Scalar scattering amplitudes: zero loci, factorization, and the double copy

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¹Work in collaboration with Yang Li and Diederik Roest; arXiv:2403.12939 [hep-th].K ロ ▶ K 個 ▶ K 할 ▶ K 할 ▶ 이 할 → 9 Q Q →

Introduction

- ▶ It was recently discovered by Arkani-Hamed et al that the tree-level, color-ordered partial amplitudes of $\mathsf{Tr}(\Phi^3)$ vanish on specific hypersurfaces in the space of Mandelstam variables.
- ▶ These zero loci naturally arise as 'causal diamonds' in a two-dimensional kinematic mesh.
- ▶ The zero loci give rise to factorisation channels into lower-point amplitudes
- The zero loci of $Tr(\Phi^3)$ carry over to two more color ordered theories: the chiral non-linear sigma model (NLSM), and scaffolded Yang-Mills.

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 \blacktriangleright The zero loci include and extend the soft-limit behaviour (Adler zero) of the NLSM.

▶ We identify three non-color ordered theories that naturally generalise these results

- ▶ flavor ordered: scaffolded General Relativity (GR).
- ▶ flavor ordered: multi-Dirac-Born Infeld (DBI)
- ▶ single scalar: special Galileon (SG).
- \triangleright The six theories together naturally span the set of scalar field theories that are related via the double copy procedure.
- ▶ I will discuss two of the six theories as representative examples.
- ▶ The papers Cao et al, arXiv:2403.08855 [hep-th], and Bartsch et al. arXiv:2403.10594 [hep-th] contain related work with some overlap. Also see the plenary talk by Jaroslav Trnka at this conference, Wednesday at 11.00 am.

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Color ordered theory: $\text{Tr}(\Phi^3)$

- ▶ Real scalar transforming under the adjoint representation of $SU(N)$.
- ▶ Amplitudes at all-loop and of any multiplicity follow from compact geometric rules on surfaces.

Four point color ordered partial amplitude (canonical ordering):

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$$
 Zero locus: $s_{13} = t = 0$.

- ▶ Two poles available for exchange interactions: *s* and u-channel.
- ▶ $A_4^{\Phi^3} \sim \frac{1}{s} + \frac{1}{t} \sim \frac{u}{s}$ st
- ▶ Partial amplitude is invariant under cyclic permutations.
- \triangleright Other permutations transform this partial amplitude into another one, with corresponding zero locus.

Six point color ordered partial amplitude (canonical ordering):

- ▶ Zero loci:
	- \triangleright "skinny rectangle": $s_{13} = s_{14} = s_{15} = 0$, and cyclic permutations.
	- ▶ "fat square": $s_{14} = s_{24} = s_{15} = s_{25} = 0$, and cyclic permutations.

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- ▶ factorization near "skinny rectangle:" $\mathcal{A}_6^{\Phi^3} \to (\frac{1}{X_1}$ $\frac{1}{X_T} + \frac{1}{X_l}$ $\frac{1}{X_B}$) $A_3^{\Phi^3} A_5^{\Phi^3}$
- ▶ factorization near "fat square:" $\mathcal{A}_6^{\Phi^3} \to (\frac{1}{X_1}$ $\frac{1}{X_T} + \frac{1}{X_l}$ $\frac{1}{X_B}$) $A_4^{\Phi^3} A_4^{\Phi^3}$

Single scalar: Special Galileon

 \blacktriangleright Lagrangian invariant under the non-linear transformation $\delta\phi = s_{\mu\nu}x^{\mu}x^{\nu} + s_{\mu\nu}\partial^{\mu}\phi\partial^{\nu}\phi.$

 \blacktriangleright Maximal soft degree: $\sigma = 3$.

Four point full amplitude:

▶ Zero locus: $s_{13} = t = 0$, and any permutation, thus also $s_{12} = s = 0$ and $s_{14} = u = 0$.

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- ▶ $A_4^{SG} \sim stu$
- ▶ Amplitude is invariant under any permutation.

Six point amplitude:

- ▶ Zero loci:
	- \triangleright "skinny rectangle": $s_{13} = s_{14} = s_{15} = 0$, and any permutation.
	- ▶ "fat square": $s_{14} = s_{24} = s_{15} = s_{25} = 0$, and any permutation.
- ▶ Factorization near "skinny rectangle:" $\mathcal{A}_6^{\mathrm{SG}} \rightarrow \mathsf{s}^3 \times \mathcal{A}_5^{\mathrm{ext\,SG}}$
- \triangleright The extended 5 point amplitude can be generated from the Lagrangian: $\mathcal{L}^{\text{ext}} = \sqrt{-g}(-\frac{1}{2})$ $\frac{1}{2} \partial^{\mu} \Phi \partial_{\mu} \Phi + \Phi^3$) in terms of the special Galileon covariant metric $g_{\mu\nu} = \eta_{\mu\nu} + \partial_{\mu}\partial_{\rho}\varphi \partial^{\rho}\partial_{\nu}\varphi$.

$$
\blacktriangleright \text{ factorization near "fat square:}\n \begin{aligned}\n &\quad \mathcal{A}_6^{\text{SG}} \rightarrow \left(\frac{1}{X_B} + \frac{1}{X_T}\right) A_4^{\text{up,SG}} \times A_4^{\text{down,SG}}\n \end{aligned}
$$

 \triangleright Considering again the "skinny rectangle" loci, the amplitude vanishes when setting any three of the Mandelstam variables $s_{\alpha\beta}$ from the set $\alpha = (1)$, $\beta = (\{2, 3, 4, 5, 6\})$ to zero. Accordingly, the amplitude must be trilinear in the above set. This implies that the amplitude scales with third power in the soft limit $p_1 \rightarrow 0$: $A_6^{\text{SG}} \sim p_1^3$.

Double copy

Figure: Tetrahedron indicating BCJ relations among theories. Figure taken from De Neeling et al, HEP 10 (2022) 066, 2204.11629 [hep-th].

Overview

Figure: The six theories with related zero loci, arranged in terms of their soft degree σ . The underlined theories have no odd vertices and hence the zeros imply their soft behaviours.

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Conclusions

- \triangleright Three double copy theories inherit amplitude zeros from their single copies: scaffolded GR, multi-DBI and the special Galileon.
- \triangleright For the two Goldstone theories, these zero loci underlie their enhanced soft behaviours.
- \triangleright The extended Lagrangians that generate the mixed amplitudes from the near-zero factorization as covariantly coupled colorless Φ^3 , where the minimal coupling is to different metrics for the three cases.
- \blacktriangleright The inherited zero loci of the double copy theories can be explained via the most general version of the KLT formula.

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