Parton Showers

Simon Plätzer Institute of Physics — NAWI, University of Graz Particle Physics — University of Vienna

At the QCD@LHC workshop Freiburg | 7 October 2024

UNIVERSITÄT GRAZ UNIVERSITY OF GRAZ







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Synthesis of Enantiopure Sulfoxides by Concurrent Photocatalytic Oxidation and Biocatalytic Reduction

Sarah Bierbaumer, Dr. Luca Schmermund, Alexander List, Dr. Christoph K. Winkler 🗙 Dr. Silvia M. Glueck















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https://particle.uni-graz.at/en/event-generators-anotherset Natural Sciences









Setting the scene



$d\sigma \sim L \times d\sigma_H(Q) \times PS(Q \to \mu) \times MPI \times Had(\mu \to \Lambda) \times \dots$



Setting the scene



100's of GeV I-2 GeV

few 100's MeV $d\sigma \sim L \times d\sigma_H(Q) \times PS(Q \to \mu) \times MPI \times Had(\mu \to \Lambda) \times \dots$



Setting the scene

Central for realistic description of the observed complexity.

Description of IR sensitive observables shower and hadronization factored, but intertwined





Probabilistic algorithms with QCD coherence or large-N limit



$$dS = \frac{\alpha_s}{2\pi} \frac{d\tilde{q}_i^2}{\tilde{q}_i^2} dz P(z_i) \exp\left(-\int_{\tilde{q}_i^2}^{Q^2} \frac{dq^2}{q^2} \int_{z_-(k^2)}^{z_+(k^2)} d\xi \frac{\alpha_i}{2}\right)$$

emission rate

no emission probability

 $d\sigma \sim L \times d\sigma_H(Q) \times PS(Q \to \mu) \times MPI \times Had(\mu \to \Lambda) \times \dots$

~ $\mathcal{M}_{n}^{\dagger}(\mathbf{p}_{1},...,\mathbf{p}_{n})$ T.... T. $\mathcal{T} \cdot \mathcal{T} \cdot \mathcal{T} \cdot \mathcal{T} \cdot \mathcal{M}_{n}(\mathbf{p}_{1},...,\mathbf{p}_{n})$

Exploit QCD coherence:







All probabilistic algorithms determine the effect of gluon exchange and virtual corrections by unitarity.





Favorite probabilistic algorithms



Driven by QCD coherence

- Recoil global
- Links to analytic use of coherent branching

Parton branchings order in angle.

Herwig 7

	Current release	Hard matrix	Shower algorithms	NLO Matching	Multijet merging	MPI	Hadronizat ion	Shower variations
17	Herwig 7	Internal, libraries, event files	QTilde, Dipoles	Internally automated	Internally automated	Eikonal	Clusters, (Strings)	Yes
Ð	Pythia 8	Internal, event files	Pt ordered, DIRE, VINCIA	External	Internal, ME via event files	Interleaved	Strings	Yes
	Sherpa 2	Internal, libraries	CSShower, DIRE, ALARIC	Internally automated	Internally automated	Eikonal	Clusters, Strings	Yes



Dipole branchings order in transverse momentum.

- Driven by large-N dipole pattern and colour flows
- Momentum conservation for each emission
- Advantageous for matching & merging

Herwig 7, Pythia 8, Sherpa, PanScales, Deductor

Plethora of approaches to compare — need to go beyond in understanding and controlling shower algorithms.









Main lines of current parton shower research

Shower development is a broad field, fortunately back on the agenda. Hadronziation is not exactly shower development, but enters at a similar level,

See Schumann's talk. Always needs to accompany shower development.

Matching/Merging

Control and demonstration of perturbative accuracy.

(N)NLL accuracy

Amplitude evolution

Hadronization

Genuine quantum effects: not limited to subleading colour. Comprehensive, factorized picture and construction of algorithms.



• Perturbative accuracy

- Beyond probabilistic algorithms.
- Factorisation and hadronization.



The struggle with QED, EW and other interactions

- The absence of a large-N limit forces us to question existing structures
- Accuracy from interleaving with QCD needs to be carefully addressed
- The inner workings and role of coherence is entirely unknown

Halfway safe ground in the quasi-collinear limit, should be exploring these algorithms. Spin correlations are vital — amplitude evolution will be crucial to build algorithms.

Recent examples:

Sectorised QED multipoles and electroweak splittings in VINCIA Multiscale interleaved angular ordering in Herwig



[Verheyen, Skands—'21]

Interactions beyond QCD











Dark sector showers

QCD-like dark sectors can in principle build on existing QCD showering. Hadronization and scale hierarchy can differ significantly: no safe territory.

New in Herwig and the cluster model — more investigations and pheno to follow.



[Stafford at PSR '24]

Interactions beyond QCD

Angularities, correlations ... extremely useful observables.

[See also recent Les Houches study (in progress) and Kiebacher @ PSR '24].

 $EEC^{\gamma}(\rho) =$

$$\frac{1}{N_{\gamma}}\sum_{i< j\in jet} (E_i E_j)^{\gamma} \delta(\Delta R_{ij} - \rho)$$



[Kulkarni, Masouminia, Plätzer, Stafford — '24]









Dark sector showers

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[Stafford at PSR '24]

Interactions beyond QCD

Think of this as BSM4QCD: all of QED, EW and BSM scenarios force us

to think about showers and hadronisation in an unprecedented way.

$$\frac{1}{N_{\gamma}}\sum_{i< j\in jet}(E_{i}E_{j})^{\gamma}\delta(\Delta R_{ij}-\rho)$$

[Kulkarni, Masouminia, Plätzer, Stafford — '24]

0.2

0.3









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Description of electroweak effects and BSM scenarios.

Interactions beyond QCD

Remaining focus of this talk:

- Perturbative accuracy
- Beyond probabilistic algorithms.
- Factorisation and hadronization.



Accuracy of parton showers



Fragmentation is fine if we get collinear physics right.





Accuracy of Parton Showers



Fragmentation is fine if we get collinear physics right.

Global event shapes from coherent branching — for two jets.

 $H(\alpha_s) \times \exp\left(Lg_1(\alpha_s L) + \right)$

LL — qualitative



$$g_2(\alpha_s L) + \alpha_s g_3(\alpha_s L) + \dots)$$

NLL — quantitative NNLL — precision





 $\alpha_s L \sim 1$

Accuracy of Parton Showers



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Accuracy of Parton Showers



Fragmentation is fine if we get

[Banfi, Marchesini, Smye '02]





The quest for NLL precision



(N)NLO with matching

NLL with coherent branching Issues in dipole showers

Can we push this to NLL_{global} / LL_{non-global} in one (dipole) algorithm?



Issues in coherent branching LL with dipole showers

 $\alpha_s L \sim 1 \quad \alpha_s N^2 \sim 1$



The quest for NLL precision

Demonstrate NLL accurate evolution:

• PanScales — numerical

[PanScales — Dasgupta, Monni, Salam, Soyez +]

• Forshaw/Holguin/Plätzer — analytical

[aim at improving Herwig 7 dipole shower]

Deductor — numerical/analytical

• Sherpa — numerical/analytical

[Herren, Höche, Krauss, Reichelt, Schönherr]

• Apollo — numerical

Can we push this to NLL_{global} / LL_{non-global} in one (dipole) algorithm?



 $\alpha_s L \sim 1 \quad \alpha_s N^2 \sim 1$



Improving Shower Accuracy — Tools









[van Bleekveld @ MBI '24]



Improving Shower Accuracy — Spin & Colour Correlations

Spin correlations building on Collins-Knowles algorithm





[Webster, Richardson - '20]

Dynamic colour factors in dipole showers

$$C_{iJ}(\theta_{iq},\theta_{LJ}) = \begin{pmatrix} C_F \delta_i^{(q)} + \frac{C_A}{2} \delta_i^{(g)} \end{pmatrix} \theta(\theta_{iq} < \theta_{LJ}) \qquad \text{Track} \\ + \begin{pmatrix} \frac{C_A}{2} \delta_J^{(g)} + C_F \delta_J^{(q)} \end{pmatrix} \theta(\theta_{iq} < \theta_{LJ}) \qquad \text{to reduce to the set of the set of$$

[Forshaw, Holguin, Plätzer — '21] [Hamilton, Medves, Salam, Scyboz, Soyez — '21]



[Karlberg, Salam, Scyboz, Verheyen –

k angular extent of evolution produce colour factors as ted by coherence.







Where it (also) matters

Coherent branching jet mass including mass effects:

$$z(1-z)\tilde{q}^{2} = -m_{\tilde{i}\tilde{j}}^{2} + \frac{m_{i}^{2}}{z} + \frac{m_{j}^{2}}{1-z} - \frac{p_{\perp}^{2}}{z(1-z)}$$

using [Gieseke, Stephens, Webber - '03]

NLL accurate for global observables with massive quarks.

Top mass definition from coherent branching.

 $m_t^{\mathrm{MC}} =$

[Hoang, Plätzer, Samitz — '18]

 $m_t^{\mathrm{CB}}(Q_0)$

Take home message: hadronization and mass scheme compensate for shower cutoff dependence.



$$m_t^{\text{pole}} + \Delta_m^{\text{pert}} + \Delta_m^{\text{non-pert}}$$
$$0) = m_t^{\text{pole}} - \frac{2}{3}Q_0 \ \alpha_s(Q_0) + \mathcal{O}(\alpha_s^2)$$



Beyond NLL_{global} — more differential, NLL_{nonglobal}, NNLL_{global}, ...

Structure of second order ingredients and dedicated resummations: uncover algorithms and benchmarks.



Kinematic and colour structure of emission kernels beyond LO

[Dasgupta, El-Menoufi — '21]

Effective couplings and collinear fragmentation

$$\frac{\alpha_s}{2\pi}K(z_g) = \frac{V_{q\bar{q}g}}{B_{q\bar{q}g}} - \frac{V_{q\bar{q}}}{B_{q\bar{q}}} + \int_0^{\tilde{v}_g} \frac{d\Phi_{q\bar{q}ij}}{d\Phi_{q\bar{q}g}} \frac{B_{q\bar{q}ij}}{B_{q\bar{q}g}} - \int_0^{v_g} \frac{d\Phi_{q\bar{q}g'}}{d\Phi_{q\bar{q}}} \frac{B_{q\bar{q}g'}}{B_{q\bar{q}}}$$

[Van Beekveld, Dasgupta, El-Menoufi, Halliwell, Karlberg, Monni — '24]

Double emission matrix element corrections

Towards second-order showers: unordered contributions

- sector showers allow to include direct $2 \rightarrow 4$ branchings in a simple way
- divide phase space into strongly-ordered and unordered region
 - s.o. region: only single-unresolved limits
 - u.o. region: only double-unresolved limits
- $2 \rightarrow 4$ branchings important ingredient to NNLO+PS $(+ \text{ virtual corrections to } 2 \rightarrow 3)$

[C. Preuss for Vincia — PSR 21] [PanScales double soft algorithms]



[Gelle

[Löschner — PSR '24] [Löschner, Plätzer, Simpson — '21]





Building amplitude evolution at the second order.

$$\mathbf{\Gamma}_{S,n}^{(2)} = -\hat{\mathbf{V}}_{n}^{(2)} \left[\partial_{S} \Xi_{n,2}\right] + \hat{\mathbf{V}}_{n}^{(1)} \left[\partial_{S} \Xi_{n,1}\right] \hat{\mathbf{V}}_{n}^{(1)} \left[1 + \hat{\mathbf{V}}_{n}^{(1)}\right] = -\hat{\mathbf{V}}_{n}^{(2)} \left[1 + \hat{\mathbf{V}}_{n}^{(2)}\right] = -\hat{\mathbf{V}}_{n}^{(2)} \left[1 +$$

$$\mathbf{R}_{n}^{(1,1)} \circ \mathbf{R}_{n}^{(1,0)\dagger} = \left(\hat{\mathbf{D}}_{n}^{(1,1)} \left[1 - \Xi_{n,1}\right] - \hat{\mathbf{D}}_{n}^{(1,0)} \hat{\mathbf{V}}_{n-1}^{(1)} \left[1 - \Xi_{n-1,1}\right]\right) \circ \hat{\mathbf{D}}_{n}^{(1,0)} \\ + \left(\hat{\mathbf{D}}_{n}^{(1,1)} \left[\partial_{S}\Xi_{n,1}\right] - \hat{\mathbf{V}}_{n}^{(1)} \left[\partial_{S}\Xi_{n,1}\right] \hat{\mathbf{D}}_{n}^{(1,0)}\right) \circ \hat{\mathbf{D}}_{n}^{(1,0)\dagger} \left(\Xi_{n-1,1}\right) \right)$$

$$\mathbf{R}_{n}^{(2,0)} \circ \mathbf{R}_{n}^{(2,0)\dagger} = \hat{\mathbf{D}}_{n}^{(2,0)} \circ \hat{\mathbf{D}}_{n}^{(2,0)\dagger} \partial_{S} \Theta_{n,2} \\ - \hat{\mathbf{D}}_{n}^{(1,0)} \hat{\mathbf{D}}_{n-1}^{(1,0)} \circ \hat{\mathbf{D}}_{n-1}^{(1,0)\dagger} \hat{\mathbf{D}}_{n}^{(1,0)\dagger} (1 - \Theta_{n-1,1})$$

[Plätzer — '22] [Nagy, Soper — PSR '24]







Beyond NLLglobal — more differential, NLLnonglobal, NNLLglobal, ...





Effective couplings

 $\frac{\alpha_s}{2\pi}K(z_g) = \frac{V_{q\bar{q}g}}{B_{q\bar{q}g}}$

[Van Beekveld, Da Halliwell, Karlberg

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[C. Preuss for Vincia [PanScales double s

Algorithms are more than the sum of their ingredients.

Way forward is hard, but possible we need to find systematic formulations.

Demonstrate accuracy for certain observables, but strive for most differential/flexible algorithms.

Role of data comparisons not always clear but encouraging.

See Karlberg's talk





Main lines of current parton shower research

Shower development is a broad field, fortunately back on the agenda. Hadronziation is not exactly shower development, but enters at a similar level,

See Schumann's talk. Always needs to accompany shower development.

Matching/Merging

Control and demonstration of perturbative accuracy.

(N)NLL accuracy

Amplitude evolution

Hadronization

Genuine quantum effects: not limited to subleading colour. Comprehensive, factorized picture and construction of algorithms.

Description of electroweak effects and BSM scenarios.

Interactions beyond QCD

Remaining focus of this talk:

- Perturbative accuracy
- Beyond probabilistic algorithms.
- Factorisation and hadronization.



Colour reconnection and hadronization is about subleading-N. So are shower accuracy and interference terms.

Colour factor algorithms

Coherent, NLL-accurate dipole showers

[Gustafson] [PanScales '21] [Forshaw, Holguin, Plätzer '21] Colour ME corrections

Colour-exact real emissions as far as possible

> [Plätzer, Sjödahl '12, '18] [Höche, Reichelt '20]

Full amplitude evolution

Colour-exact real and virtual corrections

[Forshaw, Plätzer + ... 'I 3 ...] [Nagy, Soper '07 ...]



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Colour-exact real and virtual corrections

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The amplitude evolution equation



Markovian algorithm at the amplitude level: Iterate gluon exchanges and emission.

Different histories in amplitude and conjugate amplitude needed to include interference.

[Forshaw, Plätzer, De Angelis, Torre, ... — '18 ... '24]





The amplitude evolution equation

Markovian algorithm a Iterate gluon exchange

Different histories in a amplitude needed to



[Forshaw, Plätzer, De Angelis, Torre, ... — '18 ... '24]



[Forshaw, Plätzer, DeAngelis, Torre ...]

amplitude

conjugate amplitude





Amplitude evolution at hadron colliders







Showering of hard process interferences



Amplitude evolution and cloctroweak physics

Factorizing momentum mapp

Momentum mappings to systematically factor renormalised matrix elements.

$$- = \frac{1}{2p_i \cdot Q_{i,s}} \frac{\Psi(\Lambda p_i, M_i) \overline{\Psi}(\Lambda p_i, M_i)}{1 - \Sigma'(M_i^2)} + \mathcal{O}(\lambda)$$





Find a ba 1 stri together with isospin and colour.



Electroweak bosons now mix different chiral basis states.





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Hadronization in general purpose event generators







[Altmann, Skands — '24]

Back to an active field

- Colour reconnection
- Spin physics
- Kinematics and low energy hadronization



[Gieseke, Kiebacher, Plätzer, Priedigkeit — in progress]



[Masouminia, Richardson — '23]



[Kerbizi, Lönnblad, Martin — '24]









Why we should worry about PS x Hadronization



[Bellm, Lönnblad, Plätzer, Prestel, Samitz, Siodmok, Hoang — Les Houches 2017]







Shower cutoff has central role as factorisation scale:

- Shower/hadronization constrained by renormalisation group
- Formulated in amplitude level evolution, contains event generators in large-N limit.

Subtract IR divergencies in unresolved regions

Re-arrange to resum IR enhancements

$$\int \operatorname{Tr}_{n} \left[\mathbf{M}_{n} \mathbf{U}_{nm} \right] \mathrm{d}\phi_{m} u(\phi_{m})$$

Systematic construction of amplitude evolution algorithms.

$$\mathbf{U}_{n} = \mathcal{X}_{n} \left[\mathbf{S}(\mu_{S}), \mu_{S} \right]$$

$$\sigma = \sum_{n} \alpha_{S}^{n} \int \operatorname{Tr} \left[\mathbf{A}_{n}(\mu_{S}) \mathbf{S}_{n}(\mu_{S}) \right]$$

$$\mathbf{M}_{n} Z_{g}^{n} = \mathcal{Z}_{n} \left[\mathbf{A}(\mu_{S}), \mu_{S} \right]$$

Leads to constructive prescription for "high energy" end of hadronization.

First steps demonstrated:

Lower energy dynamics in the model not sensitive to cutoff anymore.

[Hoang, Jin, Plätzer, Samitz — '24] [Gieseke, Kiebacher, Plätzer, Priedigkeit — in progress]

shower IR cutoff

[Kiebacher @ PSR '24]

Leads to constructive prescription for "high energy" end of hadronization.

Colour reconnection dynamics implied from amplitude evolution structures:

[Gieseke, Kirchgaesser, Plätzer, Siodmok — '17, '18]

[Hoang, Jin, Plätzer, Samitz — '24] [Gieseke, Kiebacher, Plätzer, Priedigkeit — in progress]

[Kiebacher @ PSR '24]

Summary

Parton showers are central to event generators.

I emphasise that we study shower *algorithms* rather than shower models — predict QCD effects in a hierarchy of strongly ordered energy scales.

We can, and need, to assess their accuracy and strive for underlying, well-defined, construction principles.

Infrared sensitive observables are of course the prime test bed — matching will do remaining aspects for jet observables. Hadronization can not be looked at in isolation but is intertwined with parton showering.

Amplitude evolution provides us with a new paradigm which we can use to construct parton showers, as a dedicated resummation tools and eventually to study the structure of amplitudes using simulation.

 $d\sigma \sim L \times d\sigma_H(Q) \times PS(Q \to \mu) \times MPI \times Had(\mu \to \Lambda) \times \dots$

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