ALESSANDRO BACCHETTA, PAVIA U. AND INFN

TMD OVERVIEW





slide from 2018 CPHI@Yerevan It's the dawn of TMD global fits era



slide from 2018 CPHI@Yerevan It's the dawn of TMD global fits era

... but there's still a lot of climbing to be done







 TMD multiplicities for pions and kaons, off protons and deuterons, from **COMPASS** and JLab



- COMPASS and JLab
- Drell-Yan and Z measurements from CERN, RHIC, FermiLab (COMPASS with pions)

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READY TO USE EIC DATA





How "wide" is the distribution?











TMDs in **black** survive integration over transverse momentum TMDs in **red** are time-reversal odd

<u>Mulders-Tangerman, NPB 461 (96)</u> <u>Boer-Mulders, PRD 57 (98)</u>





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- Some hints about all others



AVAILABLE EXTRACTIONS (NEWEST ONLY)

Unpol. TMD	MAP 22 arXiv:2206.0759	
Helicity		
Transversity	arXiv:1505.05589, arXiv:	
Sivers	MAP20 arXiv:2004.1427 arXiv:2304.14328	
Boer-Mulders	<u>arXiv:2004.02117</u> ,	
Worm-gear g1T	<u>arXiv:2110.10253, arXiv:</u>	
Worm-gear h1L		
Pretzelosity	arXiv:1411.0580	

98, <u>ART23 2305.07473</u>

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NEW THIS YEAR

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Not mentioned: pion TMDs, TMD fragmentation functions, nuclear TMDs

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- Some knowledge of g_T x-dependence
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		U	circ	linear
pol.	U	f_1^g		$h_1^{\perp g}$
leon	L		g^g_{1L}	$h_{1L}^{\perp g}$
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TMD TABLE: GLUONS, LEADING TWIST

	gluon pol.							
		U	circ	linear				
eon pol.	U	f_1^g		$h_1^{\perp g}$				
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See talk by Cristian Pisano







TMDS IN DRELL-YAN PROCESSES



1	1	

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The analysis is usually done in Fourier-transformed space



1	1	

TMDS IN DRELL-YAN PROCESSES



The analysis is usually done in Fourier-transformed space TMDs formally depend on two scales, but we set them equal.

1	1	

TMDS IN SEMI-INCLUSIVE DIS (SIDIS)





 $\hat{f}_1^a(x, |\boldsymbol{b}_T|; \mu, \zeta) = \int d^2 \boldsymbol{k}_\perp e^{i\boldsymbol{b}_T \cdot \boldsymbol{k}_\perp} f_1^a(x, \boldsymbol{k}_\perp^2; \mu, \zeta)$



$$\hat{f}_1^a(x, |\boldsymbol{b}_T|; \boldsymbol{\mu}, \boldsymbol{\zeta}) = \int d^2 \boldsymbol{k}_\perp \, e^{i\boldsymbol{b}_T \cdot \boldsymbol{k}_\perp} \, f_1^a(x, \boldsymbol{k}_\perp^2;$$

 $\hat{f}_1^a(x, b_T^2; \mu_f, \zeta_f) = [C \otimes f_1](x, \mu_{b_*}) \ e^{\int_{\mu_{b_*}}^{\mu_f} \frac{d\mu}{\mu}} \left(\gamma_{a_*}\right) = [C \otimes f_1](x, \mu_{b_*}) \ e^{\int_{\mu_{b_*}}^{\mu_f} \frac{d\mu}{\mu}} \left(\gamma_{a_*}\right) = [C \otimes f_1](x, \mu_{b_*}) \ e^{\int_{\mu_{b_*}}^{\mu_f} \frac{d\mu}{\mu}} \left(\gamma_{a_*}\right) = [C \otimes f_1](x, \mu_{b_*}) \ e^{\int_{\mu_{b_*}}^{\mu_f} \frac{d\mu}{\mu}} \left(\gamma_{a_*}\right) = [C \otimes f_1](x, \mu_{b_*}) \ e^{\int_{\mu_{b_*}}^{\mu_f} \frac{d\mu}{\mu}} \left(\gamma_{a_*}\right) = [C \otimes f_1](x, \mu_{b_*}) \ e^{\int_{\mu_{b_*}}^{\mu_f} \frac{d\mu}{\mu}} \left(\gamma_{a_*}\right) = [C \otimes f_1](x, \mu_{b_*}) \ e^{\int_{\mu_{b_*}}^{\mu_f} \frac{d\mu}{\mu}} \left(\gamma_{a_*}\right) = [C \otimes f_1](x, \mu_{b_*}) \ e^{\int_{\mu_{b_*}}^{\mu_f} \frac{d\mu}{\mu}} \left(\gamma_{a_*}\right) = [C \otimes f_1](x, \mu_{b_*}) \ e^{\int_{\mu_{b_*}}^{\mu_f} \frac{d\mu}{\mu}} \left(\gamma_{a_*}\right) = [C \otimes f_1](x, \mu_{b_*}) \ e^{\int_{\mu_{b_*}}^{\mu_f} \frac{d\mu}{\mu}} \left(\gamma_{a_*}\right) = [C \otimes f_1](x, \mu_{b_*}) \ e^{\int_{\mu_{b_*}}^{\mu_f} \frac{d\mu}{\mu}} \left(\gamma_{a_*}\right) = [C \otimes f_1](x, \mu_{b_*}) \ e^{\int_{\mu_{b_*}}^{\mu_f} \frac{d\mu}{\mu}} \left(\gamma_{a_*}\right) = [C \otimes f_1](x, \mu_{b_*}) \ e^{\int_{\mu_{b_*}}^{\mu_f} \frac{d\mu}{\mu}} \left(\gamma_{a_*}\right) = [C \otimes f_1](x, \mu_{b_*}) \ e^{\int_{\mu_{b_*}}^{\mu_f} \frac{d\mu}{\mu}} \left(\gamma_{a_*}\right) = [C \otimes f_1](x, \mu_{b_*}) \ e^{\int_{\mu_{b_*}}^{\mu_f} \frac{d\mu}{\mu}} \left(\gamma_{a_*}\right) = [C \otimes f_1](x, \mu_{b_*}) \ e^{\int_{\mu_{b_*}}^{\mu_f} \frac{d\mu}{\mu}} \left(\gamma_{a_*}\right) = [C \otimes f_1](x, \mu_{b_*}) \ e^{\int_{\mu_{b_*}}^{\mu_f} \frac{d\mu}{\mu}} \left(\gamma_{a_*}\right) = [C \otimes f_1](x, \mu_{b_*}) \ e^{\int_{\mu_{b_*}}^{\mu_f} \frac{d\mu}{\mu}} \left(\gamma_{a_*}\right) = [C \otimes f_1](x, \mu_{b_*}) \ e^{\int_{\mu_{b_*}}^{\mu_f} \frac{d\mu}{\mu}} \left(\gamma_{a_*}\right) = [C \otimes f_1](x, \mu_{b_*}) \ e^{\int_{\mu_{b_*}}^{\mu_f} \frac{d\mu}{\mu}} \left(\gamma_{b_*}\right) = [C \otimes f_1](x, \mu_{b_*}) \ e^{\int_{\mu_{b_*}}^{\mu_f} \frac{d\mu}{\mu}} \left(\gamma_{b_*}\right) = [C \otimes f_1](x, \mu_{b_*}) \ e^{\int_{\mu_{b_*}}^{\mu_f} \frac{d\mu}{\mu}} \left(\gamma_{b_*}\right) = [C \otimes f_1](x, \mu_{b_*}) \ e^{\int_{\mu_{b_*}}^{\mu_f} \frac{d\mu}{\mu}} \left(\gamma_{b_*}\right) = [C \otimes f_1](x, \mu_{b_*}) \ e^{\int_{\mu_{b_*}}^{\mu_f} \frac{d\mu}{\mu}} \left(\gamma_{b_*}\right) = [C \otimes f_1](x, \mu_{b_*}) \ e^{\int_{\mu_{b_*}}^{\mu_f} \frac{d\mu}{\mu}} \left(\gamma_{b_*}\right) = [C \otimes f_1](x, \mu_{b_*}) \ e^{\int_{\mu_{b_*}}^{\mu_f} \frac{d\mu}{\mu}} \left(\gamma_{b_*}\right) = [C \otimes f_1](x, \mu_{b_*}) \ e^{\int_{\mu_{b_*}}^{\mu_f} \frac{d\mu}{\mu}} \left(\gamma_{b_*}\right) = [C \otimes f_1](x, \mu_{b_*}) \ e^{\int_{\mu_{b_*}}^{\mu_f} \frac{d\mu}{\mu}} \left(\gamma_{b_*}\right) = [C \otimes f_1](x, \mu_{b_*}) \ e^{\int_{\mu_{b_*}}^{\mu_f} \frac{d\mu}{\mu}} \left(\gamma_{b_*}\right) = [C \otimes f_1](x,$

$$\mu_b = \frac{2e^{-\gamma_E}}{b_T}$$

$$\mu,\zeta)$$

$$\gamma_F - \gamma_K \ln \frac{\sqrt{\zeta_f}}{\mu} \left(\frac{\sqrt{\zeta_f}}{\mu_{b_*}} \right)^{K_{\text{resum}} + g_K}$$



$$\hat{f}_1^a(x, |\boldsymbol{b}_T|; \boldsymbol{\mu}, \boldsymbol{\zeta}) = \int d^2 \boldsymbol{k}_\perp \, e^{i\boldsymbol{b}_T \cdot \boldsymbol{k}_\perp} \, f_1^a(x, \boldsymbol{k}_\perp^2;$$

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collinear PDF

 $\mu_b = \frac{2e^{-\gamma_E}}{b_T}$

matching coefficients (perturbative)





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UNPOLARIZED TMD GLOBAL FITS

	Accuracy	SIDIS HERMES	SIDIS COMPASS	DY fixed target	DY collider	N of points	χ²/N _{point}
Pavia 2017 <u>arXiv:1703.10157</u>	NLL			~		8059	1.55
SV 2019 <u>arXiv:1912.06532</u>	N ³ LL-			~		1039	1.06
MAP22 <u>arXiv:2206.07598</u>	N ³ LL-			~		2031	1.06
ART23 <u>arXiv:2305.07473</u>	N4LL	×	×	•		627	0.96
MAP24 <u>arXiv:2405.13833</u>	N ³ LL					2031	1.08







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See talks by Filippo Delcarro and Valentin Moos





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See talks by Filippo Delcarro and Valentin Moos see also Parton Branching approach, talk by Louis Moureaux











$\frac{|k_{\perp}| [\text{GeV}]}{|k_{\perp}| [\text{GeV}]} = \frac{1.00 \text{ m/s}}{|k_{\perp}| [\text{GeV}]} = \frac{1.00 \text{ m/s}}{|k_{\perp}| [\text{GeV}]}$





MAP Collaboration, arXiv:2405.13833



0.20 0.50 0.10 1.00 1.20 1.000.00 0.40 $|k_{\perp}|~[{ m GeV}]$ **FLAVOR-DEPENDENT UNPOLARIZED TMDS**





MAP Collaboration, arXiv:2405.13833

See talk by Filippo Delcarro







COMPASS multiplicities (one of many bins)



COMPASS multiplicities (one of many bins)

Scimemi, Vladimirov, arXiv:1912.06532

Scimemi, Vladimirov, arXiv:1912.06532

Also in the SV19 study, the overall decrease is evident

Scimemi, Vladimirov, arXiv:1912.06532

$|q_T| = |P_{hT}|/z \ll Q$

 GeV^2 \mathbf{Q}^2

Boglione, Diefenthaler, Dolan, Gamberg, Melnitchouk, arXiv:2201.12197

$|q_T| = |P_{hT}|/z \ll Q$

Approximate region corresponding to MAP22 cuts

Boglione, Diefenthaler, Dolan, Gamberg, Melnitchouk, arXiv:2201.12197

$|q_T| = |P_{hT}|/z \ll Q$ Approximate region corresponding to MAP22 cuts

Boglione, Diefenthaler, Dolan, Gamberg, Melnitchouk, arXiv:2201.12197

The MAP22 cut is already considered to be "generous", but the physics seems to be the same for a much wider transverse momentum

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Aslan, Boglione, Gonzalez-Hernandez, Rainaldi, Rogers, Simonelli, 2401.14266

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The paper emphasizes the relevance of prescription choices and simultaneous TMD-PDF fit, but does not provide a fit to extended data sets.

Aslan, Boglione, Gonzalez-Hernandez, Rainaldi, Rogers, Simonelli, 2401.14266

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See talk by Tommaso Rainaldi

VECTOR MESON CONTAMINATIONS

Semi-inclusive



Semi-inclusive

































HARUT'S CARTOON AT TRANSVERSITY 2024



Procrustes - Greek Mythology







Procrustes - Greek Mythology



Procrustes - Greek Mythology











EXPECTED DATA



<u>G. Angelini's talk at SarWors2021</u>





Simple Guassians or bell-like shapes are not sufficient to describe data

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- The TMDs are probably different for different quark flavors

The TMD frag. functions are probably different for different final-state hadrons



ECINEE CIONS WITH OTHER FIELDS



Bermudez Martinez, Vladimirov, arXiv:2206.01105





Bermudez Martinez, Vladimirov, arXiv:2206.01105



TMD phenomenology



Bermudez Martinez, Vladimirov, arXiv:2206.01105





Bermudez Martinez, Vladimirov, arXiv:2206.01105



Avkhadiev, Shanahan, Wagman, Zhao, arXiv:2307.12359





Bermudez Martinez, Vladimirov, arXiv:2206.01105



<u>Avkhadiev, Shanahan, Wagman, Zhao, arXiv:2307.12359</u>



See talk by Patrizio Pucci



Bermudez Martinez, Vladimirov, arXiv:2206.01105



<u>Avkhadiev, Shanahan, Wagman, Zhao, arXiv:2307.12359</u>



See talk by Patrizio Pucci



CONNECTION WITH LATTICE QCD: TMDS





LPC collaboration, arxiv:2211.02340



CONNECTION WITH LHC PHYSICS: M_W



Unc. [MeV]	Total	Stat.	Syst.	PDF	A_i	Backg.	EW	е	μ	u_{T}	Lumi	Γ_W	PS
p_{T}^{ℓ}	16.2	11.1	11.8	4.9	3.5	1.7	5.6	5.9	5.4	0.9	1.1	0.1	1.5
m _T	24.4	11.4	21.6	11.7	4.7	4.1	4.9	6.7	6.0	11.4	2.5	0.2	7.0
Combined	15.9	9.8	12.5	5.7	3.7	2.0	5.4	6.0	5.4	2.3	1.3	0.1	2.3

Overview of m_w measurements

	,,		
I L L L (Combination			

7			ATLA	AS Coll	<u>ab. arX</u>	<u>iv:2403</u>	<u>3.1508</u>
-							
-							
]							
7.00	2	 	Lunai	Г			







Unc. [MeV]	Total	Stat.	Syst.	PDF	A_i	Backg.	EW	е	μ	u_{T}	Lumi	Γ_W	PS
p_{T}^ℓ	16.2	11.1	11.8	4.9	3.5	1.7	5.6	5.9	5.4	0.9	1.1	0.1	1.5
m _T	24.4	11.4	21.6	11.7	4.7	4.1	4.9	6.7	6.0	11.4	2.5	0.2	7.0
Combined	15.9	9.8	12.5	5.7	3.7	2.0	5.4	6.0	5.4	2.3	1.3	0.1	2.3

Overview of $m_{\mu\nu}$ measurements

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I L L L (Combination			

ATLAS Collab. arXiv:2403.15085

Not taking into account the flavor dependence of TMDs can lead to errors in the determination of the W mass, of the order of a few MeVs

Bacchetta, Bozzi, Radici, Ritzmann, Signori, arXiv:1807.02101

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The coupling constant of the strong force is determined from the transverse-momentum distribution of Z bosons produced in 8 TeV proton–proton collisions at the LHC and recorded by the ATLAS experiment.



ATLAS coll., arXiv:2309.12986



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 Table 1: Summary of the uncertainties in the
determination of $\alpha_s(m_Z)$, in units of 10^{-3} .

Experimental uncertainty	± 0.44			
PDF uncertainty	± 0.51			
Scale variation uncertainties	± 0.42			
Matching to fixed order	0	-0.08		
Non-perturbative model	+0.12	-0.20		
Flavour model	+0.40	-0.29		
QED ISR	± 0	.14		
N ⁴ LL approximation	± 0	.04		
Total	+0.91	-0.88		



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account recent TMD results

ATLAS coll., arXiv:2309.12986

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We try to always impose positivity limits. We prefer rigid and physical to flexible and unphysical







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We try to always impose positivity limits. We prefer rigid and physical to flexible and unphysical The fraction of same/opposite helicities is the same at any transverse momentum














The quarks with the same helicity as the proton's have less transverse momentum













The quarks with the same helicity as the proton's have more transverse momentum



<u>Yang, Liu, Sun, Zhao, Ma, arXiv:2409.08110</u>





MAP collaboration, arXiv:2409.18078





USED DATASETS

Yang, Liu, Sun, Zhao, Ma, arXiv:2409.08110

_			
Experiment	Process	Data points	χ^2/N
HERMES[82]	$e^{\pm}p \to e^{\pm}hX$	84 (160)	0.72
$\mathrm{HERMES}[82]$	$e^{\pm}d \to e^{\pm}hX$	160(317)	0.71
CLAS[83]	$e^- p \to e^- \pi^0 X$	9(21)	1.43
Total		253 (498)	0.74

MAP collaboration, arXiv:2409.18078

Experiment	$N_{\rm dat}$	$\chi^2_{ m NLL}/N_{ m dat}$	$\chi^2_{ m NNLL}/N_{ m dat}$
HERMES $(d \to \pi^+)$	47	1.34	1.30
HERMES $(d \to \pi^-)$	47	1.10	1.08
$\left[\text{HERMES } (d \to K^+)\right]$	46	1.26	1.25
$ \text{HERMES } (d \to K^-) $	45	0.93	0.89
HERMES $(p \to \pi^+)$	53	1.17	1.21
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Total	291	1.11	1.09





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More data needed

MAP collaboration, arXiv:2409.18078

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EXPECTED DATA

Multidimensional binning needed











See talks by Tim Hayward, Bakur Parsamyan

















SIVERS FUNCTION

$$\rho_{N^{\uparrow}}^{q}(x,k_{x},k_{y};Q^{2}) = f_{1}^{q}(x,k_{T}^{2};Q^{2}) - \frac{k_{x}}{M}f_{1T}^{\perp q}(x,k_{T}^{2};Q^{2})$$

In a nucleon polarized in the +y direction,

the distribution of quarks can be distorted in the x direction

Q^{2})



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In a nucleon polarized in the +y direction, the distribution of quarks can be distorted in the x direction $\int_{a}^{b} dt dt$









Bury, Prokudin, Vladimirov, arXiv:2103.03270



3D STRUCTURE IN MOMENTUM SPACE



Q=2GeV

Bacchetta, Delcarro, Pisano, Radici, arXiv:2004.14278



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Bacchetta, Delcarro, Pisano, Radici, arXiv:2004.14278

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S-FIT OF SINGLE TRANSVERSE-3PT



Interesting work from the point of view of simultaneous use of several measurements, but still limited from other perspectives (lack of TMD evolution and knowledge of the unpolarized function)



SIVERS FUNCTION WITH NEURAL NETWORKS



Interesting work from the point of view of the use of Neural Networks, but still limited from other perspectives (lack of TMD evolution and knowledge of the unpolarized function)

Fernando, Keller, arXiv:2304.14328





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Fernando, Keller, arXiv:2304.14328



See talks by Ishara Fernando





CONNECTION WITH ANGULAR MOMENTUM?



- Diehl & Kroll, arXiv:1302.4604
- Guidal et al., PR D72 (05) 054013
- Liuti et al., PRD 84 (11) 034007
- Bacchetta & Radici, PRL 107 (11) 212001



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Diehl	&	Kroll,	arXiv:1302.4604
Diem	œ	KIOII,	arAiv.1502.4004

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arXiv:2409.17955 (lattice)



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arXiv:2409.17955 (lattice)

See also arXiv:1907.06960 for a critique





IMPACT OF RICH DATA?





Artistic view of factorization







Artistic view of factorization







Artistic view of factorization



Works for SIDIS, Drell-Yan, e⁻e⁺ annihilation





Artistic view of factorization



Works for SIDIS, Drell-Yan, e⁻e⁺ annihilation







Artistic view of factorization



Works for SIDIS, Drell-Yan, e⁻e⁺ annihilation







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TMD factorization does not work for pp to hadrons





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TMD factorization does not work for pp to hadrons

Mulders, Rogers, arXiv:1001.2977



See talk by Oleg Eyser

































In five to ten years

- TMD multiplicities for pions and kaons, off protons and deuterons, from COMPASS and JLab
- Drell-Yan and Z measurements from CERN, RHIC, FermiLab (COMPASS with pions)
- TMD multiplicities for pions and kaons in e⁺e⁻ from BELLE and BES
- Better understanding and control of higher-order QCD corrections
- More flexible functional forms, flavour dependence, at least two or three alternative extractions
- Use TMDs for something else (W mass... comparison with lattice... Wigner distributions...)

READY TO USE EIC DATA

slide from 2018 CPHI@Yerevan

51




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X X • TMD multiplicities for pions and kaons, off protons and deuterons, from





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- Progress is ongoing concerning higher-twist and gluon TMDs
- Extractions of unpolarized TMDs are reaching a good level of sophistication, but there are still several open questions and new data are needed
- For other TMDs, the study has started and there is an increasing number of new results, but more data are needed