

FASER実験の最新結果とハドロン 相互作用研究の今後の展望

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FASER Collaboration

ForwArd Search ExpeRiment (FASER) 実験

ATLAS衝突点から480m 前方に検出器を設置
Run3 2022-2025の期間にデータ取得

主な物理目標

Dark photon 探査

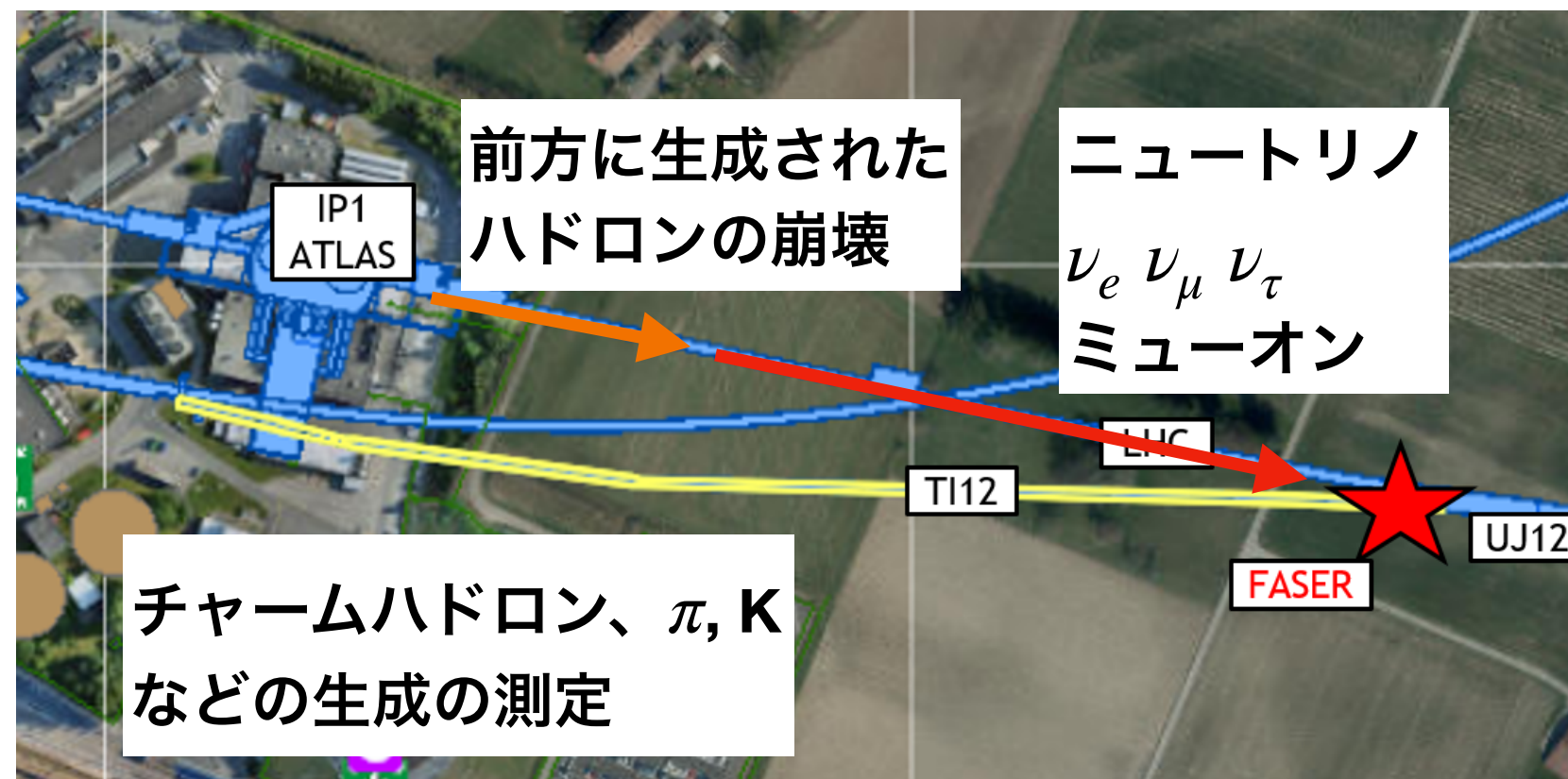
ニュートリノ測定 ($\eta > 8.8$)

三世代ニュートリノのタングステン中での
cross-section測定

Run3で想定されるイベント数:

1128 (ν_e), 5346 (ν_μ), 21.6 (ν_τ) (150fb^{-1})

超前方ハドロン生成研究

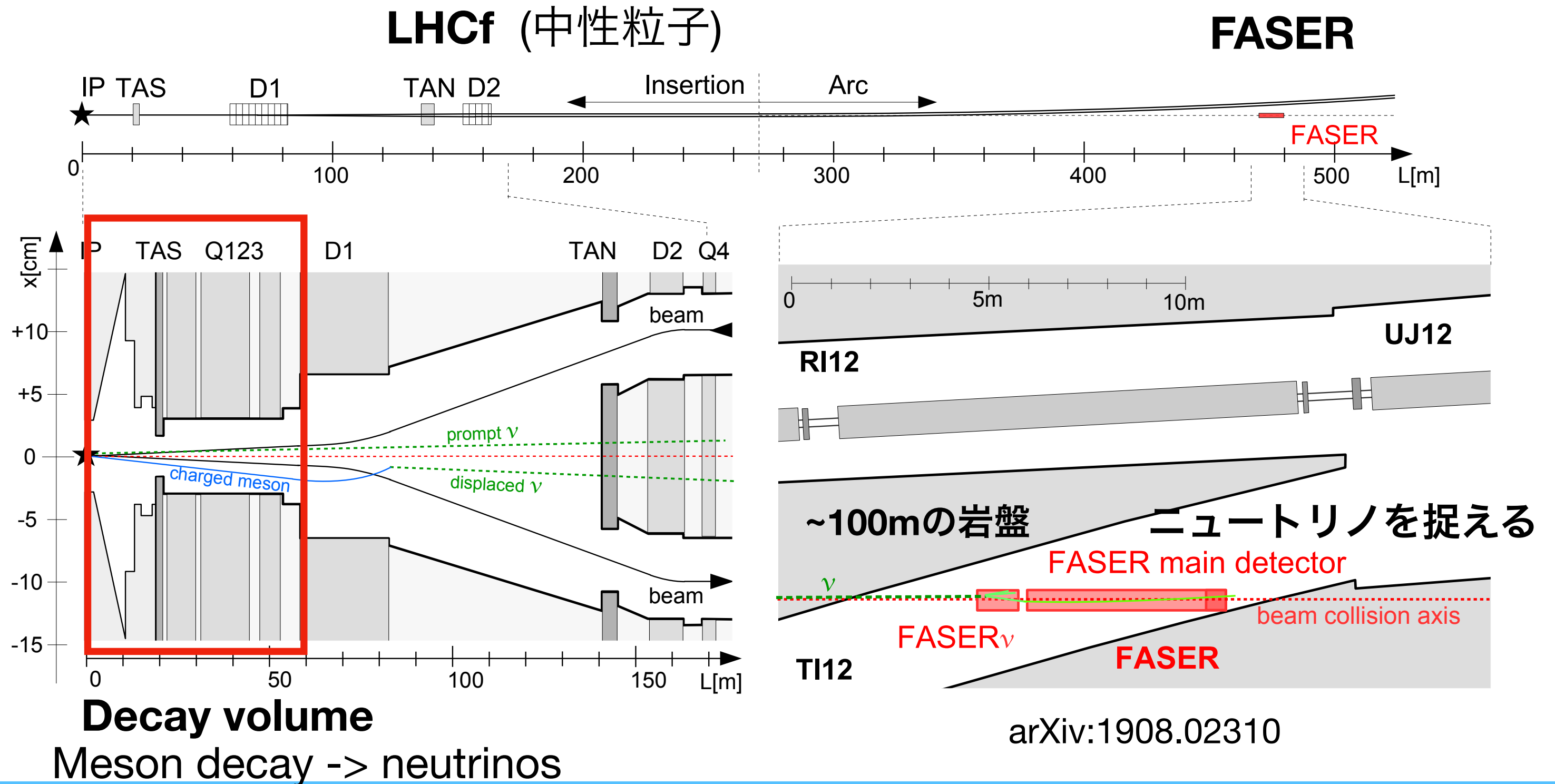


磁場 -> 運動量測定

Veto シンチレータ



荷電粒子の崩壊により生じるニュートリノを測定

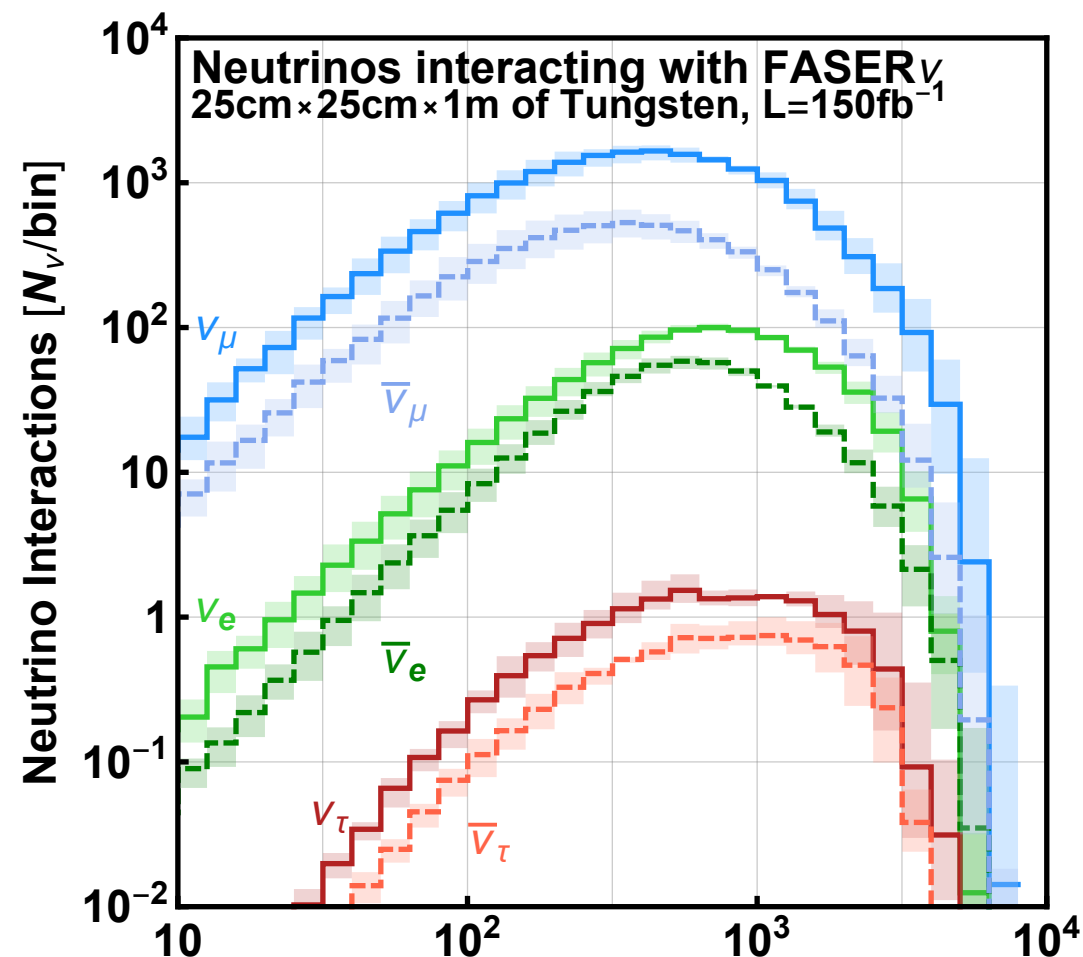


π - K の割合をニュートリノの割合から求める

電子ニュートリノは主に K から、ミューニュートリノは主に π から

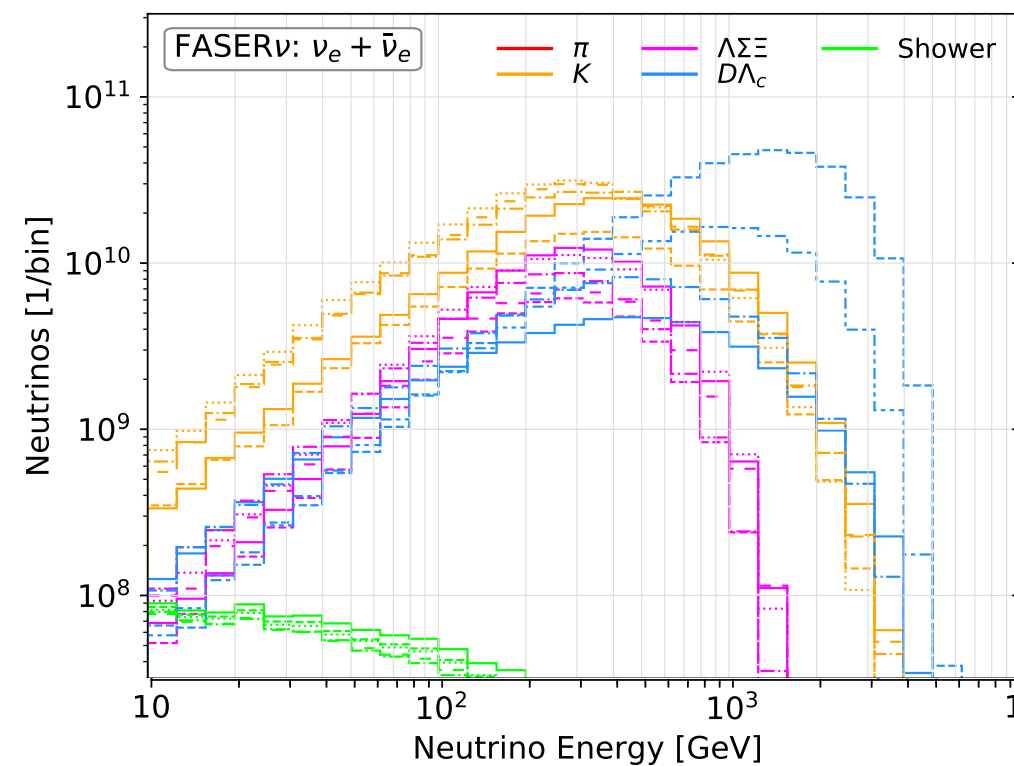
Run3で想定されるイベント数

検出器を通過するニュートリノ、親粒子ごとに分けて表示



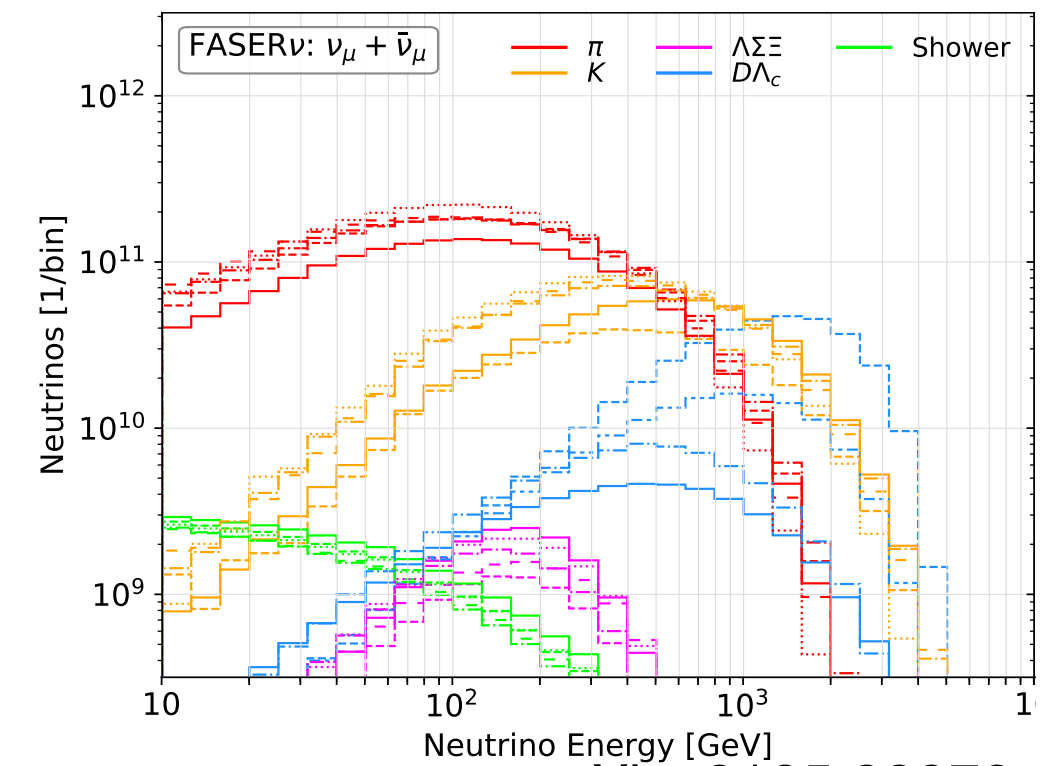
ニュートリノの E_ν [GeV]

エネルギー [GeV]



--- DPMJET 3.2017
 — SIBYLL 2.3d
 - - EPOSLHC

..... QGSJET II-04
 - - - Pythia8 (Hard)
 - - - Pythia8 (Soft)

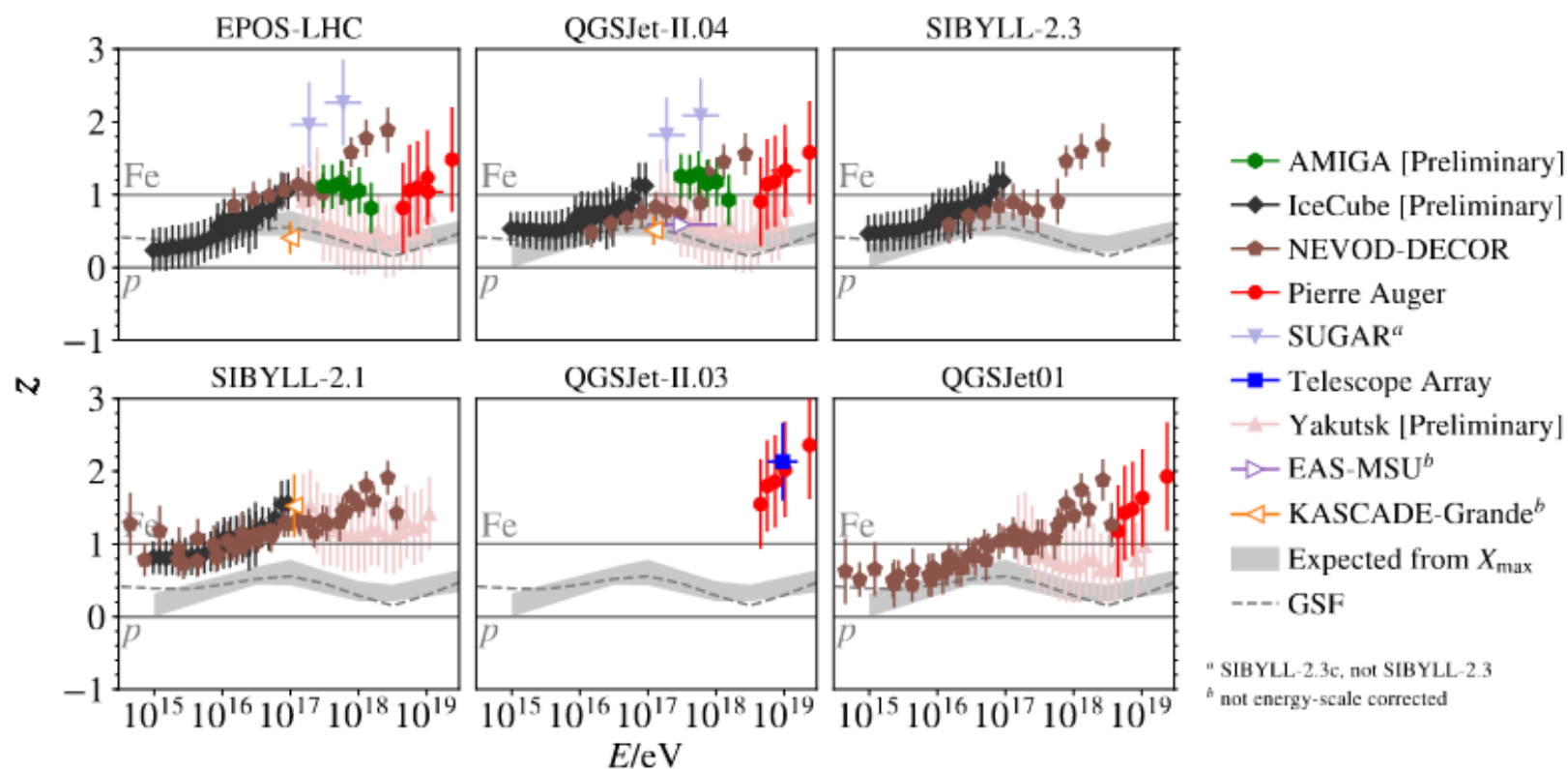


arXiv: 2105.08270

ν_e と ν_μ の割合を測定することで、 π と K の割合を測定できる!!

μ粒子超過問題

地上でのμ粒子数がシミュレーション予測よりも多い



$$z = \frac{\ln(N_{\mu}) - \ln(N_{\mu,proton})}{\ln(N_{\mu,iron}) - \ln(N_{\mu,proton})}$$

L. Cazon et al., PoS(ICRC2019) 214

Ralph Engel, Forward Physics Facility 7 meeting (March 2024)

Modification of Sibyll 2.3d to study different versions of muon enhancement

- Rho meson in pion interactions (leading particle effect only) → **not seen in FPF**
- Baryon pair production (all interactions) → **not seen in FPF**
- Kaon production (all interactions) → **directly seen in FPF**

Kaon production シナリオならFASER or Future detectorで検証できる

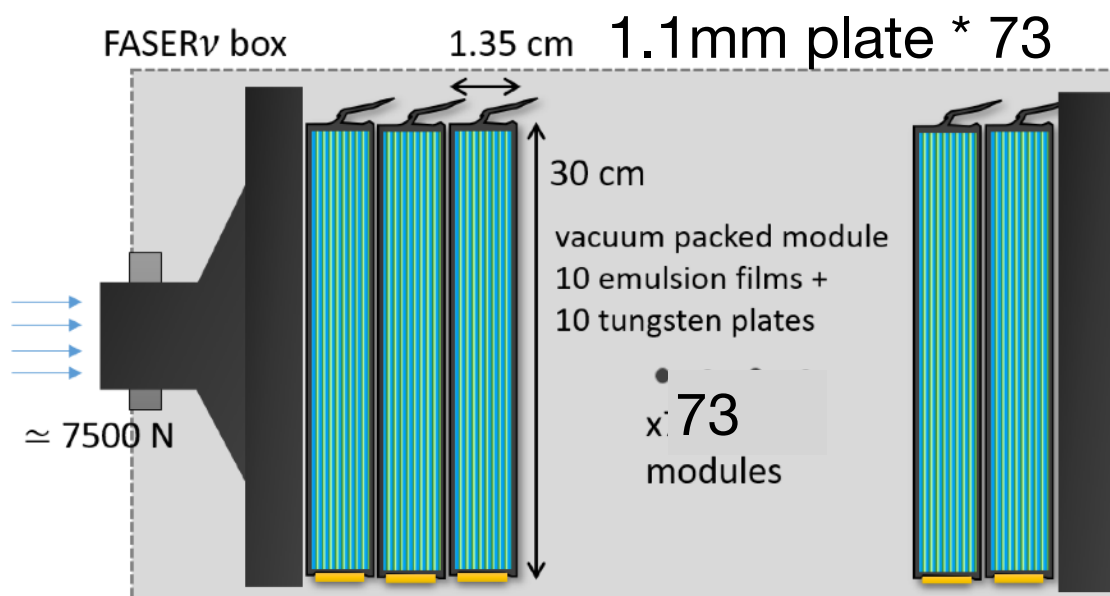
FASER ν detector

原子核乾板を用いた検出器

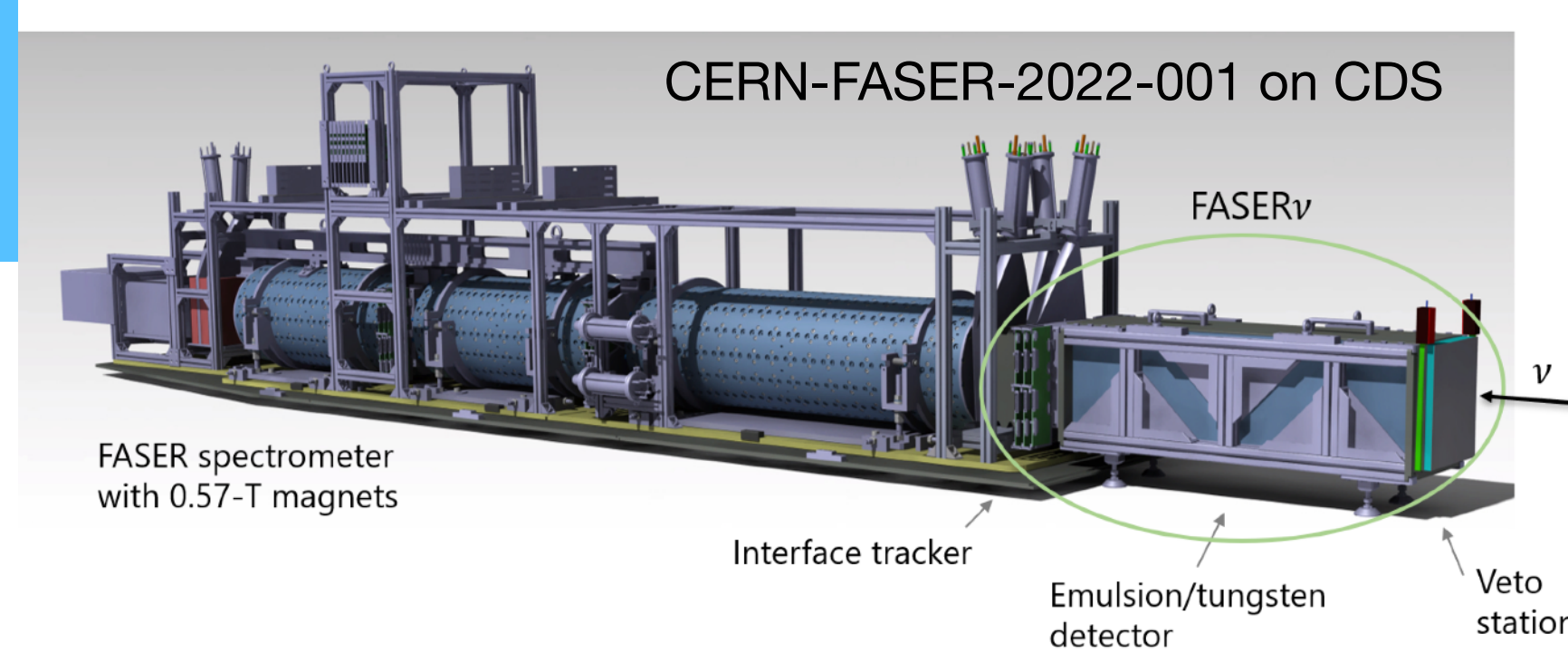
特徴

非常に高い位置分解能

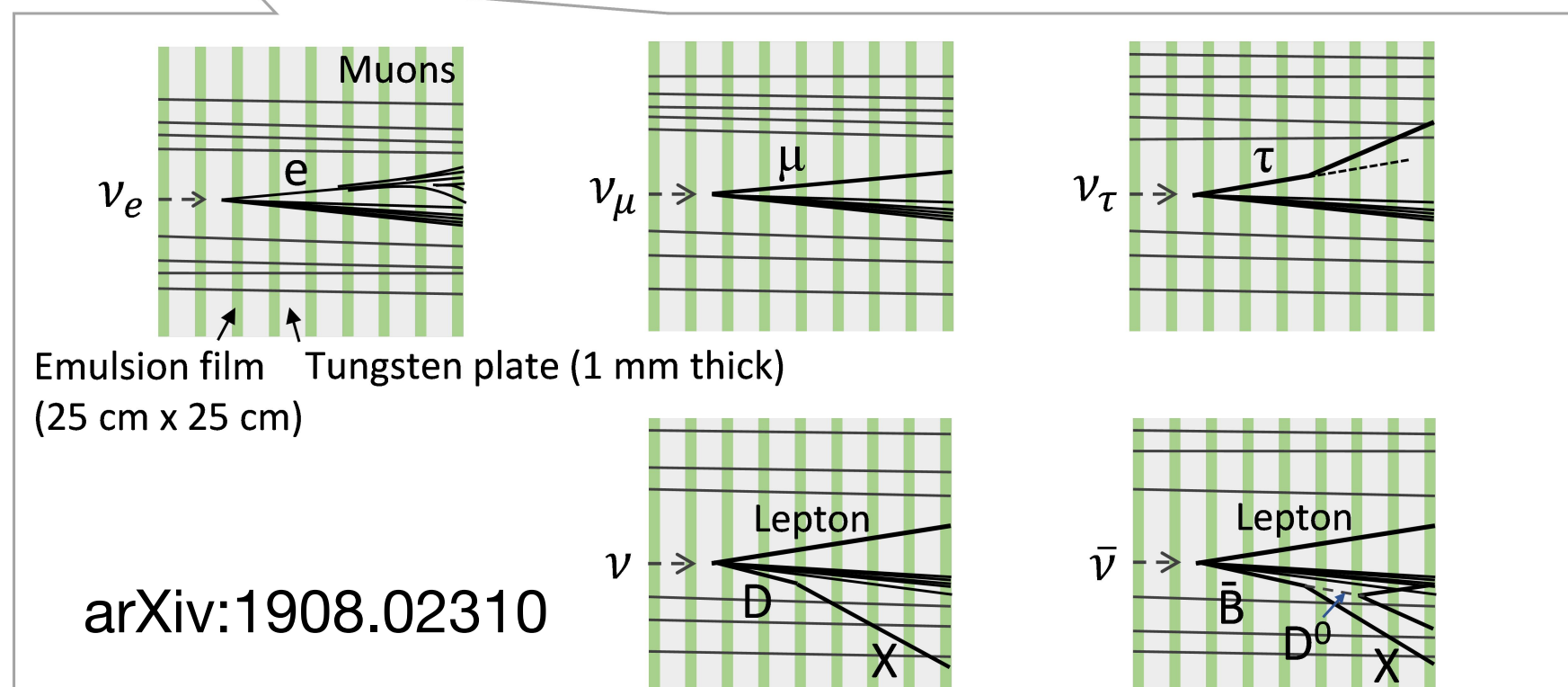
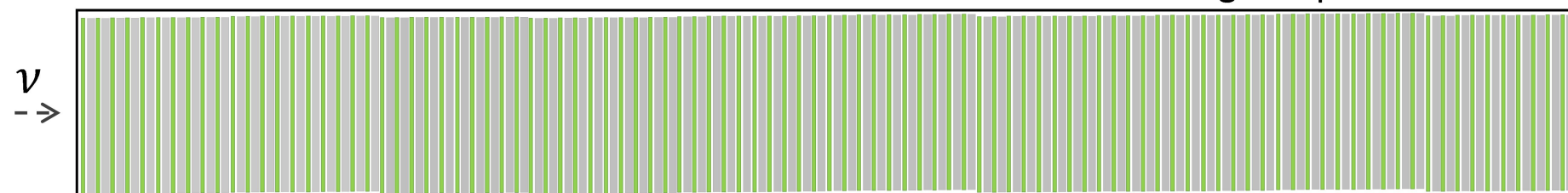
ニュートリノのフレーバーの決定



CERN-FASER-2022-001 on CDS



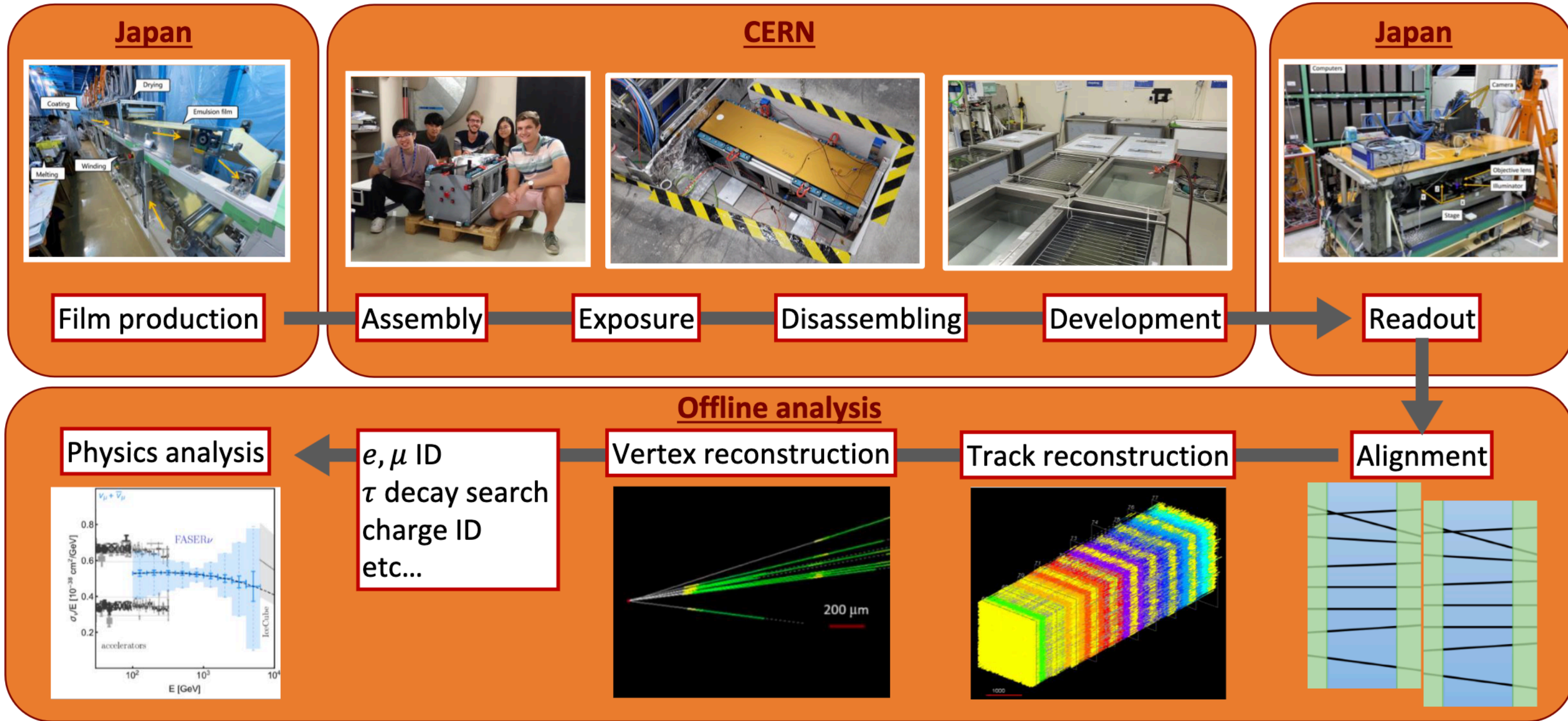
Total 1000 emulsion films interleaved with 1-mm-thick tungsten plates



arXiv:1908.02310

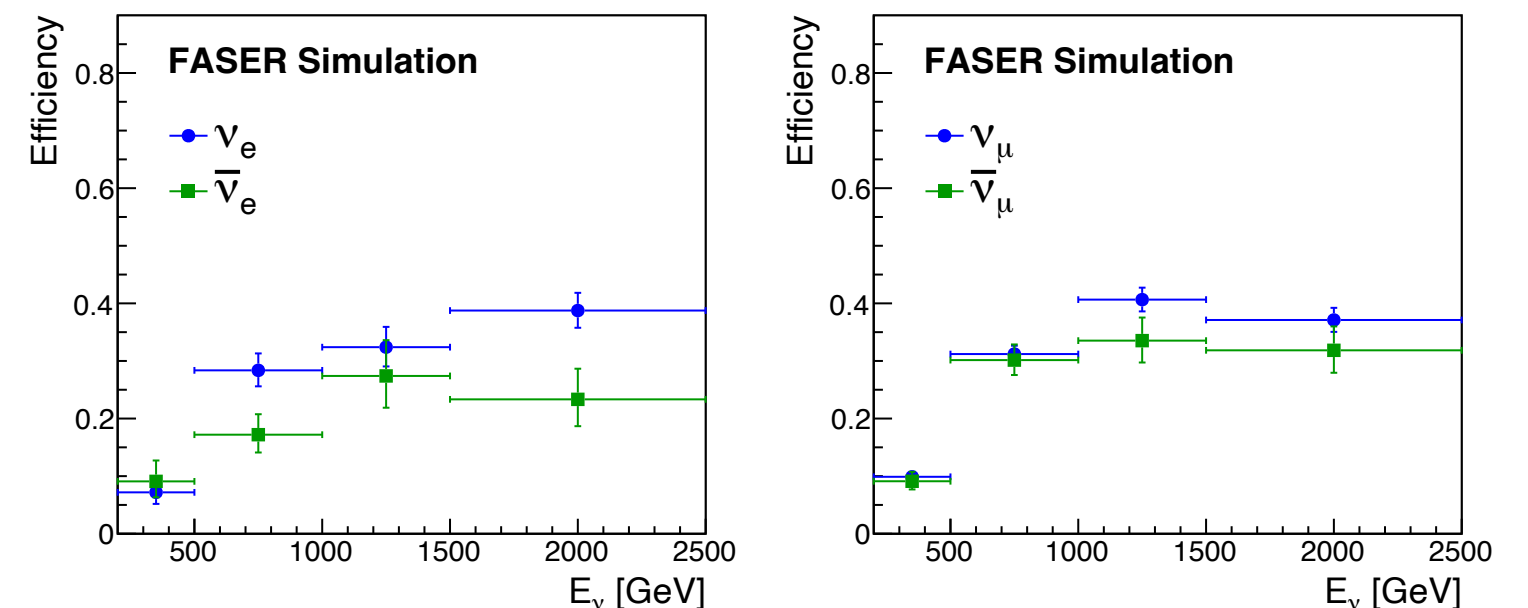
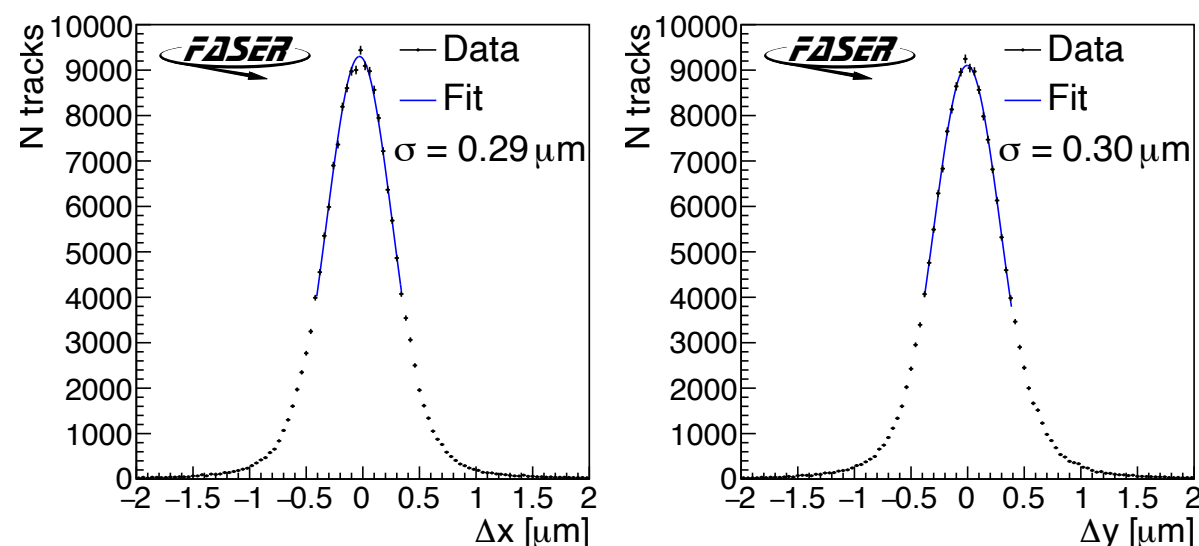
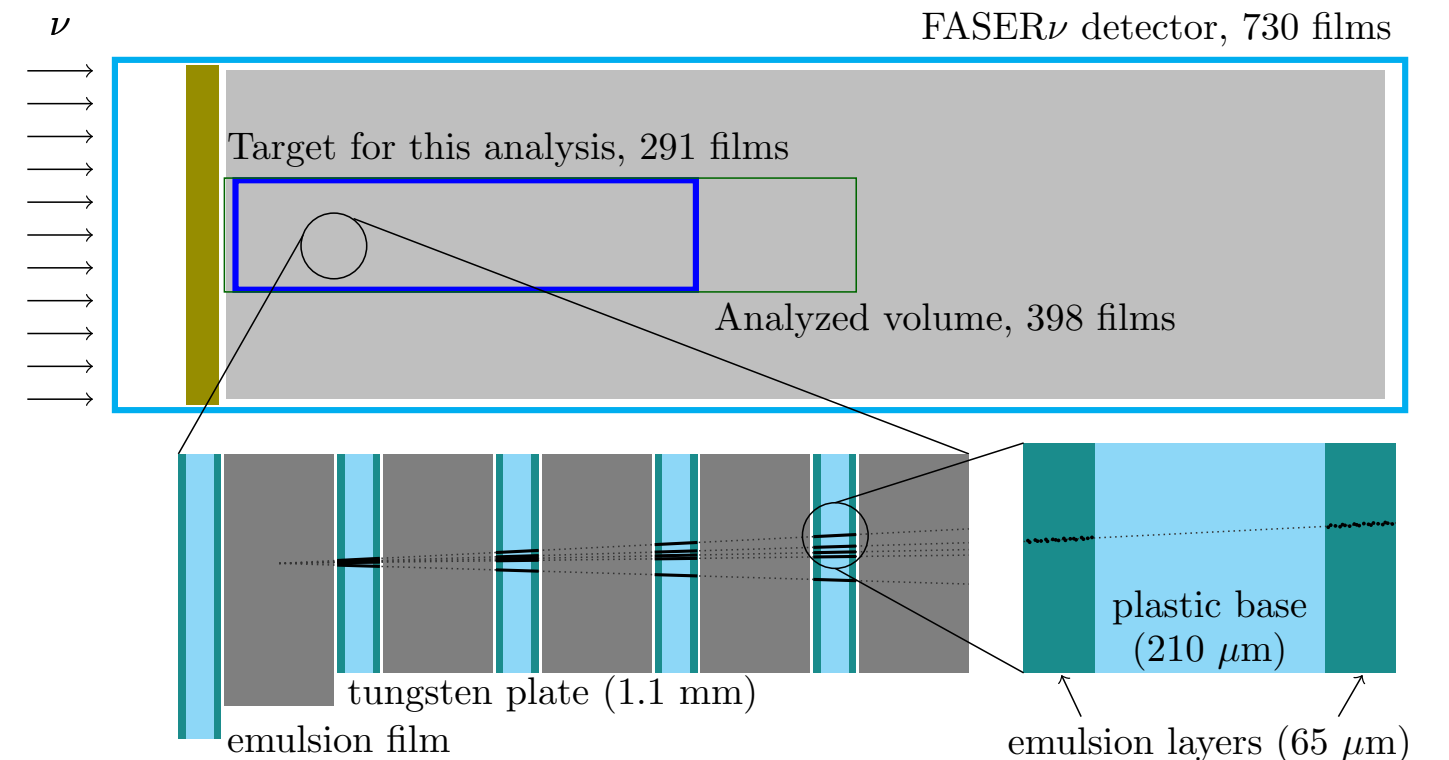
FASER ν Process

Slide By Jeremy Atkinson at NuFACT2023



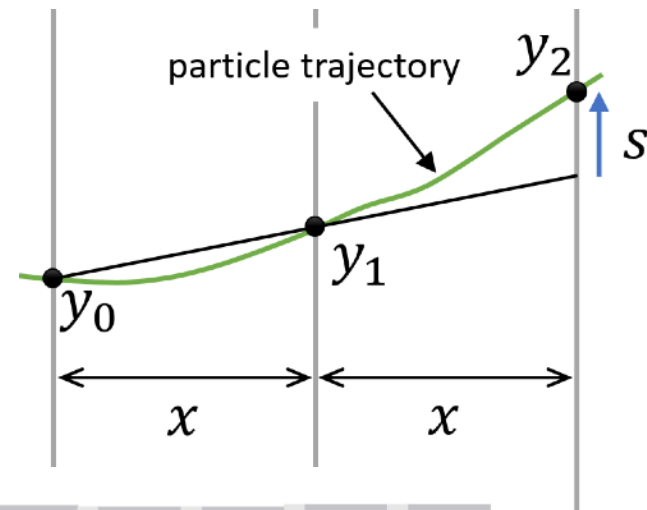
最新の物理解析結果 - ν_e & ν_μ 検出 + ν -核子反応断面積

- 2022年の2nd moduleを解析 (9.5fb^{-1})
 - Film 730枚のうち、291枚の中心部が解析領域
- 想定されるイベント数
 - 複数のハドロン相互作用モデルを考慮
 - 軽いハドロン: EPOS-LHC, QGSJET II-04, SIBYLL 2.3d, PYTHIA forward tune
 - チャームハドロン: POWHEG, pQCD のuncertaintyを考慮
- $N_{\text{expected}} = 1.68_{-0.37}^{+1.18} (\text{flux})_{-0.36}^{+0.51} (\text{others})$ for ν_e
- $N_{\text{expected}} = 8.64_{-0.76}^{+1.41} (\text{flux})_{-1.82}^{+2.65} (\text{others})$ for ν_μ

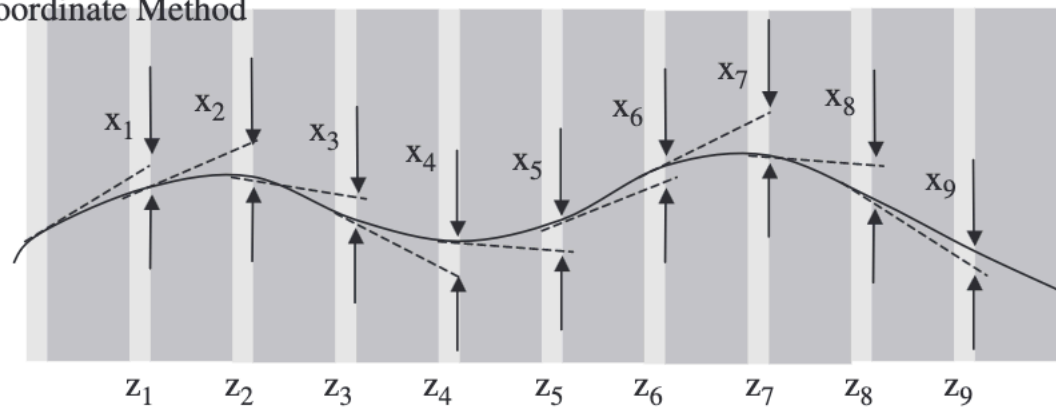


ミューオン運動量測定

Momentum dependency in the scattering angle



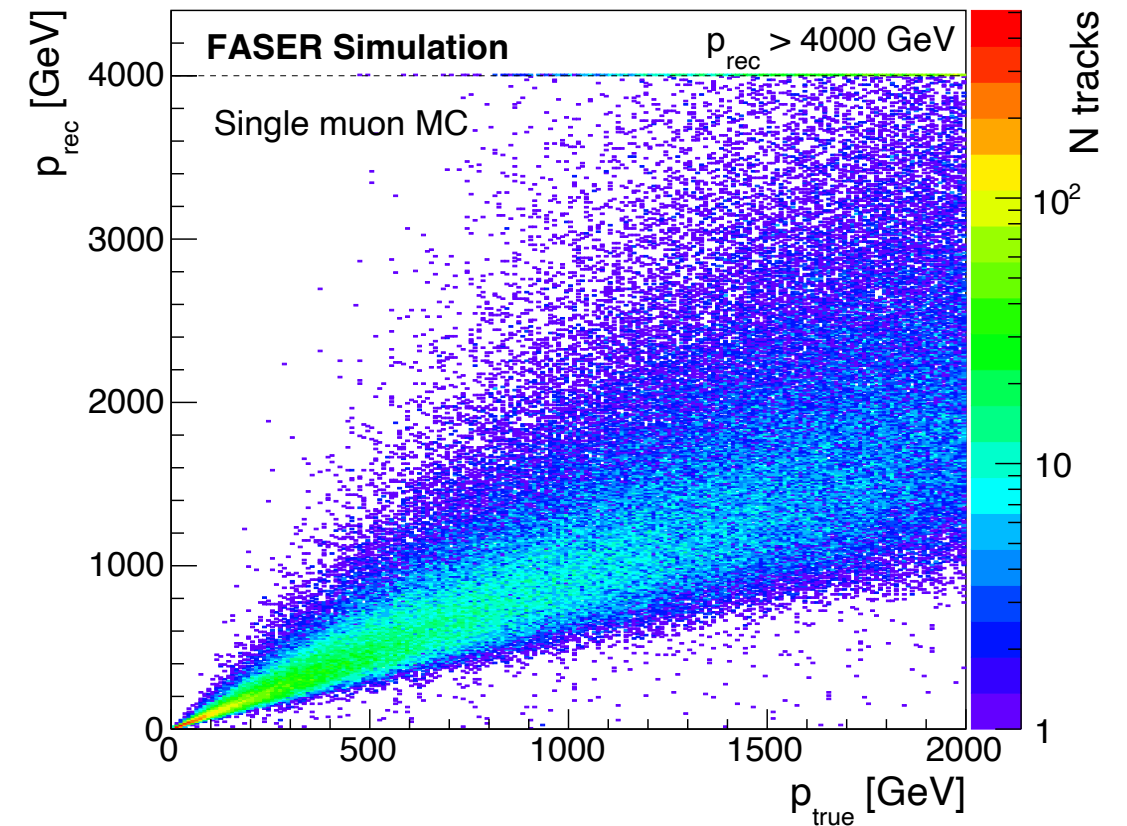
Coordinate Method



$$\theta_{plane}^{RMS} = \frac{0.0136}{\beta pc} \sqrt{\frac{z}{X_0}} \left\{ 1 + 0.038 \ln \left(\frac{z}{X_0 \beta^2} \right) \right\},$$

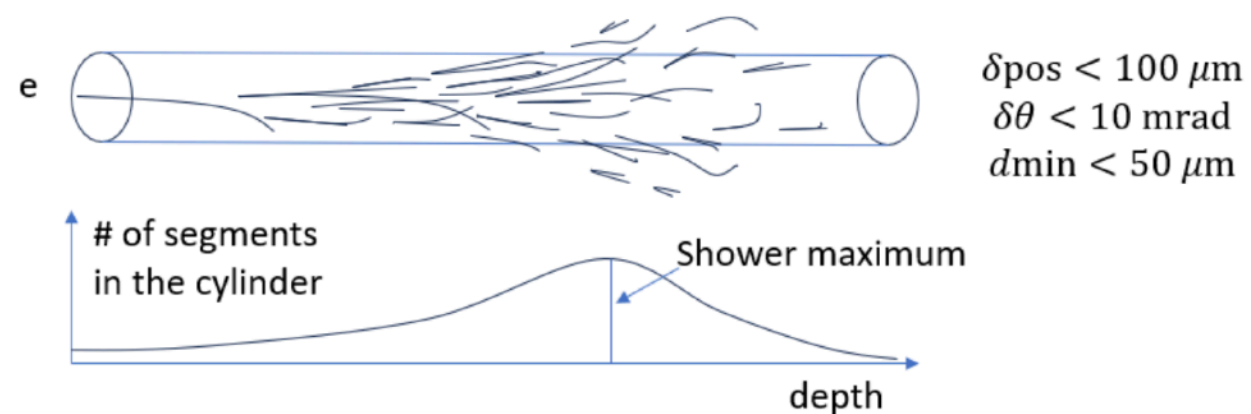
z: thickness
p: momentum

Momentum is estimated by measuring displacements for every 1 plate, every 2 plates, every 4 plates, every 8 plates, and every 16 plates and calculating RMS for each case.



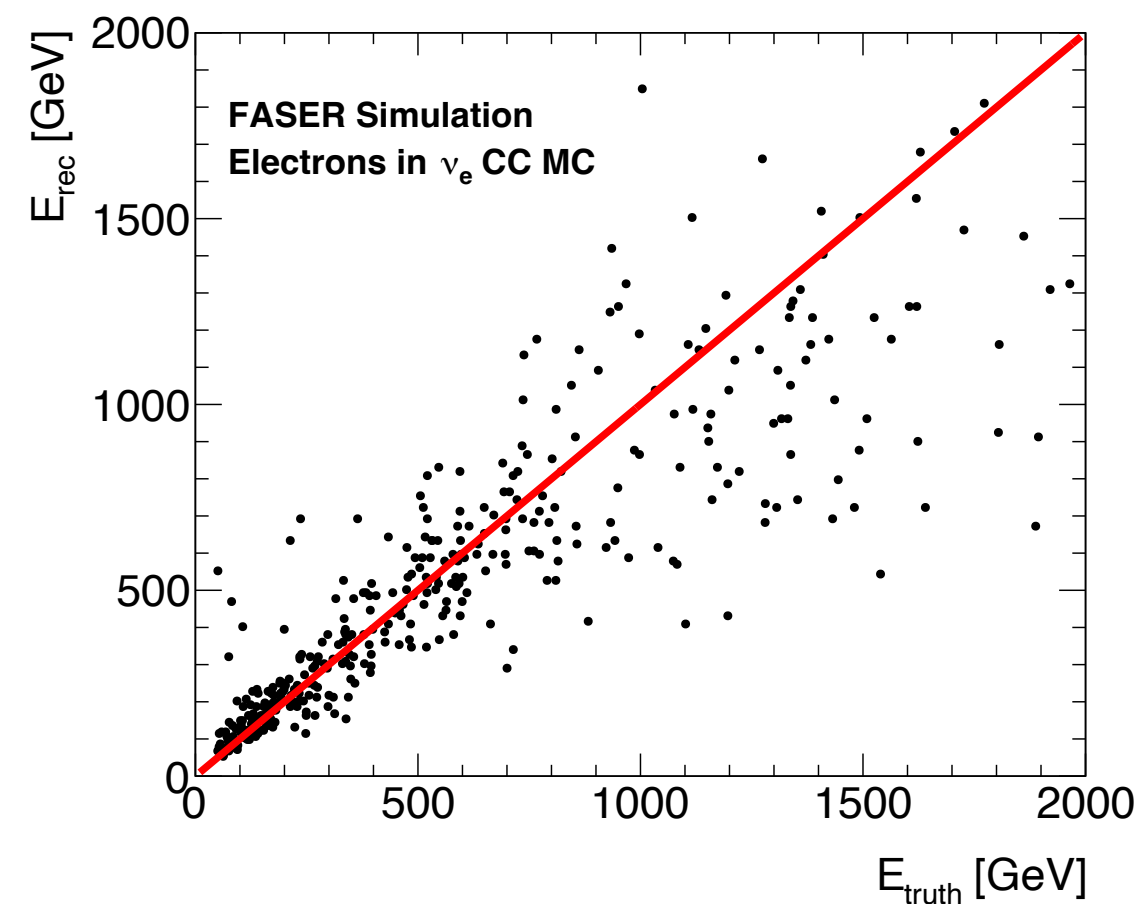
Resolution: ~20% at 200 GeV

EMシャワーエネルギー測定



Count the number of segments in ± 3 films around the shower maximum (total 7 films)

The number of backgrounds was estimated and subtracted by counting the number of segments at the cylinder randomly opened.

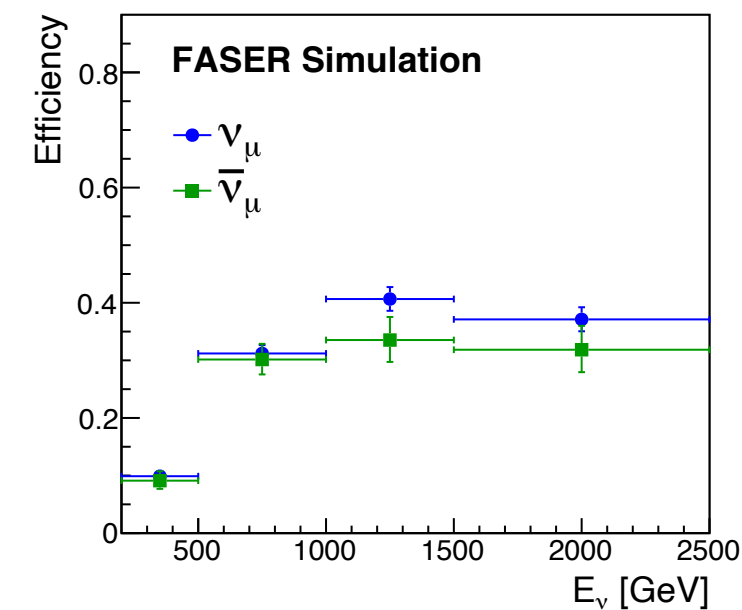
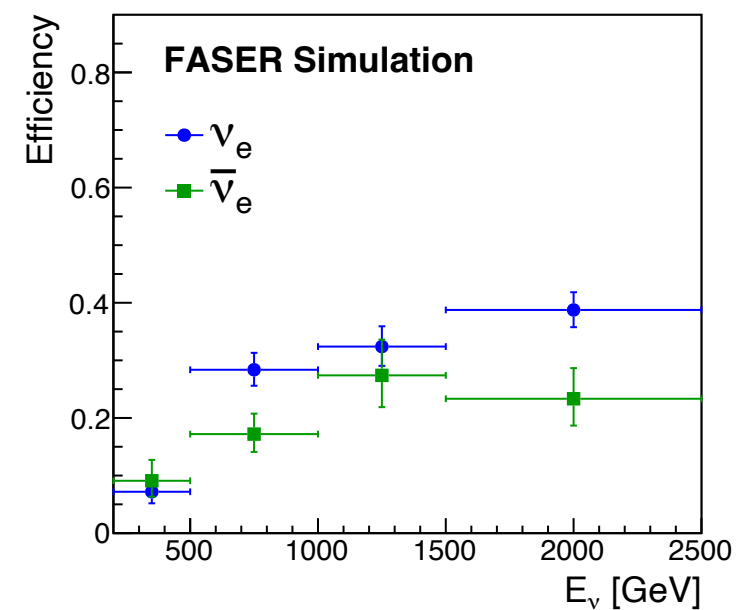
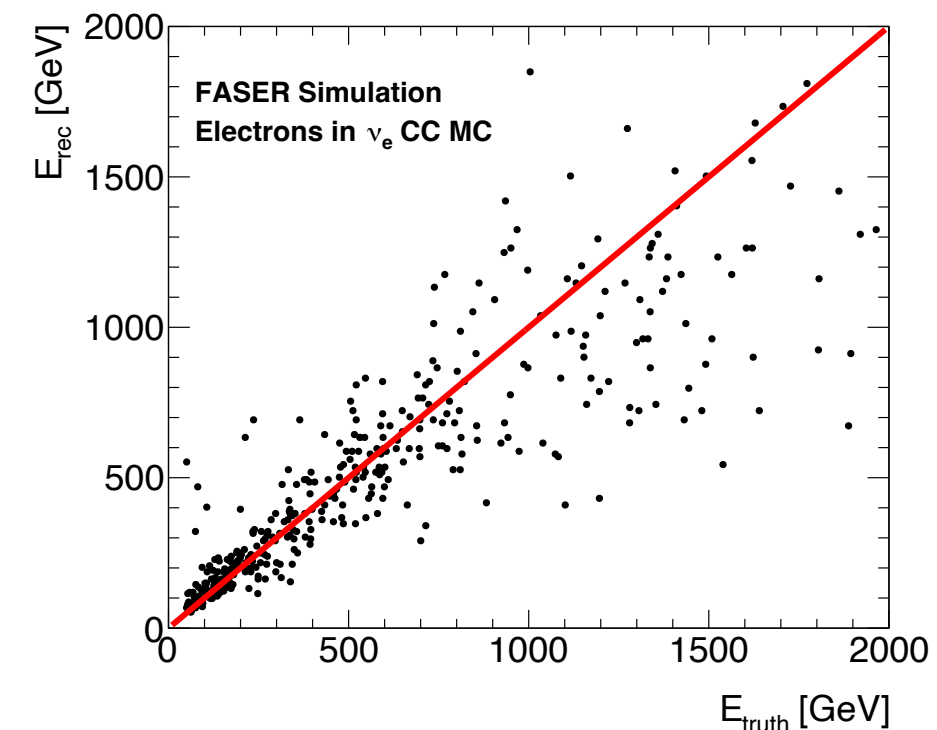
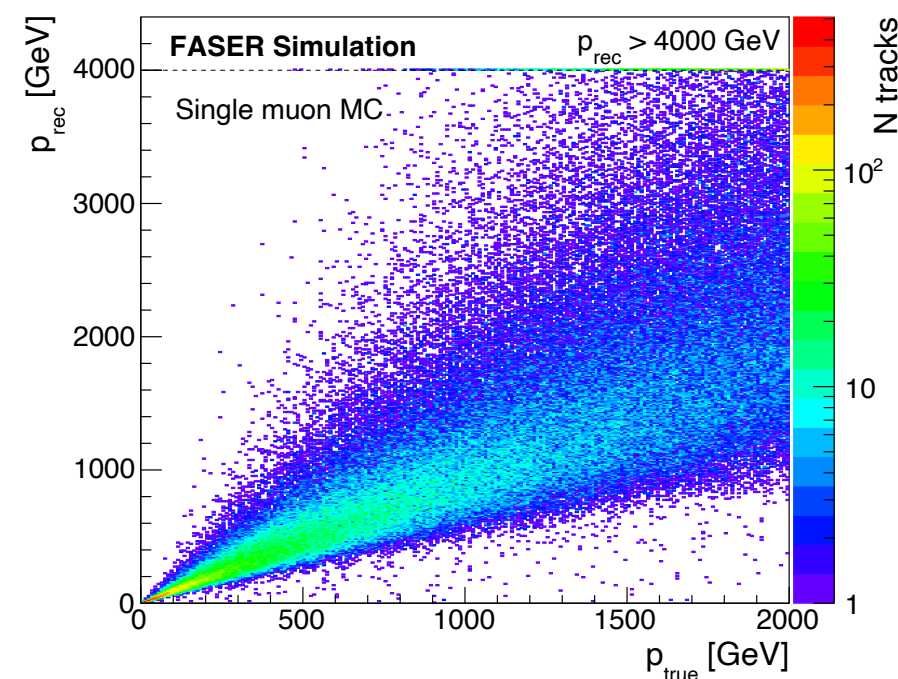
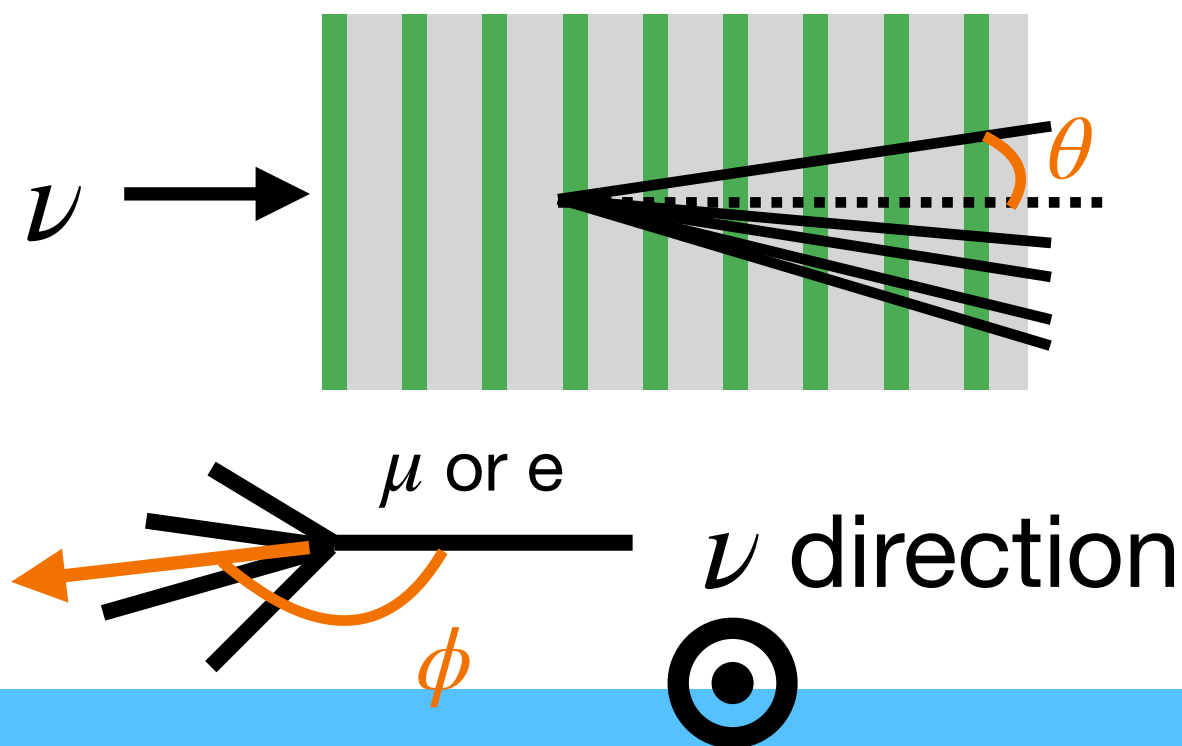


Resolution: $\sim 25\%$ at 200 GeV

イベント選別

イベント選別

- 5 or more tracks attached to a vertex
- No charged parent track
- 4 or more tracks with $\tan\theta < 0.1$
- $\tan\theta > 0.005$ for muon or EM shower
- An EM shower or a track of more than 200 GeV
- $\phi > 90^\circ$

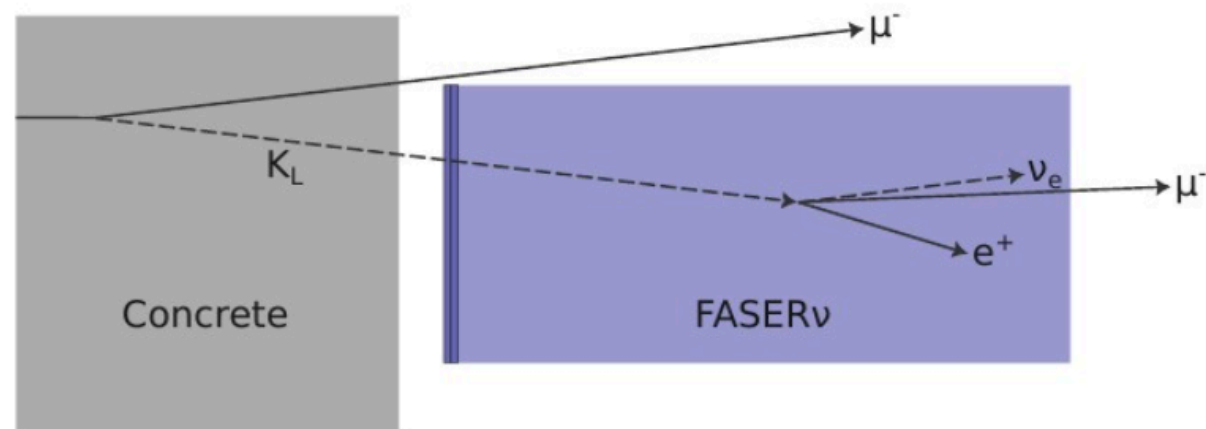


バックグラウンド

- 中性ハドロン

Hadron type	K_L	n	Λ
Events simulated ($E_h > 200$ GeV)	13497	13191	13902
Events selected as ν_e CC	0	0	0
Events selected as ν_μ CC	6	11	5
Scaling factor (data/MC)	1/232	1/256	1/423
Hadron type	K_S	\bar{n}	$\bar{\Lambda}$
Events reconstructed ($E_h > 200$ GeV)	7113	5827	5368
Events selected as ν_e CC	1	0	0
Events selected as ν_μ CC	3	3	4
Scaling factor (data/MC)	1/436	1/569	1/630

ミューオン由来の中性ハドロン生成



- ニュートリノ Neutral Current interaction

$$\nu_e: 0.008^{+0.013}_{-0.004} \text{ (flux)} \pm 0.001 \text{ (cross - section)}^{+0.007}_{-0.004} \text{ (others)}$$

$$\nu_\mu: 0.045^{+0.004}_{-0.005} \text{ (flux)} \pm 0.003 \text{ (cross - section)}^{+0.076}_{-0.024} \text{ (others)}$$

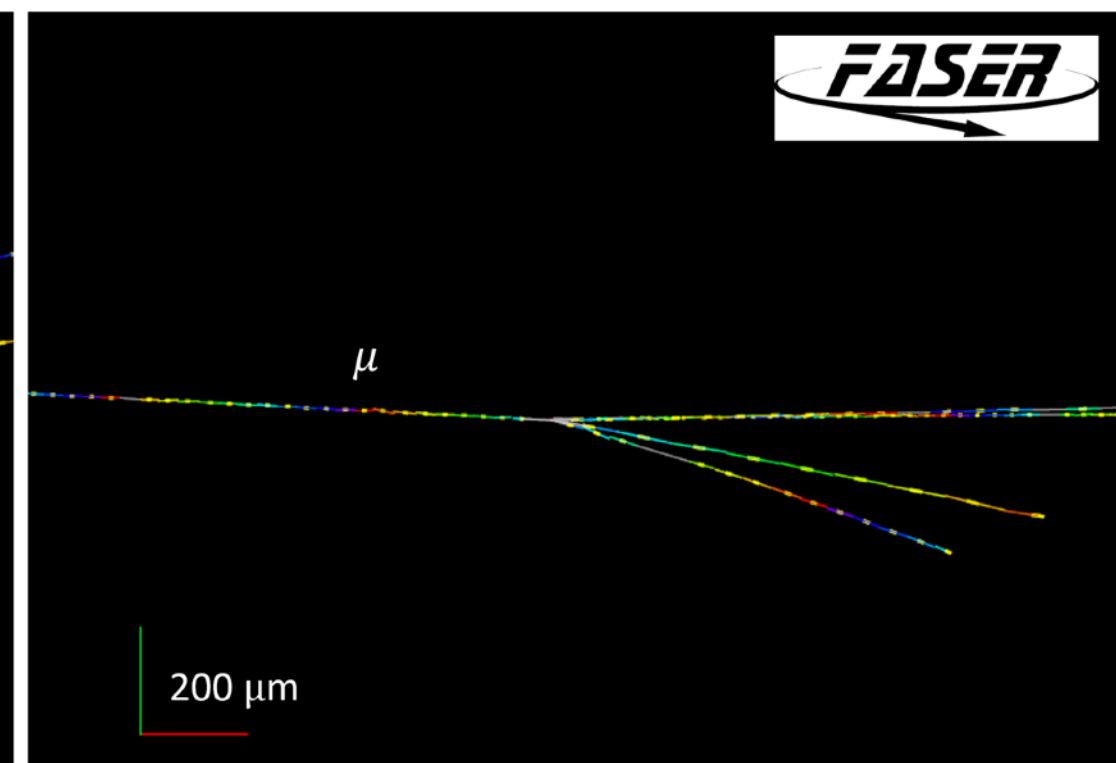
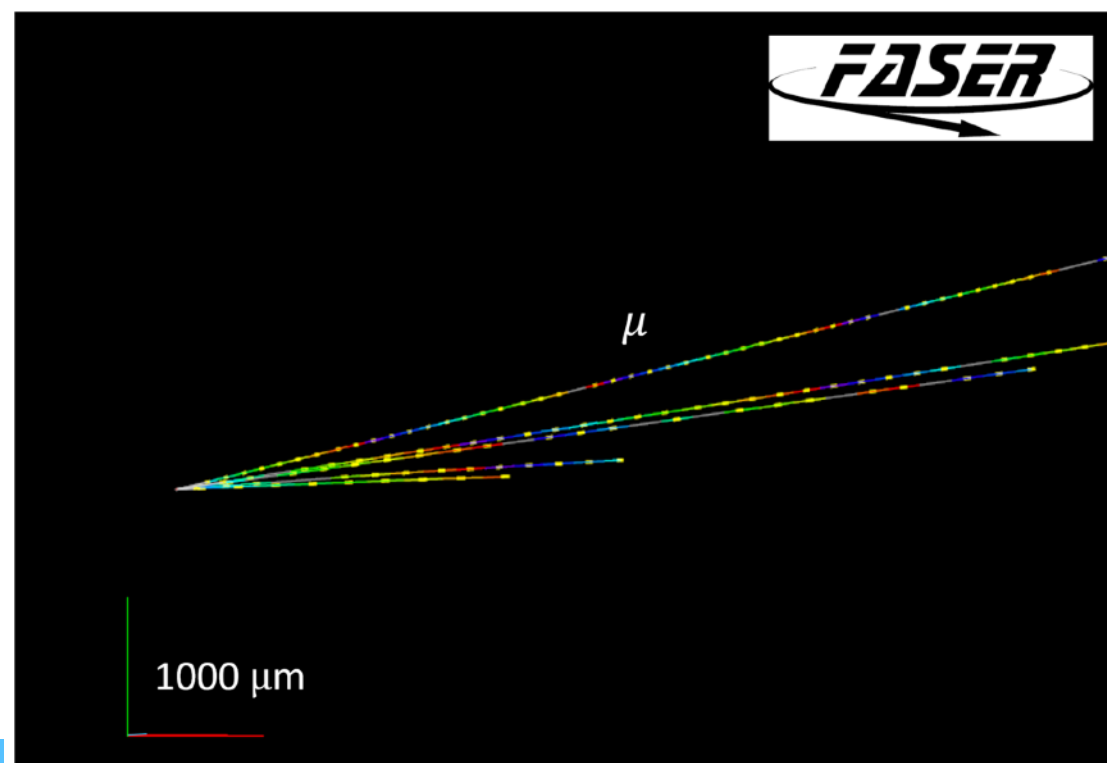
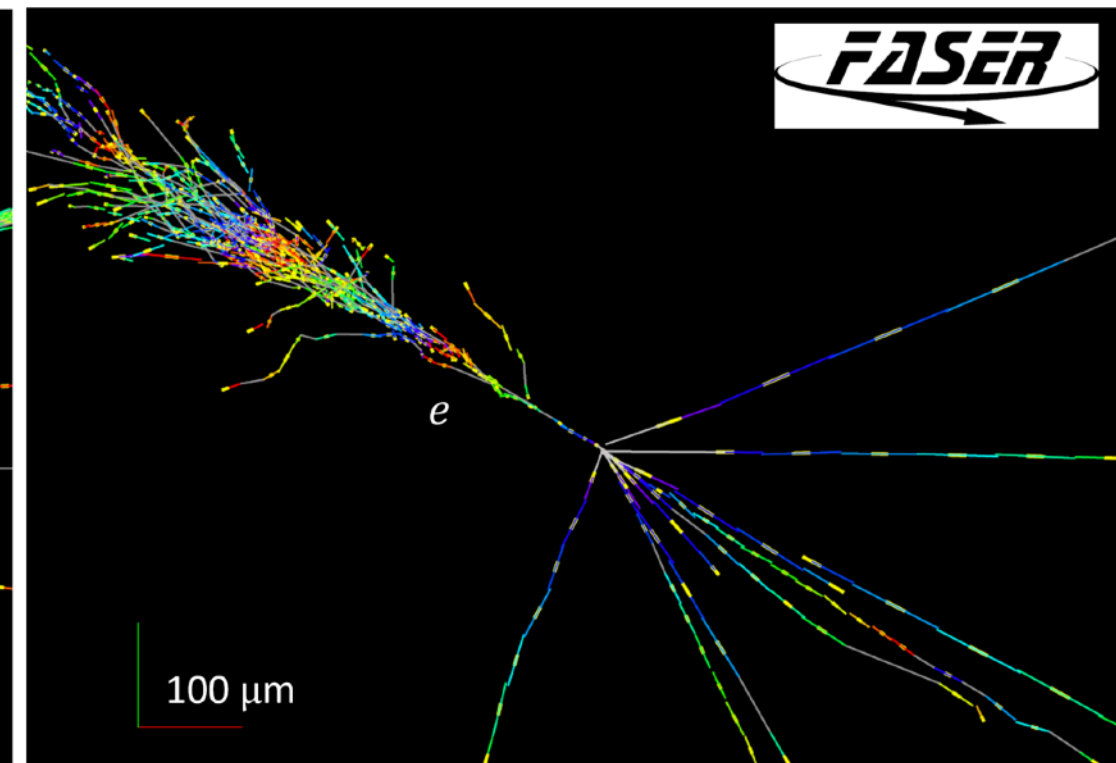
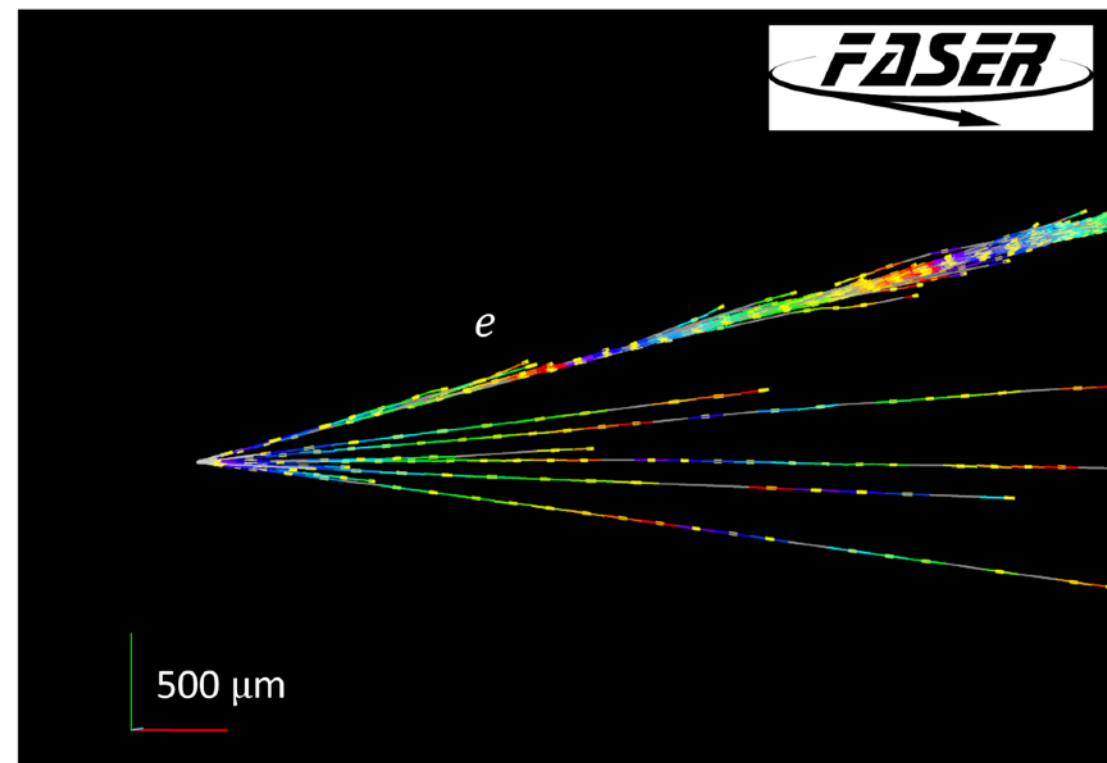
Total background

$$\nu_e: 0.025^{+0.015}_{-0.010}$$

$$\nu_\mu: 0.22^{+0.09}_{-0.07}$$

ν_e candidates, ν_μ candidates

- ν_e 候補: 4イベント
- ν_μ 候補: 8イベント
- バックグラウンド
 - $\nu_e: 0.025^{+0.015}_{-0.010}$
 - $\nu_\mu: 0.22^{+0.09}_{-0.07}$
- 10^{10} toysを振って、significanceを計算
 - $\nu_e: 5.2 \sigma$
 - $\nu_\mu: 5.7 \sigma$

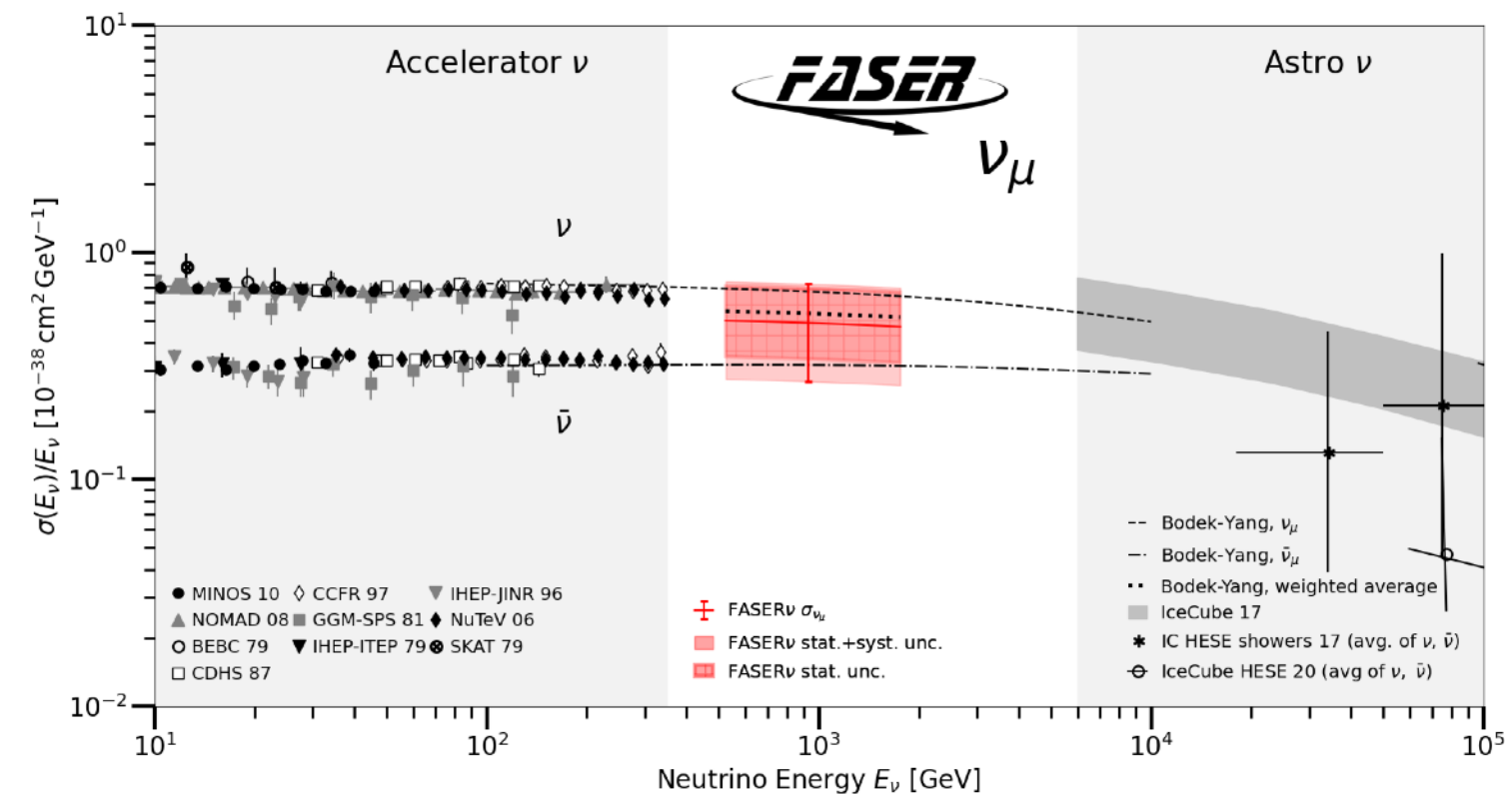
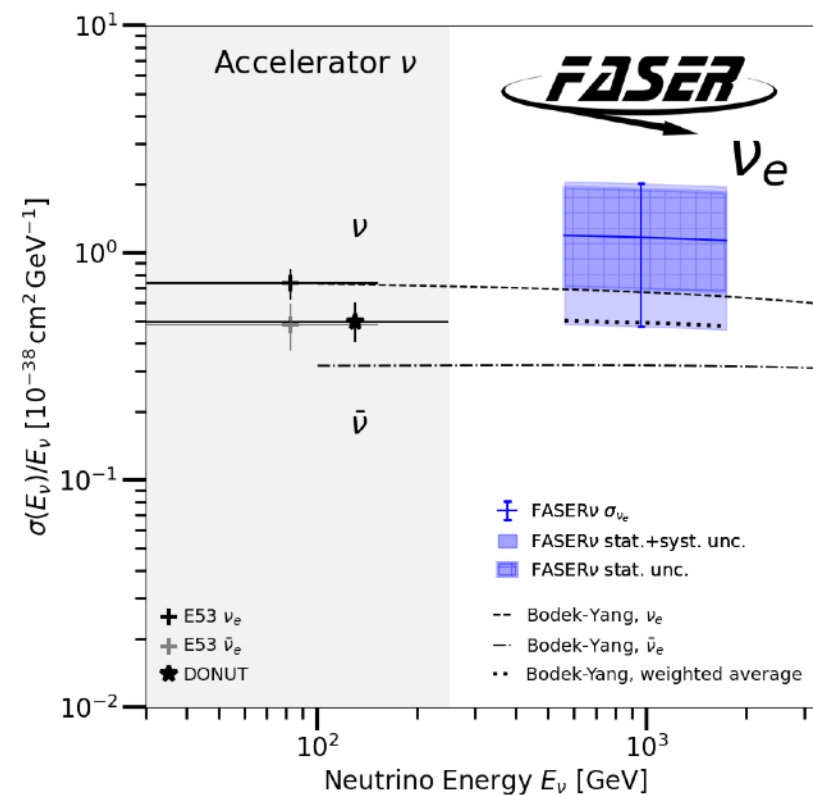


ν -核子反応断面積

- $$\sigma_{\text{obs}} = \frac{N_{\nu, \text{obs}}}{N_{\text{expected}}} \sigma_{\text{theory}}$$
- ν_e 候補: 4 イベント, ν_μ 候補: 8 イベント
- $$N_{\text{expected}} = 1.68^{+1.18}_{-0.37} (\text{flux})^{+0.51}_{-0.36} (\text{others}) \text{ for } \nu_e$$
- $$N_{\text{expected}} = 8.64^{+1.41}_{-0.76} (\text{flux})^{+2.65}_{-1.82} (\text{others}) \text{ for } \nu_\mu$$

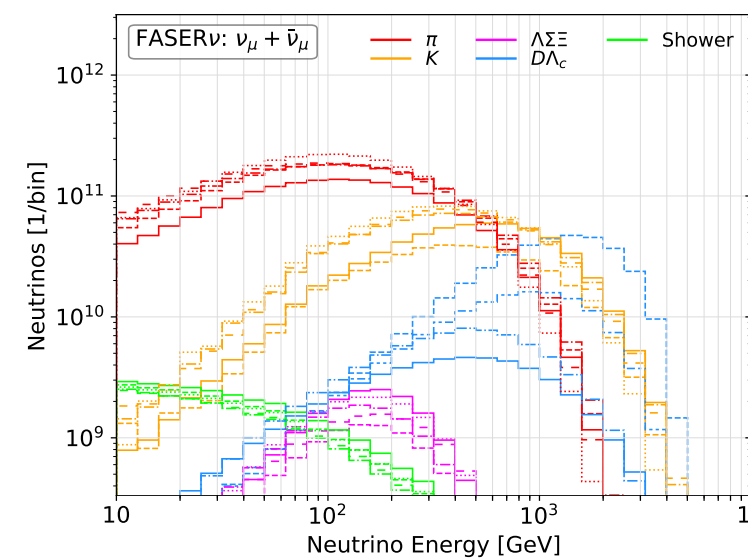
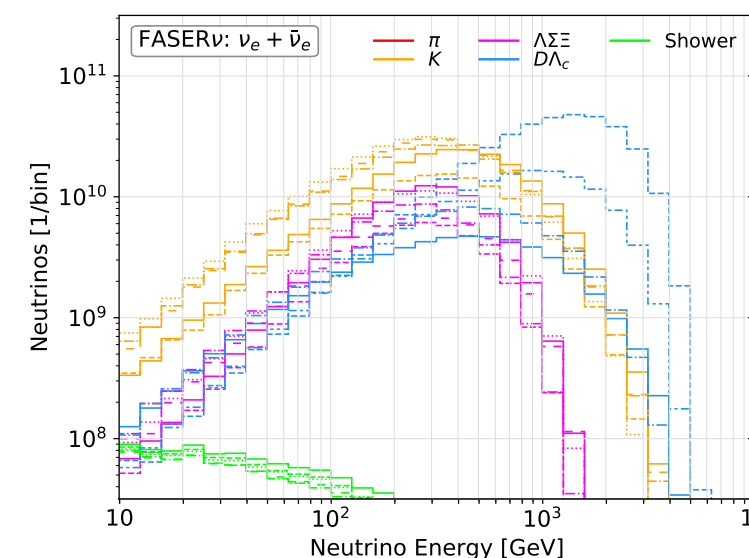
Table 2: Systematic uncertainties related to the signal expectation.

Source	Relative uncertainty	
	ν_e	ν_μ
Luminosity	2.2%	2.2%
Tungsten thickness	1%	1%
Interactions with emulsions	+3.6% -0%	+3.6% -0%
Flux uncertainty	+70% -22%	+16% -9%
Line of sight position	+2.1% -2.4%	+1.9% -2.5%
Efficiency from hadronization	+22% -5%	+23% -5%
Efficiency from reconstruction	20%	20%
Efficiency from MC statistics	4.9%	2.8%
Total	+70% (flux) -22% (flux)	+16% (flux) -9% (flux)
	+30% (other) -21% (other)	+31% (other) -21% (other)



今後の展望

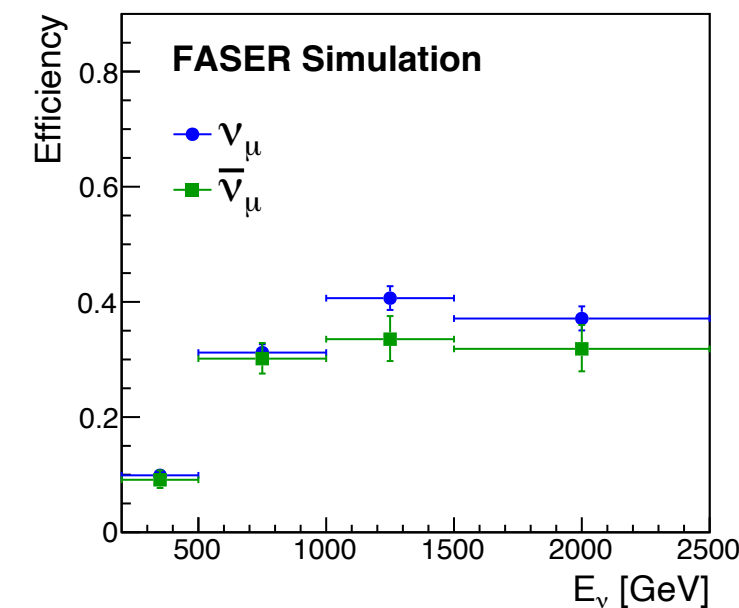
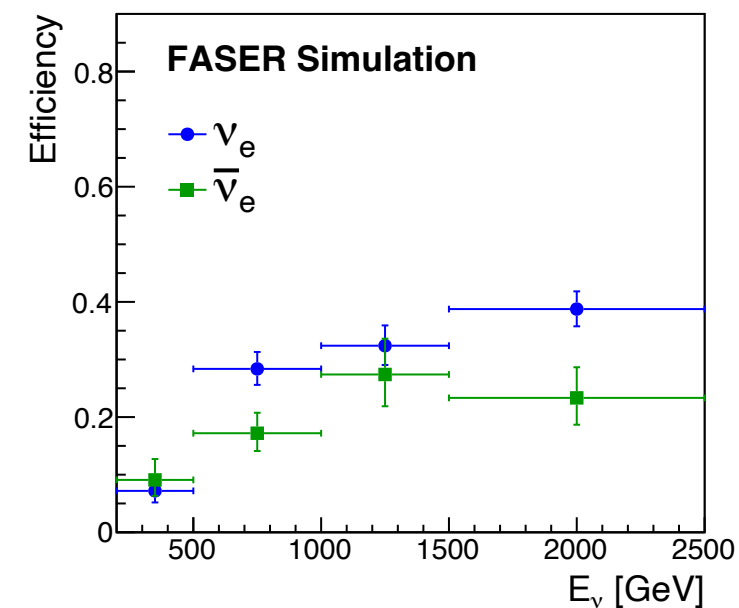
- ニュートリノエネルギー再構成
 - FASER ν 検出器での再構成手法を開発中
 - ミューオン運動量測定およびEMシャワーエネルギー測定をSPSでのテストビームにて検証
 - ハドロンの運動量やミューオン・電子のエネルギー、角度からニュートリノエネルギーを再構成する
- 2022年2nd moduleのフル解析
 - 約6倍のtarget mass
 - さらに、1 TeV以下のニュートリノの検出効率を上げる
 - 宇宙線との関係で重要なのは、100 GeV-1TeVでの ν_e と ν_μ の比
- バックグラウンドミューオン解析
 - p-p衝突で生じたハドロン由来の成分を測定できれば、ハドロン相互作用モデルの検証につながる



--- DPMJET 3.2017
 — SIBYLL 2.3d
 - - EPOS LHC

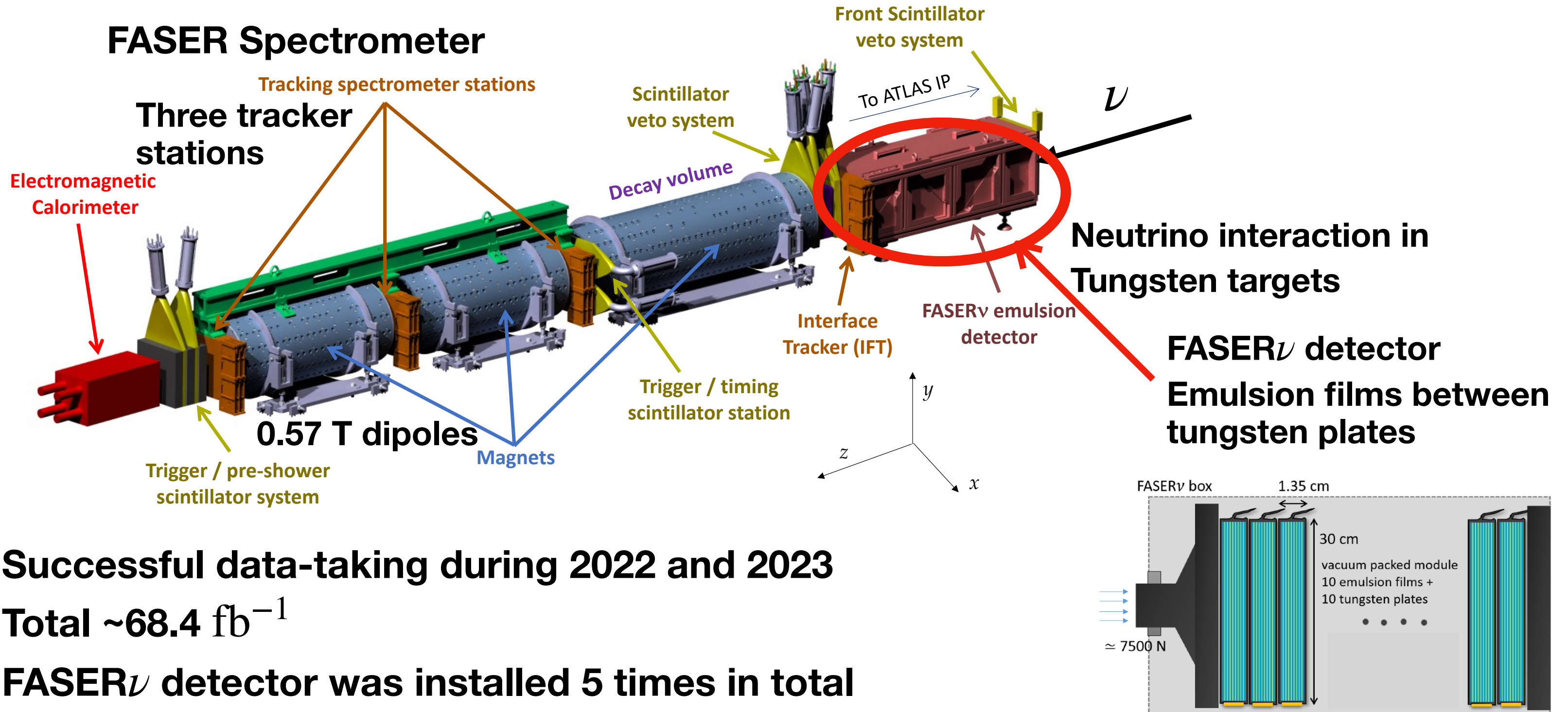
..... QGSJET II-04
 - - - Pythia8 (Hard)
 - - - Pythia8 (Soft)

arXiv:
 2105.08270



Back-up slides

The FASER detector



Successful data-taking during 2022 and 2023

Total $\sim 68.4 \text{ fb}^{-1}$

FASER ν detector was installed 5 times in total

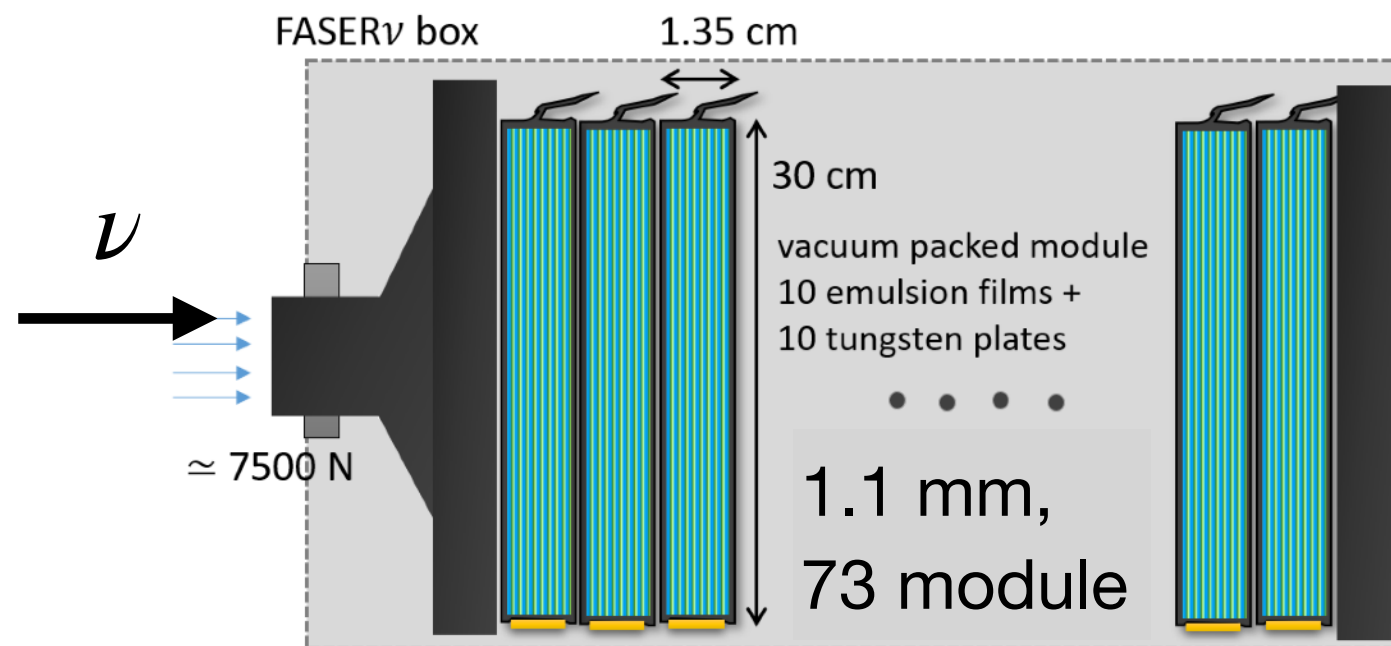
The FASER ν detector

FASER ν detector

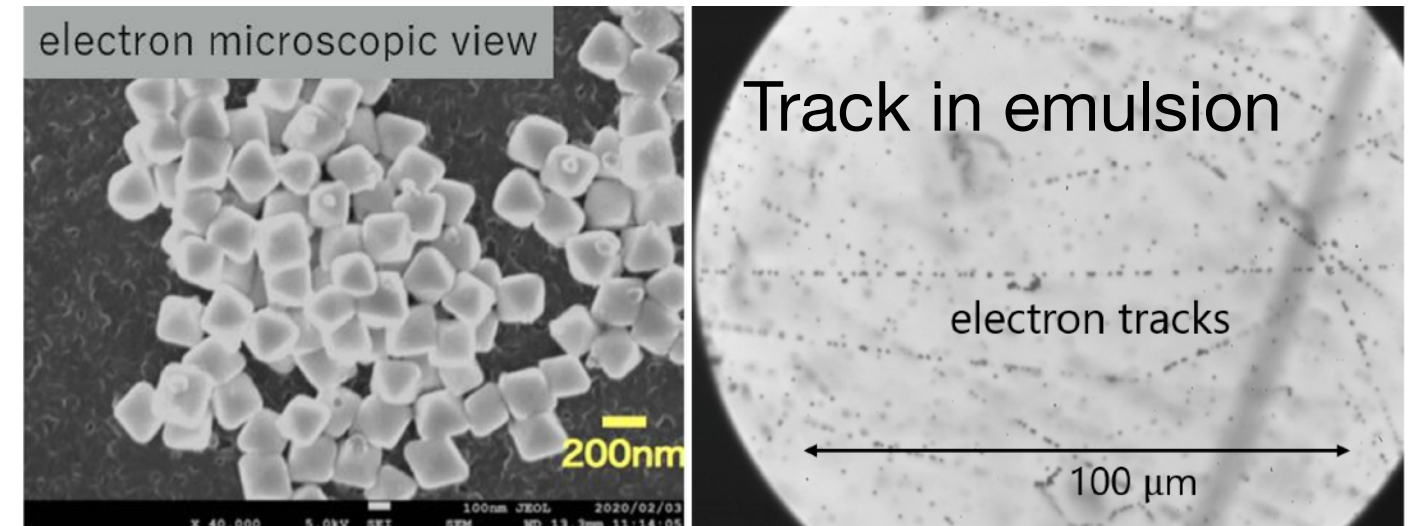
Emulsion films between tungsten plates

1.1 mm tungsten plates x 730: target

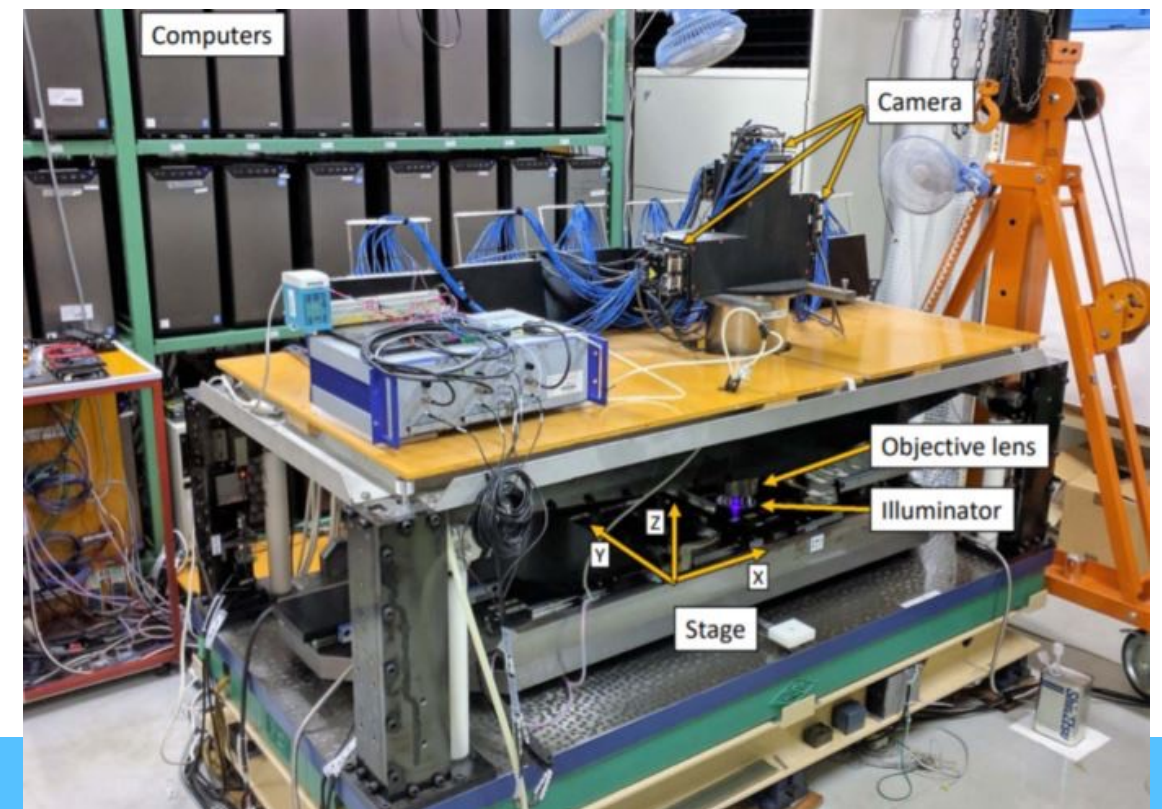
Emulsion film x 730: to measure tracks



Emulsion films (25cm x 30cm)



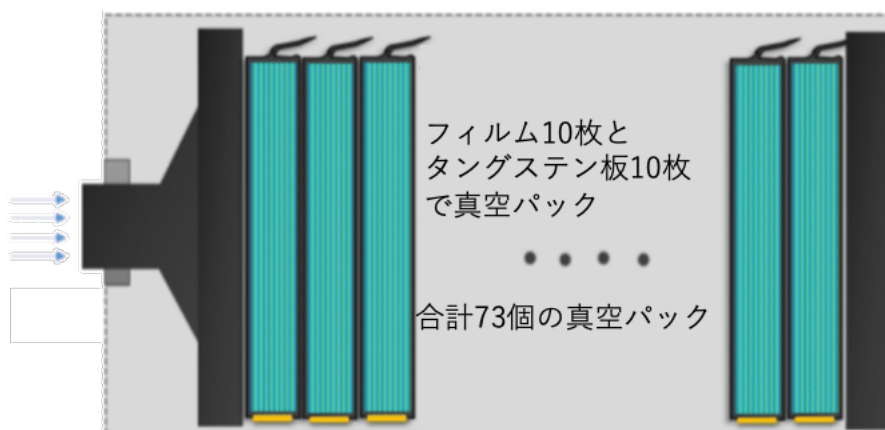
Microscope in Nagoya Univ.



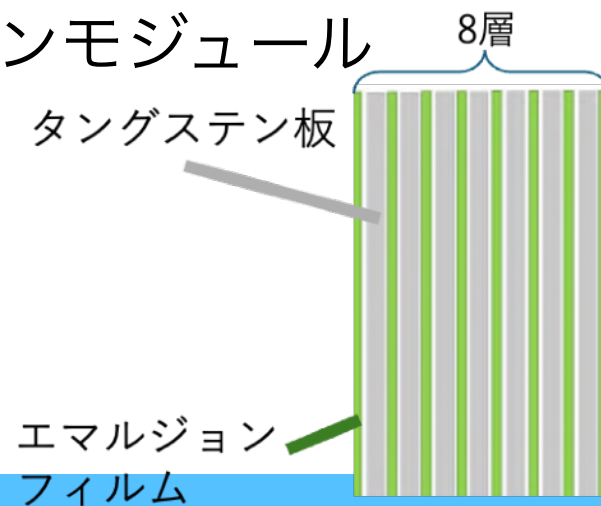
ミュオンバックグラウンドの測定

FASER ν 検出器および
ミュオンモジュールで
バックグラウンド測定
が行われている。

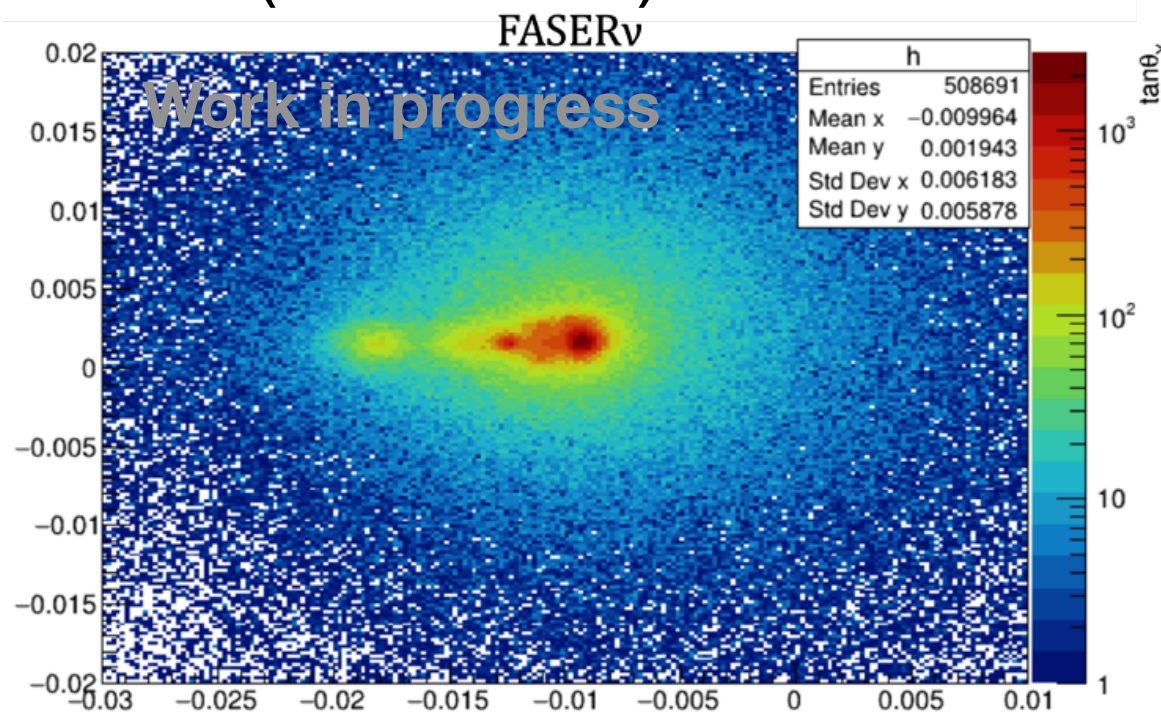
FASER ν 検出器



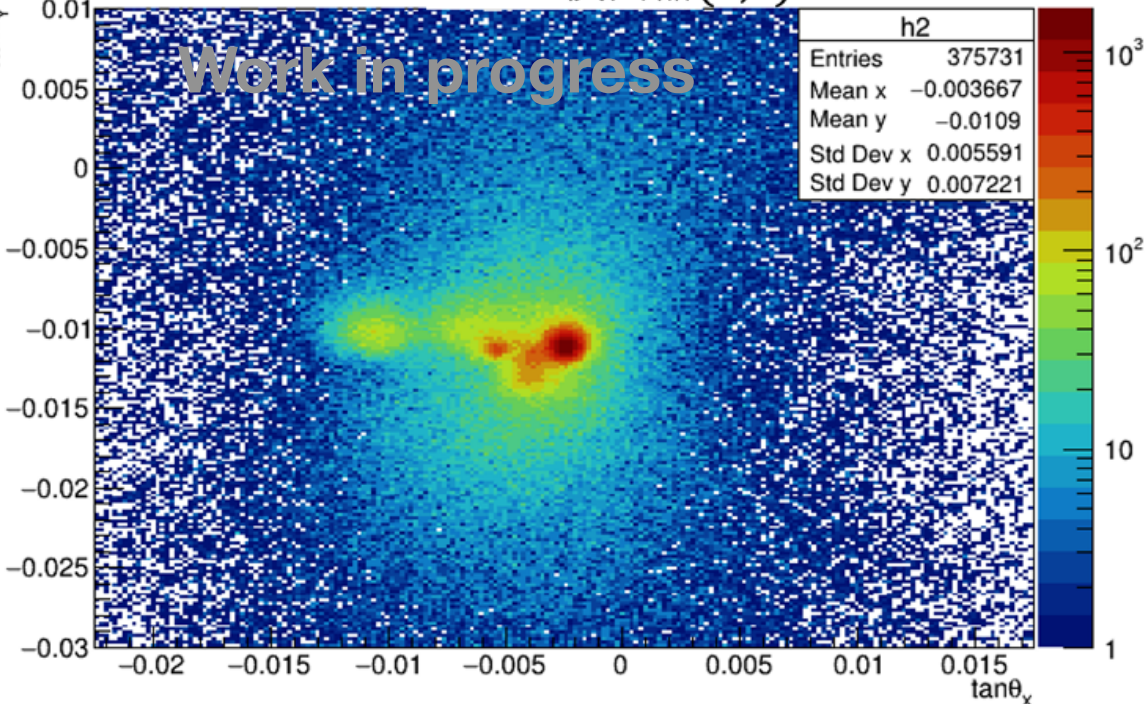
ミュオンモジュール



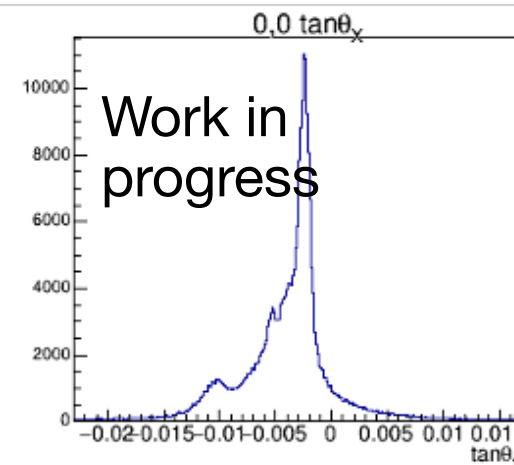
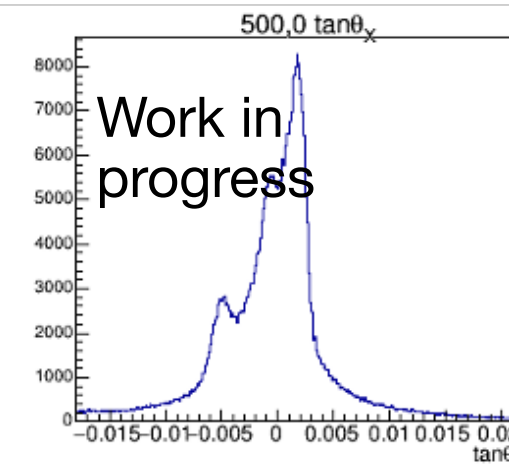
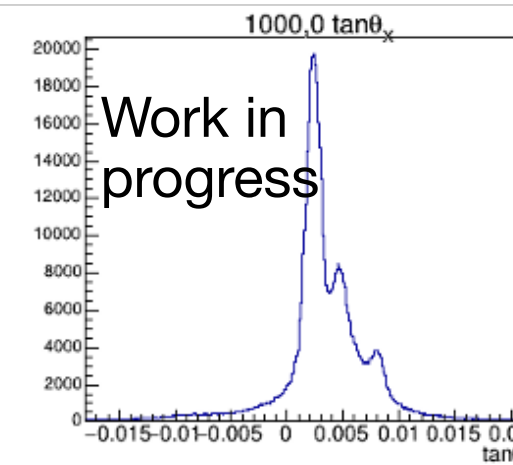
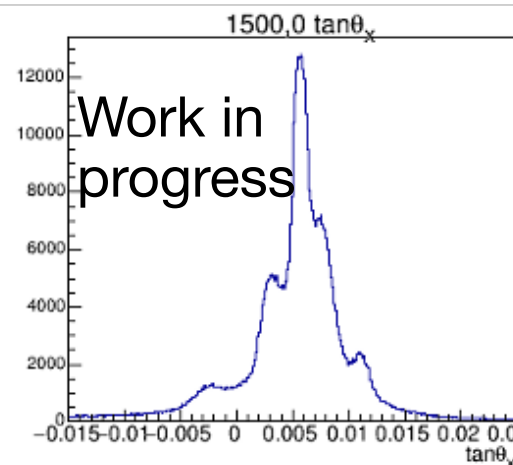
LOS(散乱角0度)での角度分布



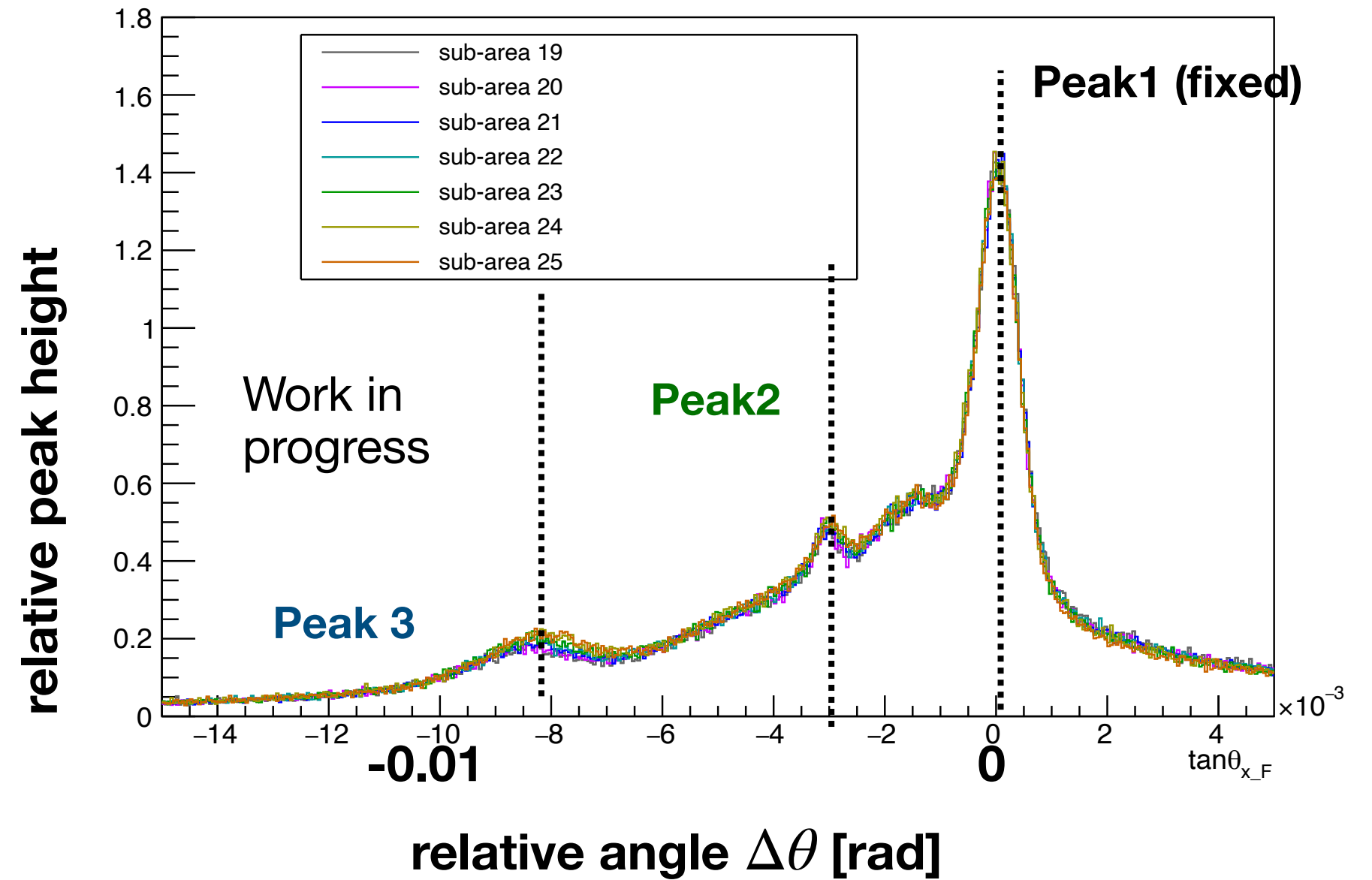
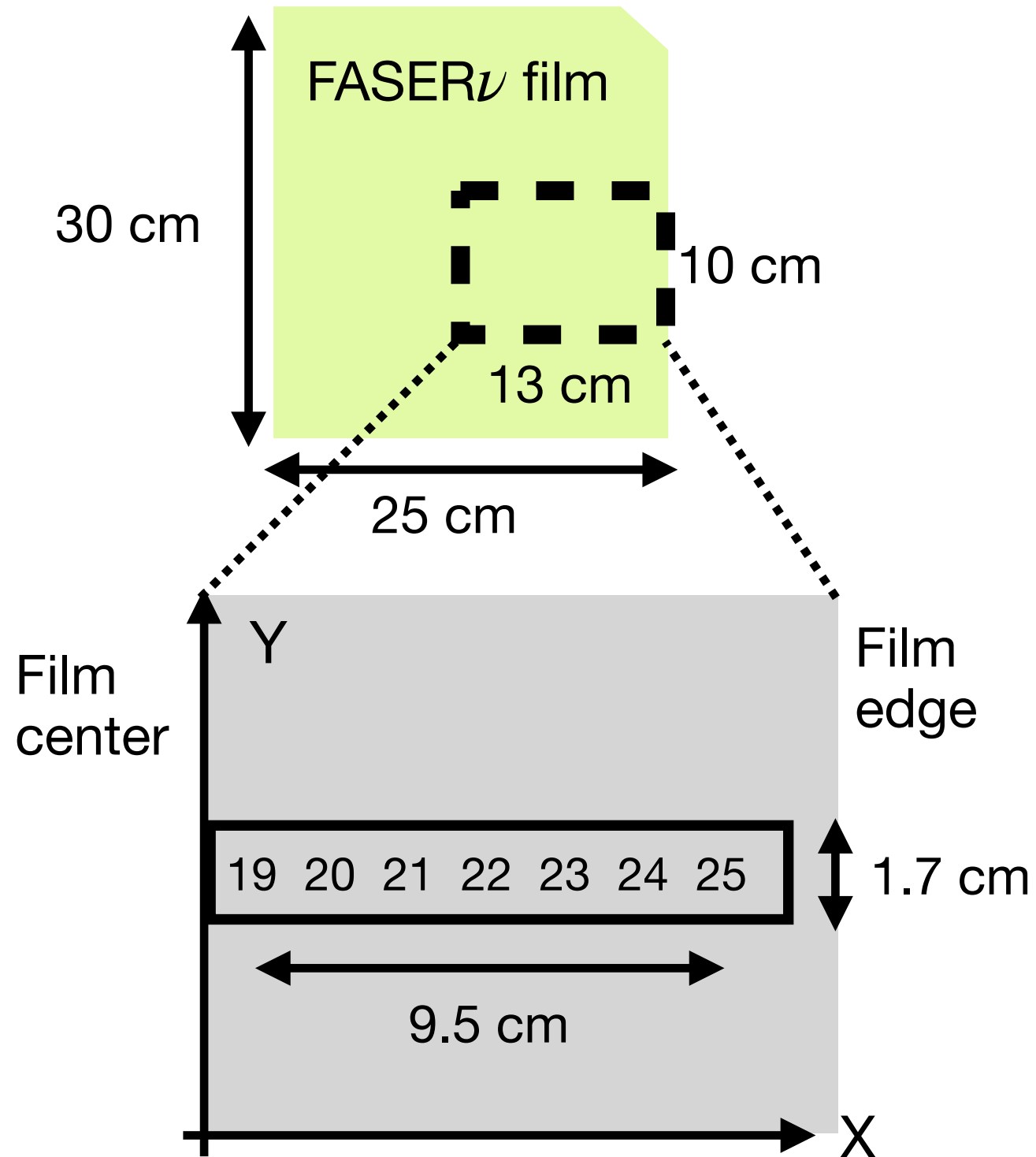
ミュオン検出器(0,0)



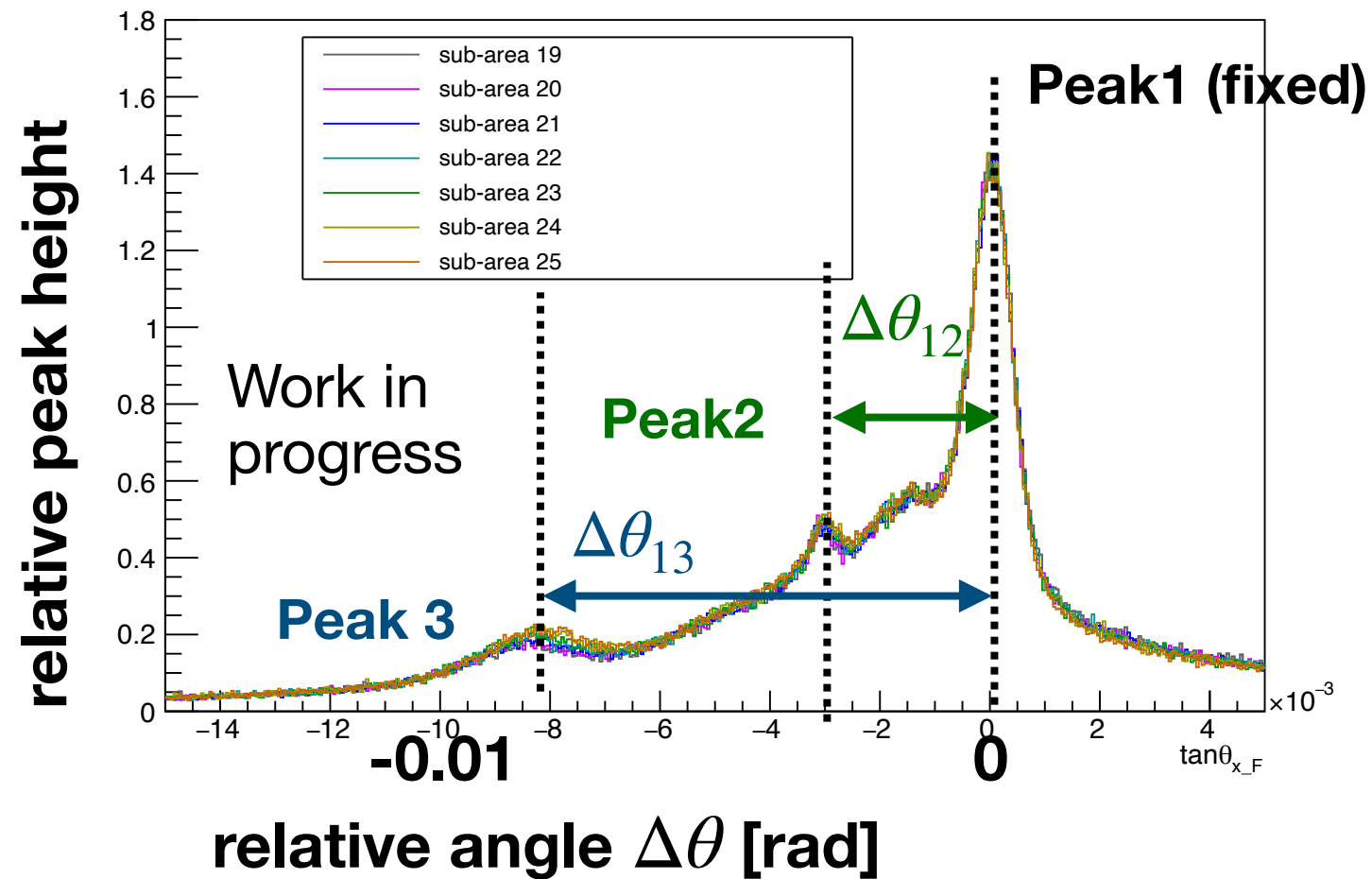
ミュオンモジュール、 $x = 1500, 1000, 500, 0$ mm



角度分布の位置依存性

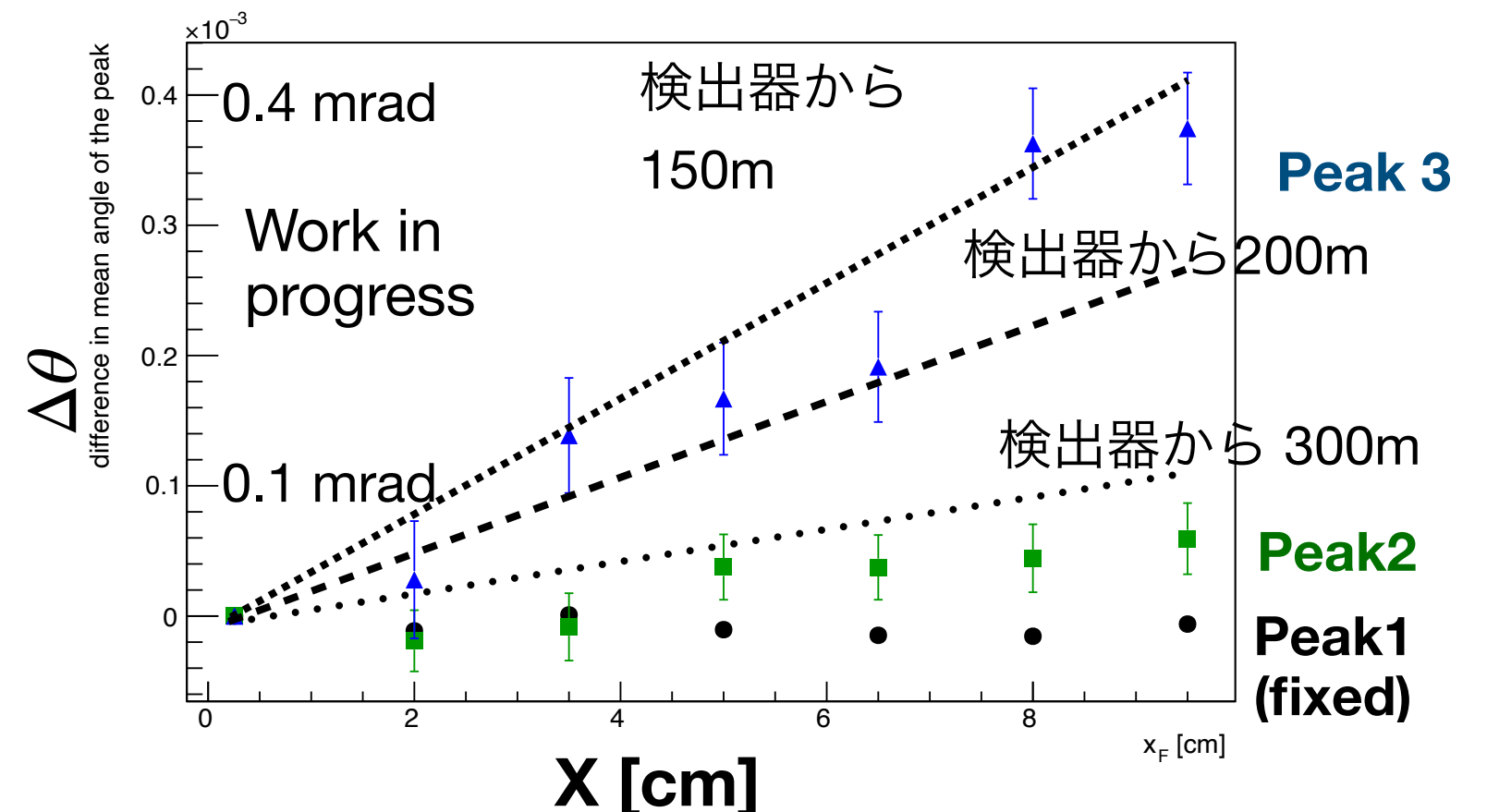


角度分布の位置依存性



角度の絶対値ではなく、一番大きなピークからの相対角度で検討

点線は、一番大きなピークのミュオンが検出器から480m先で生成したと仮定した時に検出器から150m, 200m, 300m先で生成したミュオンの角度変化の想定



今後、シミュレーションとも比較しながらさらに詳しく検討する

FASERnu first analysis

Energy range from simulation

The energy range for σ_{theory} was defined as the energy range containing 68% of reconstructed neutrinos using the baseline models, which is 560–1740 GeV and 520–1760 GeV for ν_e and ν_μ , respectively.

Number of expected events

$$N_{\text{obs}} = \frac{L\rho l}{m_{\text{nucleon}}} \int \sigma(E) \phi(E) \varepsilon(E) dA dE,$$

$$N_{\text{expected}} = 1.68_{-0.37}^{+1.18} (\text{flux})_{-0.36}^{+0.51} (\text{others}) \text{ for } \nu_e$$

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• 中性ハドロン

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• ニュートリノ Neutral Current interaction

$$\nu_e: 0.008_{-0.004}^{+0.013} (\text{flux}) \pm 0.001 (\text{cross - section})_{-0.004}^{+0.007} (\text{others})$$

$$\nu_\mu: 0.045_{-0.005}^{+0.004} (\text{flux}) \pm 0.003 (\text{cross - section})_{-0.024}^{+0.076} (\text{others})$$

FASERnu first analysis

$$L = \frac{P(N_{\text{obs}}|\mu) \cdot \prod_j G(\lambda_j|\mu' = 0., \sigma' = 1.) \cdot \prod_i P(N_{\text{MC found}}^i | n_i^{\text{background}}) G(f^i|\mu' = 1., \sigma' = 1.)}{\mu = N_\nu + \sum_i N_{\text{backgrounds}}^i (f_{\text{syst}}^i, n_i^{\text{background}}) + N_{\text{NC}}^{\text{contamination}}(\lambda_{\text{light hadron}}, \lambda_{\text{charm hadron}}, \lambda_{\text{syst.}}^{\text{NC}})}$$

$$\mu = N_\nu + \sum_i N_{\text{backgrounds}}^i (f_{\text{syst}}^i, n_i^{\text{background}}) + N_{\text{NC}}^{\text{contamination}}(\lambda_{\text{light hadron}}, \lambda_{\text{charm hadron}}, \lambda_{\text{syst.}}^{\text{NC}})$$

$$N'_{\text{expected}}(\lambda) = \begin{cases} N_{\text{center}} \left(\frac{N_{\text{center}} + \Delta N_{\text{lower}}}{N_{\text{center}}} \right)^{|\lambda|} & (\lambda < 0) \\ N_{\text{center}} \left(\frac{N_{\text{center}} + \Delta N_{\text{upper}}}{N_{\text{center}}} \right)^{|\lambda|} & (\lambda \geq 0), \end{cases}$$

$$f_{\text{syst}}(\lambda_{\text{syst}}) = \begin{cases} \frac{N_{\text{center}} + \Delta N_{\text{lower}}}{N_{\text{center}}}^{|\lambda_{\text{syst}}|} & (\lambda < 0) \\ \frac{N_{\text{center}} + \Delta N_{\text{upper}}}{N_{\text{center}}}^{|\lambda_{\text{syst}}|} & (\lambda \geq 0), \end{cases}$$

$$N_{\text{expected}}(\lambda_{\text{light hadron}}, \lambda_{\text{charm hadron}}, \lambda_{\text{syst.}}) = f_{\text{syst}}^{\text{light hadron}}(\lambda_{\text{syst.}}) N_{\text{expected}}^{\text{light hadron}}(\lambda_{\text{light hadron}}) + f_{\text{syst}}^{\text{charm hadron}}(\lambda_{\text{syst.}}) N_{\text{expected}}^{\text{charm hadron}}(\lambda_{\text{charm hadron}}),$$

Table 2: Systematic uncertainties related to the signal expectation.

Source	Relative uncertainty	
	ν_e	ν_μ
Luminosity	2.2%	2.2%
Tungsten thickness	1%	1%
Interactions with emulsions	+3.6% -0%	+3.6% -0%
Flux uncertainty	+70% -22%	+16% -9%
Line of sight position	+2.1% -2.4%	+1.9% -2.5%
Efficiency from hadronization	+22% -5%	+23% -5%
Efficiency from reconstruction	20%	20%
Efficiency from MC statistics	4.9%	2.8%
Total	+70% (flux) -22%	+16% (flux) -9%
	+30% (other) -21%	+31% (other) -21%

The posterior probability distribution of N_ν is calculated by integrating Eq. 17.

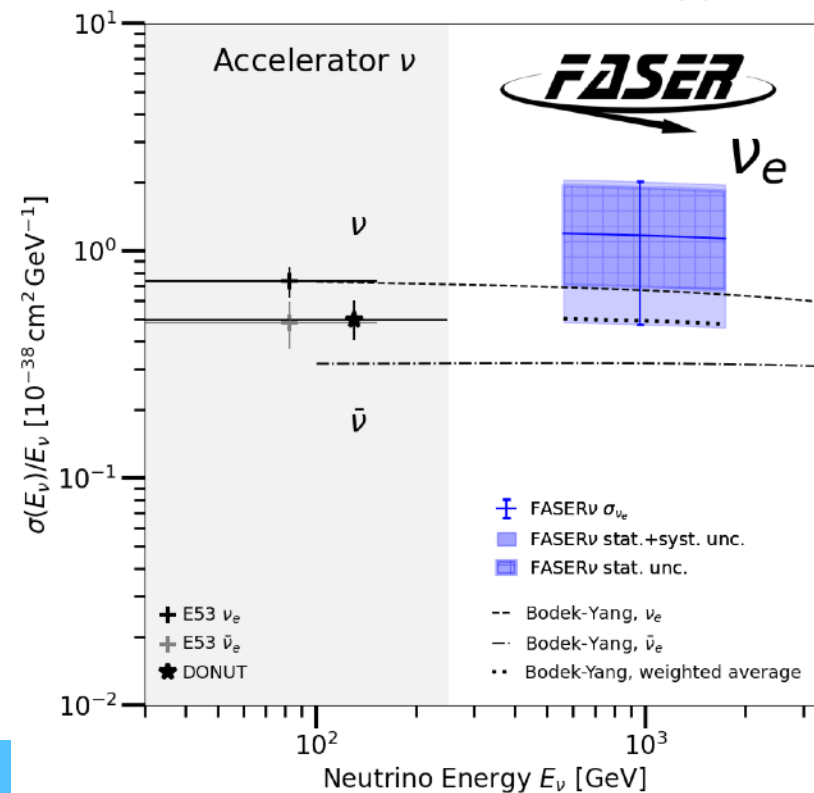
$$P(N_\nu) = \int L(N_\nu, f_{\text{syst}}^i, n_i^{\text{background}}, \lambda_{\text{light hadron}}, \lambda_{\text{charm hadron}}, \lambda_{\text{syst.}}) \pi_{\text{flat}} \prod_i df_{\text{syst}}^i dn_i^{\text{background}} d\lambda_{\text{light hadron}} d\lambda_{\text{charm hadron}} d\lambda_{\text{syst.}}$$

FASERnu first analysis

- $$\sigma_{\text{obs}} = \frac{N_{\nu, \text{obs}}}{N_{\text{expected}}} \sigma_{\text{theory}}$$

- ν_e 候補: 4 イベント, ν_μ 候補: 8 イベント

The α parameter is measured to be $2.4_{-1.3}^{+1.8}$ and $0.9_{-0.3}^{+0.5}$ for ν_e and ν_μ , respectively. The energy-independent part of the interaction cross sections per nucleon, $\sigma_{\text{obs}}/E_\nu$, is measured to be $(1.2_{-0.7}^{+0.8}) \times 10^{-38} \text{ cm}^2 \text{ GeV}^{-1}$ over the energy range of 560–1740 GeV for ν_e and $(0.5 \pm 0.2) \times 10^{-38} \text{ cm}^2 \text{ GeV}^{-1}$ over the energy range of 520–1760 GeV for ν_μ .



$$L = \frac{P(N_{\text{obs}}|\mu) \cdot \prod_j G(\lambda_j|\mu' = 0., \sigma' = 1.) \cdot \prod_i P(N_{\text{MC found}}^i | n_i^{\text{background}}) G(f^i|\mu' = 1., \sigma' = 1.)}{\dots}$$

The posterior probability distribution of N_ν is calculated by integrating Eq. 17.

$$P(N_\nu) = \int L(N_\nu, f_{\text{syst}}^i, n_i^{\text{background}}, \lambda_{\text{light hadron}}, \lambda_{\text{charm hadron}}, \lambda_{\text{syst.}}) \pi_{\text{flat}} \prod_i df_{\text{syst}}^i dn_i^{\text{background}} d\lambda_{\text{light hadron}} d\lambda_{\text{charm hadron}} d\lambda_{\text{syst.}}$$

