

# Non-Thermal Leptogenesis in Peccei-Quinn Inflation

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# What is leptogenesis?

- Leptogenesis is a model that can explain baryon asymmetry of the current universe through the seesaw model

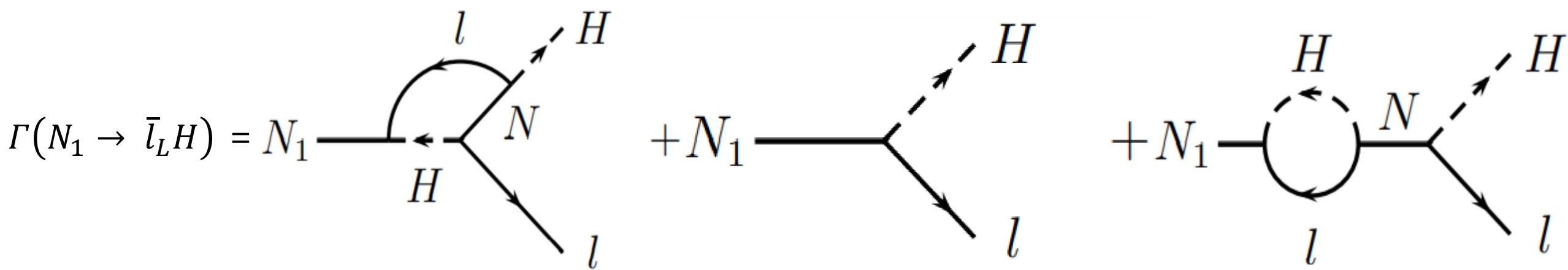
## What is different about using the PQ inflation model?

- Non zero Initial number density of RHN!

# Leptogenesis

$$\mathcal{L} \supset i\bar{N}_R \gamma^\mu \partial_\mu N_R + h\bar{l}_L H N_R + M_N \bar{N}_R^c N_R + h.c$$

Seesaw Type-1



$$i\mathcal{M} = h_{1i} + h_{ij}^* h_{kj} h_{ki} F_N$$

$$\epsilon_1 = \frac{\Gamma(N_1 \rightarrow \bar{l}_L H) - \Gamma(N_1 \rightarrow l_L H^*)}{\Gamma(N_1 \rightarrow \bar{l}_L H) + \Gamma(N_1 \rightarrow l_L H^*)} \simeq -\frac{3}{16\pi} \frac{1}{(hh^\dagger)_{11}} \sum_{j=2,3} \text{Im}[(hh^\dagger)_{1j}^2] \frac{M_{N_1}}{M_{N_j}}$$

$N_1(t_e < t < t_{RH}) \cong 0,$   
 $t_e:$  end of inflation

# Initial number density of RHN

$$N_1(t_e < t < t_{RH}) \simeq 0,$$

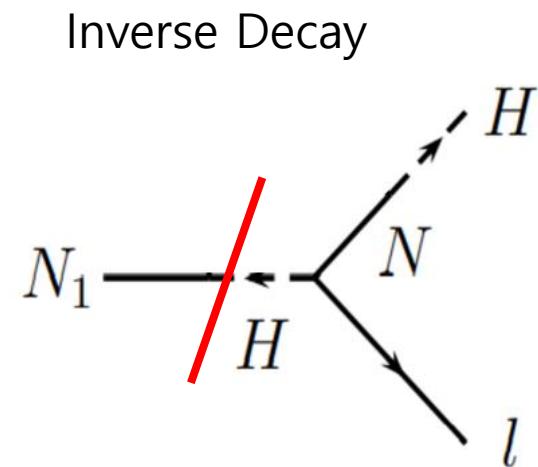
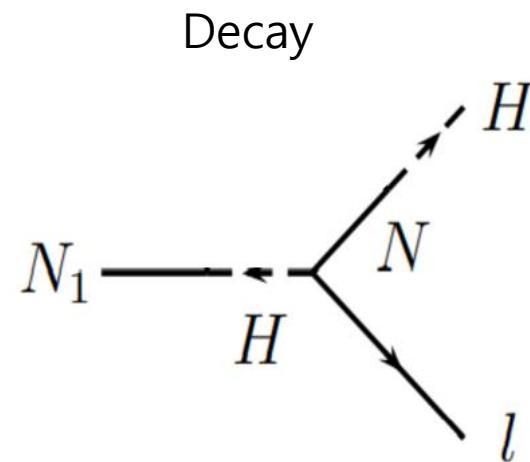
$t_e$ : end of inflation

Thermal leptogenesis scenario

$$N_1(t_e < t < t_{RH}) \neq 0,$$

$t_e$ : end of inflation

Non-Thermal leptogenesis scenario  
(Our assumption)



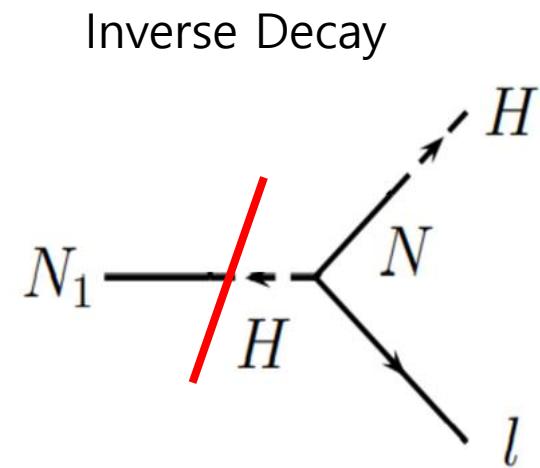
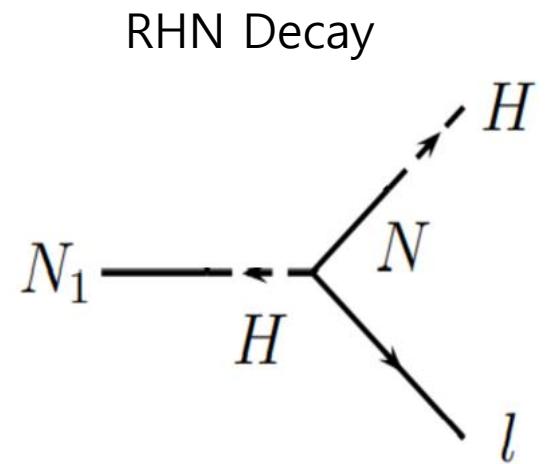
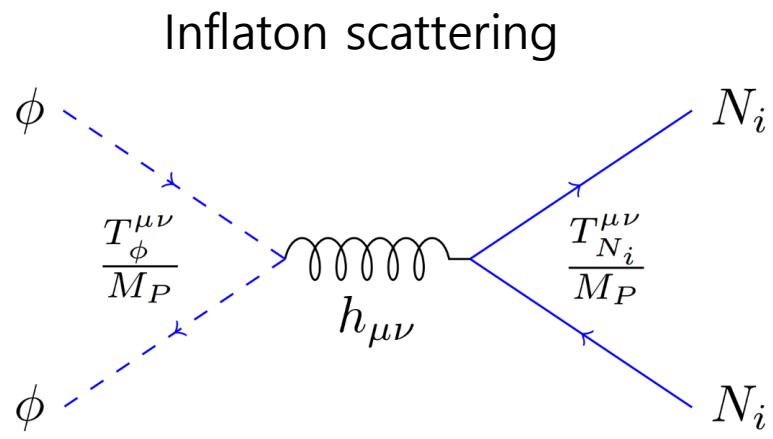
$$M_{RHN} \gg T_{RH}$$

# Initial number density of RHN

$N_1(t_e < t < t_{RH}) \neq 0,$   
 $t_e$ : end of inflation

Non-Thermal leptogenesis scenario  
(Our assumption)

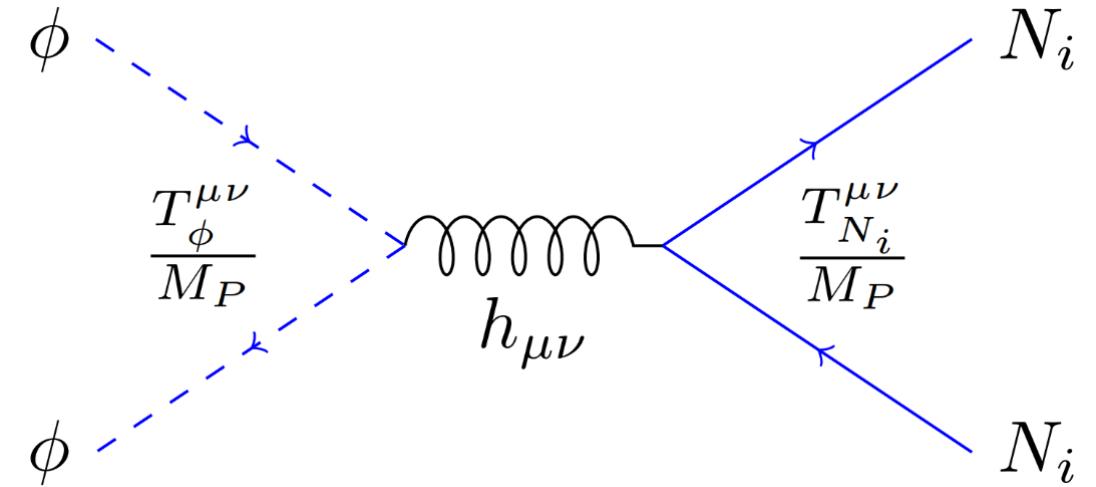
$$M_{RHN} \gg T_{RH}$$



# Gravitational Production of RHN

$$\sqrt{-g}L_{int}^1 = \frac{1}{2M_P} h_{\mu\nu} \left( T_{SM}^{\mu\nu} + T_{\phi}^{\mu\nu} + T_N^{\mu\nu} \right).$$

$$g_{\mu\nu} \cong \eta_{\mu\nu} + \frac{h_{\mu\nu}}{M_P}$$



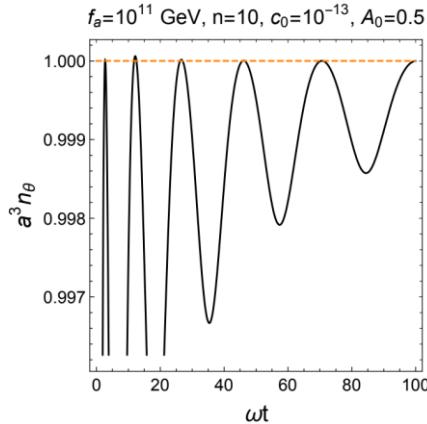
$$|\mathcal{M}_n^{\phi^k}|^2 = \frac{2\rho_\phi^2}{M_P^4} \frac{m_N^2}{s} \left[ 1 - \frac{4m_N^2}{s} \right] |(\mathcal{P}^k)_n|^2$$

$$\frac{dY_{N_1}^{\phi^k}}{da} = \frac{\sqrt{3}M_P}{\sqrt{\rho_{RH}}} a^2 \left( \frac{a}{a_{RH}} \right)^{\frac{3k}{k+2}} R_{N_1}^{\phi^k}(a)$$

# Peccei-Quinn inflation

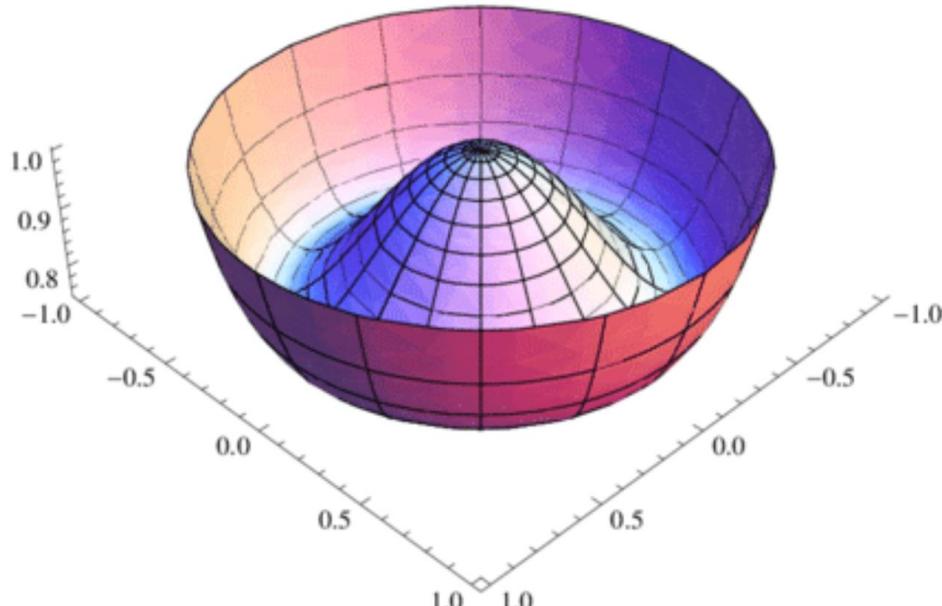
$$U(1)_{PQ}$$

$$\Phi = \frac{1}{\sqrt{2}} \rho e^{i\theta}$$



$$\frac{\mathcal{L}_E}{\sqrt{-g_E}} = -\frac{1}{2}M_P^2 R + \frac{1}{2}(\partial_\mu \phi)^2 + 3M_P^2 \sinh^2 \left( \frac{\phi}{\sqrt{6}M_P} \right) (\partial_\mu \theta)^2 - V_E(\phi, \theta)$$

$$V_E(\phi, \theta) = V_{\text{PQ}}(\phi) + V_{\text{PQV}}(\rho, \theta)$$



$$V_{\text{PQ}}(\phi) = V_0 + \frac{1}{4}\lambda_\Phi \left( 6M_P^2 \tanh^2 \left( \frac{\phi}{\sqrt{6}M_P} \right) - f_a^2 \right)^2$$

$$V_{\text{PQV}}(\rho, \theta) = 3^{n/2} M_P^4 \tanh^n \left( \frac{\phi}{\sqrt{6}M_P} \right) \sum_{k=0}^{[n/2]} |c_k| \cos \left( (n-2k)\theta + A_k \right)$$

# Interaction

Sub-dominant about Reheating

$$\mathcal{L}_{Q,\text{int}} = -y_Q \Phi \bar{Q}_R Q_L + \text{h.c.}$$

$$\mathcal{L}_{\text{gluons}} = \frac{g_s^2}{32\pi^2} \left( \bar{\theta} + \xi \frac{a}{f_a} \right) G_{\mu\nu}^a \tilde{G}^{a\mu\nu}$$

Due to Strong CP problem!

dominant about Reheating

$$\Delta V_E = \lambda_{H\Phi} |\Phi|^2 |H|^2$$

$$\mathcal{L}_{\text{int}} = -\frac{1}{2} \lambda_\Phi \phi^2 a^2$$

$$\frac{\Gamma_{\phi\phi \rightarrow aa}}{\Gamma_{\phi\phi \rightarrow HH}} \simeq \frac{\lambda_\Phi^2}{2\lambda_{H\Phi}^2}$$

$$\lambda_{H\Phi} \gtrsim \frac{1}{\sqrt{2}} \lambda_\Phi \quad \text{Reheating condition}$$

# Matching

## Constraints

1.  $M_{N_1} = 10^{13} \text{ GeV}$     2.  $\epsilon_1 = 10^{-6}$     3.  $f_a = 10^{11} \text{ GeV}$

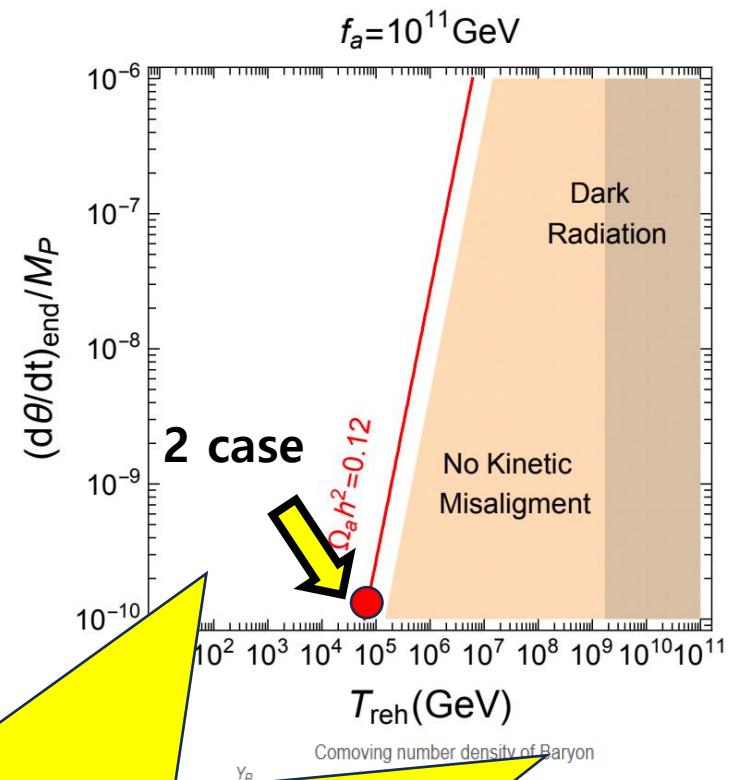
$$Y_B = \frac{n_B}{s} = \frac{28}{79} \epsilon_1 \frac{n_{N_1}(T_{RH})}{s} \quad (s = \frac{2\pi^2 g_{RH} T_{RH}^3}{45}, \text{ entropy density})$$

$$Y_B^{\text{Planck}} \cong 8.7 \times 10^{-11}$$

**1 case :**  
 $\lambda_{H\Phi} = 10^{-10.57}$   
 $T_{reh}: 6.0 * 10^4 \text{ GeV}$   
initial axial velocity  
 $: 10^{-6} M_p$

**2 case :**  
 $\lambda_{H\Phi} = 10^{-10.6}$   
 $T_{reh}: 5.5 * 10^4 \text{ GeV}$   
initial axial velocity  
 $10^{-10} M_p$

baryon asymmetry, dark matter abundance, strong CP problem could be solved!



Thank you