

Advances in Resummed Calculations

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In processes involving disparate scales $Q \gg Q_0$, higher-order corrections are enhanced by large logarithms

 $\alpha_s^n \ln^m Q/Q_0$

which can spoil perturbative expansion. Maximum power of logarithms depends on problem

- Single logarithmic: $m \leq n$
- Sudakov (soft + collinear): $m \le 2n$

Resum enhanced contributions to all orders.

- Count $\ln(Q/Q_0) \sim 1/\alpha_s$
- Systematic expansion: LL, NLL, NNLL, ...



Classic example is Z- or W-production at low transverse momentum $q_T \ll M_Z$: $L = \ln(M_Z/q_T)$ Need all-order resummations for reliable extractions of α_s and M_W from the spectrum.

→ talks by Florencia Castillo, Valentina Guglielmi, Oleg Kuprash, Giulia Marinelli

Exponentiation

One can show that cross section has the form

 $\Sigma(p_T) = \exp\left(L g_1(\alpha_s L) + g_2(\alpha_s L) + \alpha_s g_3(\alpha_s L) + \alpha_s^2 g_4(\alpha_s L) + \dots\right)$



Accuracy:

• LL: *g*₁; NLL: *g*₁, *g*₂; NNLL: *g*₁, *g*₂, *g*₃

Expand in α_s but count $\alpha_s L$ as O(1)



Many types of scale hierarchies, many different types of resummations ... and by now many different EFTs



Large logarithms arise due to soft and collinear emissions.

Resummations are based on the factorization of cross sections in **soft** and **collinear** limits.

This factorization is at the heart of collider physics and has implications in all its areas.



Insights into all-order structure of cross sections have lead to progress in all of the above areas.



These connections are exemplified by the many scientific achievements of Stefano Catani ('58-'24).

A pioneer of resummation, but his deep insights into soft and collinear dynamics led him to contribute to all areas of collider QCD.

Outline

- High-precision resummations
 - *q*_T spectrum in Drell-Yan production; *α_s* and *M_W* determinations from ATLAS and CMS
- New observables for jet substructure
 - energy-energy correlators and α_s and m_t determination
- Resummation of jet observables and parton showers
 - non-global logs, clustering logs, super-leading logs
- Back to the basics
 - all-order structure of wide-angle scattering
 - collinear factorization violation vs. PDF factorization



 α_s and m_t from precision resummations of transverse momentum spectra

$pp \rightarrow$ "EW bosons" + X at low q_T



hard function: Born + virtual corrections

- Ingredients known to very high accuracy
 - three-loop beam functions Ebert, Mistlberger, Vita '20
 - three-loop hard functions for Z/W/γ (with singlet contributions Gehrmann, Primo '21 with top mass Chen, Czakon, Niggetiedt '21), two-loop for diboson processes
 - four-loop hard anomalous dimensions Manteuffel, Panzer, and Schabinger '20 and anomaly exponent Duhr, Mistlberger, Vita '22; Moult, Zhu, Zhu '22
 - fixed-order matching to α_s^3 from MCFM and NNLOJet

Almost N4LL accuracy

hard function: Born + virtual corrections



- Missing for full N⁴LL + accuracy
 - four-loop PDF evolution, i.e. N³LO PDFs. Note: approximate N³LO PDF exist \rightarrow talks by Sven Moch and Tongzhi Yang
 - five-loop cusp anomalous dimension (likely numerically irrelevant)



- aN⁴LL resummations from several groups with different formalisms
- Results include α_{s^3} fixed order from MCFM

Comparison and uncertainties

Resummed computations are performed in a variety of (equivalent) formalisms and with different of scheme choices

- Scale setting in momentum space (CuTe, Radish) versus impact parameter space (everyone else)
- Different formalisms for rapidity logs (CSS, collinear anomaly, RRG) and associated uncertainty
- Different matching schemes / transition to fixed order

Uncertainty estimates are much less standardized than for fixedorder computations!

• Ongoing comparison/benchmark efforts by LHC EW subgroup (since 2018, to be completed this year!)

ATLAS α_s extraction



- Reconstruct inclusive spectrum rate from angular coefficients
- α_s from fit to **DYTurbo**
- N³LO fixed order MCFM and NNLOJet
- MSHT20 approximate N³LO PDFs
 - cross checks with NNLO sets
- Non-perturbative effects based on two-parameter ansatz by Collins Rogers '14



With these high-order resummed and matched computations, we have entered a new regime of precision collider calculations.

Unprecedented precision, but also difficult to be sure the uncertainties are reliably estimated ... no previous experience with this level precision at hadron colliders!

CMS Mwextraction



- W-mass from Jacobian peak in charged-lepton p_T spectrum and rapidity distribution.
- Reweigh MiNNLO_{PS} with resummation from SCETlib at N³LL, α_s² fixed-order matching through DYTurbo
- Higher-order coefficients in resummation as theory nuisance parameters Tackmann, unpublished → talk by Giulia Marinelli
 - use statistical model for their distribution; can capture correlations since same coefficient enters different observables
 - values extracted from fit to data



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Theoretical uncertainties before and after fit of nuisance parameters to data using W measurements or W and Z measurements.

- Crucial new element of analysis is that theoretical predictions are fit to data to significantly reduce their uncertainties.
 - CMS validates this procedure on Z-production, but does not use Z-data for W-mass extraction
 - Effects not accounted for in the original theory prediction (higher-order electroweak corrections, quark masses,...)?

N-jettiness event shape

- 3-loop soft function for 0-jettiness Baranowski, Delto, Melnikov, Pikelner, Wang '24
 - Also 3-loop beam functions are known Ebert, Mistlberger, Vita '20, Baranowski, Behring, Melnikov, Tancredi, Wever '22 → all ingredients for 3-loop jettiness slicing available
- New precise representations of the two-loop *N*-jettiness soft function Bell, Dehnadi, Mohrmann, Rahn'23; Agarwal, Melnikov, Pedron '24
- N³LL resummation of 1-jettiness for *Z*-boson plus jet in Geneva Alioli, Bell, Billis, Broggio, Dehnadi, Lim, Marinelli, Nagar, Napoletano, Rahn '23

Goal of is not comparison with measurements, but use for fixed-order computations and simulations of jet processes (slicing,...)



illustration from Kyle Lee

Energy-energy correlators and jet substructure



A lot of new interesting developments in using these energy $\frac{dz}{\mathcal{E}(n)} = \int_{0}^{\infty} \frac{d^2}{dt} \lim_{t \to 0} r^2 n^i T_{0i}(t) \int_{0}^{\infty} \frac{d^2}{\mathcal{E}(n)} \int_{0}^{\infty} \frac{d^2}{dt} \lim_{t \to 0} r^2 n^i T_{0i}(t) \int_{0}^{\infty} \frac{d^2}{dt} \int_{0}^{\infty} \frac{d^2}$

- weighted by energy: insensitive to soft radiation: for some source operator \mathcal{O} . $\frac{dz}{\mathcal{E}(\vec{n})} = \int_{\vec{n}} \frac{dz}{dt} \lim_{\vec{n}} r^2 n^i T_{0i}(t, t)$
- factorization, light-ray OPE, CFortecian for the source operator of the program of the source operator operator of the source operator operator operator of the source operator operator of the source operator op
 - receivent ides dispendents in the study of ANEC operators,
 - concrete situation for studying the behavior of $ANEC_{n_1} \mathcal{E}(n_2) \mathcal{E}(n_1) \mathcal{E}(n_2) \mathcal{E}(n_1) \mathcal{E}(n_2) \mathcal{E}($

There has recently been $\operatorname{significant}$ progress in the number of different directions. For generic angles, the for some source presentor \mathcal{O} of this provides a connection by

Energy-Energy Correlator

Simplest example is the two-point function

$$\operatorname{EEC}(\chi) = \sum_{a,b} \int \mathrm{d}\sigma_{e^+e^- \to a+b+X} \, \frac{E_a E_b}{Q^2} \, \delta(\cos \chi_{ab} - \cos \chi)$$

large logarithms both for small and large angles. For large angle N⁴LL is known!



Duhr, Mistlberger, Vita '22

Ingredients:

- 3-loop jet functions Ebert, Mistlberger, Vita '20
- 4-loop rapidity anomalous dimension Duhr, Mistlberger, Vita '22, Moult, Zhu, Zhu '22.
- four-loop hard anomalous dimensions Manteuffel, Panzer, and Schabinger '20; Lee, Manteuffel, Schabinger, Smirnov, Smirnov, and M. Steinhauser '22.
- four-loop cusp Henn, Korchemsky, Mistlberger '19; Manteuffel, Panzer, and Schabinger '20 + ... 5-loop cusp is missing, estimated to have very small effect.



TOP $G^{\text{FIG. 3}}_{\text{num with }m} = 2.14 \text{ GeV and } n_t^{\text{gole}} = 3.34 \text{ GeV.}$

approach at 1-loop. Decent fits are





 $8 \times 10^{29} \ c_{\tau} \tau_{0.001\xi_{t}} = 0.0001\xi_{t}^{2}$

(12)

FIG. 4. Boundaries of absolute stability (lower band, NLO) and metastability (upper line, LO). The thickness of the lower boundary indicates perturbative and α_s uncertainty.

goes to zero, (2) that Eq. (5) have no solution, and (3)

that $V_{\rm min}$ goes positive when $\Lambda_{\rm NP} = M_{\rm Pl}$ all give nearly identical boundaries in the $m_h^{\rm pole}/m_t^{\rm pole}$ plane. Know-

ing that quantum gravity is relevant at $M_{\rm Pl}$, we should

therefore be cautious about giving too strong of an in-



y is 99% m_W^2 m_t^2 equilateral iant valı = p_t^2 tted squeezed e by $2.88 \times$ $2.40 \times$ can onds to a us, the lard ιbilby quai ded in o tted tability res- $> M_{\rm Pl}$ [dly, is gau ⁷min

h. An alternative criterion is $\mathcal{O}_6 \equiv \frac{1}{\Lambda_{\rm NP}^2} h^6$ be comparable

to V_{\min} when $h = \langle h \rangle$. Although \mathcal{O}_6 and V_{\min} are gaugeinvariant, the value of \mathcal{O}_6 at the field value h where the minimum occurs *is* gauge dependent, so this condition

Proposals to the Aver and Satisfactory of the Proposals to the Aver and Satisfactory of the Aver and Satisfactory of the Average and the prevaled to the preva terpretation of the perturbative absolute stability bound in the St. Weald show in this filt to net stability added to the classical theory and its effect on $V_{\min}^{\forall \forall}$ evalbound, that the lifetime of our vacuum be larger than the age of the universe. At lowest order this translates to $\lambda(\frac{1}{R}) = \sqrt{-14.53} + 0.153$ by 1000 and 1000 by 1000 b top decays Hode up to the Moult the that the ake, Schwarz '23 tential is still negative at its minimum in the SM even taking $\Lambda_{\rm NP}' = M_{\rm Pl} = 1.22 \times 10^{19}$ GeV, we find that $\lambda(\mu)$ is gauge invariant, so is this criterion. Although for the Standard Model this approximation is probably sufficient, it has not been demonstrated that the bound can be systematically improved in a guage-invariant way [31]. $\mu_X^{\text{min}} = 6.0 \times 10^{17} \text{ GeV} \text{ and } V_{\text{min}} = -(1.1 \times 10^{17} \text{ GeV})^4.$ In this paper, we have only discussed a single physical Comparing to Eq. (13) we see that the energy of the true feature of the effective action: the value of the effective

vacuum is very Planck-sensitive.

More generally, a good fit is given by

 $V_{\rm min} = -(0.01 \,\Lambda_{\rm NP})^4, \qquad \Lambda_{\rm NP} \gtrsim 10^{12} \,{\rm GeV}$ (14)

When $\Lambda_{\rm NP} < 3.6 \times 10^{12}$ GeV, $V_{\rm min}$ becomes positive and for $\Lambda_{\rm NP} < 3.1 \times 10^{12}$ GeV the maximum and minimum potential at its extrema. There is of course much more content in the effective action especially othen temperature dependence is included. Unfortunately, many uses of the effective action involve evaluating it for particular field configurations, a procedure that has repeatedly been shown to be gauge-dependent. For example, the

Many new ideas and results

- EECs for *b* and *c*-quarks Lee, Mecaj Moult '22
- Non-Gaussianities in collider energy flux Chen, Moult, Thaler, Zhu '22
- Nucleon energy correlators Liu, Zhu '22, Cao, Liu, Zhu '23
- TMDs from Semi-inclusive Energy Correlators Liu, Xhu '24
- EECs for nuclear matter at the electron-ion collider (EIC) Devereaux, Fan, Ke, Lee, Moult '23
- EECs for studying the quark-gluon plasma Andres, Dominguez, Holguin, Marquet, Moult '23, '24; Liu, Liu, Pan, Yuan and Zhu '23
- Non-perturbative effects in EECs Schindler, Stewart, Sun '23; Lee, Pathak, Stewart, Sun '24; Chen, Liu, Ma '24, Chen, Monni, Xu, Zhu '24;
- v-point energy correlators Budhraja, Chen, Waalewijn '24
- Higgs decay EECs Yang, Zhang '24
- EECs on tracks Lee, Moult '23; Jaarsma, Moult, Waalewijn, Zhu '23
- N³LL for transverse EEC in back-to-back limit Gao, Li, Moult, Zhu '23

• ...



Resummation for jet processes: non-global, clustering and super-leading logarithms

Traditional resummation methods limited to a small set of simple, inclusive (``global'') observables.

Any observable with angular cuts is non-global:

- isolation cones (e.g. in photon production)
- exclusive jet cross sections
- gaps beween jets, veto regions

A lot of progress in resumming logarithms in more complicated observables.







Non-global logs (NGLs)

- soft gluons from secondary emissions
- QCD only

Dasgupta, Salam '00

Clustering logs (CLs)

- phase-space constraints of new emissions depends on all existing partons
- even in QED

Appleby, Seymour '02

Super-leading logs (SLLs)

- Glauber phases spoil collinear cancellations
- QCD only
- hadron-hadron colliders only

Forshaw, Seymour '06

Super-Leading Logs (SLLs)

Forshaw, Kyrieleis, Seymour '06 '08

Consider gap between jets at hadron collider, cone around beam direction



Large logarithms $\alpha_s^n L^m$ with $L = \ln(Q/Q_0)$

- e^+e^- : $m \le n$, leading logs m = n
- $pp: \alpha_s L, \ \alpha_s^2 L^2, \ \alpha_s^3 L^3, \ \alpha_s^4 L^5 \dots, \ \alpha_s^{3+n} L^{3+2n}$

LL effect, but vanishes in large- N_c limit!

By now, we have all-order factorization theorems, which enable resummation for

- NGLs in e^+e^- TB, Neubert, Rothen, Shao '16
- SLLs TB, Neubert, Shao '21 + Stillger '23
- CLs TB, Haag '23
- New results
 - First resummations of SLLs TB, Neubert, Shao '21
 + Stillger '23 and Glauber phases Böer, Neubert, Stillger '23 + Hager, Xu '23, '24
 - Resummation of subleading NGLs TB, Schalch, Xu '23

Factorization for gaps between jets



Hard functions m hard partons along fixed directions {n₁, ..., n_m} $\mathcal{H}_m \propto \mathcal{M}_m \rangle \langle \mathcal{M}_m |$

Soft + collinear function squared amplitude for *m* Wilson lines +collinear fields

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RG evolution

Renormalized hard functions fulfill RG equation

$$\frac{d}{d\ln\mu} \mathcal{H}_m = -\sum_{l=m_0}^m \mathcal{H}_l \Gamma_{lm}^H$$
matrix in multiplicity

and color space

One-loop hard anomalous dimension:



In this framework resummation is obtained by solving the associated RG equations.

Challenge: $\Gamma^{(1)}$ is infinite matrix in the space of particle multiplicities and colors!

- NGLs: implemented
 I⁽¹⁾ and
 I⁽²⁾ in the large N_c limit into MC framework MARZILI to solve RG equation numerically
- SLLs: need full color to see effect. Compute leading SLLs order by order, sum up series → talk by Philipp Böer



Monte Carlo code to solve the RG equation in the large N_c limit

- LL equivalent to Dasgupta-Salam shower. Full-color LL is available Hatta, Ueda '13; De Angelis, Forshaw and Plätzer '20
- NLL has one insertion of $\Gamma^{(2)}$, which includes double-real, realvirtual, and purely virtual terms

resummation with a parton shower

Subleading NGLs at the LHC

TB, Schalch, Xu '23

Z-production with veto $p_T < Q_0$ on radiation in rapidity slice around Z. Compute fraction of $R(Q_0)$ of events which pass the veto.



Corrections scale as $\mathcal{O}(\alpha_s^2)$ or $\mathcal{O}(\alpha_s/N_c^2)$ First NGL resummation at this accuracy level!

First resummations for **SLLs**



- Small effects for $pp \rightarrow Z/H$, $pp \rightarrow Z/H + j$, but sizable for dijet production TB, Neubert, Shao '21 + Stillger '23
 - Similar-size effects in hadronic cross sections, but will need to combine SLL+NGLs to assess phenomenolgical significance
- By now also resummation including higher Glauber terms Böer, Neubert, Stillger '23 + Hager, Xu '24 and running coupling Böer, Hager, Neubert, Stillger, Xu '24 is available.

→ talk by Philipp Böer



Connection to parton showers

In the past, not much cross talk between parton shower MCs and resummation

- analytical
- very simple observables
- NⁿLL accuracy (by now up to n=4!)
- exact color
- non-perturbative matrix elements, fits

- numerical
- fully general
- LL + many subleading effects + tuning
- large- N_c limit + some
- hadronization models

As discussed, resummation is being extended to more complicated observables and is using MC methods...

... and parton shower are moving to higher accuracy!



PanScales shower: van Beekveld, Dasgupta, El-Menoufi, Ferrario Ravasio, Hamilton, Helliwell, Karlberg, Monni, Salam, Scyboz, Soto-Ontoso, Soyez '24

Development of parton showers which systematically include higher-log effects PanScales, Alaric, $\dots \rightarrow talk$ by Alexander Karlberg

parton showers with resummation



De Angelis, Forshaw, Plätzer '21

Amplitude evolution, development of full color showers. Deductor Nagy, Soper, CVolver Plätzer, Sjodahl, De Angelis, Forshaw, Holguin, ...

→ talks by Simon Plätzer and Fernando Torre González



Resummation of subleading soft logarithms in jet processes using MC method: Gnole Banfi, Dreyer, Monni '21 Marzili TB, Schalch, Xu, '23 ...

... and same result from PanScales MC Ferrario Ravasio, Hamilton, Karlberg, Salam, Scyboz, Soyez '23.

Numerical agreement to better than 1% among the three approaches.



- 1. Momentum regions in wide-angle scattering
- 2. Collinear factorization violation and PDF factorization?

1.) New insights into the MoR

Soft-Collinear Effective theory is based on method of regions (MoR) expansion of loop (and phase-space) integrals.

Long-standing open question whether usual soft and collinear regions are sufficient to all orders?

- Yes, for massless wide-angle scattering! Gardi, Herzog, Jones, Ma, Schlenk, '22; proof: Ma '23
- but proof only applies to "facet regions" and Gardi, Herzog, Jones, Mao '24 have identified a set of "inside regions" for special topologies.

Gardi, Herzog, Jones, Mao '24 conjecture



Hidden inside region correspond to multiple hard scattering. Compatible with SCET and power suppressed in QCD.

Hidden regions in forward scattering correspond to Glauber modes.

2.) PDF Factorization vs SLLs



Scale separation

- perturbative hard-scattering $\hat{\sigma}_{ij}$ at scale Q
- non-perturbative PDFs $f_i(x)$ at scale Λ_{QCD}
- No low-energy interactions between incoming hadrons
 - cancellation of soft and Glauber physics CSS '85 (for DY)
- Purely collinear, single logarithmic DGLAP evolution

Collinear Factorization Violation

Catani, de Florian, Rodrigo '11; Forshaw, Seymour, Siodmok '12;



New results for **Sp** Henn, Ma,Xu, Yan, Zhang, Zhu '24 Guan, Herzog, Ma, Mistlberger, Suresh '24

For space-like collinear limit $1 \parallel j$ the splitting amplitude Sp depends on the colors and directions of the partons not involved in the splitting!

• Related to non-cancellation due to soft phases

Implications for PDF factorization?

SLLs vs DGLAP

- Double-logarithmic SLLs directly related to collinear factorization breaking
- SLLs generated from double logarithmic running
- If PDF factorization holds, something interesting must happen at scale Q₀ which converts between the two evolutions



EFT analysis

TB, Hager, Jaskiewicz, Neubert, Schwienbacher '24

In our effective theory framework questions about PDF factorization on previous slides can be formulated concisely and answered

- Use RG to predict the form of the low-energy matrix elements \mathcal{W}_m required for consistency with PDF factorization
 - Show that form of \mathscr{W}_m requires soft-collinear interactions at three-loop order
- Perform a systematic method of region analysis of the relevant diagrams contributing to \mathcal{W}_m
 - Identify a hidden active-active Glauber region, well defined in dimensional regularization
 - Computing this Glauber contribution, we find that it indeed has the required form to convert the double logarithmic SLL to single-log running

Collinear factorization breaking at $\mu = Q$

soft-collinear factorization breaking by Glauber modes at $\mu = Q_0$

PDF factorization for $\mu < Q_0$

"factorization restoration"

Х

TB, Hager, Jaskiewicz, Neubert, Schwienbacher '24

Note: Analysis is a consistency check at 4-loop order, not a factorization proof. Nevertheless remarkable that factorization survives; all elements for the breaking are present at this order.

Summary

- High-precision resummations up to N⁴LL for q_T spectra, but experiment is far ahead of theory!
- Energy-energy correlators are a promising new class of LHC observables
 - many new ideas and new measurements
- Resummations for logarithms in jet observables are becoming available: CLs, NGLs, SLLs
 - new insights into soft-collinear interplay and factorization
- New partons showers with higher-logarithmic resummations

Extra Slides



Source of uncertainty	Impact (MeV)	
	Nominal	Global
Muon momentum scale	4.8	4.4
Muon reco. efficiency	3.0	2.3
W and Z angular coeffs.	3.3	3.0
Higher-order EW	2.0	1.9
$p_{\rm T}^{\rm V}$ modeling	2.0	0.8
PDF	4.4	2.8
Nonprompt background	3.2	1.7
Integrated luminosity	0.1	0.1
MC sample size	1.5	3.8
Data sample size	2.4	6.0
Total uncertainty	9.9	9.9

