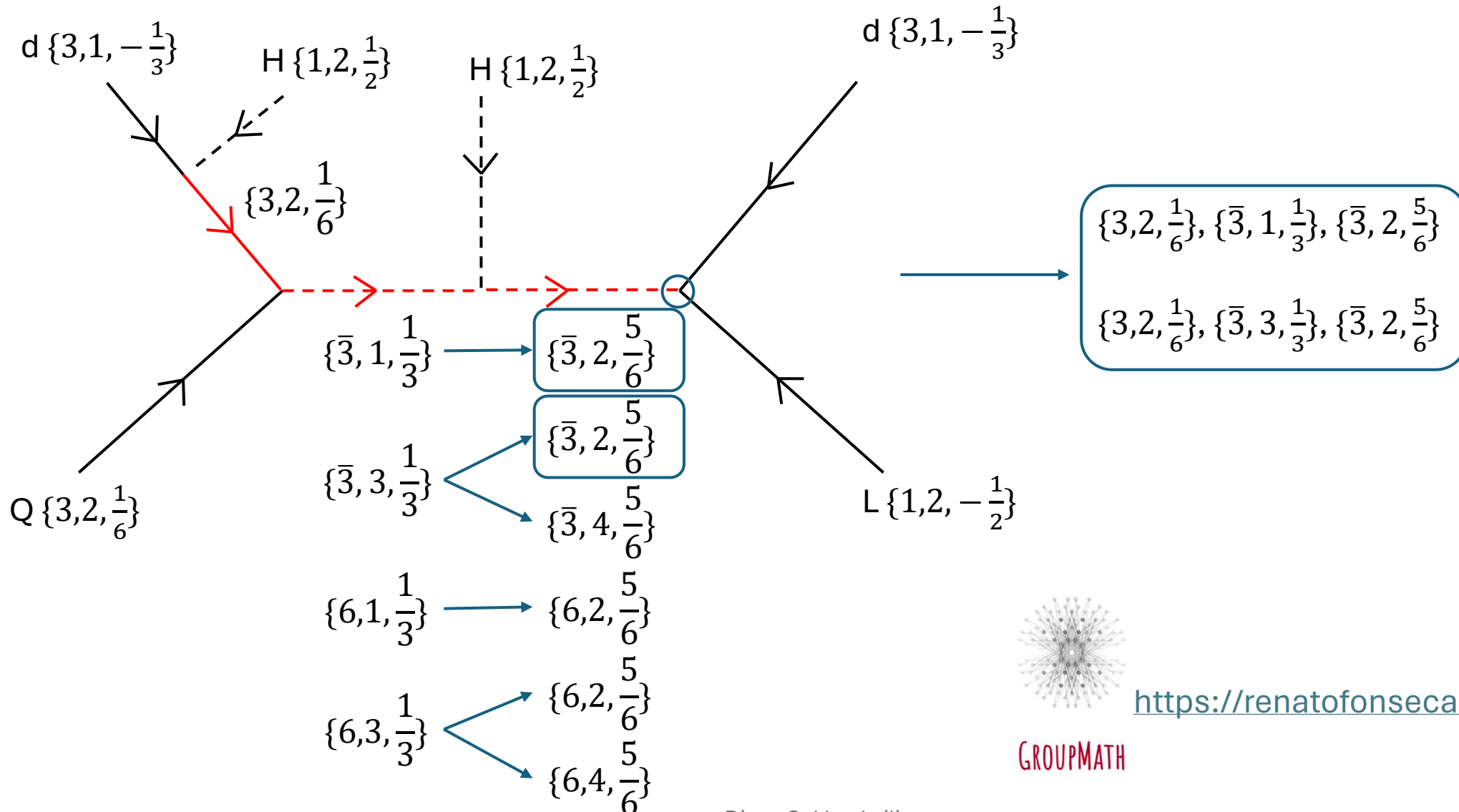
Internal particles



Internal particles



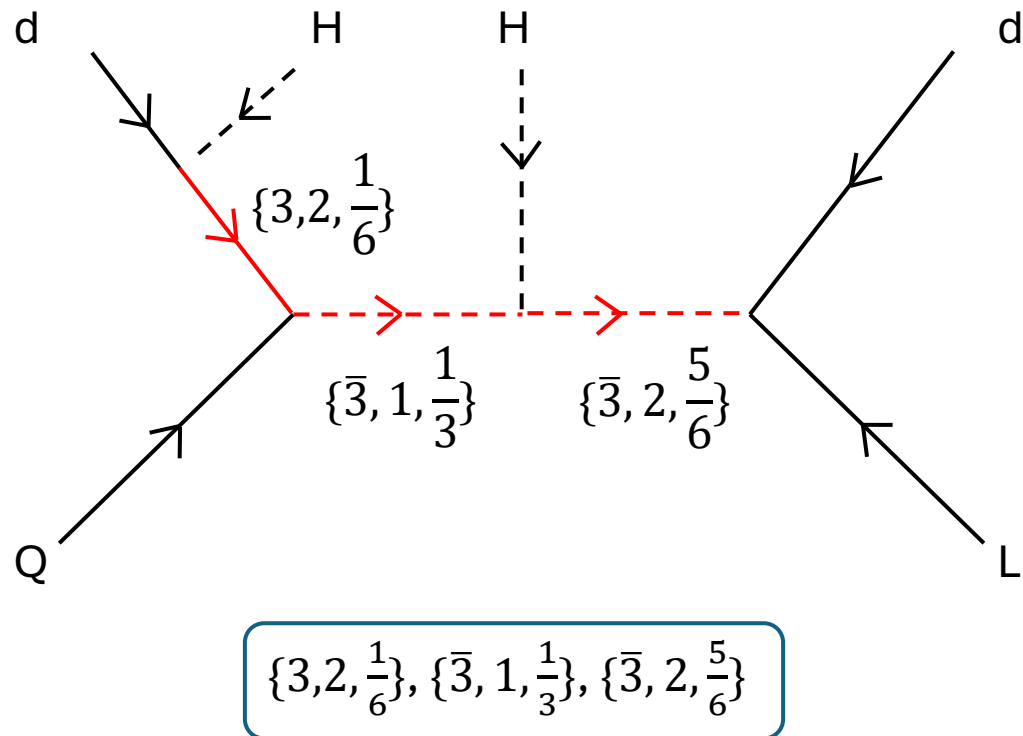
Internal particles



GROUPMATH

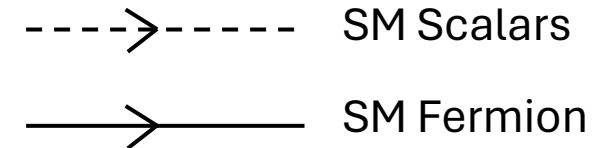
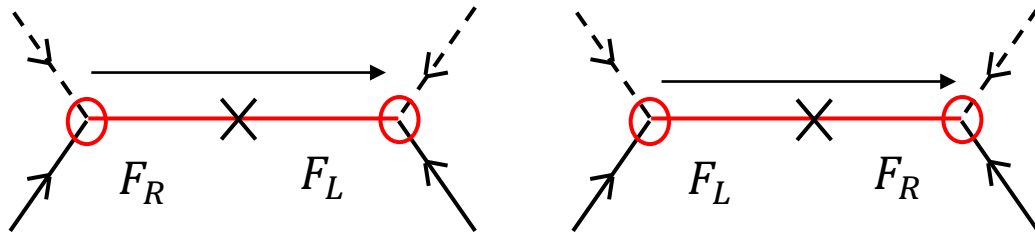
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Internal particles

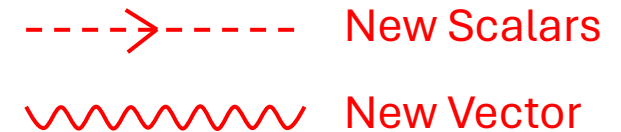
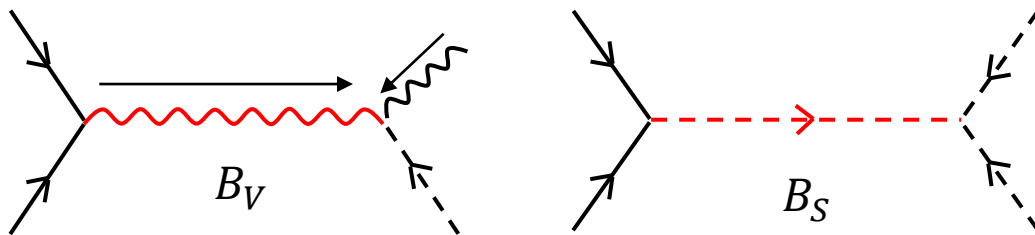


Lorentz group (SU2×SU2)

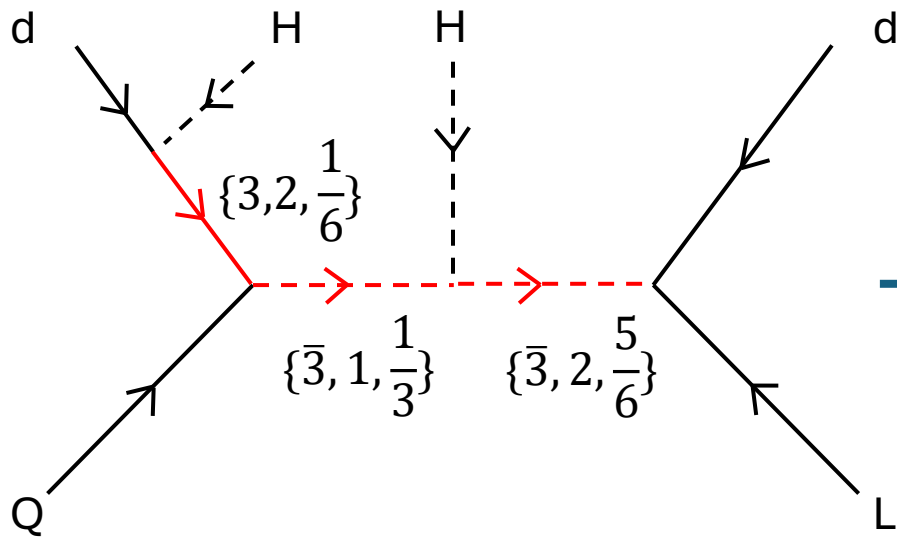
- For internal fermions we need to make sure that chirality is flipped to ensure that these fermions are massive.



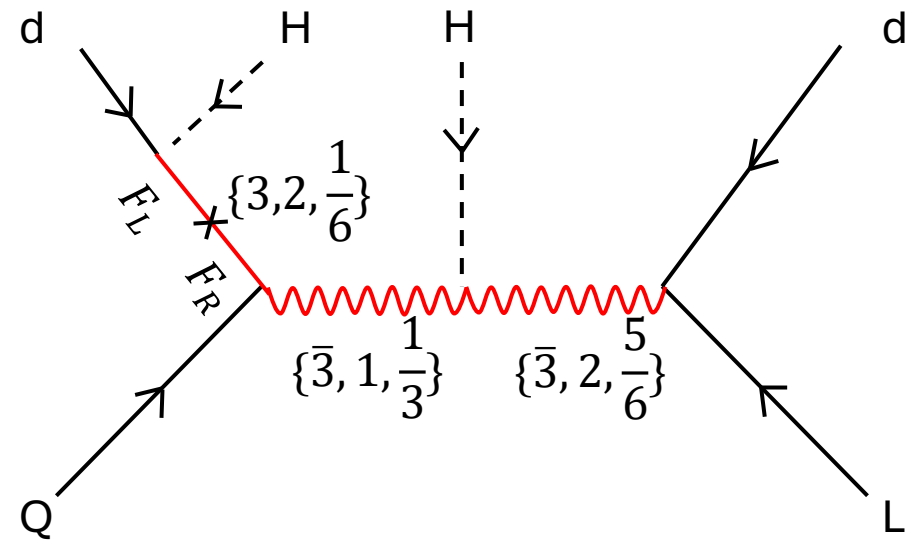
- For internal bosons we determine whether they are vector or scalar.



Internal particles



$$\left\{ 3, 2, \frac{1}{6} \right\}, \left\{ \bar{3}, 1, \frac{1}{3} \right\}, \left\{ \bar{3}, 2, \frac{5}{6} \right\}$$



$$\left\{ 3, 2, \frac{1}{6} \right\}_F, \left\{ \bar{3}, 1, \frac{1}{3} \right\}_V, \left\{ \bar{3}, 2, \frac{5}{6} \right\}_V$$



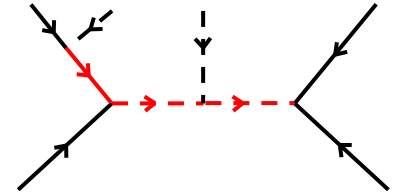
SYM2INT

<https://renatofonseca.net/sym2int>

Permuting fields around topology

Operator: ddLQHH

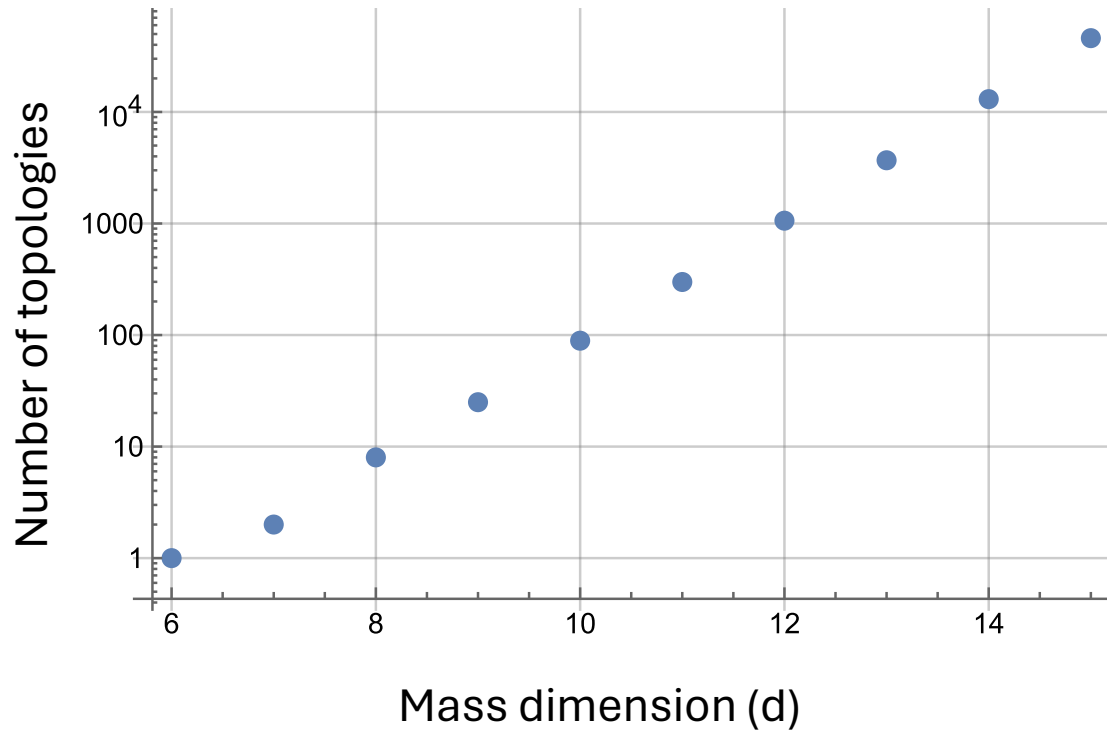
Topology:



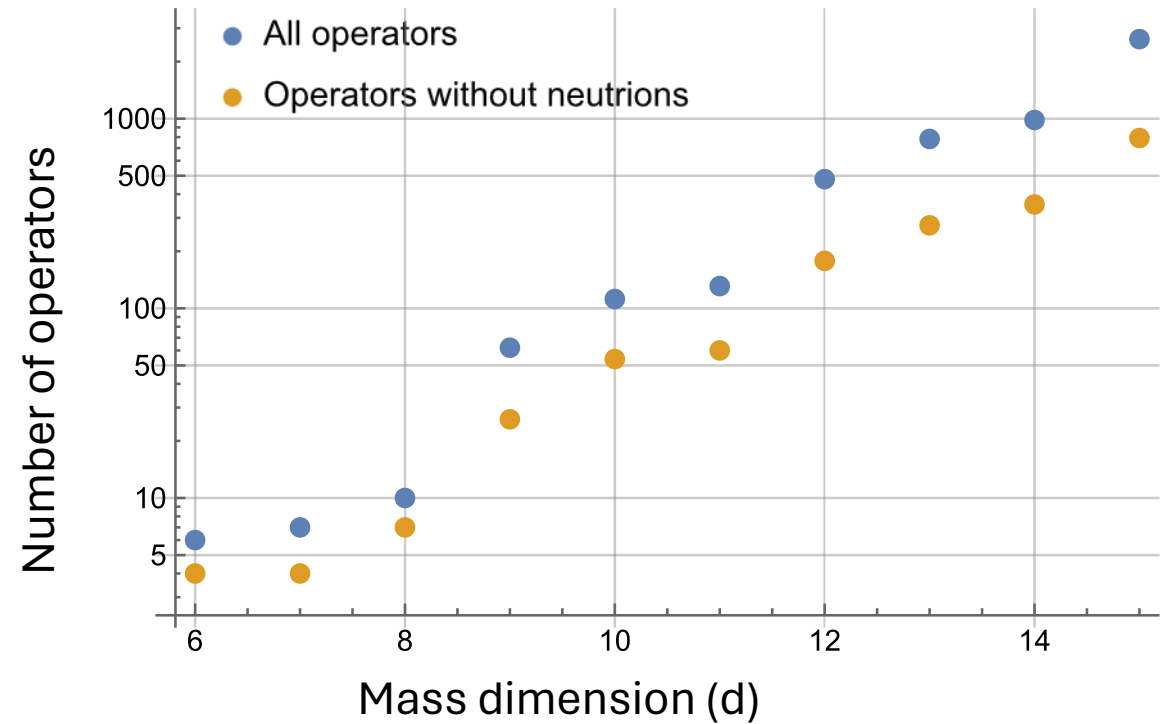
$$\begin{aligned}
 & \{ \{ \{ H, H, d, Q, L, d, \{ F, \{ 3, 2, \frac{1}{6} \} \}, \{ V, \{ \bar{3}, 3, \frac{1}{3} \} \}, \{ V, \{ \bar{3}, 2, \frac{5}{6} \} \} \} \}, \{ \{ H, H, d, Q, L, d, \{ F, \{ 3, 2, \frac{1}{6} \} \}, \{ V, \{ \bar{3}, 1, \frac{1}{3} \} \}, \{ V, \{ \bar{3}, 2, \frac{5}{6} \} \} \} \}, \\
 & \{ \{ H, H, d, L, Q, d, \{ F, \{ 3, 2, \frac{1}{6} \} \}, \{ V, \{ 3, 3, -\frac{1}{3} \} \}, \{ V, \{ 3, 2, \frac{1}{6} \} \} \} \}, \{ \{ H, H, d, L, Q, d, \{ F, \{ 3, 2, \frac{1}{6} \} \}, \{ V, \{ 3, 1, -\frac{1}{3} \} \}, \{ V, \{ 3, 2, \frac{1}{6} \} \} \} \}, \\
 & \{ \{ H, H, d, d, Q, L, \{ F, \{ 3, 2, \frac{1}{6} \} \}, \{ S, \{ \bar{3}, 2, -\frac{1}{6} \} \}, \{ S, \{ \bar{3}, 1, \frac{1}{3} \} \} \} \}, \{ \{ H, H, d, d, Q, L, \{ F, \{ 3, 2, \frac{1}{6} \} \}, \{ S, \{ \bar{3}, 2, -\frac{1}{6} \} \}, \{ S, \{ \bar{3}, 3, \frac{1}{3} \} \} \} \}, \\
 & \{ \{ H, H, Q, d, L, d, \{ F, \{ 3, 3, \frac{2}{3} \} \}, \{ V, \{ \bar{3}, 3, \frac{1}{3} \} \}, \{ V, \{ \bar{3}, 2, \frac{5}{6} \} \} \} \}, \{ \{ H, H, Q, d, L, d, \{ F, \{ 3, 1, \frac{2}{3} \} \}, \{ V, \{ \bar{3}, 1, \frac{1}{3} \} \}, \{ V, \{ \bar{3}, 2, \frac{5}{6} \} \} \} \}, \\
 & \{ \{ H, H, Q, L, d, d, \{ F, \{ 3, 3, \frac{2}{3} \} \}, \{ S, \{ 3, 2, \frac{1}{6} \} \}, \{ S, \{ 3, 1, \frac{2}{3} \} \} \} \}, \{ \{ H, H, Q, L, d, d, \{ F, \{ 3, 1, \frac{2}{3} \} \}, \{ S, \{ 3, 2, \frac{1}{6} \} \}, \{ S, \{ 3, 1, \frac{2}{3} \} \} \} \}, \\
 & \{ \{ H, H, L, d, Q, d, \{ F, \{ 1, 3, 0 \} \}, \{ V, \{ 3, 3, -\frac{1}{3} \} \}, \{ V, \{ 3, 2, \frac{1}{6} \} \} \} \}, \{ \{ H, H, L, d, Q, d, \{ F, \{ 1, 1, 0 \} \}, \{ V, \{ 3, 1, -\frac{1}{3} \} \}, \{ V, \{ 3, 2, \frac{1}{6} \} \} \} \}, \\
 & \{ \{ H, H, L, Q, d, d, \{ F, \{ 1, 3, 0 \} \}, \{ S, \{ 3, 2, \frac{1}{6} \} \}, \{ S, \{ 3, 1, \frac{2}{3} \} \} \} \}, \{ \{ H, H, L, Q, d, d, \{ F, \{ 1, 1, 0 \} \}, \{ S, \{ 3, 2, \frac{1}{6} \} \}, \{ S, \{ 3, 1, \frac{2}{3} \} \} \} \}
 \end{aligned}$$

Number of operators and topologies

Number of topologies VS Mass dim



Number of operators VS Mass dim



Conclusion:

- Baryon number violation is uniquely sensitive to heavy new physics, probing mass dimensions $\gg 6$.
- We found all UV completions for operators up to $d=15$ in terms of scalars, fermions and vectors including right-handed neutrinos.
- Previously, only incomplete results up to dimension 9.
- Useful for model building and in the event of a discovery in Super-Kamiokande, Hyper-Kamiokande, JUNO or DUNE.
- Our method can be easily applied to other operators, for example violating lepton number, only limiting factor is presentation of actual results due to a number of solutions.

Thank You For Your Attention!

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