

Lattice investigations of the chimera baryon spectrum in the $Sp(4)$ gauge theory

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Outline

- Introduction:
 - ▶ ~~Sp(4) gauge theory: A Composite Higgs model~~
 - ▶ ~~Lattice method~~
 - ▶ Chimera baryon
- Results
 - ▶ Projections
 - ▶ Mass hierarchy of chimera baryons
 - ▶ Chiral EFT and AIC
- Summary and Outlook

Our choice of model

- Sp(4) gauge theory with $2F+3AS$ Dirac fermions



- Breaking pattern: $4F+6AS$ 2-component Weyl fermions

$$G/H = \underline{SU(4) \times SU(6)} / Sp(4) \times SO(6)$$

Enhanced global symmetry due to the (pseudo-) reality

● $SU(4)/Sp(4)$ gives 5 goldstone bosons.

- ▶ 4: SM Higgs doublet
- ▶ 1: made heavy in model building

● SU(3) embedded in antisymmetric representation:

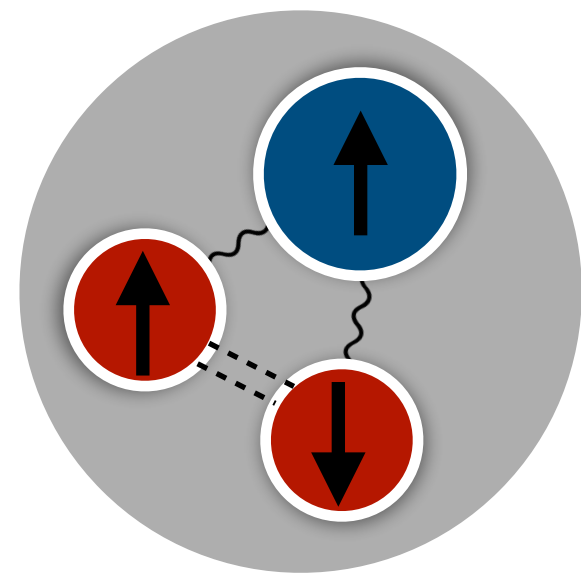
$$SU(6) \rightarrow SO(6) \supset SU(3)$$

↳ QCD colour SU(3)

Chimera Baryon

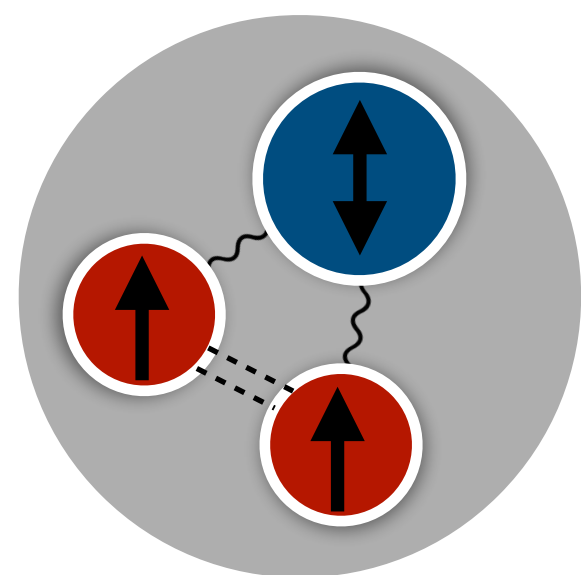
- Interpolating operators

- Λ type: $\mathcal{O}_{\text{CB},\gamma^5} = (\bar{\psi}^{1a}\gamma^5\psi^{2b})\Omega_{bc}\chi^{kca}$

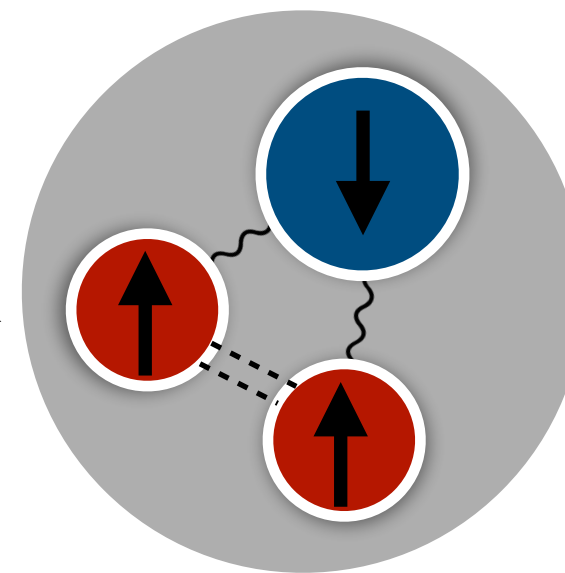


$(J, R) = (1/2, 5)$
*top partner

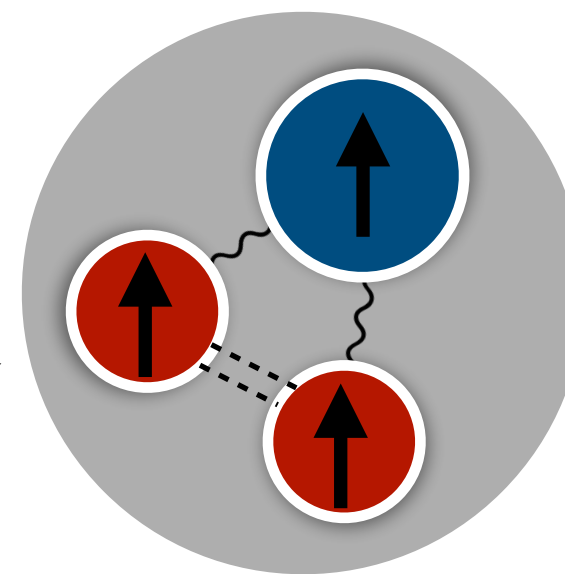
- Σ type: $\mathcal{O}_{\text{CB},\gamma^\mu} = (\bar{\psi}^{1a}\gamma^\mu\psi^{2b})\Omega_{bc}\chi^{kca}$



Spin projection



$\Sigma: (J, R) = (1/2, 10)$
*top partner



$\Sigma^*: (J, R) = (3/2, 10)$

a, b, c : hypercolour

Ω : 4×4 symplectic matrix

J : spin

R : irreducible rep. of the fundamental sector


Results

Quenched approximation

- ▶ Projections
- ▶ Mass hierarchy of chimera baryons
- ▶ Chiral EFT and AIC

\hat{m}_{PS} : fundamental
 \hat{m}_{ps} : Antisymmetric

$\hat{a} \equiv a/\omega_0$ and $\hat{m} \equiv \omega_0 m$



Ensemble	β	$N_t \times N_s^3$	$\langle P \rangle$	ω_0/a
QB1	7.62	48×24^3	0.60192	1.448(3)
QB2	7.7	60×48^3	0.608795	1.6070(19)
QB3	7.85	60×48^3	0.620381	1.944(3)
QB4	8.0	60×48^3	0.630740	2.3149(12)
QB5	8.2	60×48^3	0.643228	2.8812(21)

Results

Projection-CB two-point function

► Interpolating operator

$$\mathcal{O}_{\text{CB}}^\gamma(x) \equiv \left(Q^{ia}{}_\alpha(x) \Gamma^{1\alpha\beta} Q^{jb}{}_\beta(x) \right) \Omega_{ad} \Omega_{bc} \Gamma^{2\delta\gamma} \Psi^{kcd}{}_\gamma(x)$$

► two-point function

At large Euclidean time

$$C^{\gamma\gamma'}(t) \equiv \sum_{\vec{x}} \langle \mathcal{O}_{\text{CB}}^\gamma(x) \overline{\mathcal{O}_{\text{CB}}^{\gamma'}}(0) \rangle \rightarrow P_e [c_e e^{-m_e t} + c_o e^{-m_o(T-t)}] - P_o [c_o e^{-m_o t} + c_e e^{-m_e(T-t)}]$$

$$= - \sum_{\vec{x}} \left(\Gamma^2 S_{\Psi}^{kcd}{}_{c'd'}(x,0) \overline{\Gamma^2} \right)_{\gamma\gamma'} \Omega_{cb} \Omega^{b'c'} \Omega_{ad} \Omega^{d'a'} \quad P_e \equiv \frac{1}{2}(1 + \gamma^0) \text{ and } P_o \equiv \frac{1}{2}(1 - \gamma^0)$$

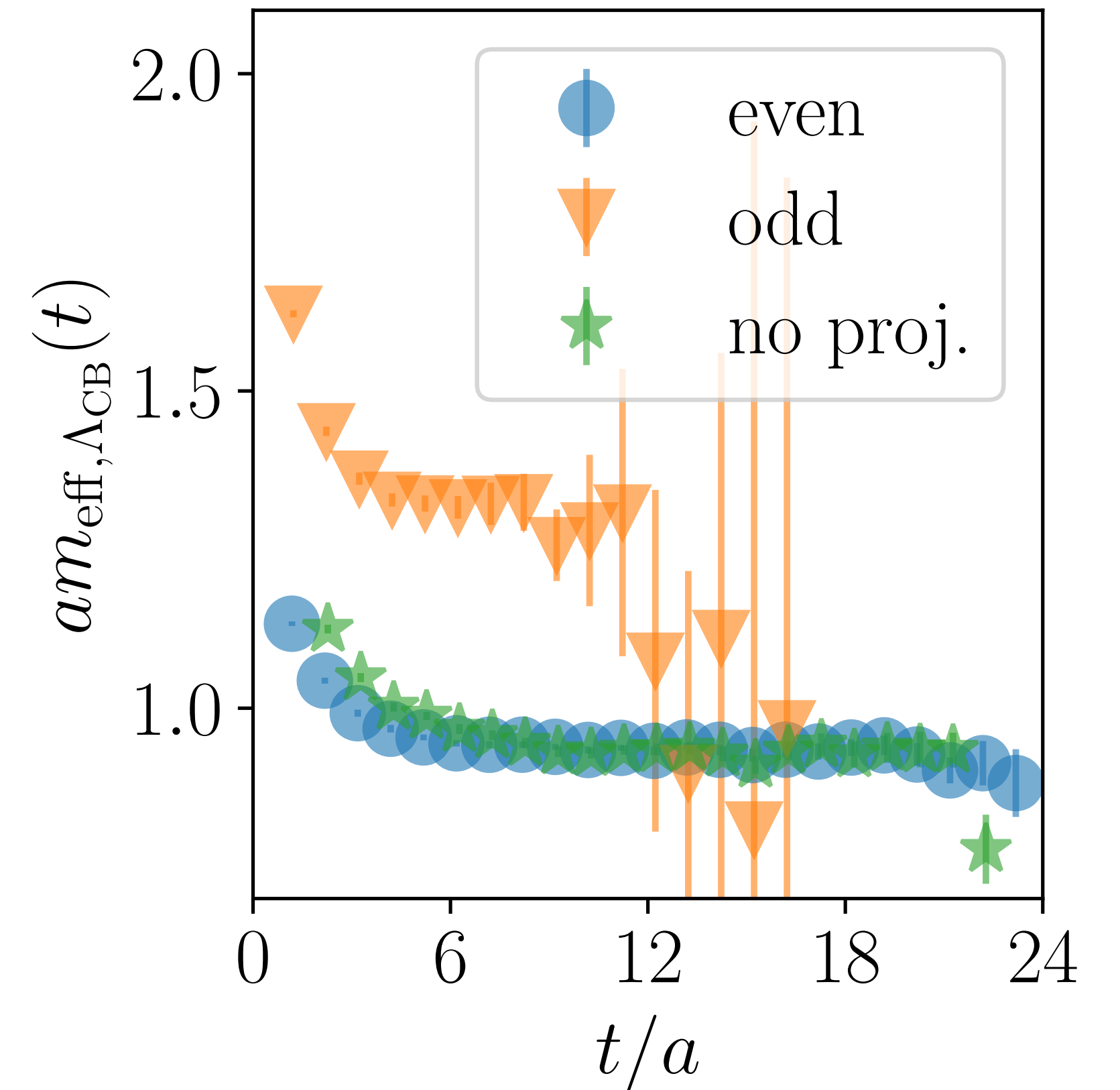
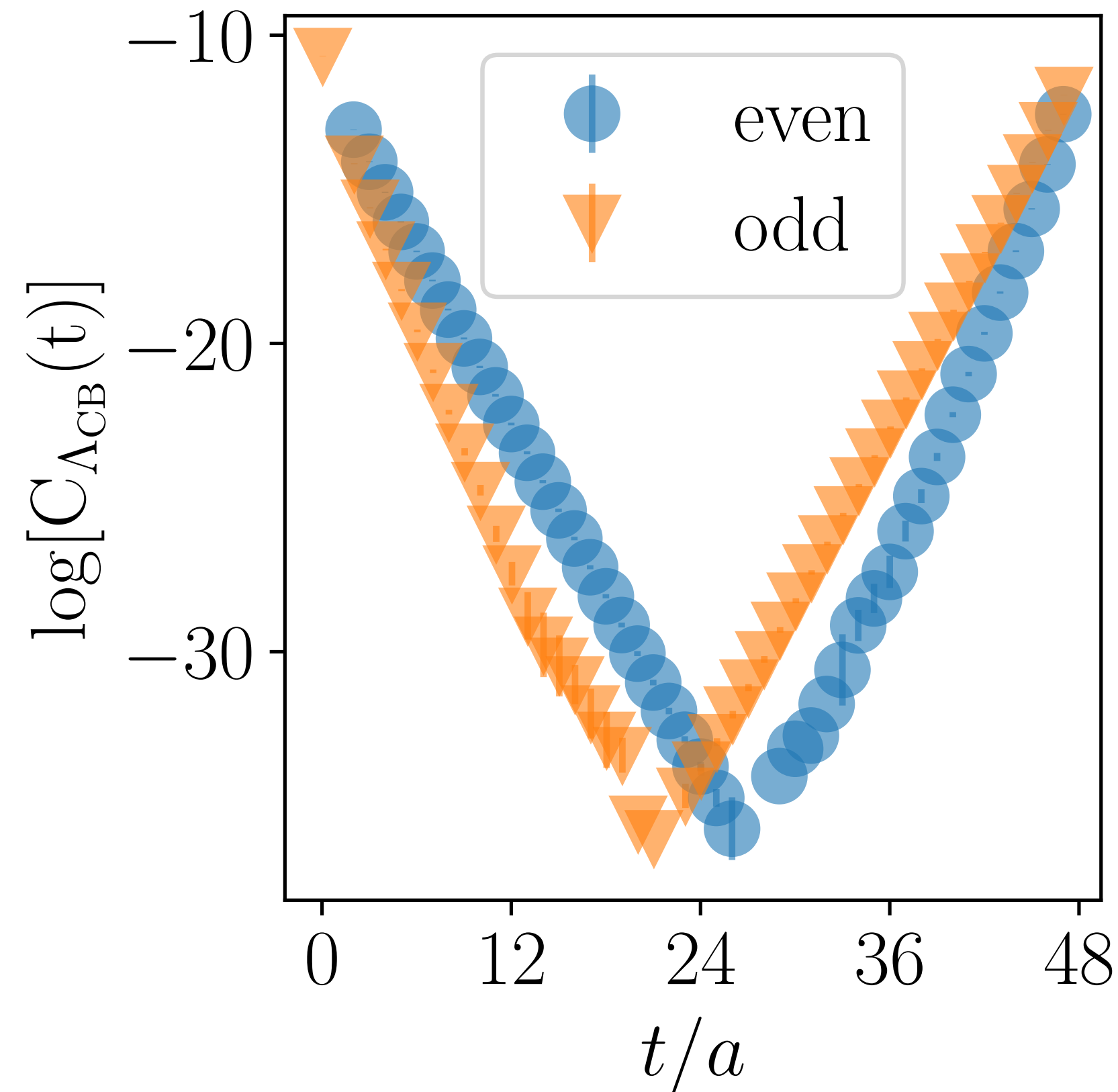
$$\times \text{Tr} \left[\Gamma^1 S_Q^b{}_{b'}(x,0) \overline{\Gamma^1} S_Q^a{}_{a'}(x,0) \right]$$

Results

Projection-Parity

► The log plot of the chimera baryon correlators (left) and their effective mass plot (right) with the parity projection.

$$C_{\text{CB}}(t) \rightarrow P_e [c_e e^{-m_e t} + c_o e^{-m_o(T-t)}] - P_o [c_o e^{-m_o t} + c_e e^{-m_e(T-t)}]$$



Chimera Baryon

- Spin projector for Σ -type baryon:

$$(P^{3/2})^{ij} = \delta^{ij} - \frac{1}{3}\gamma^i\gamma^j$$

$$(P^{1/2})^{ij} = \frac{1}{3}\gamma^i\gamma^j$$

- Two-point function

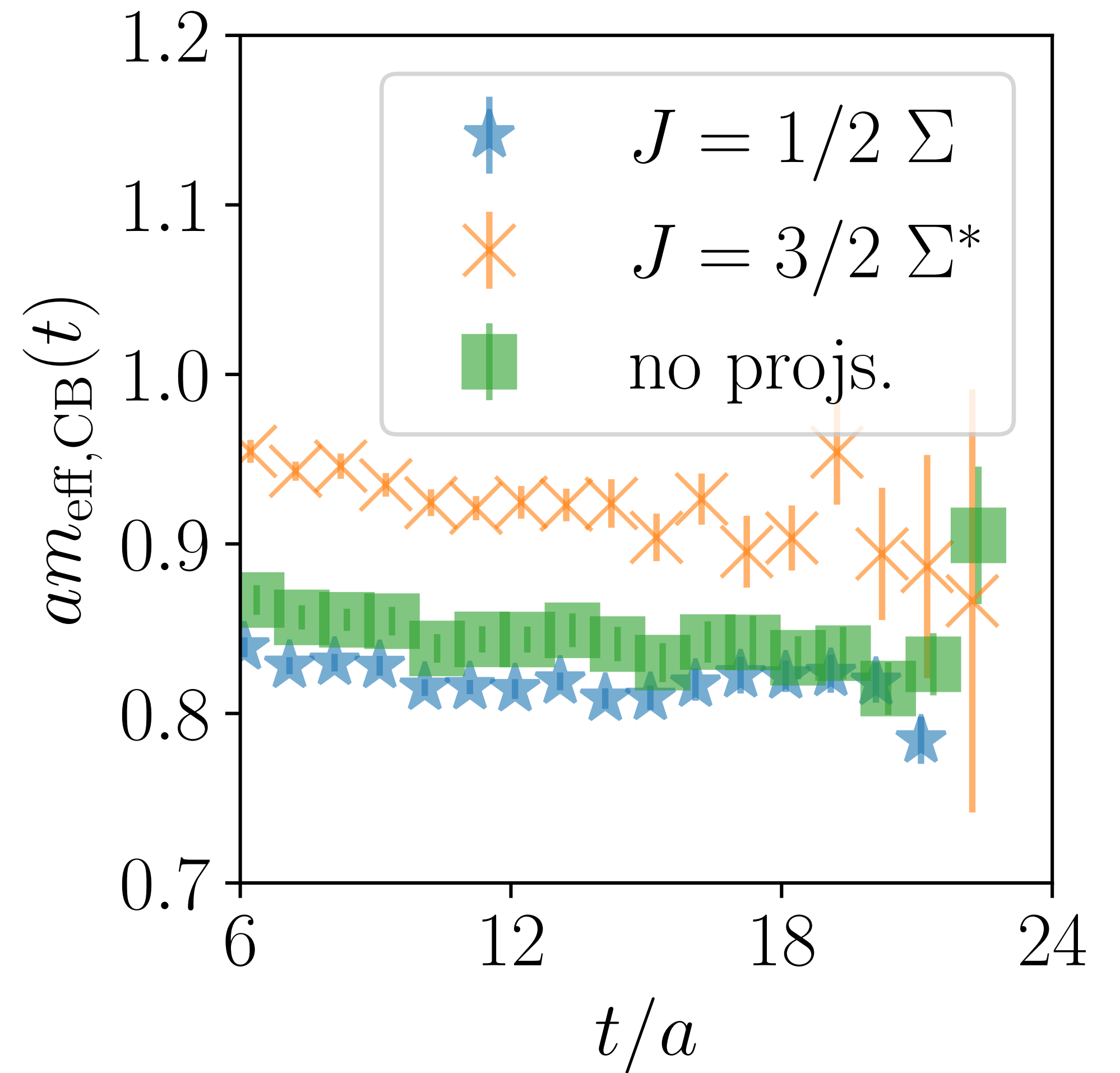
$$C_{ij}(t) = \sum_{\vec{x}} \left\langle \mathcal{O}_{\text{CB}}^i(x) \bar{\mathcal{O}}_{\text{CB}}^j(0) \right\rangle \text{ with } \mathcal{O}_{\text{CB}}^i = (\bar{\psi}\gamma^i\psi)\chi$$

$$\rightarrow C_{\Sigma}^{1/2}(t) = \text{Tr} \left[(P^{1/2})^{ij} C_{jk}(t) \right]$$

Results

Projection-Spin

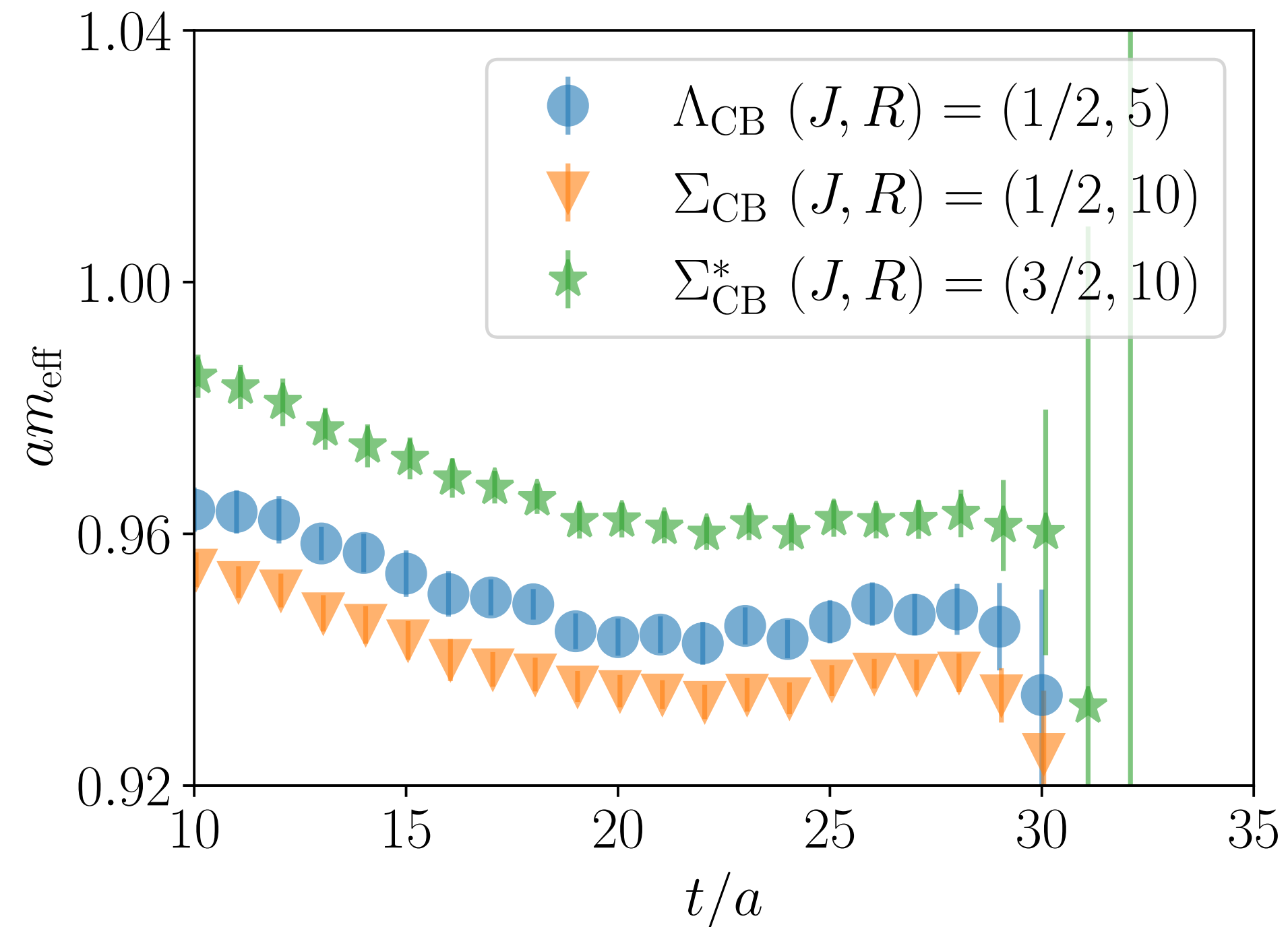
- Comparison of effective mass plot between two spin projected states and the state without spin projection.



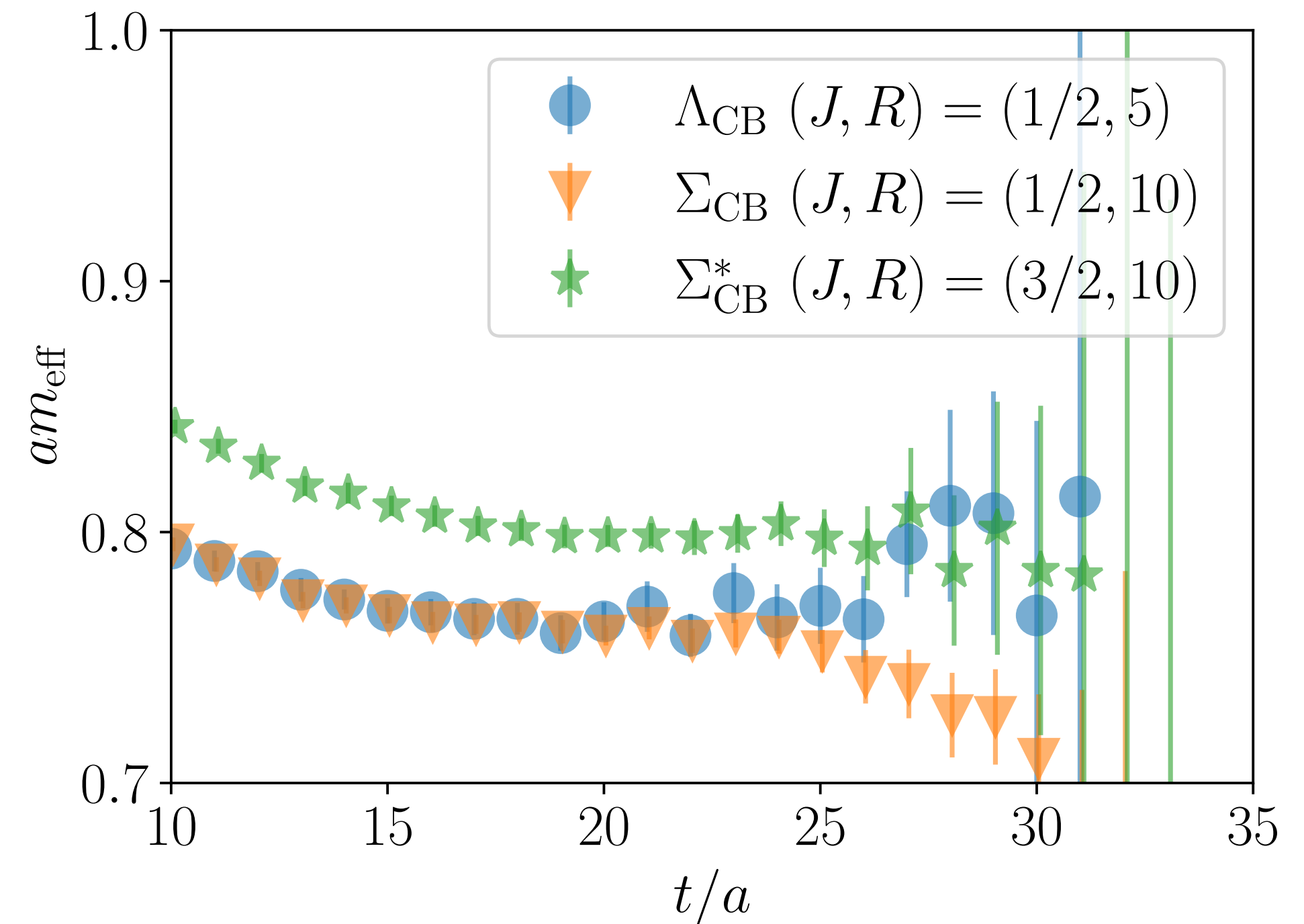
Results

Mass hierarchy

heavy F fermion mass



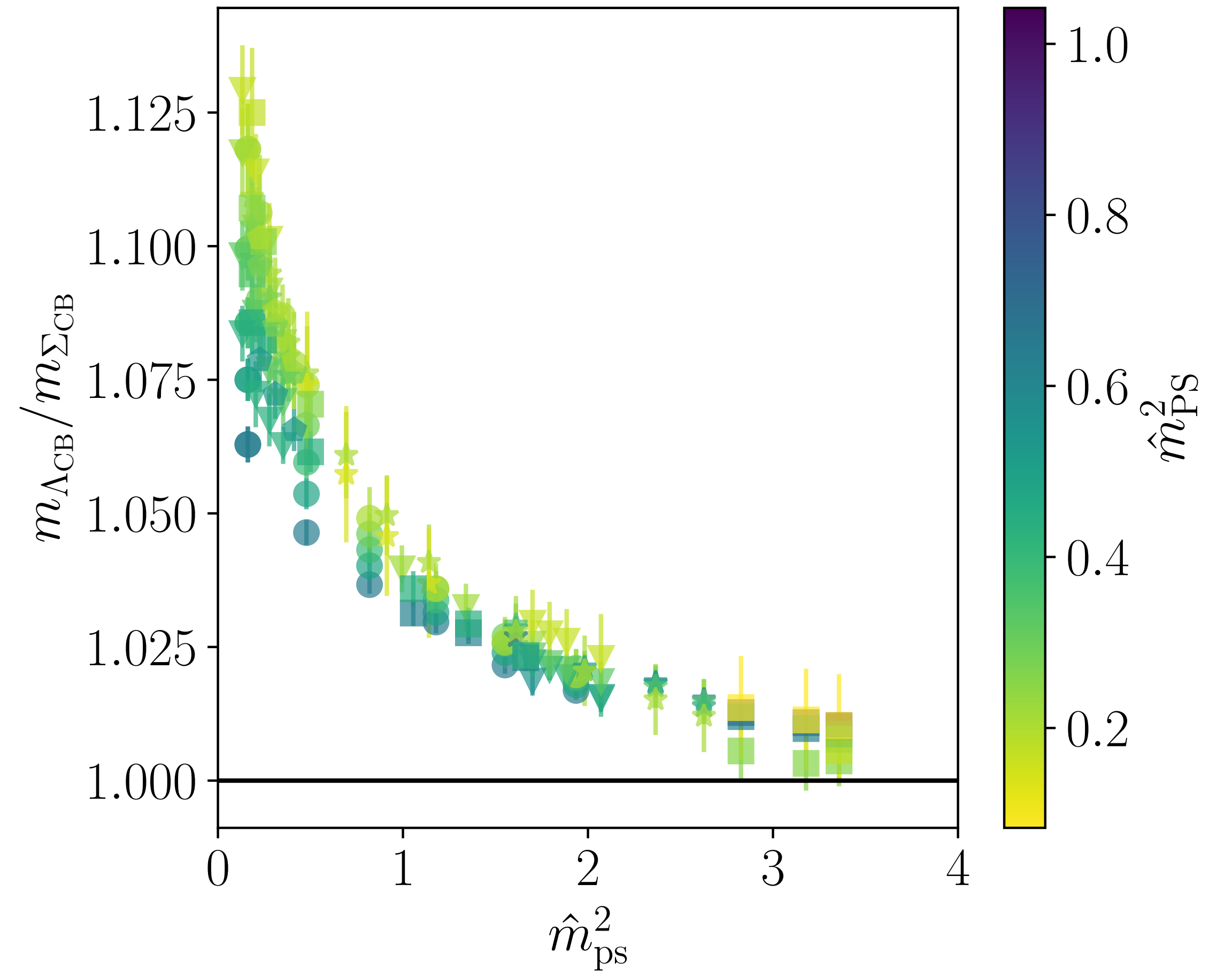
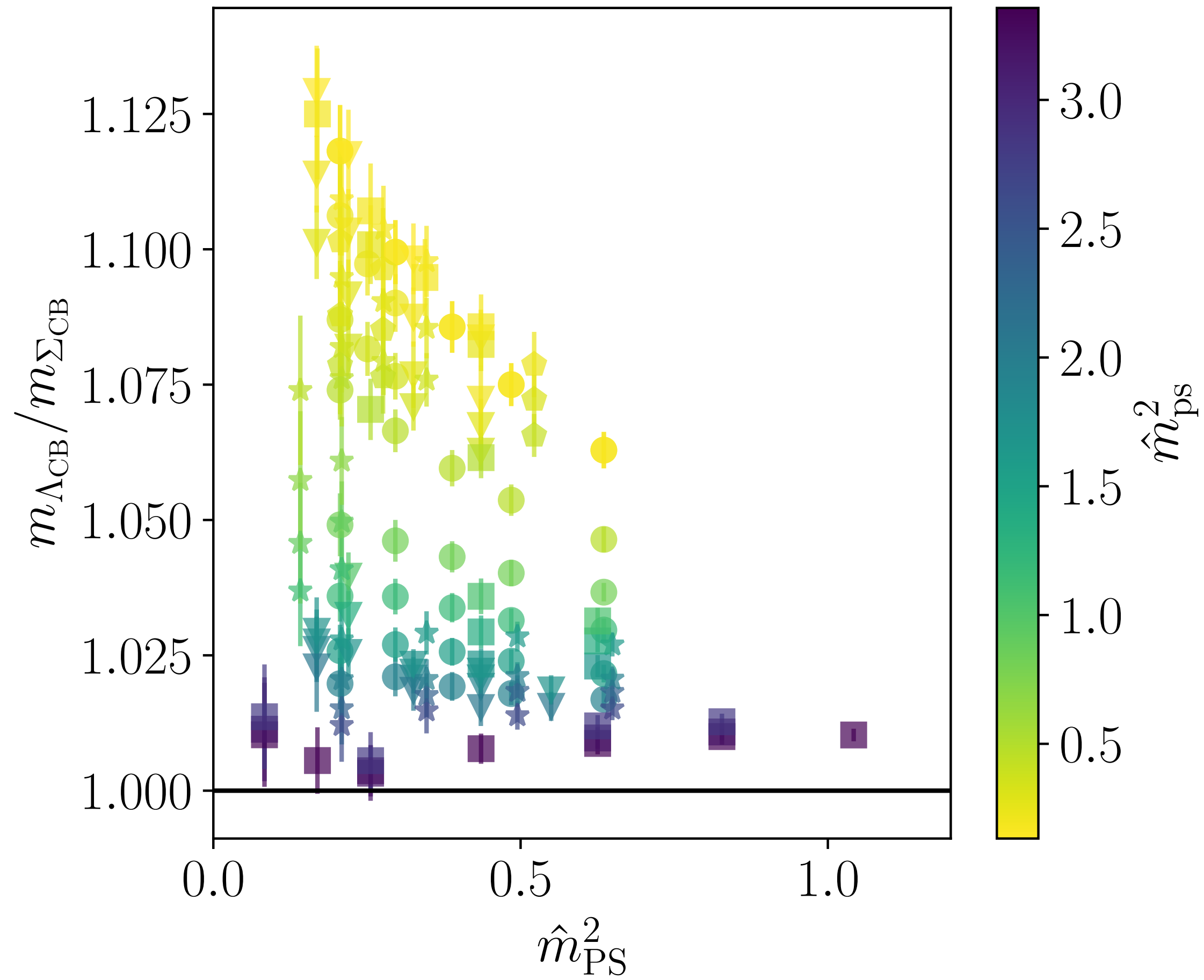
light F fermion mass



Effective mass plot of chimeron baryons calculated with different F fermion masses, at fixed AS fermion mass. The lattice size is 60×48^3 with $\beta = 8.0$.

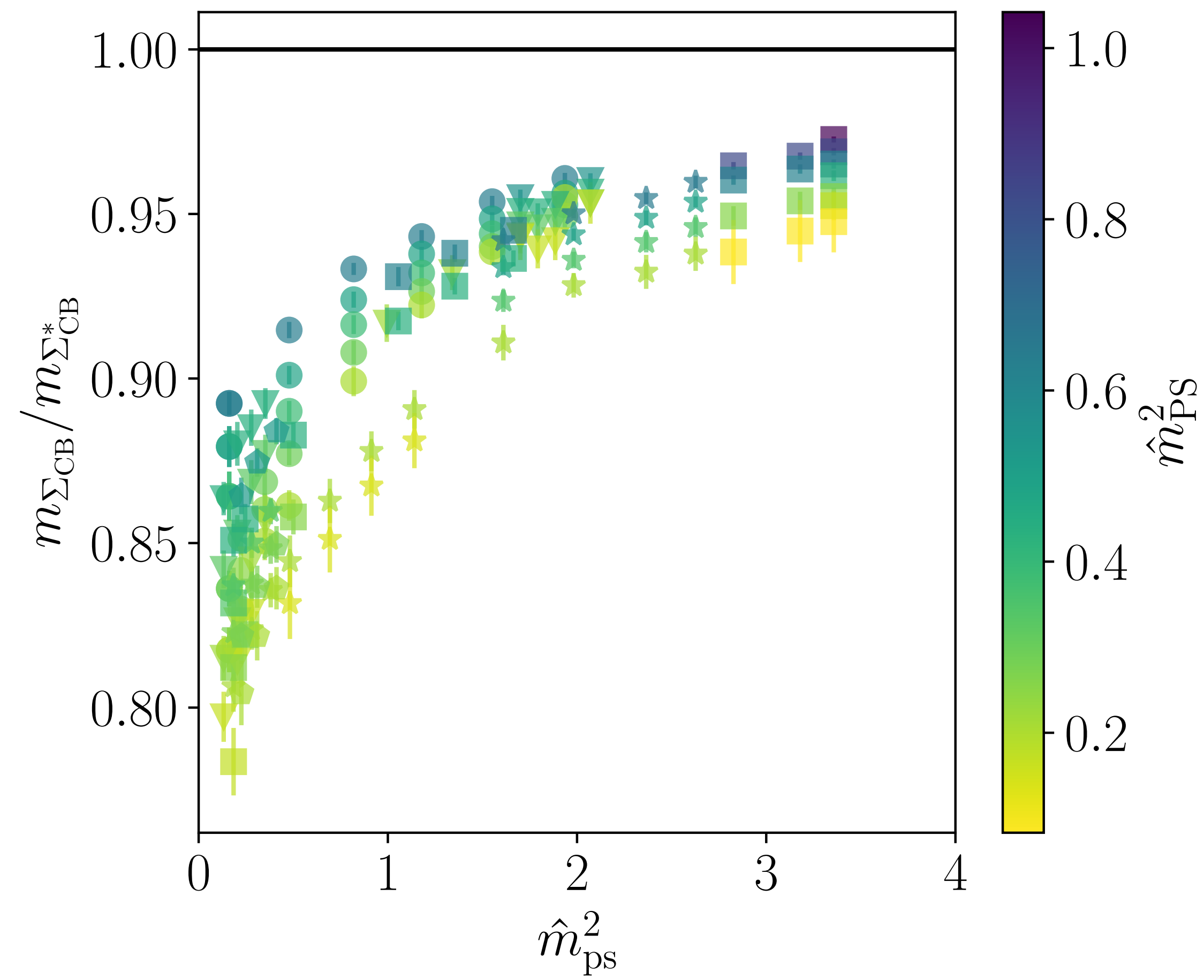
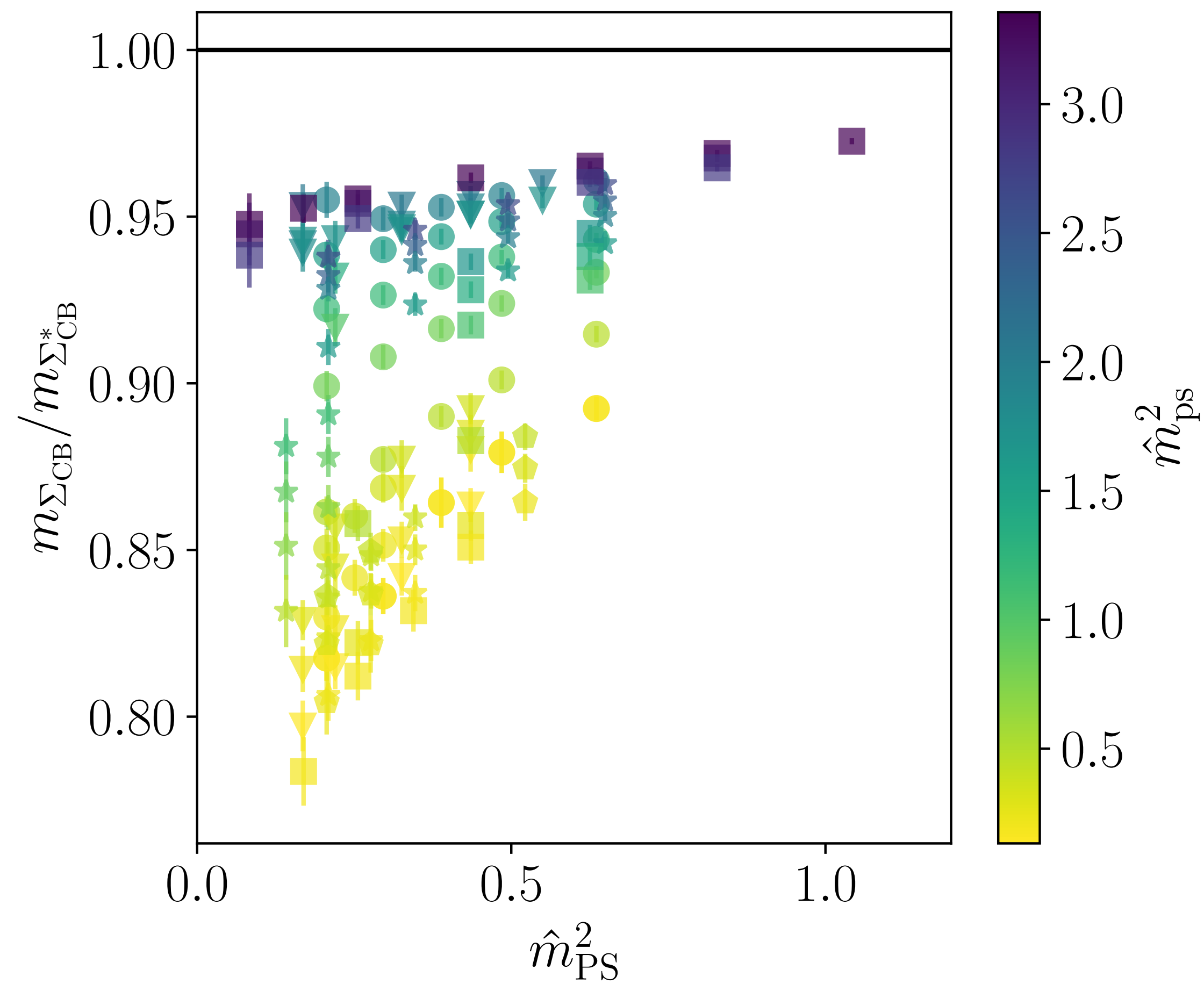
Results

Mass hierarchy



Results

Mass hierarchy



Results

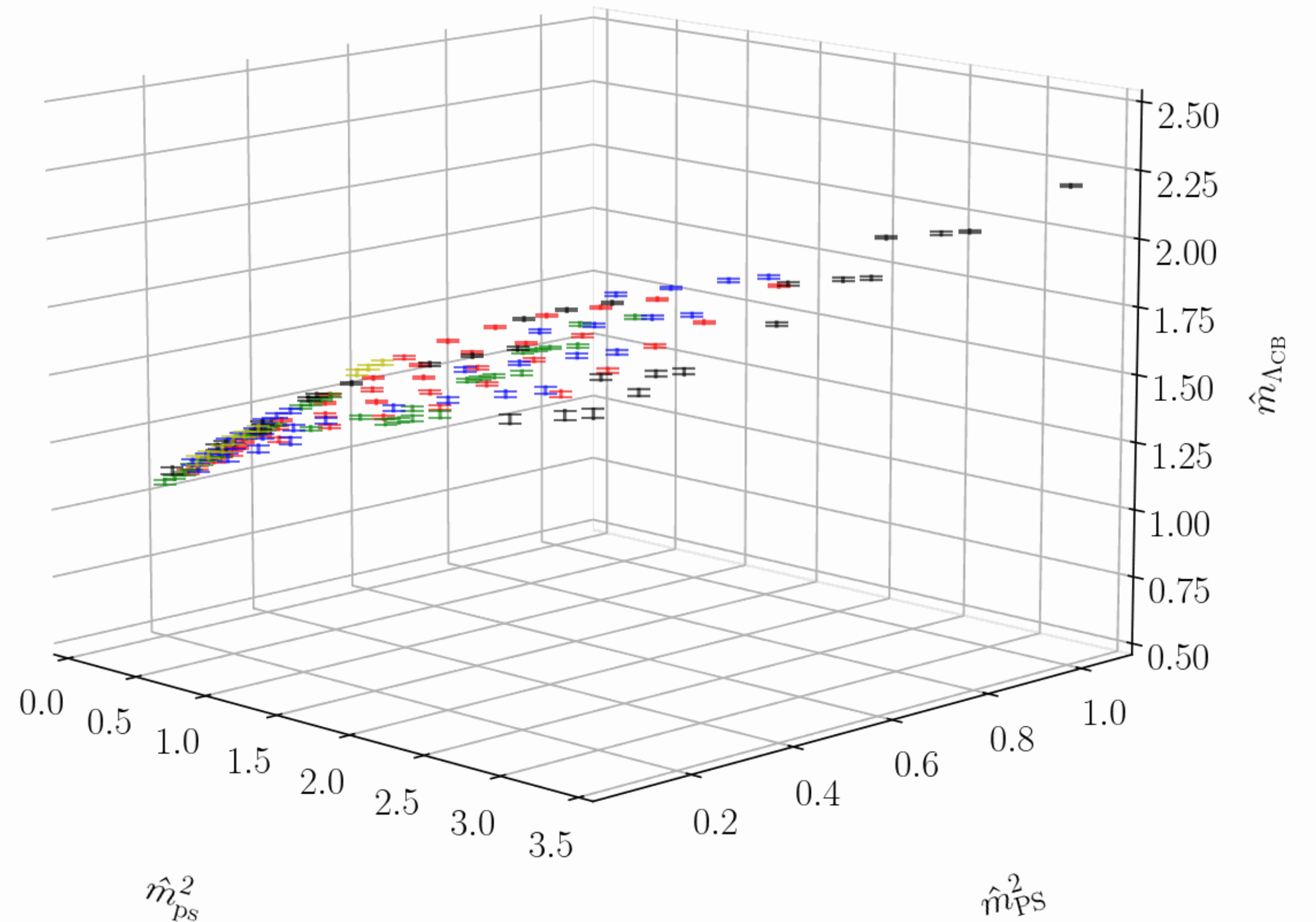
Fitting

- Apply tree-level baryon chiral perturbation theory

$$\begin{aligned} m_{\text{CB}} = & m_{\text{CB}}^{\chi} + F_2 \hat{m}_{\text{PS}}^2 + A_2 \hat{m}_{\text{ps}}^2 + L_1 \hat{a} \\ & + F_3 \hat{m}_{\text{PS}}^3 + A_3 \hat{m}_{\text{ps}}^3 + L_{2F} \hat{a} \hat{m}_{\text{PS}}^2 + L_{2A} \hat{a} \hat{m}_{\text{ps}}^2 \\ & + F_4 \hat{m}_{\text{PS}}^4 + A_4 \hat{m}_{\text{ps}}^4 + C_4 \hat{m}_{\text{PS}}^2 \hat{m}_{\text{ps}}^2 \end{aligned}$$



Returning large $\chi^2/N_{\text{d.o.f.}}$



Results

Optimal search

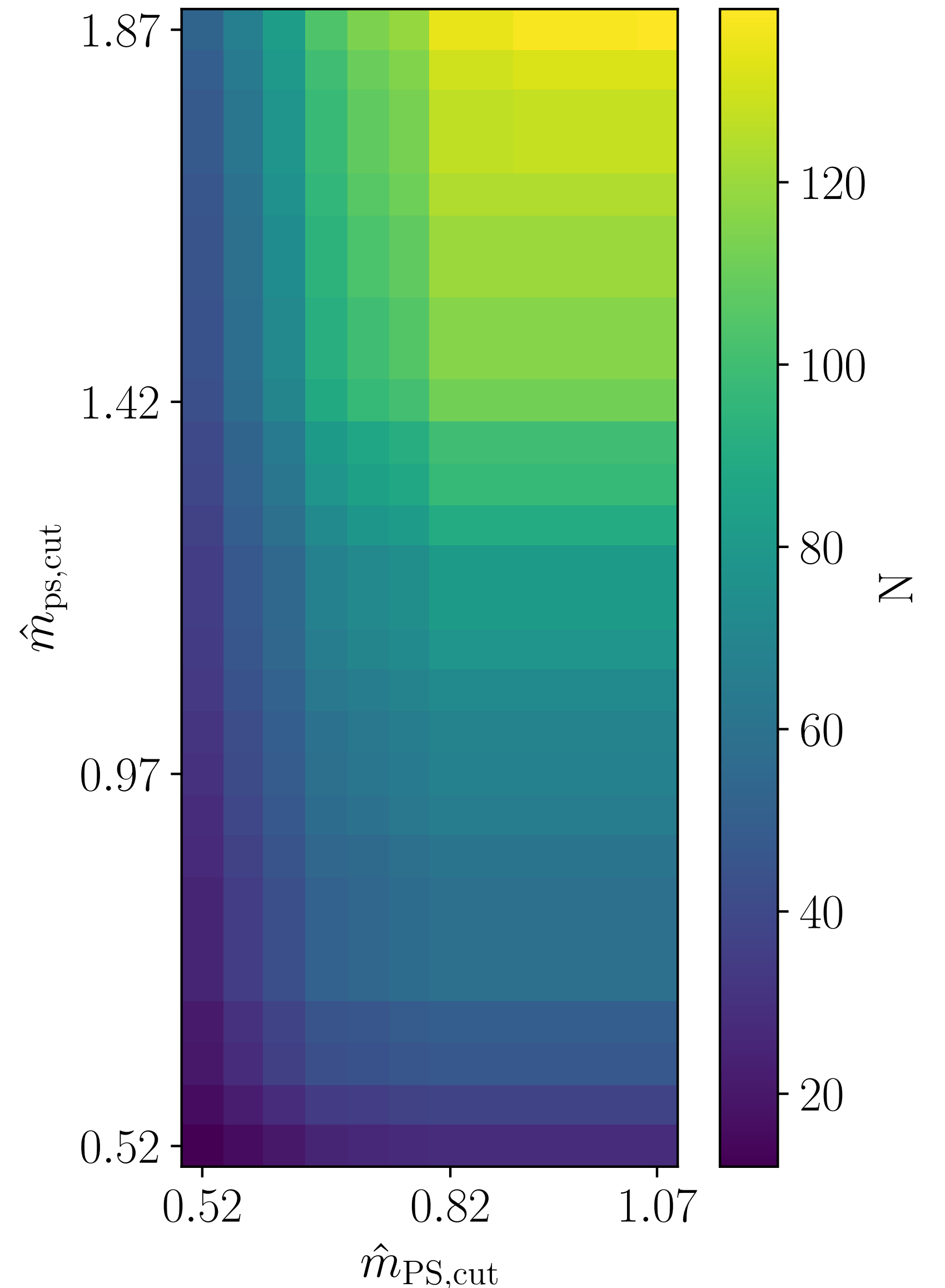
- ▶ Try including different order of corrections
- ▶ Calculate AICs for each data set, and scan through all the possible cuts:
 - ➔ Fix the cut value for \hat{m}_{PS} and vary \hat{m}_{ps}
 - ➔ Increase the fixed value of \hat{m}_{PS}
- ▶ Goodness of a fit: Akaike information criterion (AIC)

William I. Jay and Ethan T. Neil [2008.01069]

$$\text{AIC}(\text{M}, N_{\text{cut}}) \equiv \chi^2 + 2k + 2N_{\text{cut}}$$

- ▶ Probability weight

$$W(\text{M}, N_{\text{cut}}) = \frac{1}{\mathcal{N}} \exp \left[-\frac{1}{2} \text{AIC}(\text{M}, N_{\text{cut}}) \right]$$



Results

William I. Jay and Ethan T. Neil [2008.01069]

Optimal search

Try

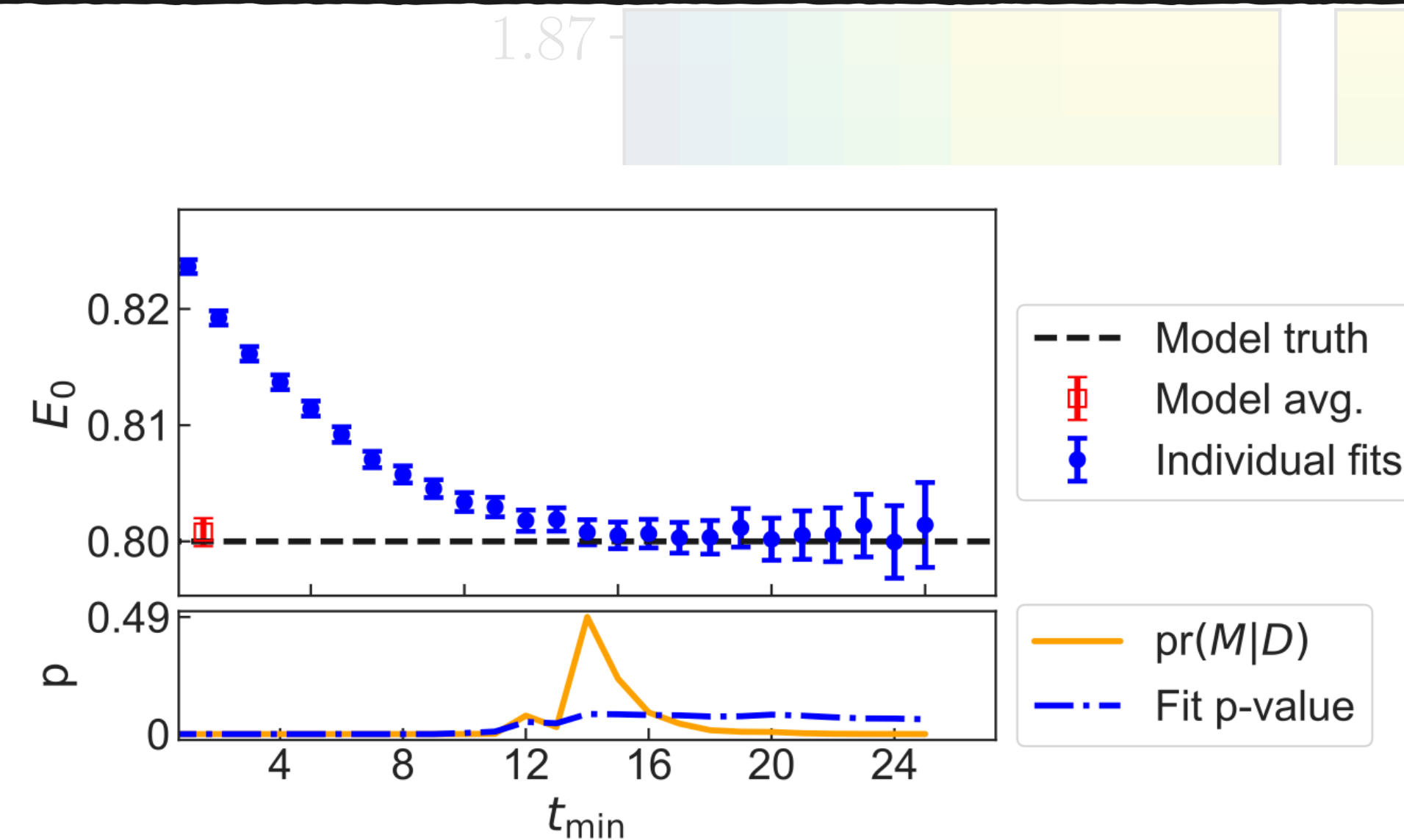
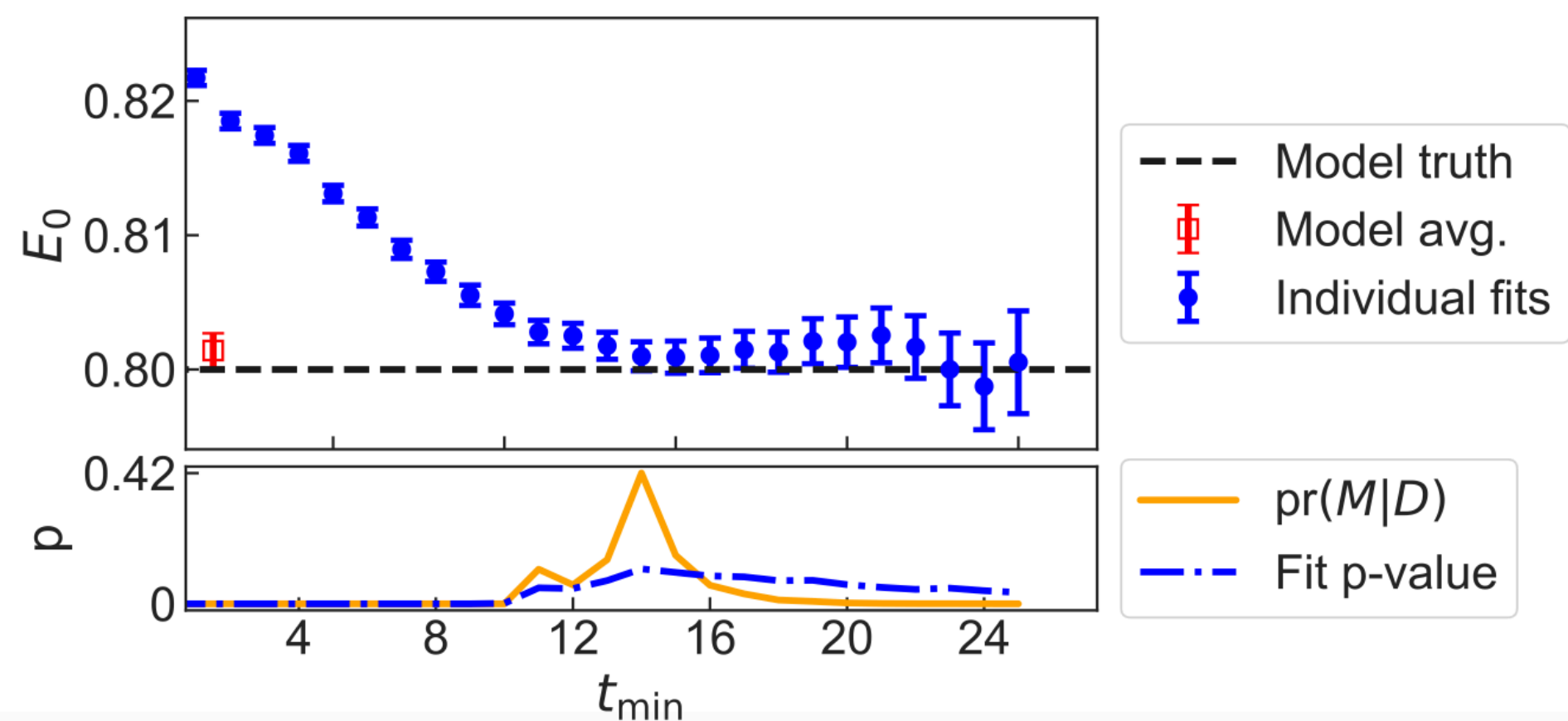
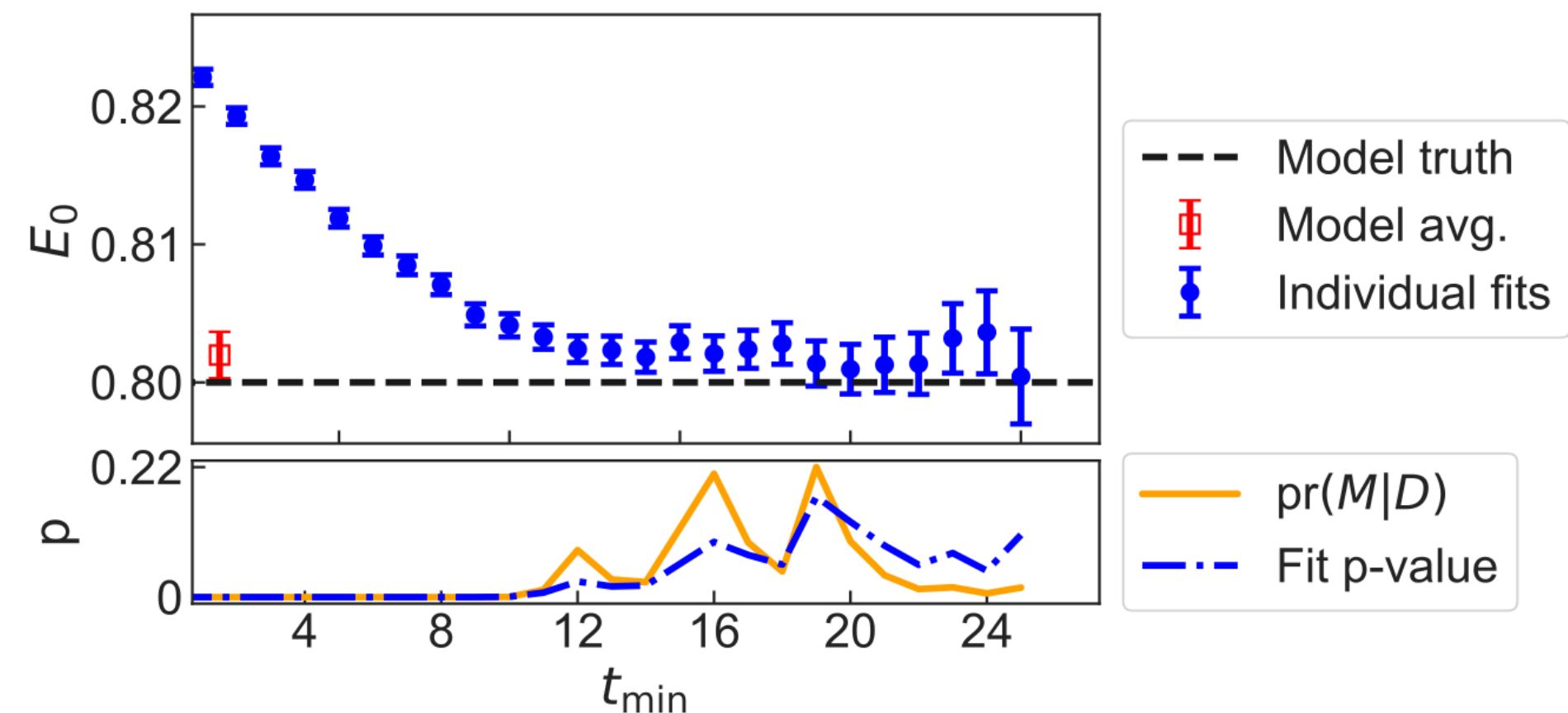
Cal
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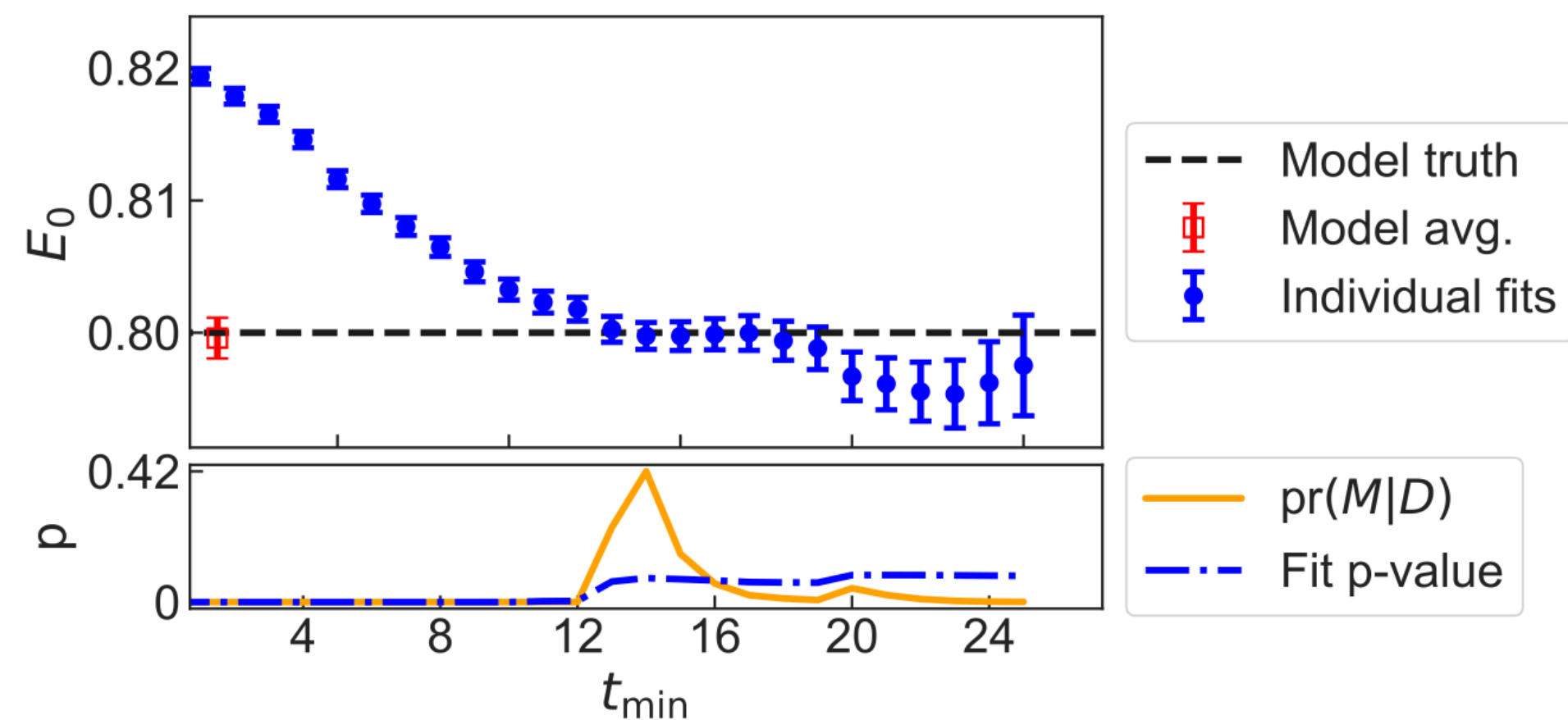
Goc

AIC

Prol



(A)



$$W(M, N_{\text{cut}}) = \frac{1}{\mathcal{N}} \exp \left[-\frac{1}{2} \text{AIC}(M, N_{\text{cut}}) \right]$$

Results

Optimal search

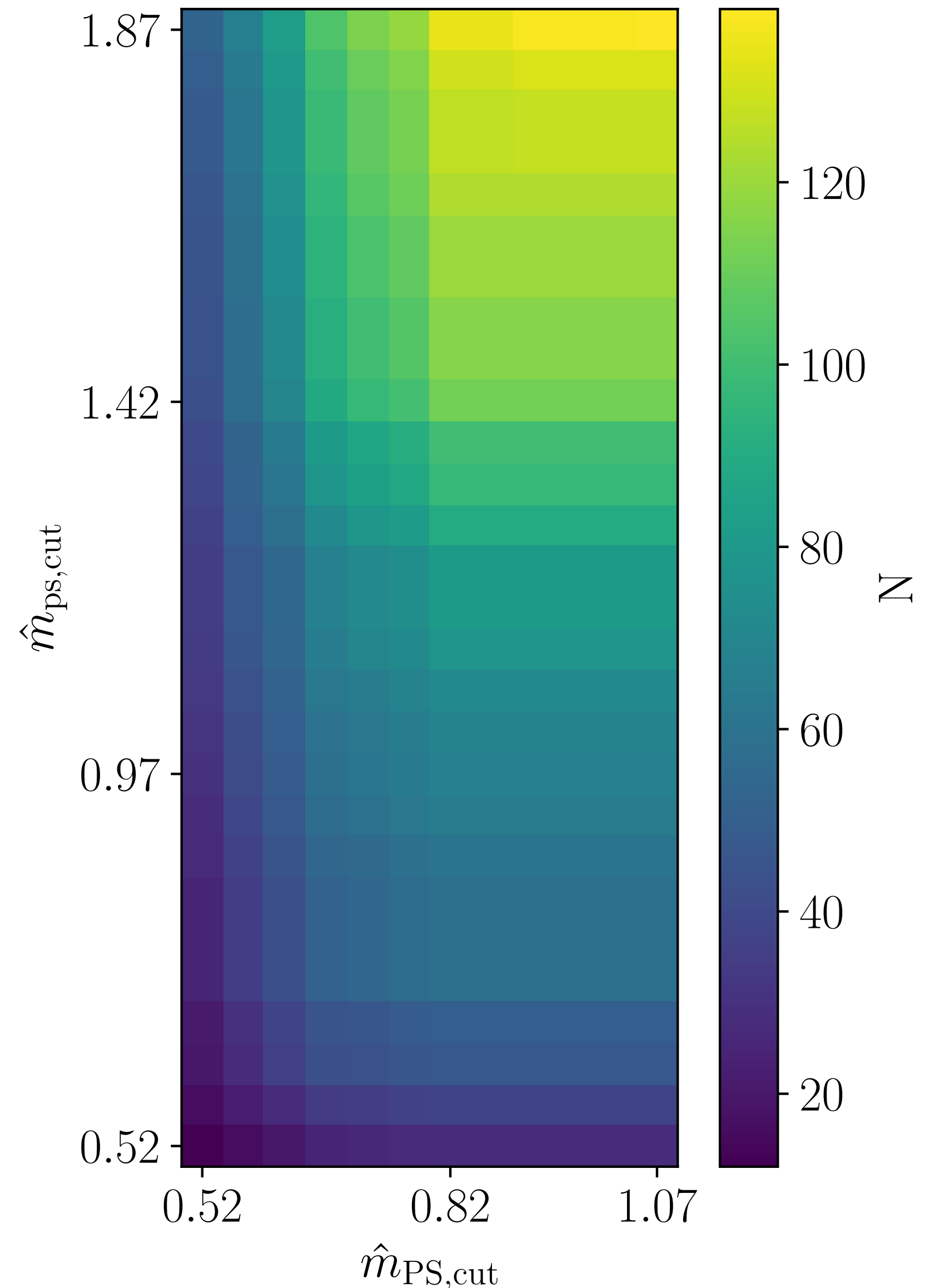
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Results

Fitting

- Apply tree-level baryon chiral perturbation theory

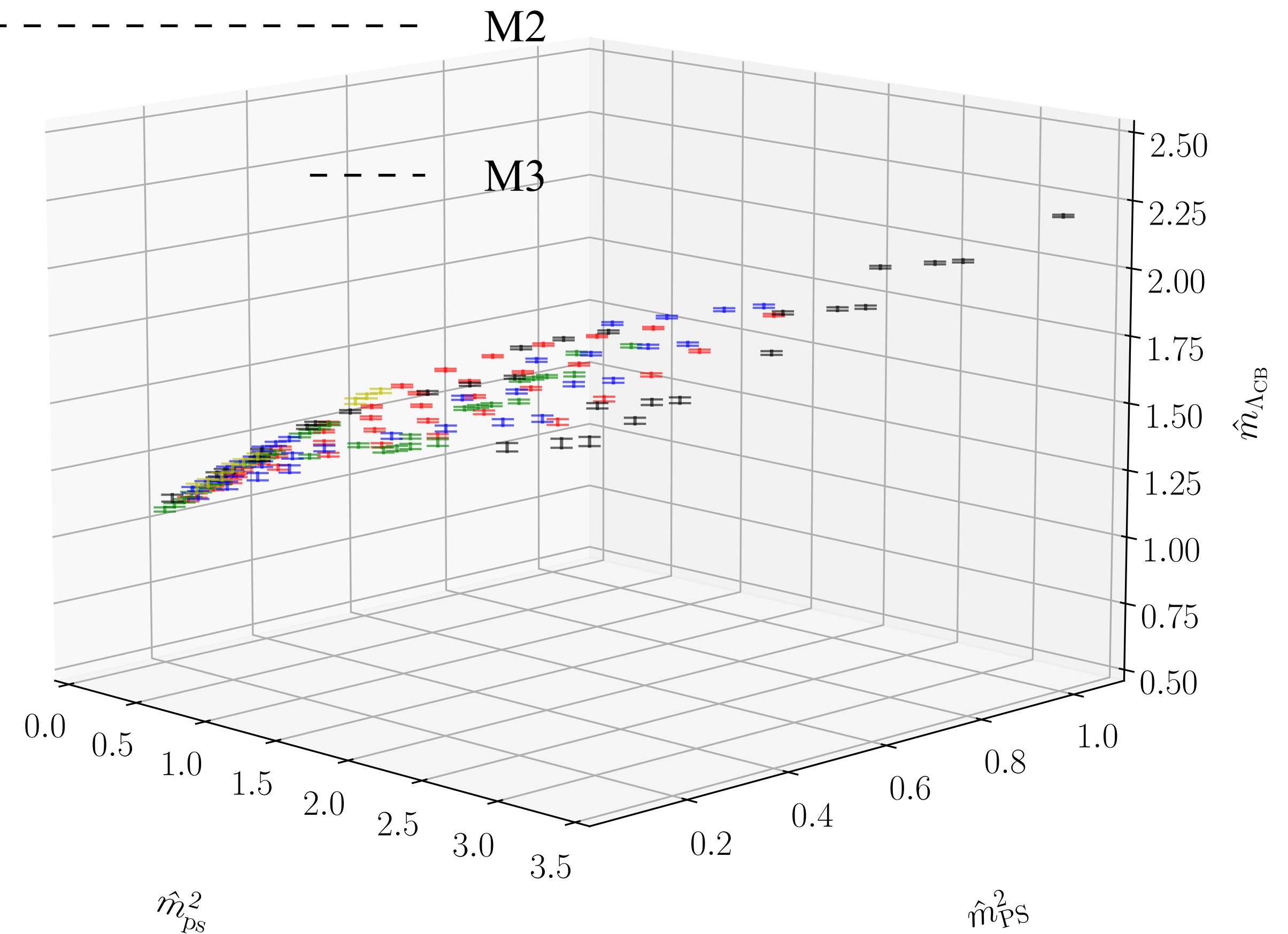
$$\begin{aligned}
 m_{\text{CB}} = & m_{\text{CB}}^{\chi} + F_2 \hat{m}_{\text{PS}}^2 + A_2 \hat{m}_{\text{ps}}^2 + L_1 \hat{a} \\
 & + F_3 \hat{m}_{\text{PS}}^3 + A_3 \hat{m}_{\text{ps}}^3 + L_{2F} \hat{a} \hat{m}_{\text{PS}}^2 + L_{2A} \hat{a} \hat{m}_{\text{ps}}^2 \\
 & + F_4 \hat{m}_{\text{PS}}^4 + A_4 \hat{m}_{\text{ps}}^4 + C_4 \hat{m}_{\text{PS}}^2 \hat{m}_{\text{ps}}^2
 \end{aligned}$$

MF4

MA4

MC4

Still returning large $\chi^2/N_{\text{d.o.f.}}$



Results

Fittings of Λ_{CB}

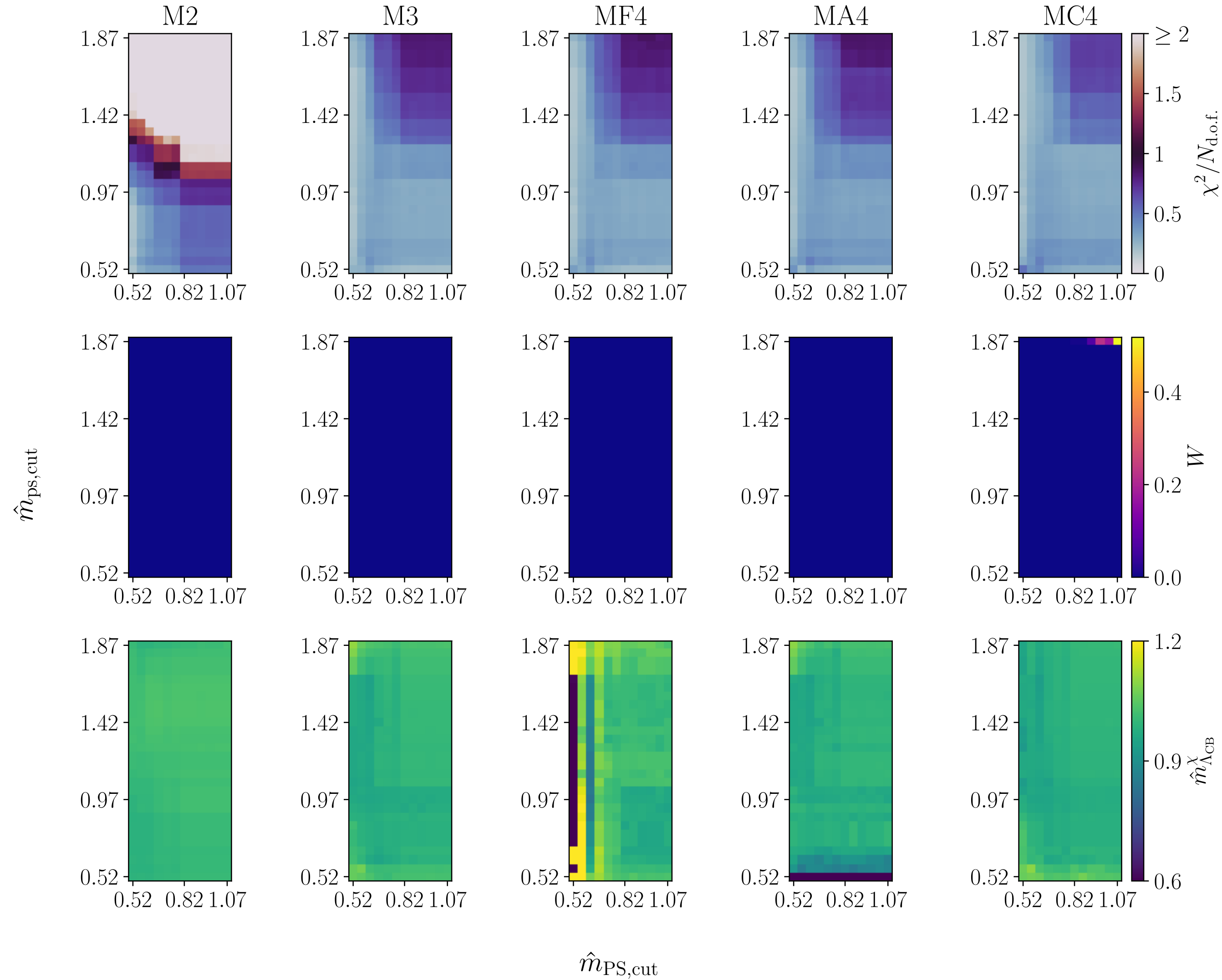
► Apply tree level baryon chiral perturbation theory

$$\begin{aligned}
 \text{M2} \quad m_{\text{CB}} &= m_{\text{CB}}^{\chi} + F_2 \hat{m}_{\text{PS}}^2 + A_2 \hat{m}_{\text{ps}}^2 + L_1 \hat{a} \\
 \text{M3} \quad &+ F_3 \hat{m}_{\text{PS}}^3 + A_3 \hat{m}_{\text{ps}}^3 + L_{2F} \hat{a} \hat{m}_{\text{PS}}^2 + L_{2A} \hat{a} \hat{m}_{\text{ps}}^2 \\
 &+ F_4 \hat{m}_{\text{PS}}^4 + A_4 \hat{m}_{\text{ps}}^4 + C_4 \hat{m}_{\text{PS}}^2 \hat{m}_{\text{ps}}^2
 \end{aligned}$$

MF4
MA4
M2C

► probability weight

$$W(\text{M}, N_{\text{cut}}) = \frac{1}{\mathcal{N}} \exp \left[-\frac{1}{2} \text{AIC}(\text{M}, N_{\text{cut}}) \right]$$



Results

Fittings of Σ_{CB}

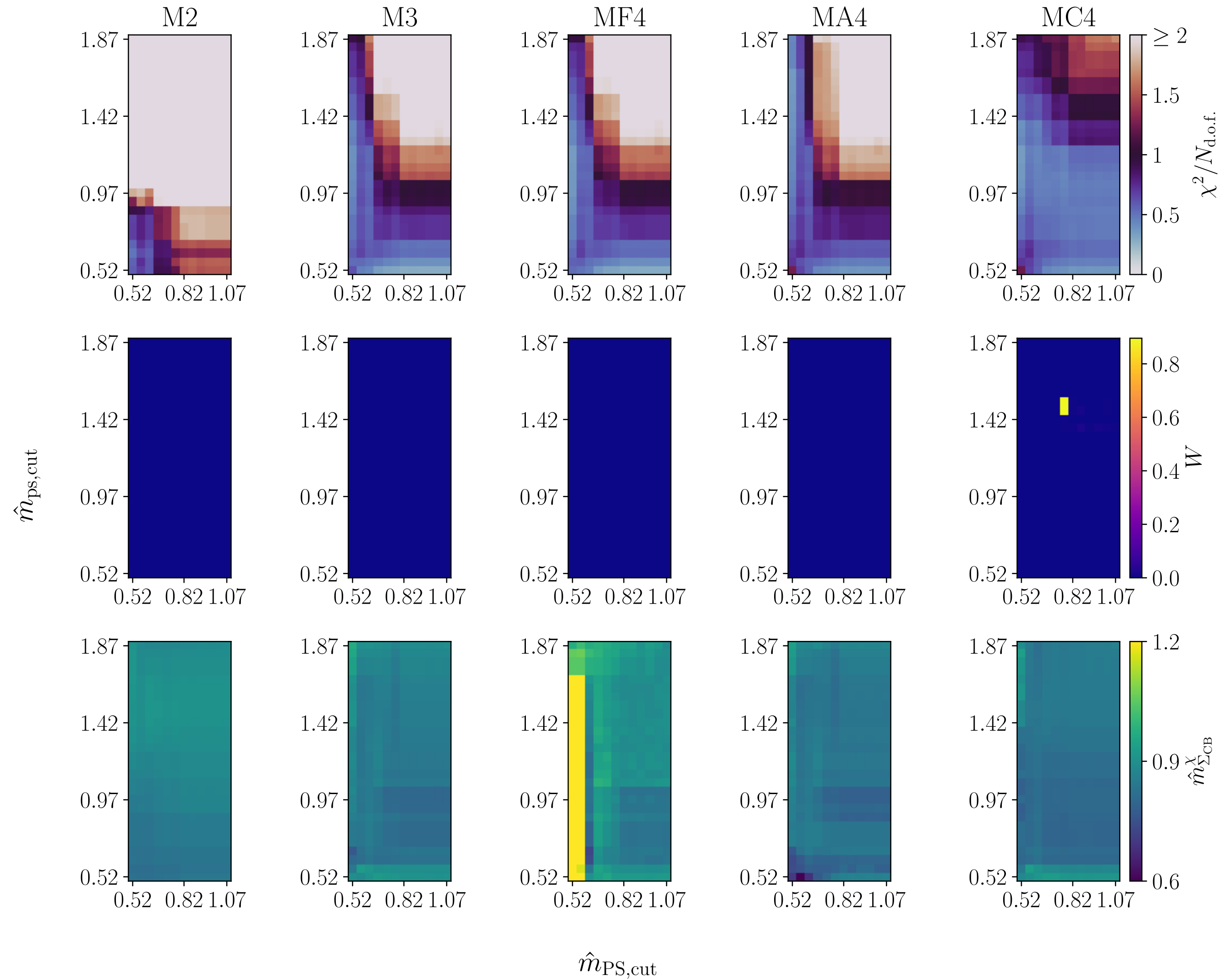
► Apply tree level baryon chiral perturbation theory

$$\begin{aligned}
 \text{M2} \quad m_{\text{CB}} &= m_{\text{CB}}^x + F_2 \hat{m}_{\text{PS}}^2 + A_2 \hat{m}_{\text{ps}}^2 + L_1 \hat{a} \\
 \text{M3} \quad &+ F_3 \hat{m}_{\text{PS}}^3 + A_3 \hat{m}_{\text{ps}}^3 + L_{2F} \hat{a} \hat{m}_{\text{PS}}^2 + L_{2A} \hat{a} \hat{m}_{\text{ps}}^2 \\
 &+ F_4 \hat{m}_{\text{PS}}^4 + A_4 \hat{m}_{\text{ps}}^4 + C_4 \hat{m}_{\text{PS}}^2 \hat{m}_{\text{ps}}^2
 \end{aligned}$$

MF4

MA4

M2C



Results

Fittings of Σ_{CB}^*

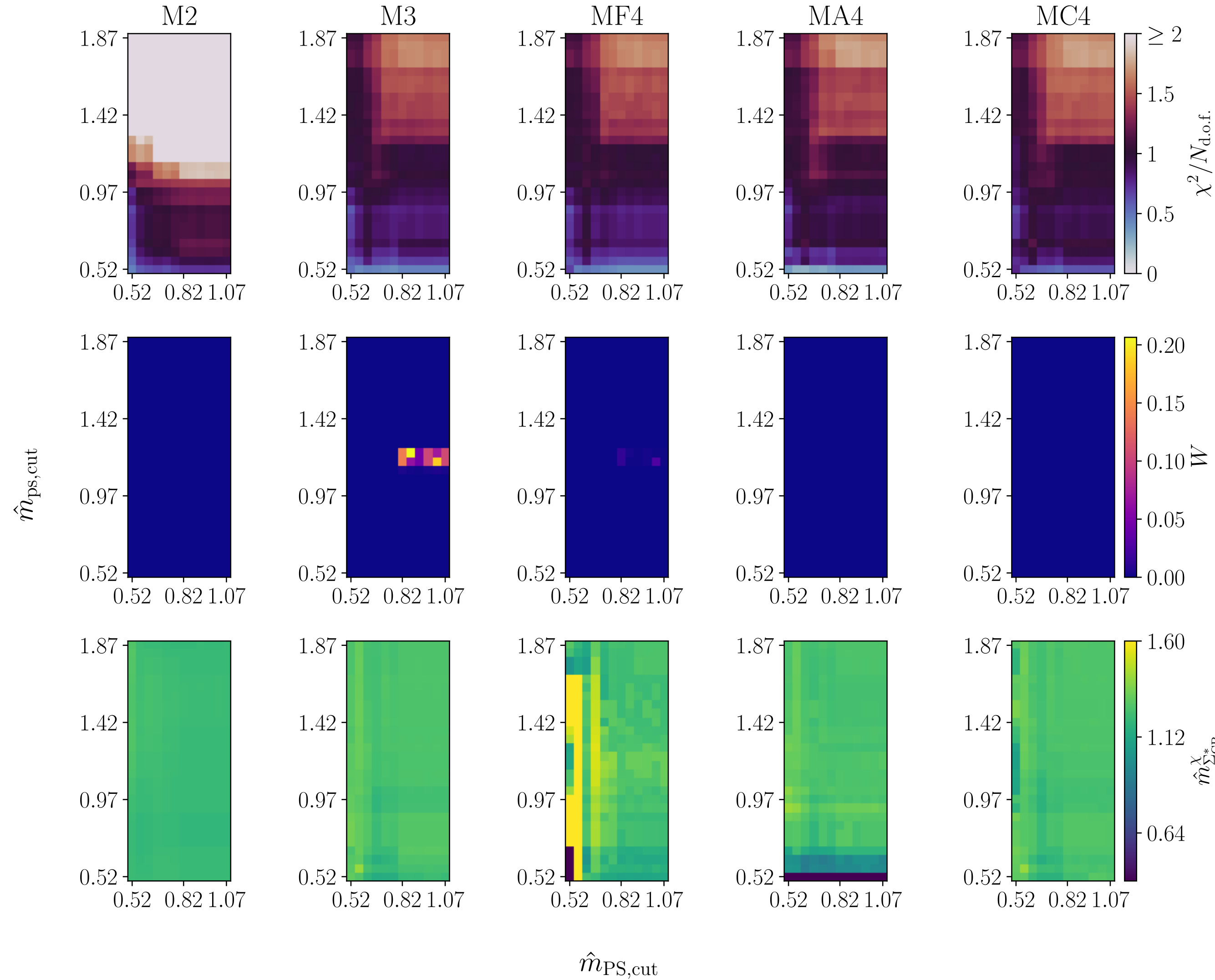
► Apply tree level baryon chiral perturbation theory

$$\begin{aligned}
 \text{M2} \quad m_{\text{CB}} &= m_{\text{CB}}^x + F_2 \hat{m}_{\text{PS}}^2 + A_2 \hat{m}_{\text{ps}}^2 + L_1 \hat{a} \\
 \text{M3} \quad &+ F_3 \hat{m}_{\text{PS}}^3 + A_3 \hat{m}_{\text{ps}}^3 + L_{2F} \hat{a} \hat{m}_{\text{PS}}^2 + L_{2A} \hat{a} \hat{m}_{\text{ps}}^2 \\
 &+ F_4 \hat{m}_{\text{PS}}^4 + A_4 \hat{m}_{\text{ps}}^4 + C_4 \hat{m}_{\text{PS}}^2 \hat{m}_{\text{ps}}^2
 \end{aligned}$$

MF4

MA4

M2C



Results

Cross check

► Apply tree-level baryon chiral perturbation theory

$$\begin{aligned} m_{\text{CB}} = & m_{\text{CB}}^{\chi} + F_2 \hat{m}_{\text{PS}}^2 + A_2 \hat{m}_{\text{ps}}^2 + L_1 \hat{a} \\ & + F_3 \hat{m}_{\text{PS}}^3 + A_3 \hat{m}_{\text{ps}}^3 + L_{2F} \hat{a} \hat{m}_{\text{PS}}^2 + L_{2A} \hat{a} \hat{m}_{\text{ps}}^2 \\ & + F_4 \hat{m}_{\text{PS}}^4 + A_4 \hat{m}_{\text{ps}}^4 + C_4 \hat{m}_{\text{PS}}^2 \hat{m}_{\text{ps}}^2 \end{aligned}$$

Results

Cross check

► At a fixed \hat{m}_{PS}^{as} , the fitting function becomes

$$\begin{aligned} m_{\text{CB}} = & m_{\text{CB}}^{\chi} + F_2 \hat{m}_{\text{PS}}^2 + A_2 \hat{m}_{\text{ps}}^2 + L_1 \hat{a} \\ & + F_3 \hat{m}_{\text{PS}}^3 + A_3 \hat{m}_{\text{ps}}^3 + L_{2F} \hat{a} \hat{m}_{\text{PS}}^2 + L_{2A} \hat{a} \hat{m}_{\text{ps}}^2 \\ & + F_4 \hat{m}_{\text{PS}}^4 + A_4 \hat{m}_{\text{ps}}^4 + C_4 \hat{m}_{\text{PS}}^2 \hat{m}_{\text{ps}}^2 \end{aligned}$$

Results

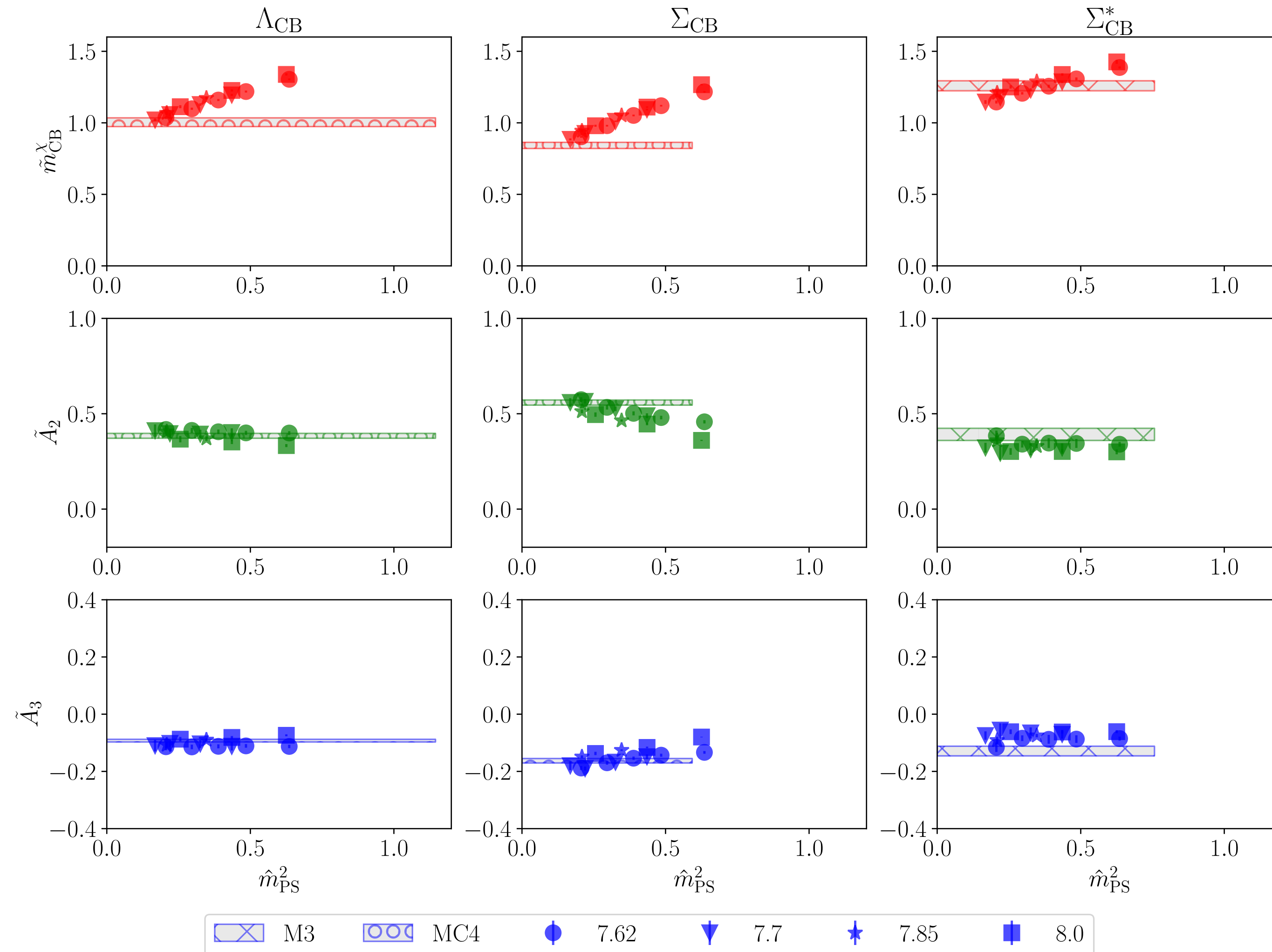
Cross check

► At a fixed $\hat{m}_{\text{PS}}^{\text{as}}$, the fitting function becomes

$$\begin{aligned} m_{\text{CB}} &= m_{\text{CB}}^{\chi} + A_2 \hat{m}_{\text{PS}}^{\text{as}^2} + L_1 \hat{a} + A_3 \hat{m}_{\text{PS}}^{\text{as}^3} + L_{2A} \hat{a} \hat{m}_{\text{PS}}^{\text{as}^2} + A_4 \hat{m}_{\text{PS}}^{\text{as}^4} \\ &\quad + F_2 \hat{m}_{\text{PS}}^{\text{f}^2} + C_4 \hat{m}_{\text{PS}}^{\text{f}^2} \hat{m}_{\text{PS}}^{\text{as}^2} + L_{2F} \hat{a} \hat{m}_{\text{PS}}^{\text{f}^2} \\ &\quad + F_3 \hat{m}_{\text{PS}}^{\text{f}^3} + F_4 \hat{m}_{\text{PS}}^{\text{f}^4} \\ &\Rightarrow \tilde{m}_{\text{CB}}^{\chi}(\hat{m}_{\text{ps}}, A, L, \hat{a}) + \tilde{F}_2(\hat{m}_{\text{ps}}, C, L, \hat{a}) \hat{m}_{\text{PS}}^2 + \tilde{F}_3 \hat{m}_{\text{PS}}^3 + F_4 \hat{m}_{\text{PS}}^{\text{f}^4} \end{aligned}$$

Results

Cross check

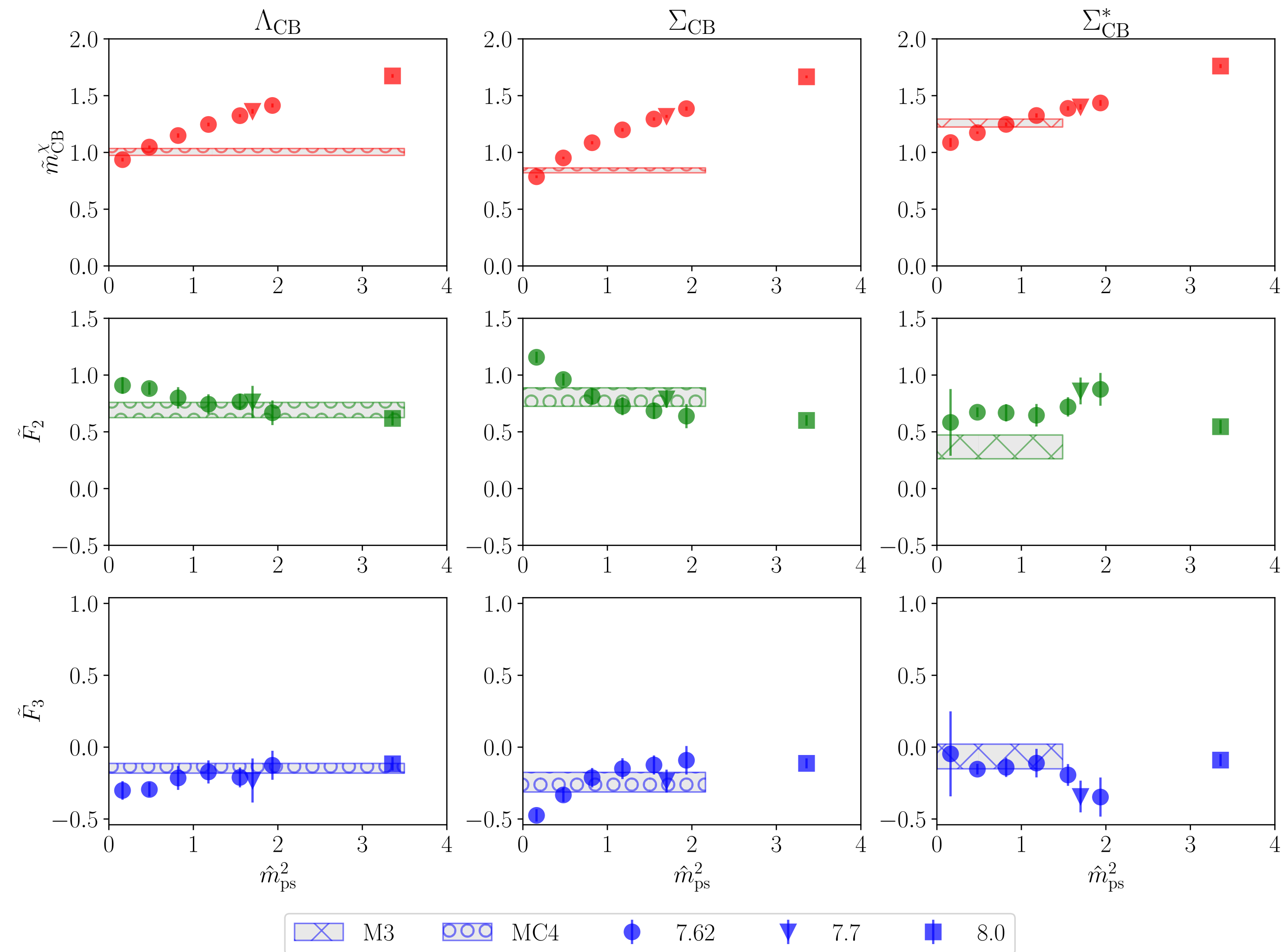


► At a fixed \hat{m}_{ps}

$$\begin{aligned}
 m_{\text{CB}} = & \tilde{m}_{\text{CB}}^x(\hat{m}_{\text{ps}}, A, L, \hat{a}) \\
 & + \tilde{F}_2(\hat{m}_{\text{ps}}, C, L, \hat{a}) \hat{m}_{\text{PS}}^2 + \tilde{F}_3 \hat{m}_{\text{PS}}^3
 \end{aligned}$$

Results

Cross check



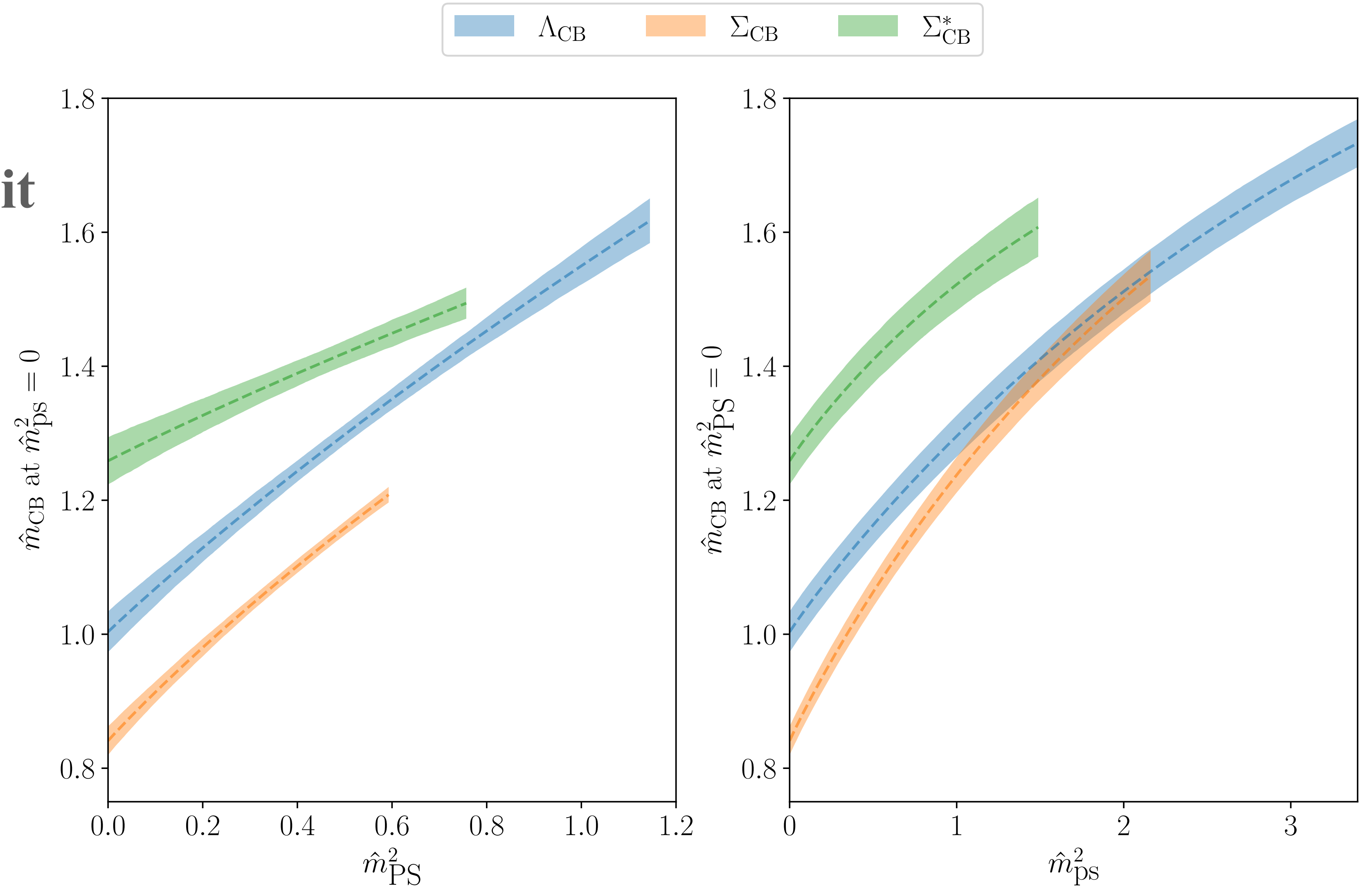
► At a fixed \hat{m}_{PS}

$$m_{\text{CB}} = \tilde{m}_{\text{CB}}^x(\hat{m}_{\text{PS}}, F, L, \hat{a})$$

$$+ \tilde{A}_2(\hat{m}_{\text{PS}}, C, L, \hat{a}) \hat{m}_{\text{ps}}^2 + \tilde{A}_3 \hat{m}_{\text{ps}}^3$$

Results

Massless-continuum limit

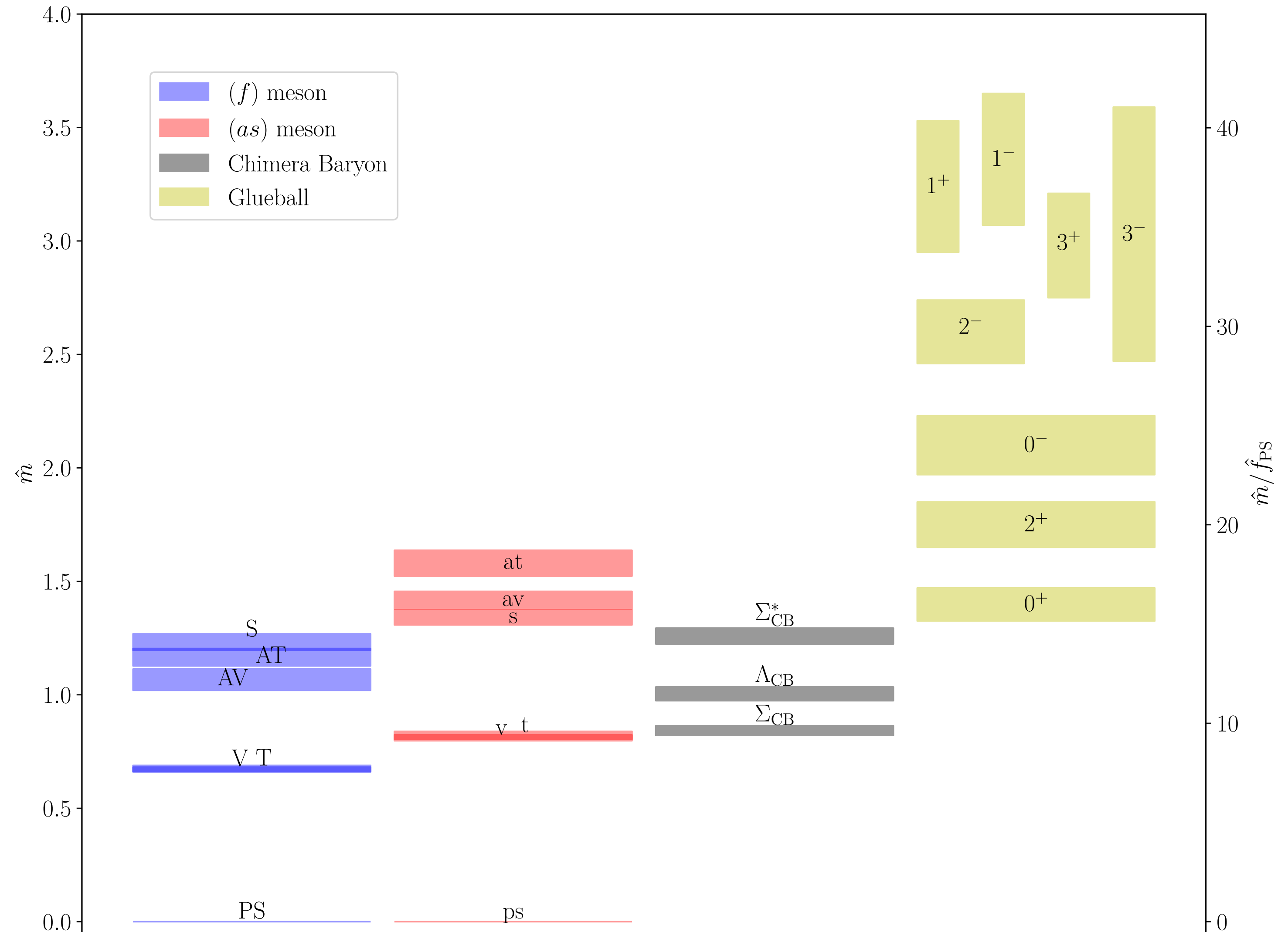


CB	Ansatz	\hat{m}_{CB}^x	F_2	A_2	L_1	F_3	A_3	L_{2F}	L_{2A}	C_4
Λ_{CB}	MC4	1.004(30)	0.692(67)	0.384(12)	-0.14(46)	-0.14(33)	-0.092(46)	0.091(76)	0.003(13)	-0.024(60)
Σ_{CB}	MC4	0.842(21)	0.806(81)	0.558(13)	-0.14(33)	-0.24(68)	-0.162(77)	0.193(62)	-0.01(16)	-0.079(62)
Σ_{CB}^*	M3	1.258(35)	0.36(10)	0.391(31)	-0.33(53)	-0.06(85)	-0.12(16)	0.335(86)	0.006(30)	-

Results

Massless-continuum limit

Comparison with masses of mesons in quenched approximation for fermions in the fundamental (blue bands) and antisymmetric (red bands) representation of $Sp(4)$, and glueballs (yellow at massless-continuum limit).



Summary and Outlook

☑ Chimera baryons

- Λ and Σ : Top partner candidates in our model
- Σ^* with spin-3/2

☑ Projection (Spin and Parity)

☑ The mass hierarchy of chimera baryons

☑ Chiral effective field theory

END
Thank you