

Composite Dark Matter from $Sp(2N)$ gauge theory

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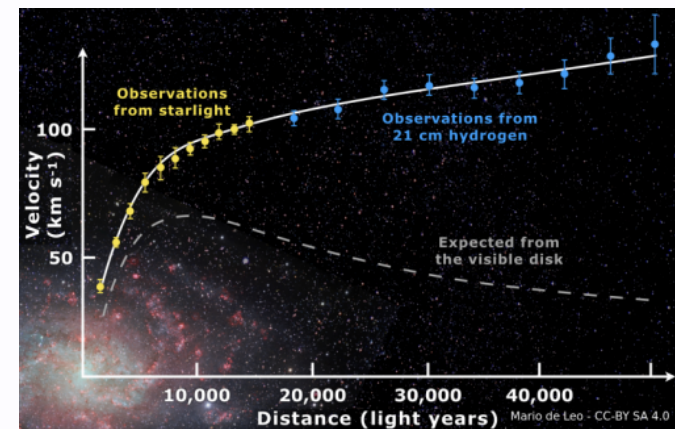
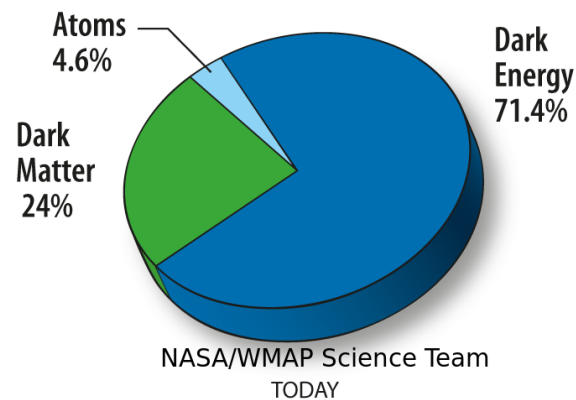
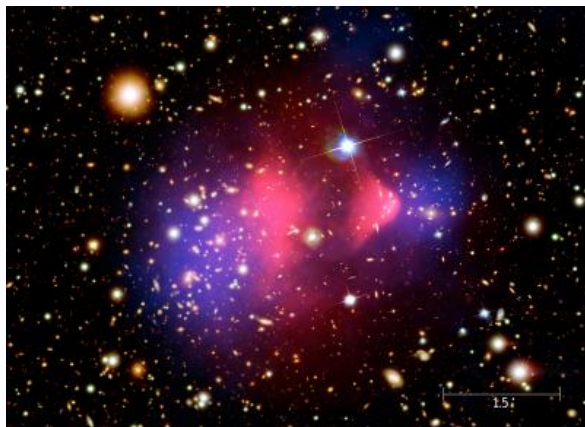
Busan, February 20, 2024

PNU workshop on Composite Higgs Physics

1. QCD inspired Dark Matter models

Dark Matter - Why?

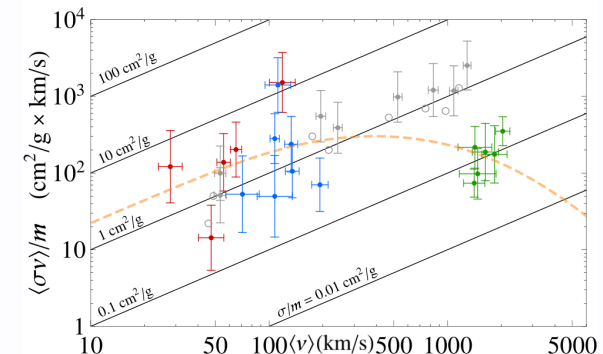
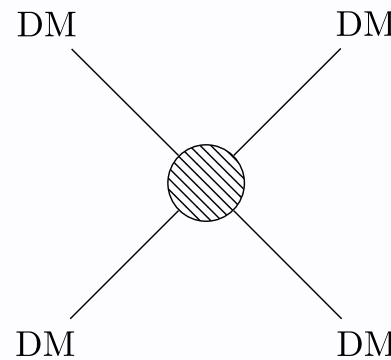
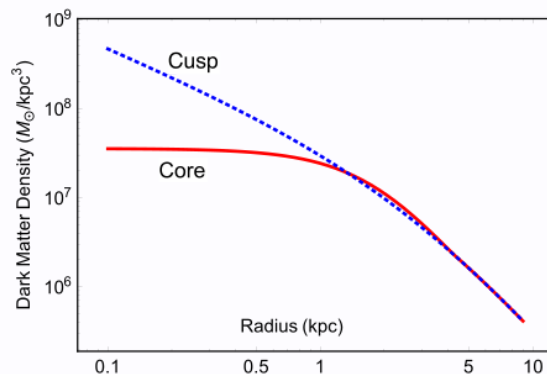
- Strong observational evidence at many scales! ^[1]
- Modified Gravity ^[2] is a potential alternative
- New particles beyond the Standard Model (BSM) promising!



[1] see e.g. Bullock, Boylan-Kolchin [1707.04256], Tulin, Yu [1705.02358]

Dark Matter properties

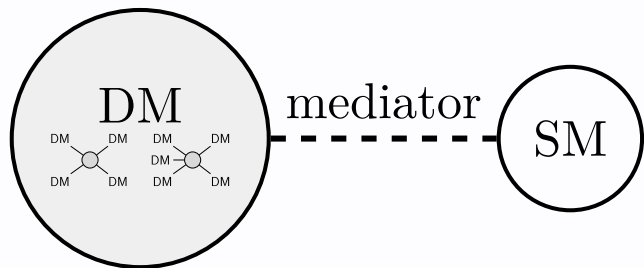
- DM self-interaction phenomenologically allowed^[1] and potentially relevant for small-scale structure problems
 - non-vanishing scattering cross-sections $\sigma_{2\text{DM}\rightarrow 2\text{DM}}$
 - velocity dependence of $\sigma_{2\text{DM}\rightarrow 2\text{DM}}$ preferred



QCD-like Dark Matter can those provide self-interactions!

Strongly Interacting Gauge Theories in DM Models

- With fermions: Global symmetries make DM stable
- With mediator: Dark sector coupled to SM


$$\mathcal{L}_{\text{DM}} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + \bar{\psi}_f(i\not{D} + m_f)\psi_f$$

- Non-vanishing self-scattering cross-section arise

$$\langle v\sigma_{\pi\pi\rightarrow\pi\pi} \rangle \neq 0$$

- Relic density driven by strong processes

Dark meson scattering: Determine DM relic density

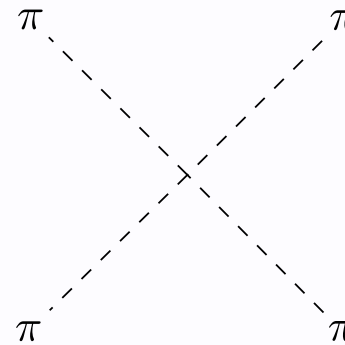
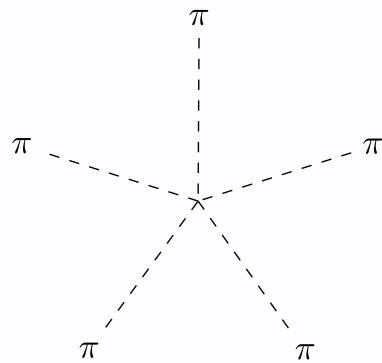
- Any model must predict the current density of DM correctly
 - number density n can be calculated using Boltzmann equations

$$\partial_t n + 3Hn = f(\langle v\sigma_{\text{number changing}} \rangle)$$

- Cross-sections $\langle \sigma v \rangle$ are input for Boltzmann equations
 - describe non-equilibrium dynamics
 - H is the Hubble rate

Relevant pion scattering channels

- $3\pi \rightarrow 2\pi$ (semi-annihilation) [1]
- $2\pi \rightarrow 2\pi$ (self-scattering)
 - self-scattering among DM [2]
 - resonant enhancements [3]
- $2n\pi \rightarrow 2\pi$ (multi-hadron bound states) [4]



Strongly Interacting Massive Particles (SIMPs)

- Depletion via $3\text{DM} \rightarrow 2\text{DM}$ ^[1], i.e. $3\pi \rightarrow 2\pi$
 - same as $KK \rightarrow 3\pi$ in QCD ^[2]
 - Early universe: SM \rightleftharpoons DM equilibrium
 - Dark matter depletion process: **freeze-out**
- LO ChiPT matches relic density at

$$m_\pi \approx \mathcal{O}(100)\text{MeV} - \mathcal{O}(1)\text{GeV}$$

Dark Matter with $3\text{DM} \rightarrow 2\text{DM}$ depletion and self-interactions

Other mass scales than QCD are relevant!

- Lagrangian has two free parameters: g^2 and m_f
 - one overall energy scale
 - one intrinsic strong scale
- Overall scale should allow sufficiently heavy DM
- m_f should lead to parametrically light m_π
 - both scales can deviate strongly from QCD!

Lattice investigations of a larger parameter space are useful!

[1] Hochberg et. al. [[1411.3727](#)] [[1512.07917](#)] [2] Choi et.al. [[1801.07726](#)] Bernreuther et.al. [[2311.17157](#)]

[3] Kulkarni et.al. [[2202.05191](#)]

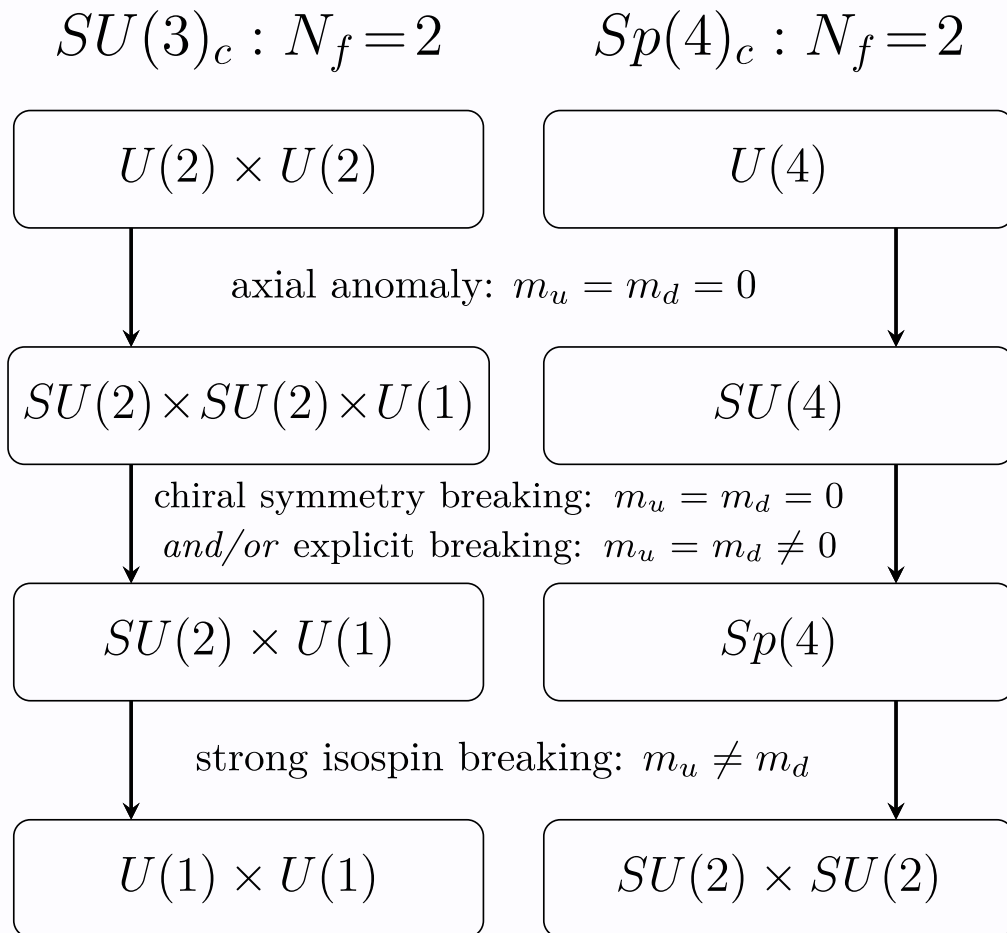
see also Hansen et.al. [[1507.01590](#)]

Other relevant channels

- decay to Standard Model: $2\pi \rightarrow SM$ ^[1]
- involvement of vector mesons: $\pi\pi \rightarrow \pi\rho$, $3\pi \rightarrow \pi\rho$ ^[2]
- influence of light singlets: $\eta'\eta' \rightarrow \pi\pi$, $\pi\pi \rightarrow \eta'\pi$, ... ^[3]
- The relevance depends on the spectrum
 - investigation of the meson spectrum important
 - lattice investigations inform EFT construction

[1] Kosower (Phys.Lett.B. 1984)

[2] Hochberg et. al. [1411.3727] [1512.07917]



SIMPs from $Sp(4)$ gauge theory

- Pseudo-real representation: ^[1]
 - \Rightarrow more pseudo-Goldstones
 - \Rightarrow no fermionic bound states
- $N_f = 2$: exactly 5 Goldstones
 - Allows 3DM \rightarrow 2DM ^[2]

$Sp(4)$ with two fermions is a minimal SIMP DM realisation

Lagrangian of $Sp(4)_c$ with fermions

$$\mathcal{L}_{Sp(4)} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + \sum_{f=u,d} \bar{\psi}_f (i\not{D} + m_f)\psi_f$$

- Higher symmetry than QCD-like theories

$$\Psi = \begin{pmatrix} u_L \\ d_L \\ -SCu_R^* \\ -SCd_R^* \end{pmatrix} = \begin{pmatrix} u_L \\ d_L \\ \tilde{u}_R \\ \tilde{d}_R \end{pmatrix} \quad \begin{array}{l} C \dots \text{charge conj.} \\ S \dots \text{colour matrix} \end{array}$$

$$\mathcal{L}_{Sp(4)} = i\bar{\Psi}\not{D}\Psi - \frac{1}{2}(\Psi^T SCM\Psi + h.c.) - \frac{1}{4}F_{\mu\nu}F^{\mu\nu}$$

- generators τ_a in fundamental repr. : $S\tau_a S = -\tau_a^T$

Meson multiplets of $Sp(4)_c$ with $N_f = 2$

- Extra gauge invariant states: $q^T \dots q$ and $\bar{q} \dots \bar{q}^T$
- $Sp(2N_f)$ flavour symmetry

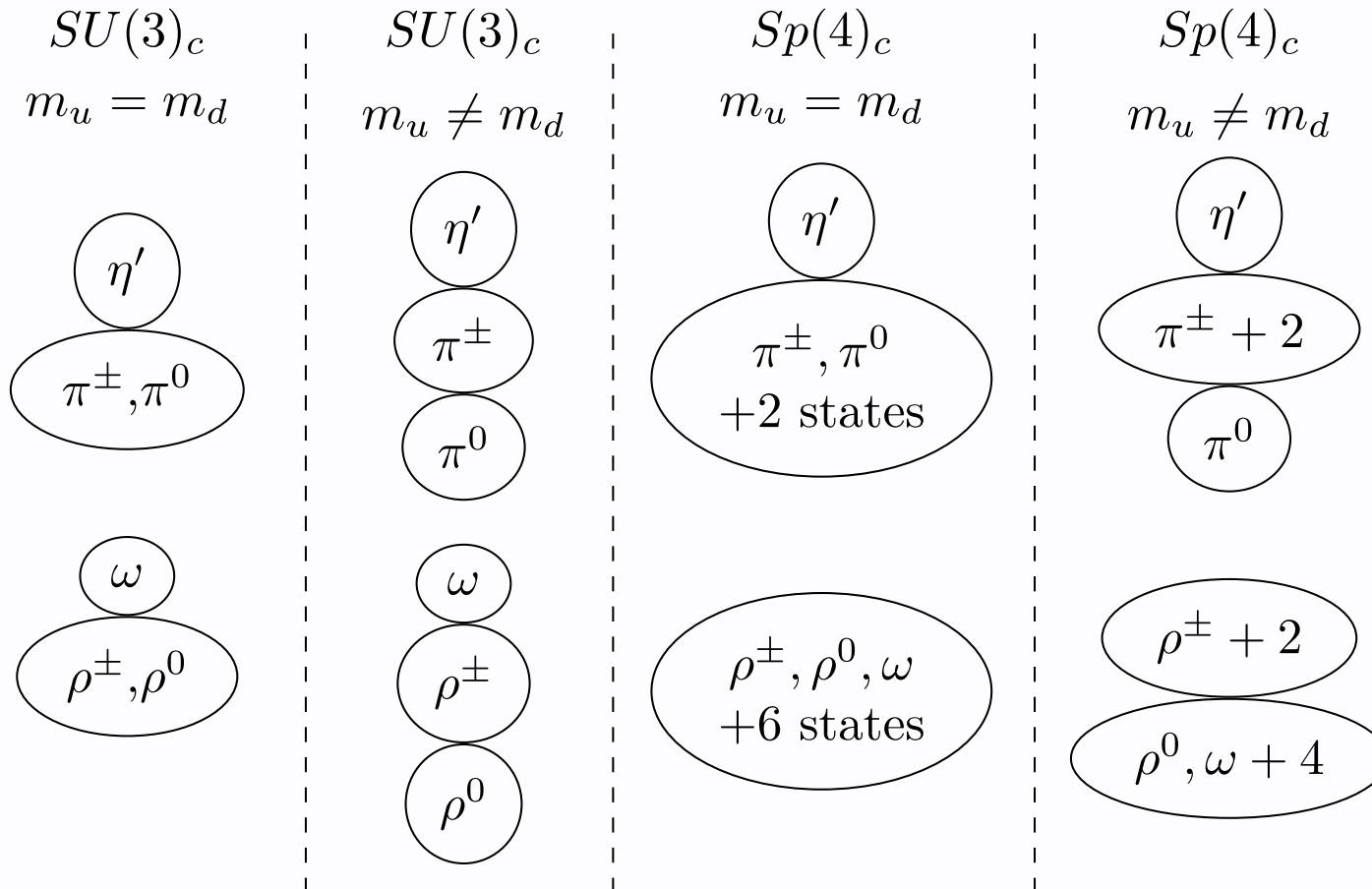
$$Sp(4)_F : \quad 4 \otimes 4 = 1 \oplus 5 \oplus 10$$

- Different multiplets depending on J^P number

$$5\pi, 10\rho, 5a_0, 5a_1, 1\eta', 1\sigma, 0\omega$$

The global symmetries lead to a richer meson multiplet structure!

Pseudoscalar (PS) and vector (V) multiplets



The same patterns persist for other channels.

BSM wishlist from the lattice

1. Masses and decay constants of dark hadrons
 - Non-singlet and singlet mesons, glueballs
2. Scattering of dark pions
 - $2\pi \rightarrow 2\pi$ for self-interaction crosssection
 - $3\pi \rightarrow 2\pi$ for SIMP semi-annihilation
3. Applicability of χ PT and related EFTs

[1]Nogradi,Szikszai[2107.05996][2]Bennett et.al.[2010.15781][3]Bennett et.al.[2202.05516],Drach et.al.[2107.09974]

[4] e.g. Boz. et.al. [1912.10975], [5] Mason et.al. [2310.02145]

Applications of $Sp(2N)$ gauge theory beyond SIMP DM

- Generic features of non-Abelian confining gauge theories
 - Hadron masses as functions of N_f and N_c [1]
 - large N_c limit [2]
- Higgs compositeness, partial top compositeness [3]
 - Mixed fermion representations: near conformal behaviour?
- Model theory for finite density calculations (no sign problem) [4]
- Finite temperature behaviour: Deconfinement and chiral symmetry
 - Potential first order phase transitions? [5]

Composite Higgs studies can be repurposed

- Coset spaces for Higgs physics are large enough for SIMP DM
 - applies also to different fermion reps. (e.g. $Sp(4)$, $N_f^{as} \geq 2$)
 - or mixed representation theories
- Particle spectrum determines relevant hadronic states
- Scattering studies in the context of WW scattering

Lattice Investigations:

Quantitative Insights

Lattice spectroscopy: Getting meson masses

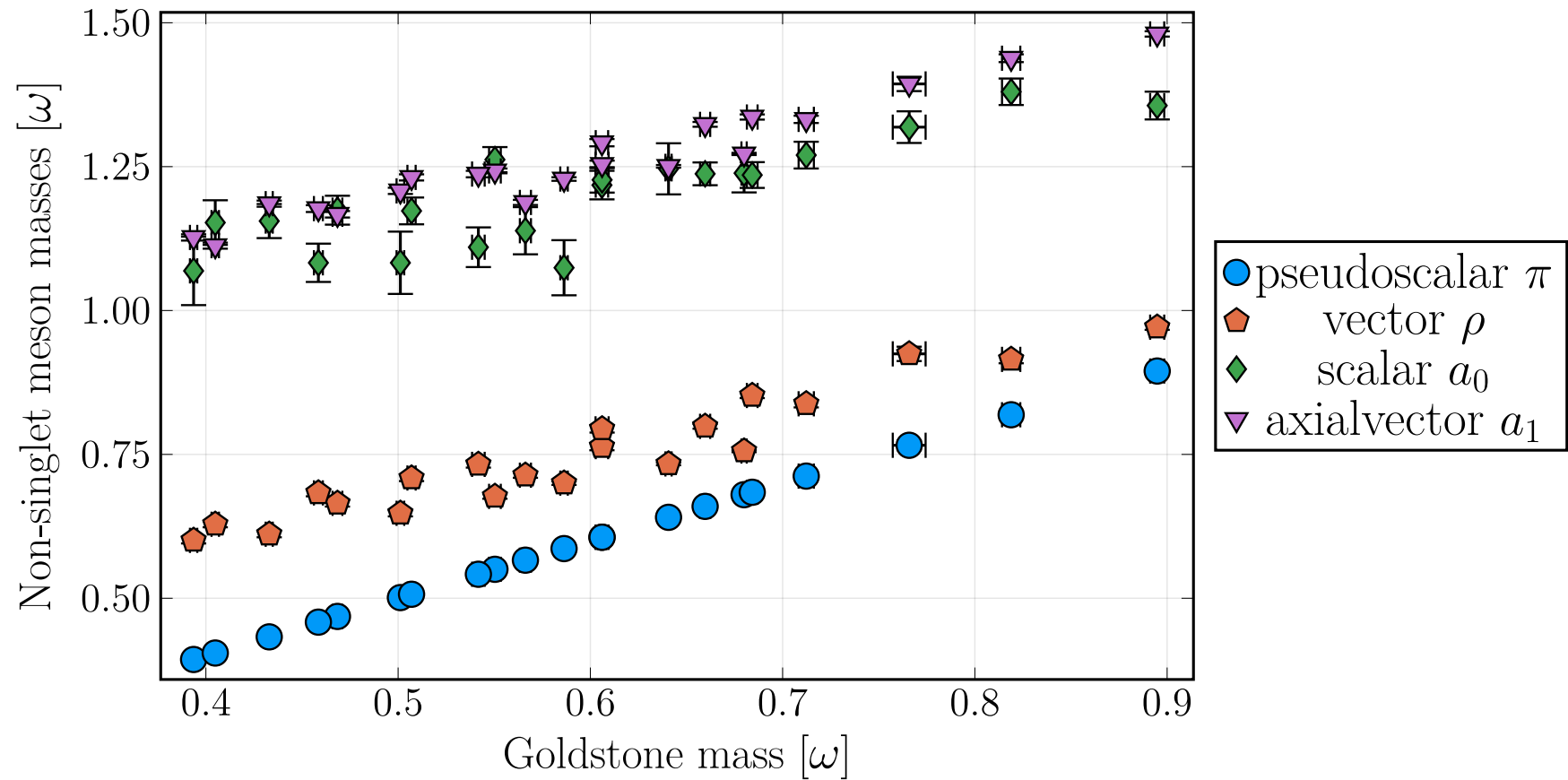
- Construct operator with same quantum numbers
- Energy levels from Euclidean correlator

$$C_{\mathcal{O}}(t) = \sum_n \frac{1}{2E_n} \langle 0 | \mathcal{O} | n \rangle^* \langle n | \mathcal{O} | 0 \rangle e^{-E_n t}.$$

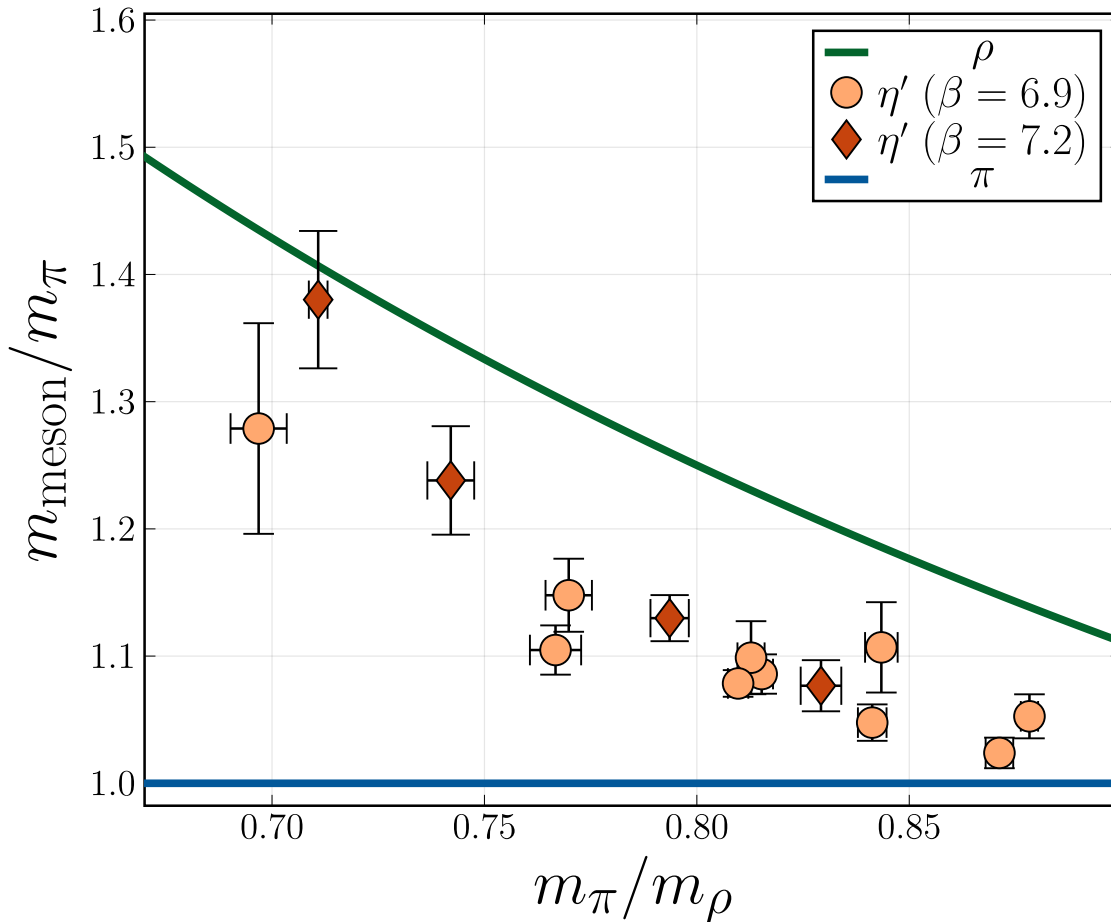
- For mesons a generic correlator

$$C(t - t') = \sum_{\vec{x}, \vec{y}} \left(\begin{array}{c} \text{Diagram 1: } \vec{x}, t \text{ and } \vec{y}, t' \text{ with two loops} \\ \text{Diagram 2: } \vec{x}, t \text{ and } \vec{y}, t' \text{ with two arcs} \end{array} \right) + \underbrace{\text{const.}}_{=|\langle 0 | \mathcal{O} | 0 \rangle|^2}$$

Non-singlet spectrum



The pseudoscalar and vector mesons are the lightest non-singlets.²⁰



The pseudoscalar singlet η' is surprisingly light!

- Phenomenologically relevant:
 - $m_\rho > m_{\eta'}$ different from QCD
 - relevant low-energy dof
 - η' relevant for $\pi\pi$ scattering
 - more accessible channels for decays into SM

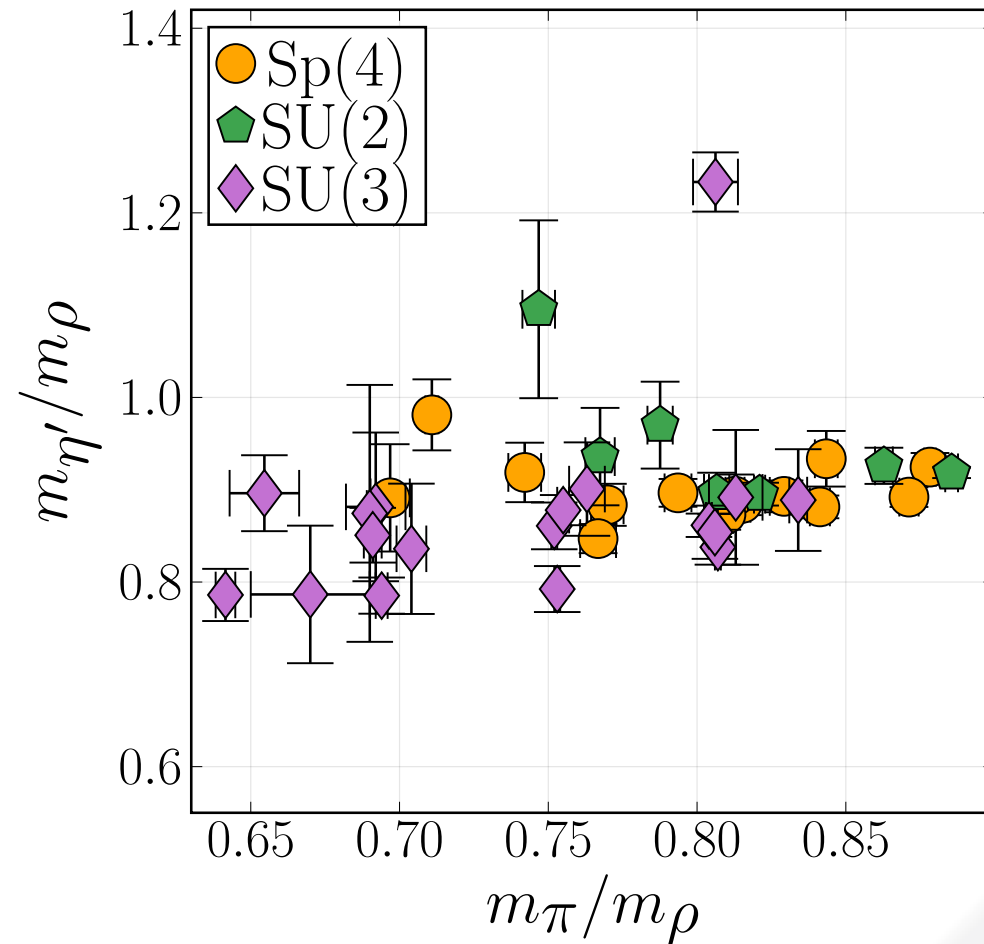
Interesting! Is this surprising?

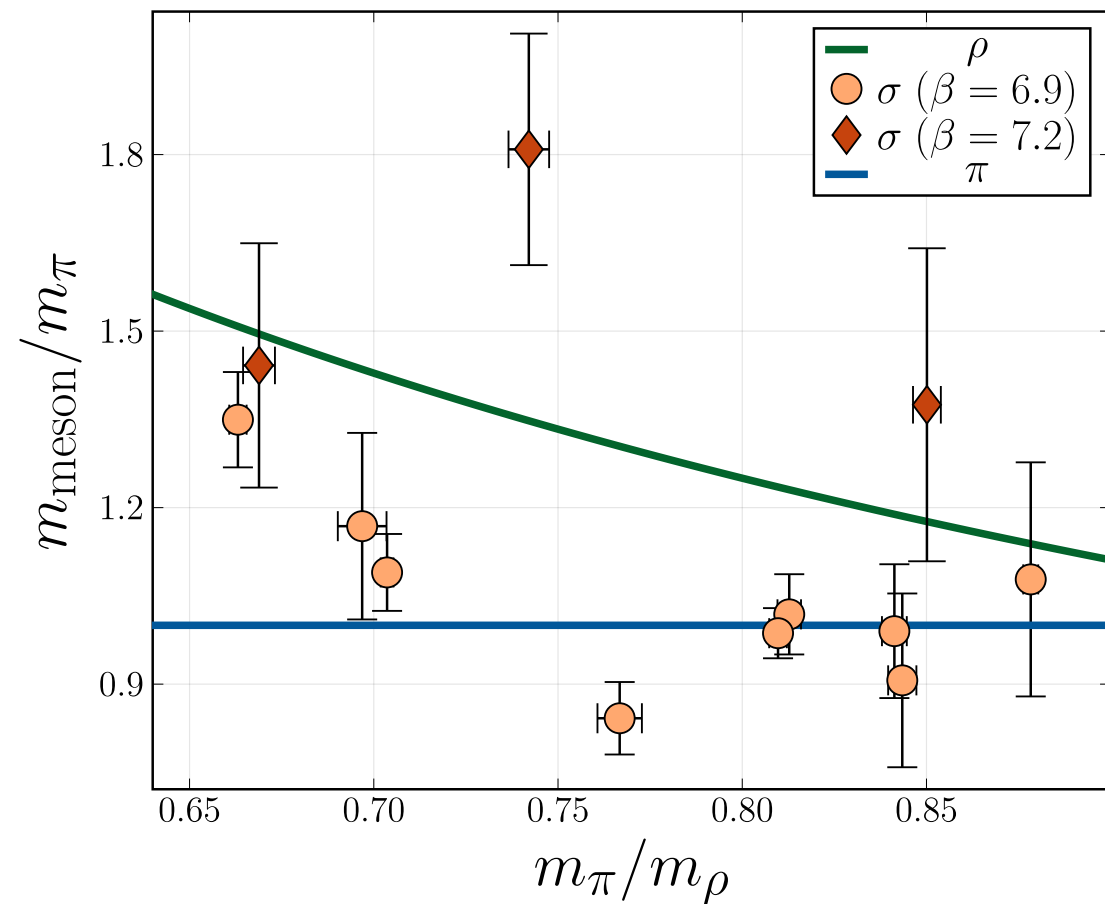
Consider different theories:

- Large N_c : $m_{\eta'} - m_\pi \propto N_f / N_c$
 - $N_f = 2$ could be "small"
 - $N_c = 4$ could be "large"

SU(2) and SU(3) comparison:

- Similarities: generic $N_f = 2$ feature?
- QCD: strong N_f dependence
- Differences may arise $m_\pi / m_\rho \rightarrow 0$
mass driven by flavour content!





The scalar singlet σ

- Strong finite spacing effects!
- Potentially light state below $2m_\pi$ threshold
- Unclear systematics

Overall, inconclusive results.

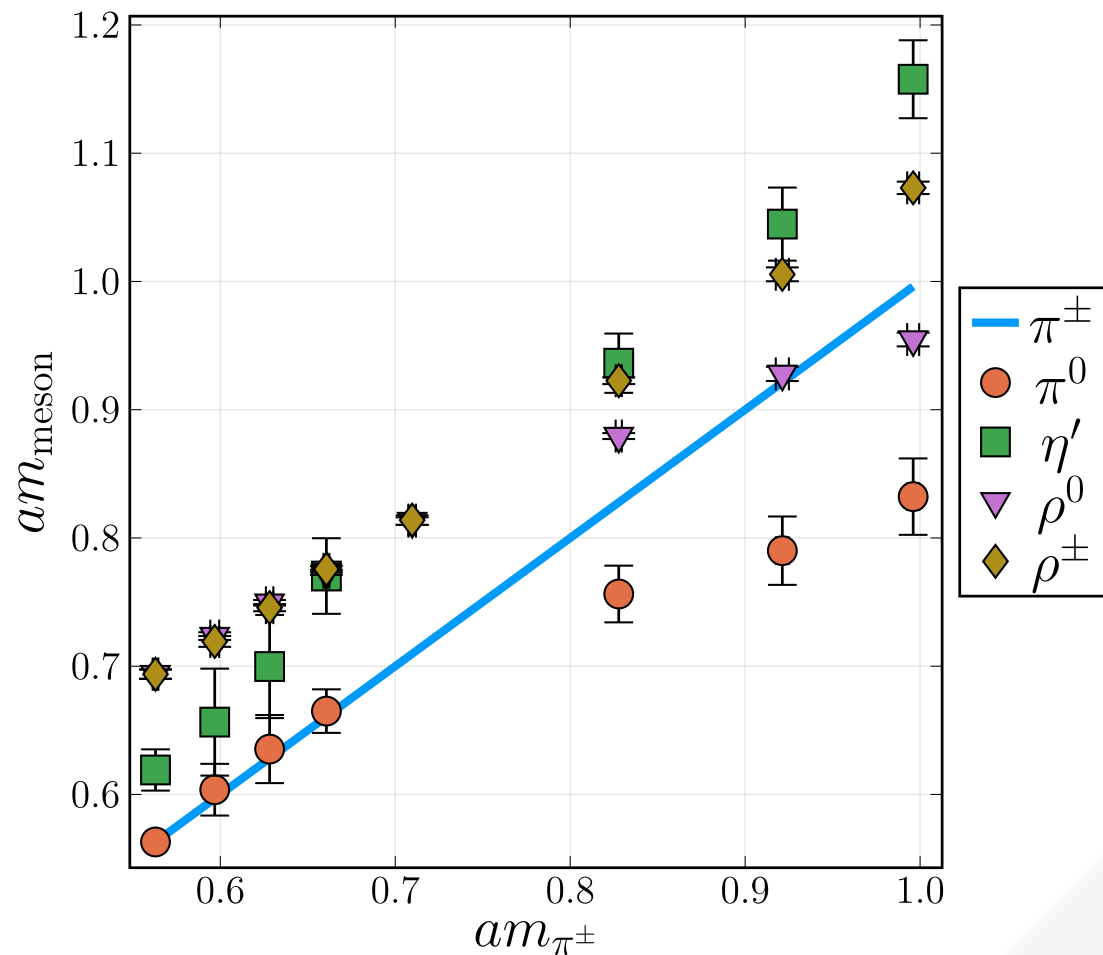
[1] Bennett et. al. [2304.07191]

Kulkarni et. al. [2202.05191]

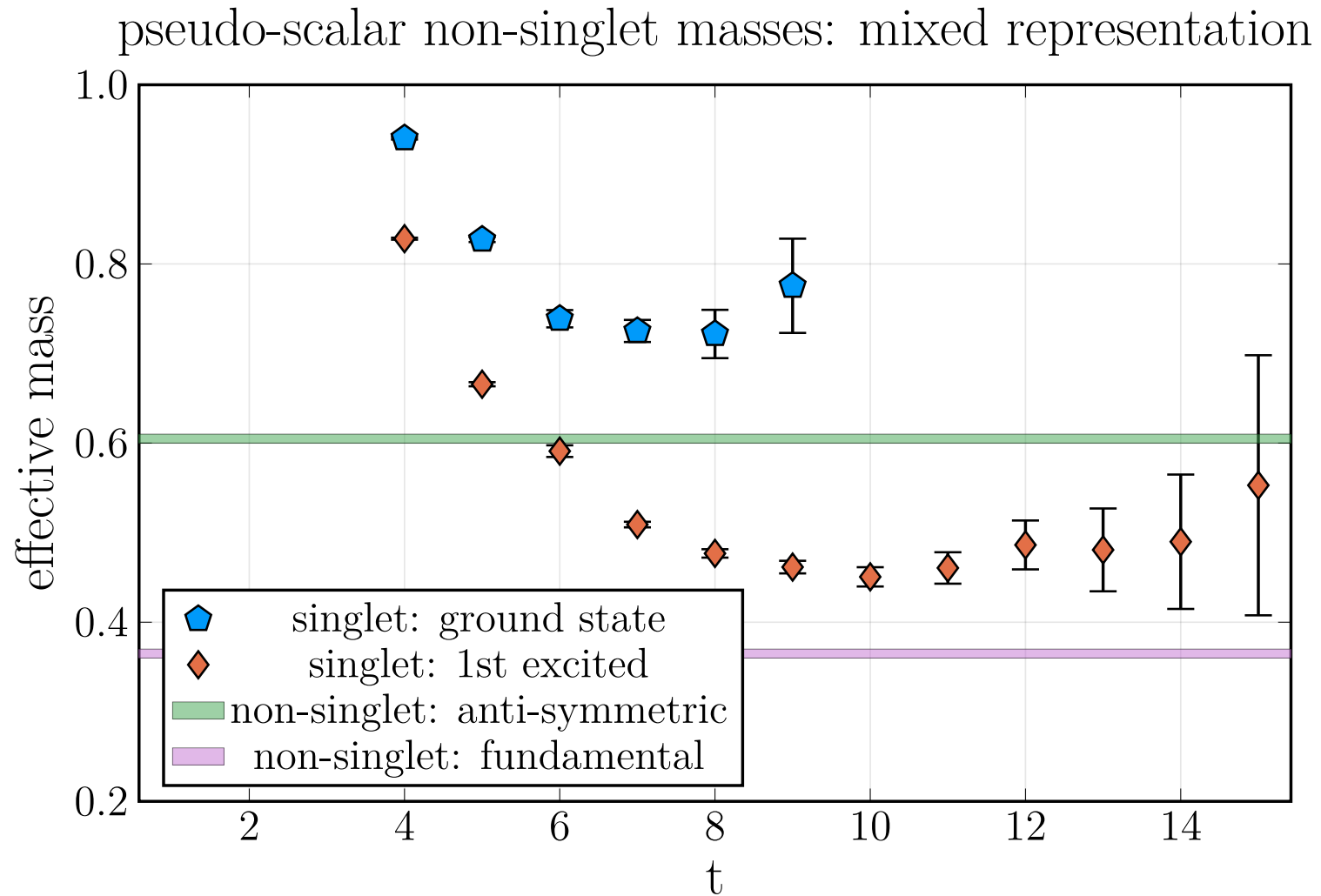
Non-degenerate fermions

- Different mass hierarchies
- Transition from a flavour-symmetric theory to a heavy-light system
- One light and one heavy pseudoscalar flavour-singlet

$Sp(4)$, $m_u \neq m_d$



Pseudoscalar (non-)singlets: mixed representation



Consequences for Dark Matter

- Mass hierarchies: limit χ PT validity
 - inclusion of other states than π required, e.g. η' and ρ
 - additional tests needed (fermions are too heavy)
- Light unprotected states η' , π^0 allow decays into SM
 - no protection from symmetry

Are these fermion masses phenomenologically relevant?

Dark Matter Scattering on the Lattice

- Pions are in the 5-dimensional representations
- A two pion scattering is in one of three irreps

$$5 \times 5 = 14 \oplus 10 \oplus 1$$

- Corresponds to the usual QCD channels
 - $14 \Leftrightarrow$ isospin $I = 2$ in QCD, e.g. $\pi^+ \pi^+$
 - $10 \Leftrightarrow$ isospin $I = 1$ in QCD, e.g. $\pi\pi \rightarrow \rho$
 - $0 \Leftrightarrow$ isospin $I = 0$ in QCD, e.g. $\pi\pi \rightarrow \sigma / f_0$

Scattering information from the lattice

- Scattering phase shift $\delta_0(p)$ from finite volume energy

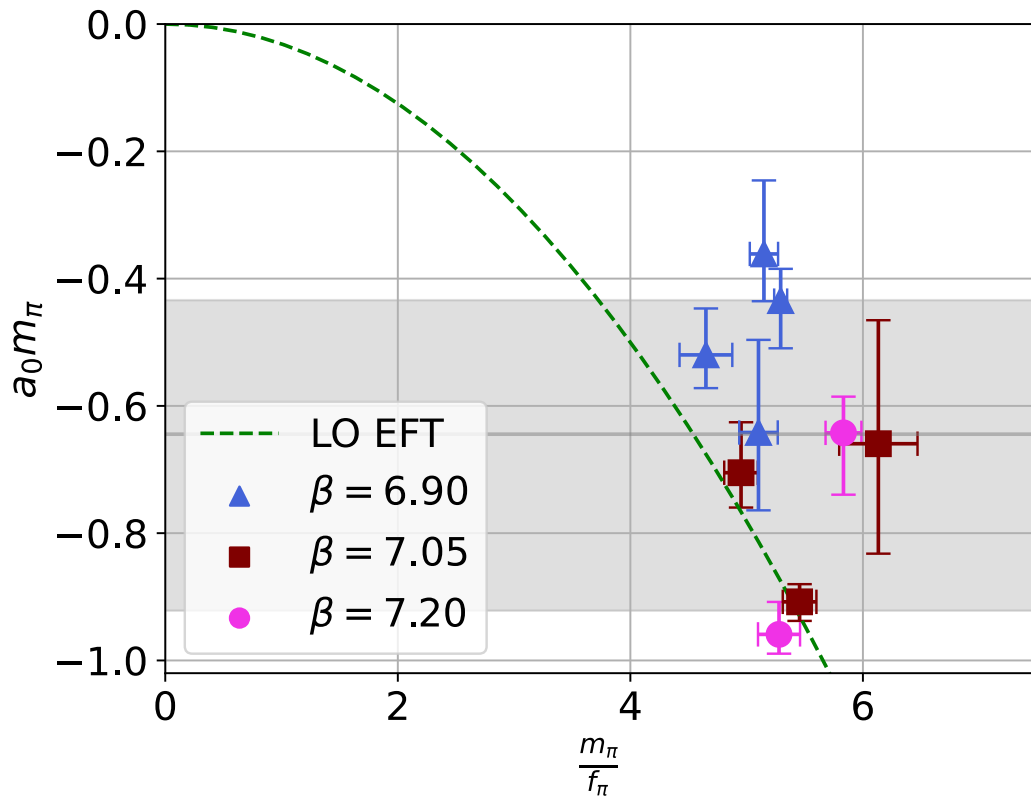
$$\tan(\delta_0(q)) = \frac{\pi^{\frac{3}{2}} q}{\mathcal{Z}_{00}^{\vec{0}}(1, q^2)}, \quad q = p^* \frac{L}{2\pi}$$

$$\cosh\left(\frac{E_{\pi\pi}}{2}\right) = \cosh(m_{\pi\pi}) + 2 \sin\left(\frac{p^*}{2}\right)^2$$

- Low-velocity behaviour: Scattering length

⇒ relation between $\pi\pi$ energy $E_{\pi\pi}$ and m_π on a lattice [1]

$$\frac{\delta E_{\pi\pi}}{m_\pi} = \frac{4\pi m_\pi a_0}{(m_\pi L)^3} \left(1 + c_1 \frac{m_\pi a_0}{m_\pi L} + c_2 \left(\frac{m_\pi a_0}{m_\pi L} \right)^2 \right)$$



First investigation of isospin-2 scattering

- repulsive $\pi\pi$ interaction
- few lattice energy levels available \Rightarrow systematics
- finite volume effects present
- roughly matches ChiPT

Summary

- Full light hadron spectrum of two-flavour $Sp(4)$
 - surprisingly light η' , potentially light σ
 - input for EFTs: masses and decay constants
 - first determination of isospin-2 $\pi\pi$ scattering

Outlook

- Full scattering analysis of $2\pi \rightarrow 2\pi$ and $3\pi \rightarrow 2\pi$
- Better understanding of singlets and scattering states:
Reduce lattice artefacts
- Spectroscopy closer to the chiral limit

Thank you