# Neutrino Mass and Gravitational Wave

Arnab Dasgupta, Tao Han, Tong Arthur Wu

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y_{\nu} L \tilde{H} \nu_R + \text{h.c.} \rightarrow \frac{y_{\nu} v}{\sqrt{2}} (\nu_L \nu_R + \nu_R \nu_L)
$$

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y_{\nu} L\tilde{H}\nu_R + \text{h.c.} \rightarrow \left(\frac{y_{\nu}v}{\sqrt{2}}\right)\nu_L \nu_R + \nu_R \nu_L
$$
\n
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m_{\nu} < 0.1 \text{eV} \rightarrow y_{\nu} \sim 10^{-13} \text{ Unnaturally small parameter!}
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Inverse seesaw mechanism:  
\n
$$
y_{\nu}L\tilde{H}\nu_R + M\nu_R S + \frac{1}{2}\mu SS + \text{h.c.} \rightarrow \frac{1}{2}(\nu_L \nu_R S) \begin{pmatrix} 0 & \frac{y_{\nu}v}{\sqrt{2}} & 0 \\ \frac{y_{\nu}v}{\sqrt{2}} & 0 & M \\ 0 & M & \mu \end{pmatrix} \begin{pmatrix} \nu_L \\ \nu_R \\ S \end{pmatrix} + \text{h.c.}
$$

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y_{\nu}L\tilde{H}\nu_{R} + M\nu_{R}S + \frac{1}{2}\mu SS + \text{h.c.} \rightarrow \frac{1}{2}(\nu_{L}\nu_{R}S)\begin{pmatrix} 0 & \frac{y_{\nu}v}{\sqrt{2}} & 0\\ \frac{y_{\nu}v}{\sqrt{2}} & 0 & M\\ 0 & M & \mu \end{pmatrix}\begin{pmatrix} \nu_{L}\\ \nu_{R}\\ S \end{pmatrix} + \text{h.c.}
$$
\n
$$
m_{\nu} = \frac{\mu}{2}(\frac{y_{\nu}v}{M})^{2} \longrightarrow \text{Natural suppression!}
$$

#### Model

#### Scalar induced mass term:



$$
\mathcal{L}_{\text{Yukawa}} \supset -\left[y_{\nu} L \tilde{H} \nu_R + (y_{\sigma} \sigma \nu_R S + \bar{y}_{\sigma} \sigma^* \nu_R \bar{S}) + (y_{\phi} \phi S S + \bar{y}_{\phi} \phi^* \bar{S} \bar{S}) + \text{h.c.}\right]
$$
\n
$$
m_{\nu} = \sqrt{2} y_{\phi} v_{\phi} \left(\frac{y_{\nu} v_H}{y_{\sigma} v_{\sigma}}\right)^2
$$













$$
V_{\text{eff}}(\phi,T) = V_{\text{tree}}(\phi) + V_{1-\text{loop}}(\phi) + V_T(\phi,T) + V_{\text{daisy}}(\phi,T)
$$





#### Relevant parameters

nucleation rate: 
$$
\frac{\Gamma(T)}{\Gamma(T)} \sim T^4 \left(\frac{S_3(\phi, T)}{2\pi T}\right)^{3/2} e^{-S_3(\phi, T)/T}
$$

$$
S_3(\phi, T) = \int d^3x \left[\frac{1}{2}(\nabla \phi)^2 + V_{\text{eff}}(\phi, T)\right]
$$

nucleation temperature:

$$
\boxed{T_n}
$$
 when  $\left. \frac{\Gamma(T)}{H(T)^4} \right|_{T=T_n} = 1$  (one bubble per Hubble radius)

vacuum energy density:

$$
\fbox{$\alpha$} \!=\! \frac{1}{\rho_{\rm rad}}\Bigl(-1+T\frac{\rm d}{{\rm d}T}\Bigr)\Delta V_{\rm eff}\biggr|_{T=T_n}
$$

$$
\frac{\beta}{H} = -\frac{T}{\Gamma} \frac{d\Gamma}{dT} \bigg|_{T=T_n}
$$

inverse timescale: bubble-wall velocity:

$$
\boxed{\frac{\mathcal{U}_w}{\mathcal{U}}}\text{=}\begin{cases}\sqrt{\frac{\Delta V_{\text{eff}}}{\alpha\rho_{\text{rad}}}} & \text{for }\frac{\Delta V_{\text{eff}}}{\alpha\rho_{\text{rad}}}<\mathcal{U}_J \\ 1 & \text{for }\frac{\Delta V_{\text{eff}}}{\alpha\rho_{\text{rad}}}> \mathcal{U}_J\end{cases}}
$$

#### Gravitational wave spectrum

$$
h^{2}\Omega_{\rm GW}(f) = h^{2}\Omega_{b}(f) + h^{2}\Omega_{s}(f) + h^{2}\Omega_{t}(f)
$$
\ncollisions of bubble walls\n
$$
\begin{array}{ccc}\n\text{collisions of bubble walls} & \text{sound waves} & \text{turbulence} \\
\parallel & & \parallel & & \parallel \\
\text{magnitude }1.67 \times 10^{-5} \left(\frac{v_{w}}{\beta/H}\right)^{2} \left(\frac{100}{g(T_{n})}\right)^{1/3} \left(\frac{\kappa_{b}\alpha}{1+\alpha}\right)^{2} \left(\frac{0.11v_{w}}{0.42 + v_{w}^{2}}\right) & 2.65 \times 10^{-6} \left(\frac{v_{w}}{\beta/H}\right) \left(\frac{100}{g(T_{n})}\right)^{1/3} \left(\frac{\kappa_{s}\alpha}{1+\alpha}\right)^{2} & 3.35 \times 10^{-4} \left(\frac{v_{w}}{\beta/H}\right) \left(\frac{100}{g(T_{n})}\right)^{1/3} \left(\frac{\kappa_{t}\alpha}{1+\alpha}\right)^{3/2} \\
\times & & \times & \times & \times \\
\text{function} & & & \left(\frac{f}{f_{b}}\right)^{2.8} \left(\frac{3.8}{1+2.8(f/f_{b})^{3.8}}\right) & & & \left(\frac{f}{f_{s}}\right)^{3} \left(\frac{7}{4+3(f/f_{s})^{2}}\right)^{7/2} & \left(\frac{f}{f_{t}}\right)^{3} \left(\frac{1}{1+(f/f_{t})}\right)^{11/3} \left(\frac{1}{1+8\pi f/h_{*}}\right)\n\end{array}
$$

 $\Box$  Obtained from simulation **Parametrized by**  $T_n$ **,**  $\alpha$ **,**  $\beta$ **,**  $v_w$ 

$$
\mathcal{L}_{\text{scalar}} \supset \mu_{\sigma}^{2} |\sigma|^{2} - \lambda_{\sigma} |\sigma|^{4} + \mu_{\phi}^{2} |\phi|^{2} - \lambda_{\phi} |\phi|^{4} - \lambda_{\sigma\phi} |\sigma|^{2} |\phi|^{2} - (\mu \phi^{*} \sigma^{2} + \text{h.c.})
$$
  

$$
\mathcal{L}_{\text{Yukawa}} \supset -\left[ y_{\nu} L \tilde{H} \nu_{R} + \left( y_{\sigma} \sigma \nu_{R} S + \bar{y}_{\sigma} \sigma^{*} \nu_{R} \bar{S} \right) + \left( y_{\phi} \phi S S + \bar{y}_{\phi} \phi^{*} \bar{S} \bar{S} \right) + \text{h.c.} \right]
$$

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$$

Phase transition:



$$
\mathcal{L}_{\text{scalar}} \supset \mu_{\sigma}^2 |\sigma|^2 - \lambda_{\sigma} |\sigma|^4 + \mu_{\phi}^2 |\phi|^2 - \lambda_{\phi} |\phi|^4 - \lambda_{\sigma\phi} |\sigma|^2 |\phi|^2 - (\mu \phi^* \sigma^2 + \text{h.c.})
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$$



#### Phase transition: Benchmark parameters:



$$
\mathcal{L}_{\text{scalar}} \supset \mu_{\sigma}^{2} |\sigma|^{2} - \lambda_{\sigma} |\sigma|^{4} + \mu_{\phi}^{2} |\phi|^{2} - \lambda_{\phi} |\phi|^{4} - \lambda_{\sigma\phi} |\sigma|^{2} |\phi|^{2} - (\mu \phi^{*} \sigma^{2} + \text{h.c.})
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\mathcal{L}_{\text{Yukawa}} \supset -\left[ y_{\nu} L \tilde{H} \nu_{R} + \left( y_{\sigma} \sigma \nu_{R} S + \bar{y}_{\sigma} \sigma^{*} \nu_{R} \bar{S} \right) + \left( y_{\phi} \phi S S + \bar{y}_{\phi} \phi^{*} \bar{S} \bar{S} \right) + \text{h.c.} \right]
$$



Benchmark parameters:



#### Collider signal from seesaw model





#### **Complementarity**



#### Conclusion

- Seesaw model + scalar fields ⇒ strong gravitational waves
- Gravitational wave observation is highly complementary to collider searches

Thank you!