

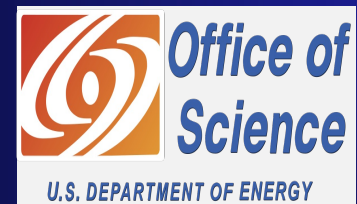
# Alternative Method for Measuring the Radius of the Deuteron

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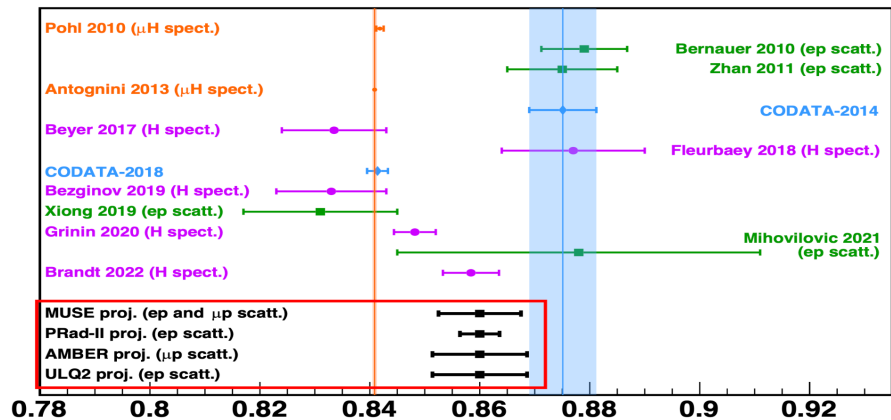


Joint 20<sup>th</sup>-IWHSS-5<sup>th</sup> CPHI  
September 30-October 4, 2024, Yerevan, Armenia

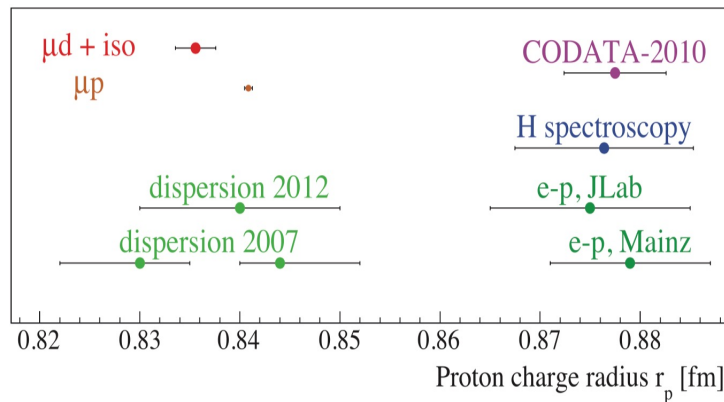


# *Spectroscopic vs e-Scattering Measurement of Radii*

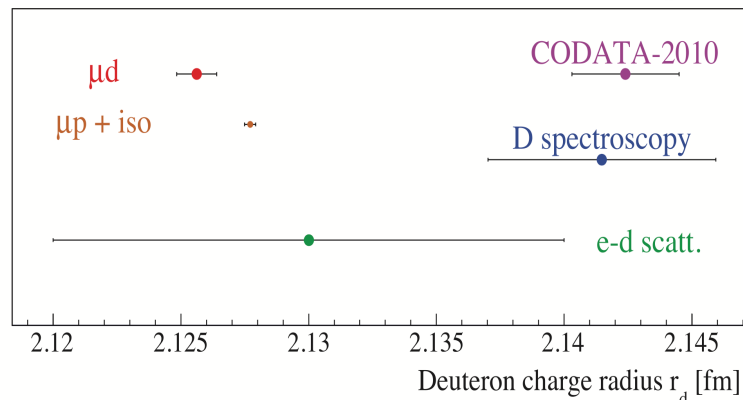
<i>Radii</i>	<b>e-Spectroscopic</b>	<b>mu-Spectroscopic</b>	<b>e-Scattering</b>
<i>Proton</i>	0.8775(51)fm	0.84087(39)fm	0.879(8)fm
<i>Deuteron</i>	2.1415(45)fm	2.12562(13)fm	2.13(12)fm
<i><sup>4</sup>He</i>	3 body problem	1.67824(13)fm	1.681(4)fm



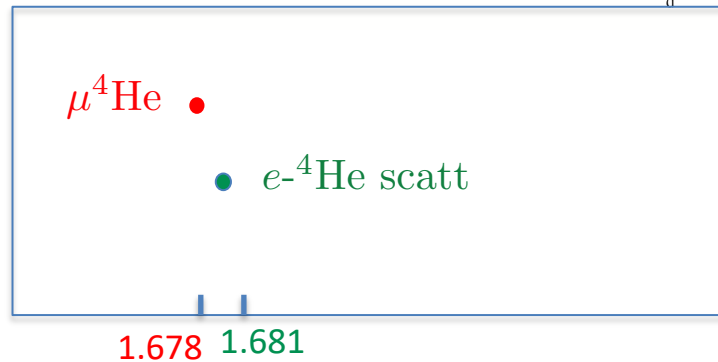
A.Gasparian, yesterday



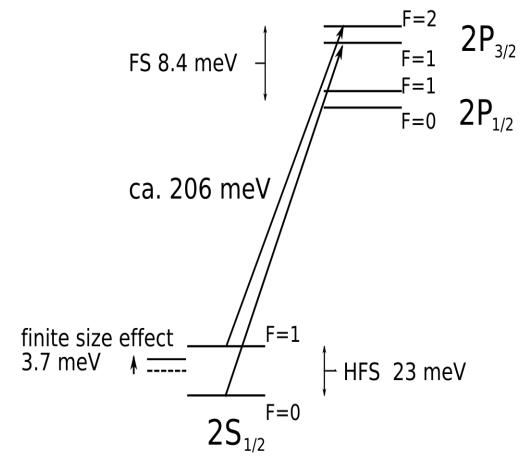
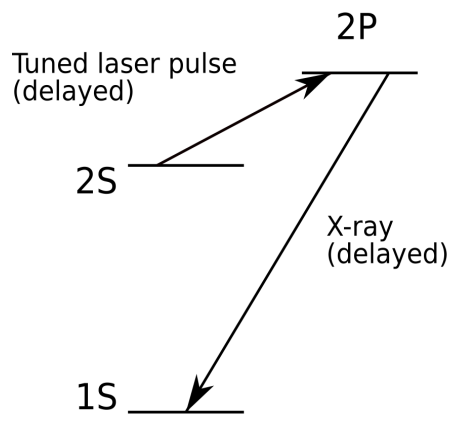
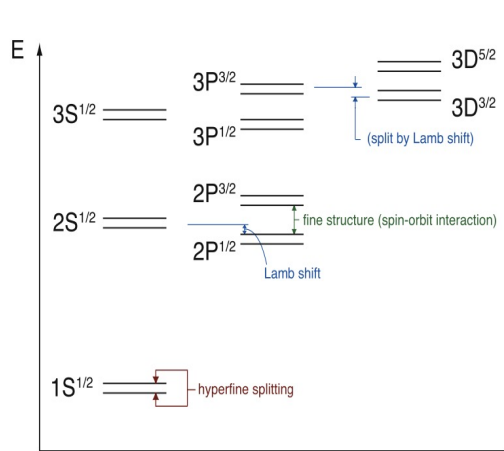
R.Pohl, Nature, 2010



R.Pohl, Science, 2016

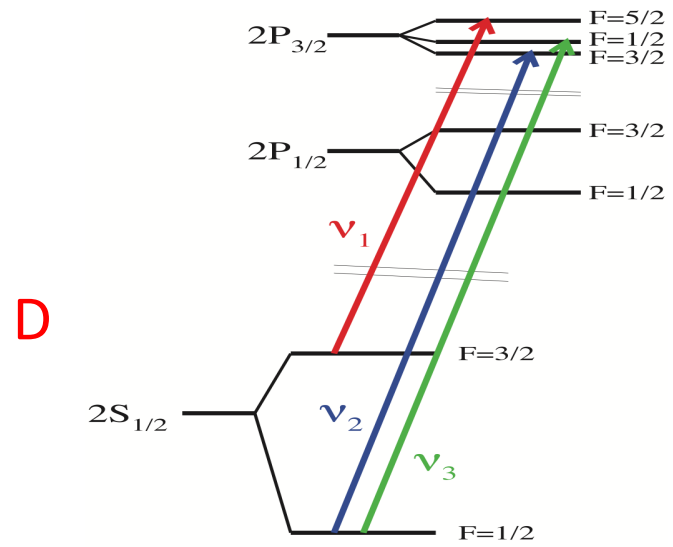


J.Krauth, Nature, 2021

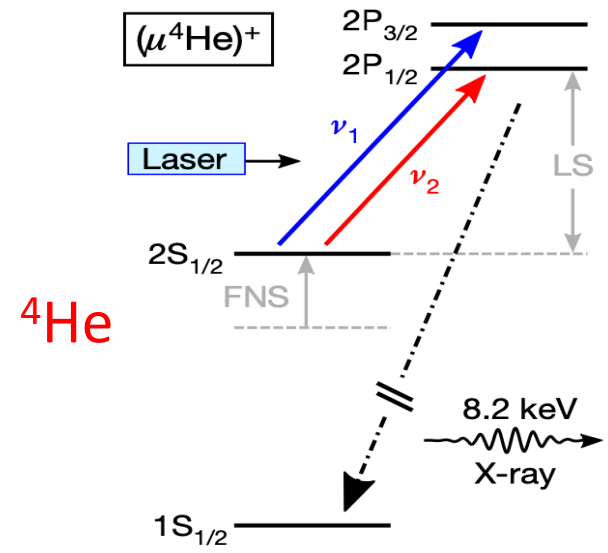


H

C. Carlson, PNP, 2015

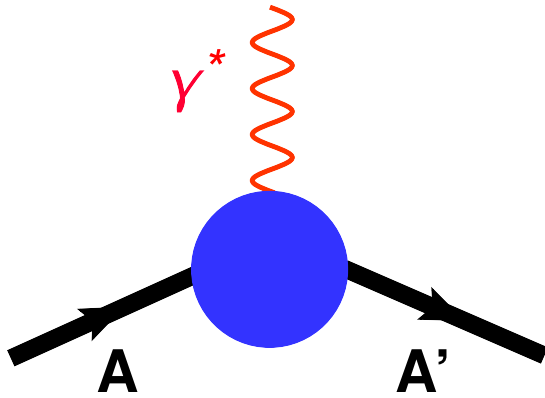


D



<sup>4</sup>He

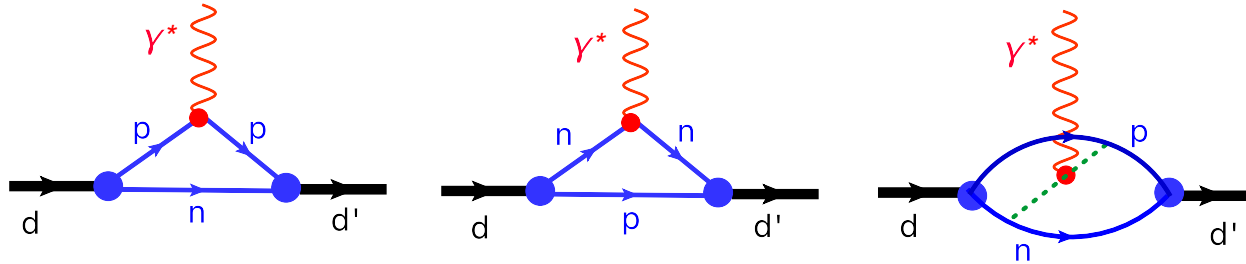
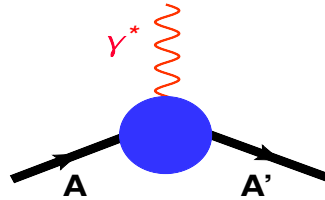
$$e + A \rightarrow e' + A'$$



$$R_{ch}^2 = -6 \frac{dG_{ch}(Q^2)}{dQ^2} \Big|_{Q^2 \rightarrow 0}$$

$$e + A \rightarrow e' + A'$$

New Method



$$T_{eA}(Q^2) = ZT_{ep}(Q^2)S_A(\frac{q}{2}) + (A - Z)T_{en}(Q^2)S_A(\frac{q}{2}) + T_{MEC} + T_{Rel}$$

$$Q^2 \rightarrow 0$$

$$r_{A,ch}^2 = Zr_{A,m}^2 + Zr_{p,ch}^2 + (A - Z)r_{n,ch}^2 + r_{MEC}^2 + r_{Rel}^2$$

$$e + d \rightarrow e' + d'$$

$$r_{d,ch}^2 = r_{d,m}^2 + r_{p,ch}^2 + r_{n,ch}^2 + r_{MEC}^2 + \frac{3}{4} \frac{1}{m_p^2}$$

Klarsfeld, et al, Nuclear Physics A456, 1986

This relation was used to evaluate **the matter radius**

Instead we try to evaluate deuteron charge radius using

$$\begin{aligned} r_{d,m} &= 1.973 fm && \text{EFT} \\ r_{p,ch} &= 0.879 fm && \text{ep-scatt} \\ r_{n,ch}^2 &= -0.106 fm^2 \\ r_{MEC}^2 &= 0.0034 fm^2 \end{aligned}$$

$$r_{d,ch} = 2.135 fm \quad \mathbf{2.13(12)fm} \quad ed - \text{scattering}$$

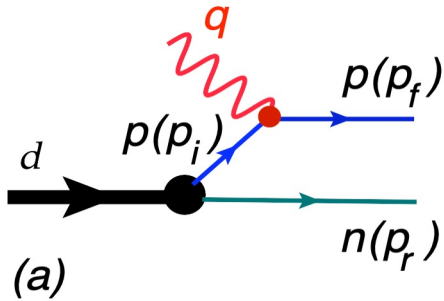
Alternative Method

# Empirical Evaluation of Deuteron Matter Radius

Consider Deuteron Electrodisintegration process



Within Plane Wave Impulse Approximation



$$\sigma_{ed} = K \sigma_{ep}(Q^2)(U^2(k) + W^2(k))$$

$$\vec{k} = \vec{p}_{miss}$$

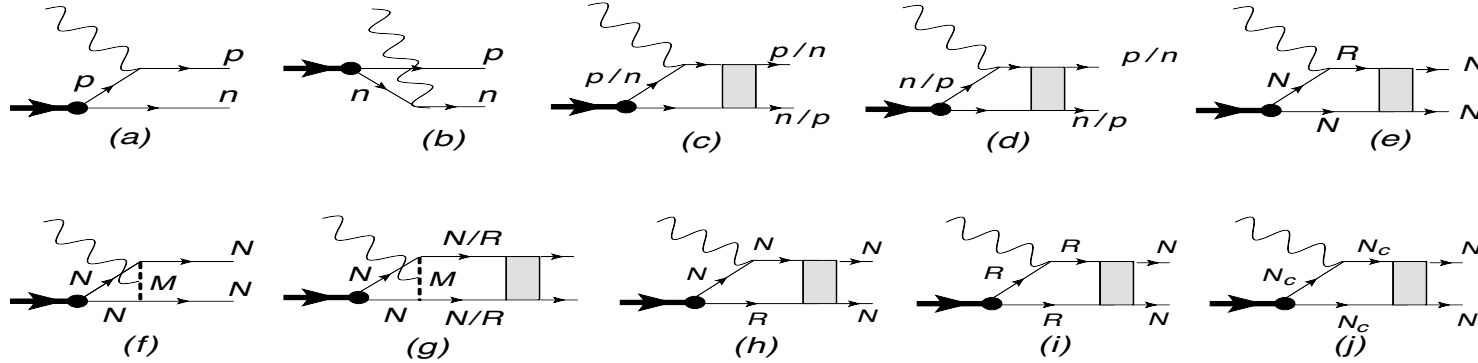
$$\frac{u(r)}{r} = \frac{1}{(2\pi)^{\frac{3}{2}}} \int U(k) j_0(kr) d^3 k$$

$$\frac{w(r)}{r} = \frac{1}{(2\pi)^{\frac{3}{2}}} \int W(k) j_2(kr) d^3 k$$

$$r_{d,m}^2 = \frac{1}{4} \int r^2 (u(r)^2 + w(r)^2) dr$$



# Reality is more complicated



GEA

M.Sargsian, PRC,2010

$$\langle s_f, s_r | A^\mu | s_d \rangle = \langle s_f, s_r | A_0^\mu | s_d \rangle + \langle s_f, s_r | A_1^\mu | s_d \rangle + \langle s_f, s_r | A_{1,chev}^\mu | s_d \rangle + \langle s_f, s_r | A_{1,R}^\mu | s_d \rangle, \quad (2)$$

$$\langle s_f, s_r | A_0^\mu | s_d \rangle = \sqrt{2} \sqrt{(2\pi)^3 2E_r} \times \sum_{s_i} J_N^\mu(s_f, p_f; s_i, p_i) \Psi_d^{s_d}(s_i, p_i, s_r, p_r),$$

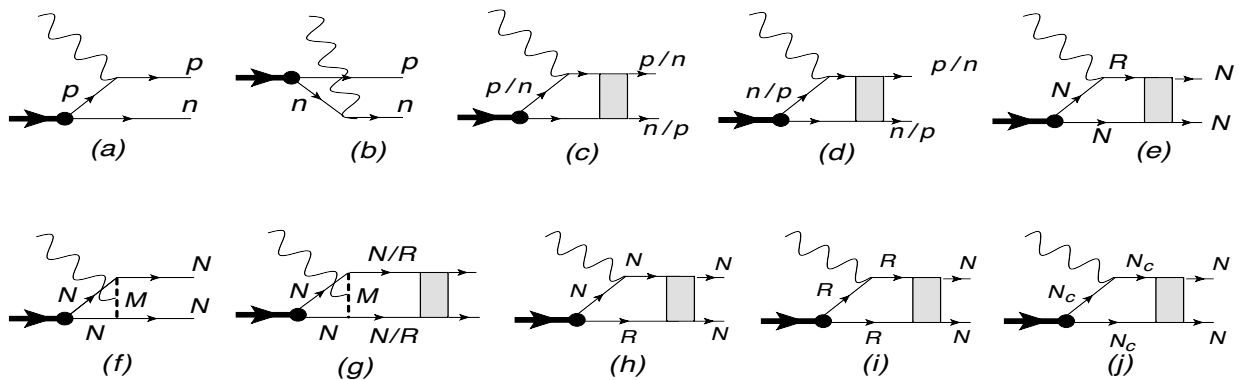
$$\begin{aligned} \langle s_f, s_r | A_1^\mu | s_d \rangle &= - \int \frac{d^4 p'_r}{i(2\pi)^4} \\ &\times \frac{\bar{u}(p_f, s_f) \bar{u}(p_r, s_r) F_{NN}[\not{p}'_r + m][\not{p}_d - \not{p}'_r + \not{q} + m]}{(p_d - p'_r + q)^2 - m^2 + i\epsilon} \\ &\times \frac{\Gamma_{\gamma^* N}[\not{p}_d - \not{p}'_r + m] \Gamma_{DNN} \chi^{s_d}}{((p_d - p'_r)^2 - m^2 + i\epsilon)(p_r'^2 - m^2 + i\epsilon)}, \end{aligned} \quad (11)$$

# Non PWIA Contributions

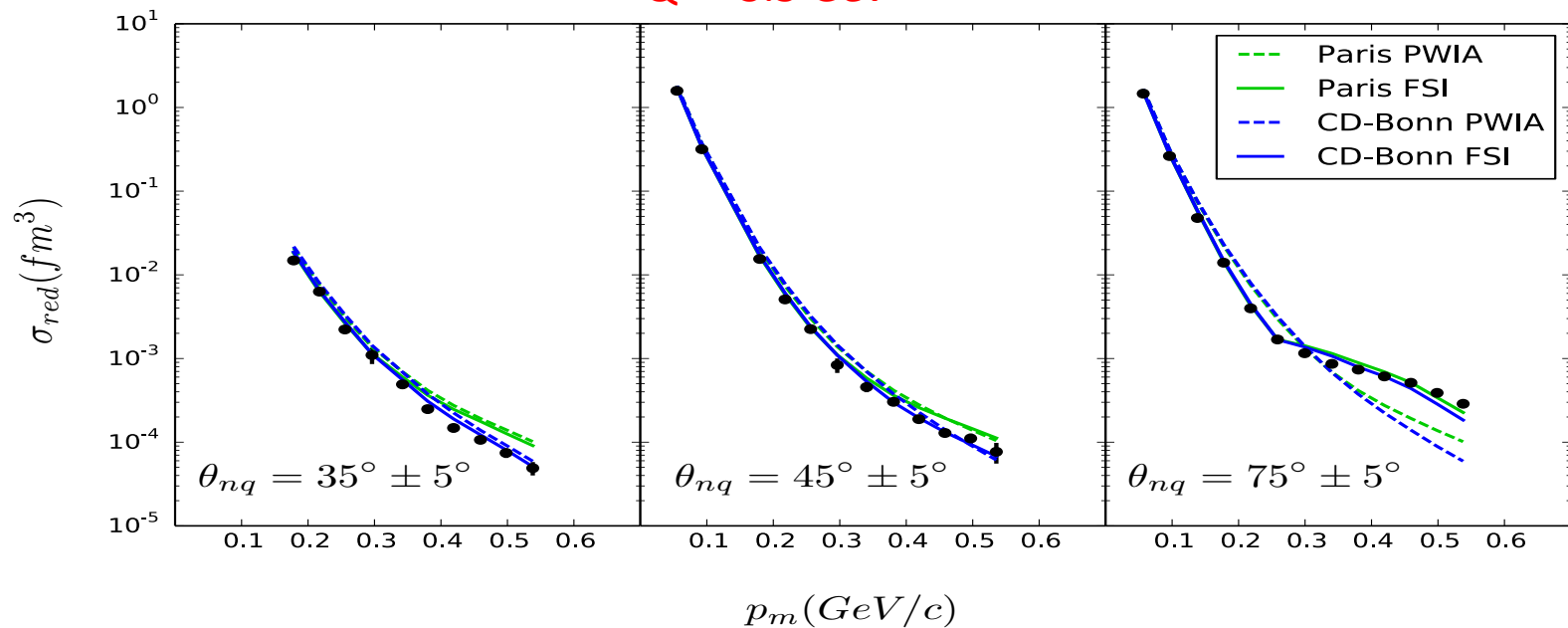
$$\begin{aligned}
 \langle s_f, s_r | A_1^\mu | s_d \rangle &= - \int \frac{d^4 p'_r}{i(2\pi)^4} \\
 &\times \frac{\bar{u}(p_f, s_f) \bar{u}(p_r, s_r) F_{NN}[\not{p}'_r + m][\not{p}_d - \not{p}'_r + \not{q} + m]}{(p_d - p'_r + q)^2 - m^2 + i\epsilon} \\
 &\times \frac{\Gamma_{\gamma^* N}[\not{p}_d - \not{p}'_r + m] \Gamma_{DNN} \chi^{s_d}}{((p_d - p'_r)^2 - m^2 + i\epsilon)(p'^2_r - m^2 + i\epsilon)}, \quad (11)
 \end{aligned}$$

$$\begin{aligned}
 \langle s_f, s_r | A_1^\mu | s_d \rangle &= \frac{i\sqrt{2}(2\pi)^{\frac{3}{2}}}{4} \sum_{s'_f, s'_r, s_i} \int \frac{d^2 p'_r}{(2\pi)^2} \frac{\sqrt{2\tilde{E}'_r} \sqrt{s(s-4m^2)}}{2\tilde{E}'_r |q|} \langle p_f, s_f; p_r, s_r | f^{NN, on}(t, s) | \tilde{p}'_r, s'_r; \tilde{p}'_f, s'_f \rangle \\
 &\times J_N^\mu(s'_f, p'_f; s_i, \tilde{p}'_i) \cdot \Psi_d^{s_d}(s_i, \tilde{p}'_i, s'_r, \tilde{p}'_r) \\
 &- \frac{\sqrt{2}(2\pi)^{\frac{3}{2}}}{2} \sum_{s'_f, s'_r, s_i} \mathcal{P} \int \frac{dp'_{r,z}}{2\pi} \int \frac{d^2 p'_r}{(2\pi)^2} \frac{\sqrt{2E'_r} \sqrt{s(s-4m^2)}}{2E'_r |\mathbf{q}|} \frac{\langle p_f, s_f; p_r, s_r | f^{NN, off}(t, s) | p'_r, s'_r; p'_f, s'_f \rangle}{p'_{r,z} - \tilde{p}'_{r,z}} \\
 &\times J_N^\mu(s'_f, p'_f; s_i, p'_i) \cdot \Psi_d^{s_d}(s_i, p'_i, s'_r, p'_r), \quad (13)
 \end{aligned}$$

# Reality is more complicated



$Q^2 = 3.5 \text{ GeV}^2$



W.Boeglin, et al, PRL 2011

Jlab Experiment

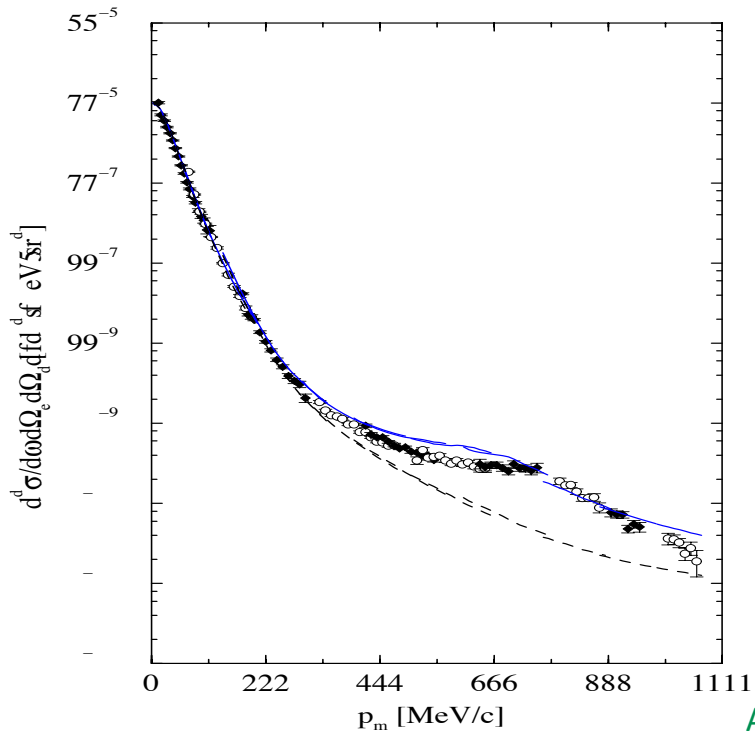
# Anomalous Distribution at Very Small Missing Momenta



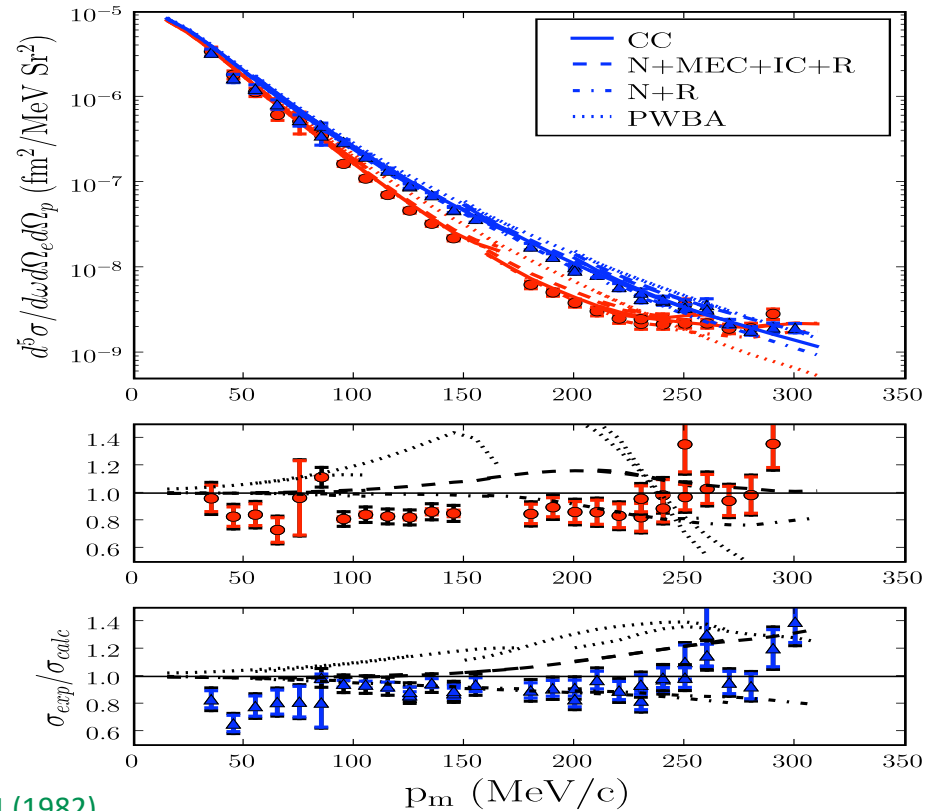
W.Boeglin et al Phys.Rev.C 78 (2008) 054001

Mainz experiment

Mainz Measurement at  $Q^2 = 0.33 \text{ GeV}^2$



Arenhoevel  
Nucl.Phys.A 384 (1982)

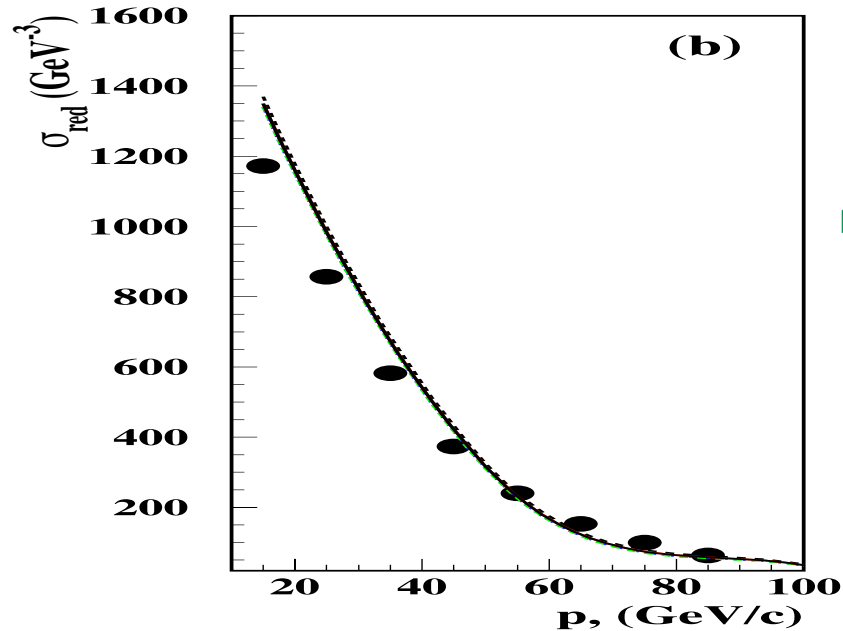
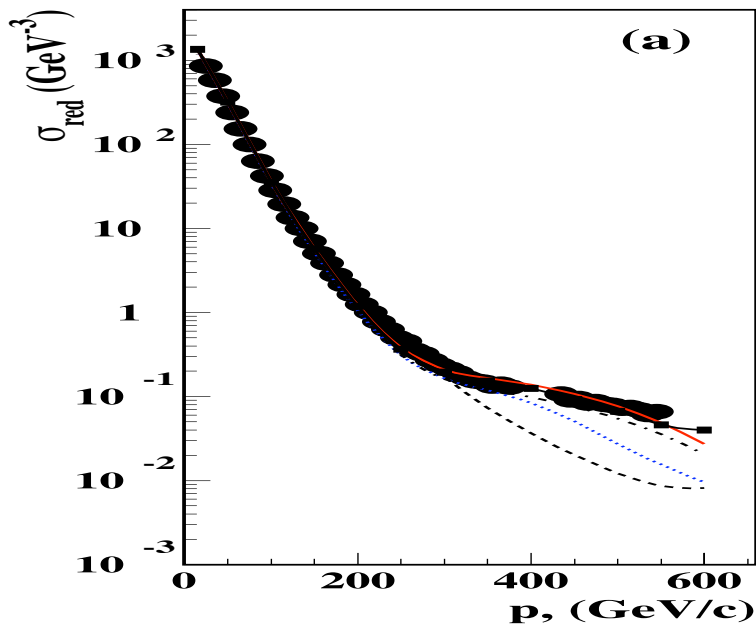


# Anomalous Distribution at Very Small Missing Momenta



JLab Measurement at  $Q^2 = 0.66 \text{ GeV}^2$

M.Sargsian, *Phys.Rev.C* 82 (2010) 014612



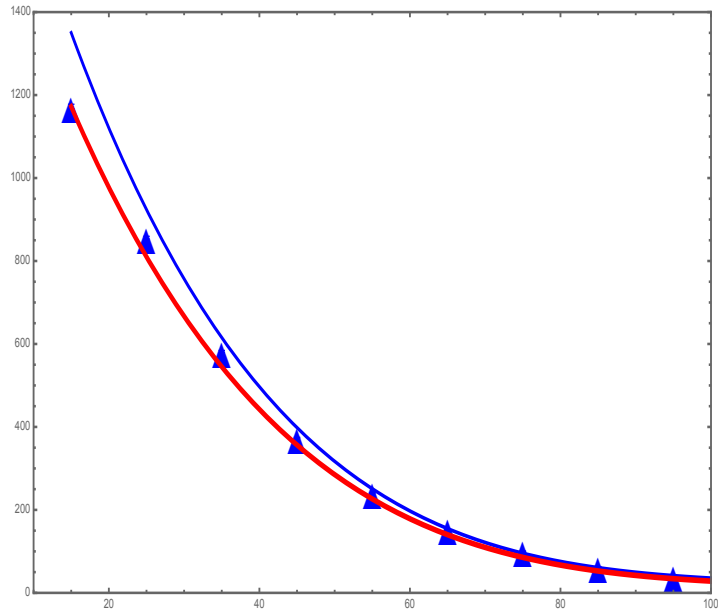
P.Ulmer, et al, PRL 2002

JLab experiment

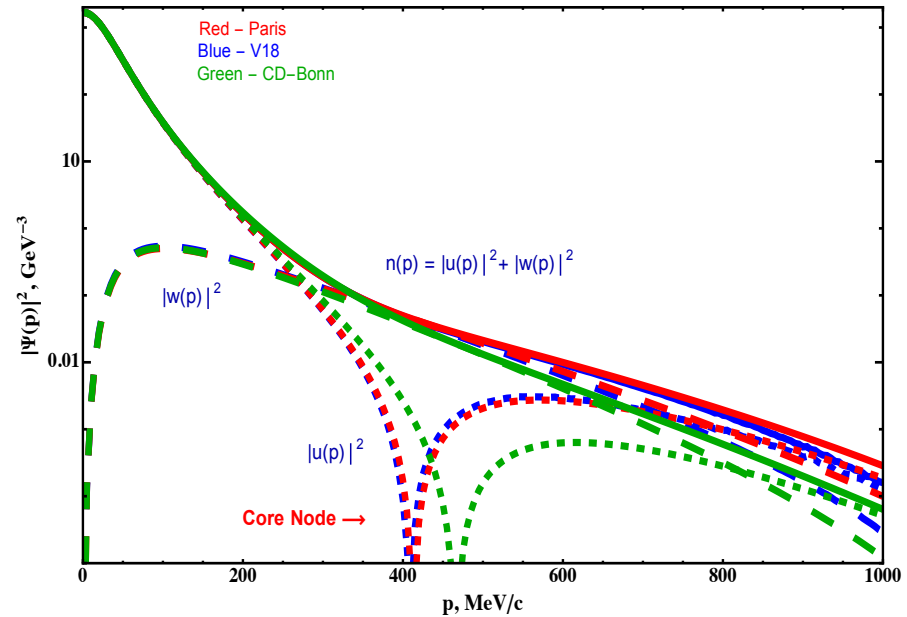
# Fitting Deuteron Wave Function to the data



JLab Measurement at  $Q^2 = 0.66 \text{ GeV}^2$



$$r_{d,m} = 1.957 \text{ Fm}$$



Only S-state needs to be modified

$$r_{d,m} = 1.957 fm$$

$$r_{p,ch} = 0.879 fm$$

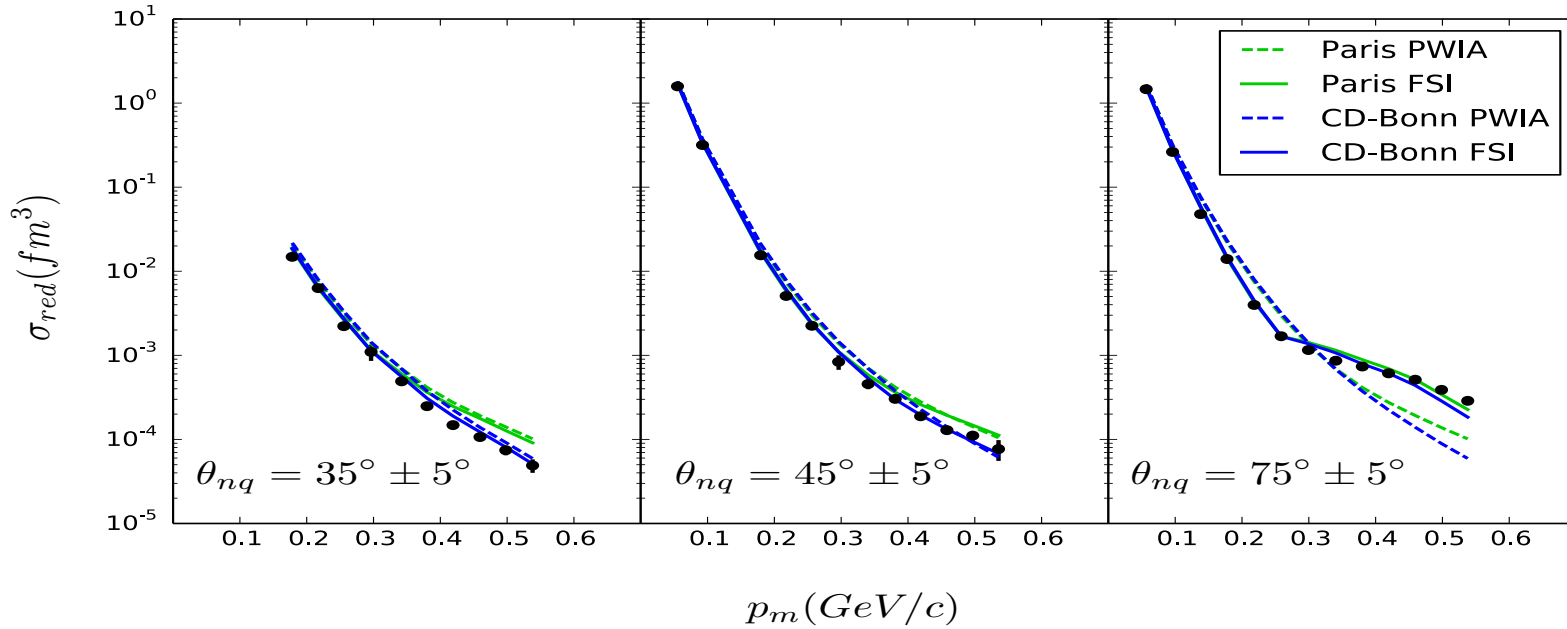
$$r_{n,ch}^2 = -0.106 fm^2$$

$$r_{MEC}^2 = 0.0034 fm^2$$

$$r_{d,ch} = 2.121 fm \quad \mathbf{2.12562(13)fm} \quad \mu d - \text{spectrometry}$$

Alternative Method

$Q^2 = 3.5 \text{ GeV}^2$



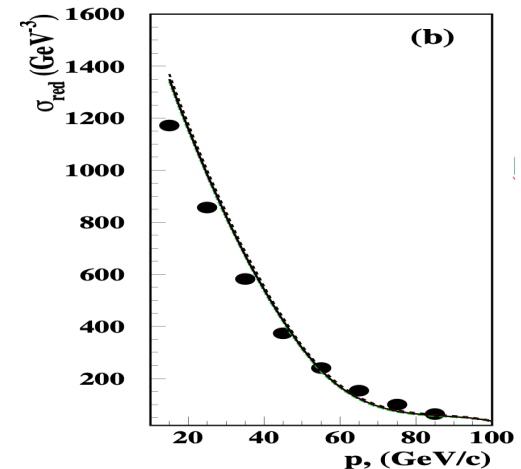
W.Boeglin, et al, PRL 2011

Jlab Experiment



# Conclusion/Outlook

- “Alternative” method is suggested to evaluate charge radius of the deuteron
- ed – scattering results can be reconciled with ep scattering better than ed spectroscopy
- the same may be true for  $^4\text{He}$  nucleus
- Initial analysis indicates that the deuteron matter radius may be smaller
- it may indicate that BSM effects may be related to electron spectroscopy rather than to muon spectroscopy
- careful analysis of deuteron break-up reaction at small missing momenta may be needed (*this will require new experiments with short beam time*)
- as well as reliable measurement of neutron charge radius square and MEC



POETIC XI  
Miami, Feb 24-28



*The road to EIC, as seen from South Florida...*

## Physics Opportunities at an Electron-Ion Collider XI

Feb 24 – 28, 2025

Miami, Florida International University, Modesto Maidique Campus

US/Eastern timezone

Overview

Scientific Program

Call for Abstracts

Registration

Conference Fee Payment

**Ignore any e-mails not coming from the organizers or indico related to this meeting! They are phis attempts with promises of booked hotels etc.**

POETIC XI, the eleventh international conference on the 'Physics Opportunities at an Electron-Ion Collider', follows POETIC X which took place at the ICTP-SAIFR, São Paulo, Brazil in May 2023.

The primary goal of the conference is to continue the advancement of the field of the future Electron-Ion Collider (EIC) physics which was granted Critical Decision 3A (CD-3A) by the U.S. Department of Energy.

