Selected

Recent Advances in PDFs From Lattice QCD

Martha Constantinou



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Disclaimer

Mistakes attributed to Wine Tasting Excursion in Breisach



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Before...





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OUTLINE

- A. Lattice QCD @ LHC
- **B.** Methods to access PDFs (+ GPDs, TMDs) from lattice QCD
- C. Selected recent results on:
 - Quark PDFs
 - Gluon PDFs
 - Novel Developments
- **D.** Synergy with phenomenology
- E. Concluding remarks



Lattice QCD @ LHC

Inputs for interpreting high-energy collisions at LHC for SM and Beyond.

M. Seidel, Thu 9:00 am

M. Muskinja, Thu 9:30 am

J. M. Cruz Martinez, Thu 10:00 am

Predictions of the Higgs boson production cross section in pp collisions

G. Heinrich, Tue 9:30 am

- ★ Lattice QCD-derived PDFs complement phenomenological fits to experimental data sets. Better constraints on the isovector proton PDF for high-x region at ~10% level will impact predictions for new-physics searches at ATLAS and CMS [2017 PDFLattice Report, Prog. Part. Nucl. Phys. 100, 107 (2018)]
- ★ Lattice can complement experimental data in kinematic regions where data are sparse or unavailable. Essential for heavy ion collisions, small-x physics, and precision Higgs and top-quark physics
- ★ Lattice results on spin-dependent PDFs interest for experiments at the LHC and other spin-polarized collision experiments.
- Strange PDFs have important role in precision EW physics can be coupled with W + c data from LHC (unpolarized PDFs).
 T. Hobbs, Tue 2 pm

Accessing PDFs/GPDs from lattice QCD



Nucleon Characterization

Wigner distributions

Т

- ★ Fully characterize partonic structure of hadrons
- ★ Provide multi-dim images of the parton distributions in phase space



Correlations between momenta, positions, spins

★ Information on the hadron's mechanical properties (OAM, pressure, etc.)

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★ Partons contain information on
 x: longitudinal momentum fraction
 k_T: transverse momentum
 *b*_⊥: impact parameter



Correlations between momenta, positions, spins

Information on the hadron's mechanical properties (OAM, pressure, etc.)

Hadron Structure

★ Structure of hadrons explored in high-energy scattering processes





★ Processes cross-section contains information on hadron $\sigma_{\text{DIS}}(x, Q^2) = \sum_{i} \left[H^i_{\text{DIS}} \otimes f_i \right](x, Q^2) \qquad [a \otimes b](x) \equiv \int_x^1 \frac{d\xi}{\xi} a\left(\frac{x}{\xi}\right) b(\xi)$

 Hadron structure expressed in terms of distribution functions of partonic constituents (PDFs, GPDs, TMDs)





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Perturb. part (process dependent)

 ★ Hadron structure expressed in terms of distribution functions of partonic constituents (PDFs, GPDs, TMDs)





Hadron Structure

Structure of hadrons explored in high-energy scattering processes





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Perturb. part
(process dependent) Non-Perturb. part
(process "independent")

 Hadron structure expressed in terms of distribution functions of partonic constituents (PDFs, GPDs, TMDs)





- In parton model, physical picture valid for infinite momentum frame [R. P. Feynman, Phys. Rev. Lett. 23, 1415 (1969)]
- **PDFs** parameterized via matrix elements of nonlocal light-cone operators

$$f(x) = \frac{1}{4\pi} \int dy^{-} e^{-ixP^{+}y^{-}} \langle P, S | \bar{\psi}_{f} \gamma^{+} \mathcal{W} \psi_{f} | P, S \rangle$$



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$$\langle N(P') | \mathcal{O}_{V}^{\mu\mu_{1}\cdots\mu_{n-1}} | N(P) \rangle \sim \sum_{\substack{i=0 \\ \text{even}}}^{n-1} \left\{ \gamma^{\{\mu} \Delta^{\mu_{1}} \cdots \Delta^{\mu_{i}} \overline{P}^{\mu_{i+1}} \cdots \overline{P}^{\mu_{n-1}\}} A_{n,i}(t) - i \frac{\Delta_{\alpha} \sigma^{\alpha\{\mu}}{2m_{N}} \Delta^{\mu_{1}} \cdots \Delta^{\mu_{i}} \overline{P}^{\mu_{i+1}} \cdots \overline{P}^{\mu_{n-1}\}} B_{n,i}(t) \right\} \\ + \frac{\Delta^{\mu} \Delta^{\mu_{1}} \cdots \Delta^{\mu_{n-1}}}{m_{N}} C_{n,0}(\Delta^{2}) \Big|_{n \text{ even}} \rangle \left\{ \left. + \frac{\Delta^{\mu} \Delta^{\mu_{1}} \cdots \Delta^{\mu_{n-1}}}{m_{N}} C_{n,0}(\Delta^{2}) \right|_{n \text{ even}} \right\} \right\} \\ + \frac{\Delta^{\mu} \Delta^{\mu_{1}} \cdots \Delta^{\mu_{n-1}}}{m_{N}} C_{n,0}(\Delta^{2}) \Big|_{n \text{ even}} \right\} \\ + \frac{\Delta^{\mu} \Delta^{\mu_{1}} \cdots \Delta^{\mu_{n-1}}}{m_{N}} C_{n,0}(\Delta^{2}) \Big|_{n \text{ even}} \right\} \\ + \frac{\Delta^{\mu} \Delta^{\mu_{1}} \cdots \Delta^{\mu_{n-1}}}{m_{N}} C_{n,0}(\Delta^{2}) \Big|_{n \text{ even}} \right\} \\ + \frac{\Delta^{\mu} \Delta^{\mu_{1}} \cdots \Delta^{\mu_{n-1}}}{m_{N}} C_{n,0}(\Delta^{2}) \Big|_{n \text{ even}} \right\} \\ + \frac{\Delta^{\mu} \Delta^{\mu_{1}} \cdots \Delta^{\mu_{n-1}}}{m_{N}} C_{n,0}(\Delta^{2}) \Big|_{n \text{ even}} \right\} \\ + \frac{\Delta^{\mu} \Delta^{\mu_{1}} \cdots \Delta^{\mu_{n-1}}}{m_{N}} C_{n,0}(\Delta^{2}) \Big|_{n \text{ even}} \right\} \\ + \frac{\Delta^{\mu} \Delta^{\mu_{1}} \cdots \Delta^{\mu_{n-1}}}{m_{N}} C_{n,0}(\Delta^{2}) \Big|_{n \text{ even}} \right\} \\ + \frac{\Delta^{\mu} \Delta^{\mu_{1}} \cdots \Delta^{\mu_{n-1}}}{m_{N}} C_{n,0}(\Delta^{2}) \Big|_{n \text{ even}} \right\} \\ + \frac{\Delta^{\mu} \Delta^{\mu_{1}} \cdots \Delta^{\mu_{n-1}}}{m_{N}} C_{n,0}(\Delta^{2}) \Big|_{n \text{ even}} \right\} \\ + \frac{\Delta^{\mu} \Delta^{\mu_{1}} \cdots \Delta^{\mu_{n-1}}}{m_{N}} C_{n,0}(\Delta^{2}) \Big|_{n \text{ even}} \right\} \\ + \frac{\Delta^{\mu} \Delta^{\mu_{1}} \cdots \Delta^{\mu_{n-1}}}{m_{N}} C_{n,0}(\Delta^{2}) \Big|_{n \text{ even}} \right\} \\ + \frac{\Delta^{\mu} \Delta^{\mu_{1}} \cdots \Delta^{\mu_{n-1}}}{m_{N}} C_{n,0}(\Delta^{2}) \Big|_{n \text{ even}} \right\} \\ + \frac{\Delta^{\mu} \Delta^{\mu_{1}} \cdots \Delta^{\mu_{n-1}}}{m_{N}} C_{n,0}(\Delta^{2}) \Big|_{n \text{ even}} \right\} \\ + \frac{\Delta^{\mu} \Delta^{\mu_{1}} \cdots \Delta^{\mu_{n-1}}}{m_{N}} C_{n,0}(\Delta^{2}) \Big|_{n \text{ even}} \right\} \\ + \frac{\Delta^{\mu} \Delta^{\mu_{1}} \cdots \Delta^{\mu_{n-1}}}{m_{N}} C_{n,0}(\Delta^{2}) \Big|_{n \text{ even}} \right\} \\ + \frac{\Delta^{\mu} \Delta^{\mu_{1}} \cdots \Delta^{\mu_{n-1}}}{m_{N}} C_{n,0}(\Delta^{2}) \Big|_{n \text{ even}} \right\}$$



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$$\wedge \text{ Mellin moments}_{\text{(local OPE expansion)}} \bar{q}(-\frac{1}{2}z) \gamma^{\sigma} W[-\frac{1}{2}z, \frac{1}{2}z] q(\frac{1}{2}z) = \sum_{n=0}^{\infty} \frac{1}{n!} z_{\alpha_{1}} \dots z_{\alpha_{n}} \left[\bar{q} \gamma^{\sigma} \overleftrightarrow{D}^{\alpha_{1}} \dots \overleftrightarrow{D}^{\alpha_{n}} q \right]$$

$$\text{Reconstruction of PDFs/GPDs very challenging}$$

$$\langle N(P') | \mathcal{O}_{V}^{\mu\mu_{1}\cdots\mu_{n-1}} | N(P) \rangle \sim \sum_{i=0}^{n-1} \left\{ \gamma^{(\mu}\Delta^{\mu_{1}}\cdots\Delta^{\mu_{i}}\overline{P}^{\mu_{i+1}}\cdots\overline{P}^{\mu_{n-1}}A_{n,i}(t) - i\frac{\Delta_{\alpha}\sigma^{\alpha(\mu}}{2m_{N}}\Delta^{\mu_{1}}\cdots\Delta^{\mu_{i}}\overline{P}^{\mu_{i+1}}\cdots\overline{P}^{\mu_{n-1}}B_{n,i}(t) \right\} + \frac{\Delta^{\mu}\Delta^{\mu_{1}}\cdots\Delta^{\mu_{n-1}}}{m_{N}} C_{n,0}(\Delta^{2})|_{n \text{ even}} \right\} + \frac{\Delta^{\mu}\Delta^{\mu_{1}}\cdots\Delta^{\mu_{n-1}}}{m_{N}} C_{n,0}(\Delta^{2})|_{n \text{ even}} \right] U(P)$$



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★ Matrix elements of nonlocal operators (quasi-GPDs, pseudo-GPDs, …)

 $\langle N(P_f) | \overline{\Psi}(z) \Gamma \mathcal{W}(z,0) \Psi(0) | N(P_i) \rangle_{\mu}$

Nonlocal operator with Wilson line

$$\langle N(P')|O_V^{\mu}(x)|N(P)\rangle = \overline{U}(P') \left\{ \gamma^{\mu}H(x,\xi,t) + \frac{i\sigma^{\mu\nu}\Delta_{\nu}}{2m_N} E(x,\xi,t) \right\} U(P) + \text{ht},$$

$$\langle N(P')|O_A^{\mu}(x)|N(P)\rangle = \overline{U}(P') \left\{ \gamma^{\mu}\gamma_5 \widetilde{H}(x,\xi,t) + \frac{\gamma_5\Delta^{\mu}}{2m_N} \widetilde{E}(x,\xi,t) \right\} U(P) + \text{ht},$$

$$\langle N(P')|O_T^{\mu\nu}(x)|N(P)\rangle = \overline{U}(P') \left\{ i\sigma^{\mu\nu}H_T(x,\xi,t) + \frac{\gamma^{[\mu}\Delta^{\nu]}}{2m_N} E_T(x,\xi,t) + \frac{\overline{P}^{[\mu}\Delta^{\nu]}}{m_N^2} \widetilde{H}_T(x,\xi,t) + \frac{\gamma^{[\mu}\overline{P}^{\nu]}}{m_N} \widetilde{E}_T(x,\xi,t) \right\} U(P) + \text{ht},$$



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$$\langle N(P')|O_A^{\mu}(x)|N(P)\rangle = \overline{U}(P') \left\{ \gamma^{\mu}\gamma_5 \widetilde{H}(x,\xi,t) + \frac{\gamma_5\Delta^{\mu}}{2m_N}\widetilde{E}(x,\xi,t) \right\} U(P) + \text{ht},$$

$$\langle N(P')|O_T^{\mu\nu}(x)|N(P)\rangle = \overline{U}(P') \left\{ i\sigma^{\mu\nu}H_T(x,\xi,t) + \frac{\gamma^{[\mu}\Delta^{\nu]}}{2m_N}E_T(x,\xi,t) + \frac{\overline{P}^{[\mu}\Delta^{\nu]}}{m_N^2}\widetilde{H}_T(x,\xi,t) + \frac{\gamma^{[\mu}\overline{P}^{\nu]}}{m_N}\widetilde{E}_T(x,\xi,t) \right\} U(P) + \text{ht},$$



Novel Approaches

Hadronic tensor Auxiliary scalar quark Fictitious heavy quark Auxiliary scalar quark Higher moments Quasi-distributions (LaMET) Compton amplitude and OPE Pseudo-distributions Good lattice cross sections PDFs without Wilson line Moments of PDFs of any order

[K.F. Liu, S.J. Dong, PRL 72 (1994) 1790, K.F. Liu, PoS(LATTICE 2015) 115]
[U. Aglietti et al., Phys. Lett. B441, 371 (1998), arXiv:hep-ph/9806277]
[W. Detmold, C. J. D, Lin, Phys. Rev. D73, 014501 (2006)]
[V. Braun & D. Mueller, Eur. Phys. J. C55, 349 (2008), arXiv:0709.1348]
[Z. Davoudi, M. Savage, Phys. Rev. D86, 054505 (2012)]
[X. Ji, PRL 110 (2013) 262002, arXiv:1305.1539; Sci. China PPMA. 57, 1407 (2014)]
[A. Chambers et al. (QCDSF), PRL 118, 242001 (2017), arXiv:1703.01153]
[A. Radyushkin, Phys. Rev. D 96, 034025 (2017), arXiv:1705.01488]
[Y-Q Ma & J. Qiu, Phys. Rev. Lett. 120, 022003 (2018), arXiv:1709.03018]
[Y. Zhao Phys.Rev.D 109 (2024) 9, 094506, arXiv:2306.14960]
[A. Shindler, Phys.Rev.D 110 (2024) 5, L051503, arXiv:2311.18704]



Novel Approaches

Hadronic tensor
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 Auxiliary scalar quark
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 K.F. Liu, S.J. Dong, PRL 72 (1)
 Auxiliary scalar quark
 Magietti et al., Phys. Lett. B
 W. Detmold, C. J. D, Lin, Phys
 W. Braun & D. Mueller, Eur. Pr
 Z. Davoudi, M. Savage, Phys
 X. Ji, PRL 110 (2013) 262002,
 A. Chambers et al. (QCDSF),
 A. Radyushkin, Phys. Rev. D 109
 Compton applications
 A. Shindler, Phys.Rev.D 109 (2022)

[K.F. Liu, S.J. Dong, PRL 72 (1994) 1790, K.F. Liu, PoS(LATTICE 2015) 115]
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[A. Shindler, Phys.Rev.D 110 (2024) 5, L051503, arXiv:2311.18704]

★ Reviews of methods and applications

- A guide to light-cone PDFs from Lattice QCD: an overview of approaches, techniques and results K. Cichy & M. Constantinou (invited review) Advances in HEP 2019, 3036904, arXiv:1811.07248
- Large Momentum Effective Theory X. Ji, Y.-S. Liu, Y. Liu, J.-H. Zhang, and Y. Zhao (2020), 2004.03543
- The x-dependence of hadronic parton distributions: A review on the progress of lattice QCD
 M. Constantinou (invited review) Eur. Phys. J. A 57 (2021) 2, 77, arXiv:2010.02445





Novel Approaches



$$\mathcal{M}(P_f, P_i, z) = \langle N(P_f) \, | \, \bar{\Psi}(z) \, \Gamma \, \mathcal{W}(z, 0) \Psi(0) \, | \, N(P_i) \rangle_{\mu}$$



$$\mathcal{M}(P_{f}, P_{i}, z) = \langle N(P_{f}) | \bar{\Psi}(z) \Gamma \mathcal{W}(z, 0) \Psi(0) | N(P_{i}) \rangle_{\mu}$$

$$[X. Ji, Phys. Rev. Lett. 110 (2013) 262002] [X. Ji, Sci. China Phys. M.A. 57 (2014) 1407] \qquad quasi-PDFs \qquad pseudo-ITD [A. Radyushkin, PRD 96, 034025 (2017)]$$

$$\tilde{q}_{\Gamma}^{\text{GPD}}(x, t, \xi, P_{3}, \mu) = \int \frac{dz}{4\pi} e^{-ixP_{3}z} \mathcal{M}(P_{f}, P_{i}, z) \qquad \mathfrak{M}(\nu, \xi, t; z_{3}^{2}) \equiv \frac{\mathcal{M}(\nu, \xi, t; z_{3}^{2})}{\mathcal{M}(0, 0, 0; z^{2})} \qquad (\nu = z \cdot p)$$



$$\mathcal{M}(P_{f}, P_{i}, z) = \langle N(P_{f}) | \bar{\Psi}(z) \Gamma \mathcal{W}(z, 0) \Psi(0) | N(P_{i}) \rangle_{\mu}$$

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Quark PDFs:

The unpolarized case



Collection of results



[M. Constantinou et al. (2020 PDFLattice Report), Prog.Part.Nucl.Phys. 121 (2021) 103908]

★ Several improvements:

- More calculations at physical quark masses
- Ensembles at various lattice spacings
- Addressing systematic uncertainties due to methodologies

Refining the unpolarized proton PDF (u-d)

★ Physical quark masses

- HISQ, a=0.076 fm, P~1.5 GeV
- Deep Neural Network for inverse problem
- NNLO for matching



[X. Gao et al., PRD 107 (2023) 7, 074509]



Refining the unpolarized proton PDF (u-d)

★ Physical quark masses

- HISQ, a=0.076 fm, P~1.5 GeV
- Deep Neural Network for inverse problem
- NNLO for matching



[X. Gao et al., PRD 107 (2023) 7, 074509]

★ Continuum limit

- TM&clover, a=0.09 fm, m_{π} =350 MeV
- P~1.8 GeV
- NNLO for matching





Improving evolution of PDFs

★ Continuum limit - higher twist effects

- Clover, a=0.075, 0.065, 0.048 fm
- m_{π} =440 MeV
- Jacobi polynomials for controlling finite-a & higher twist



[Karpie et al., JHEP 11 (2021) 024]

Improving evolution of PDFs

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ר'

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[Karpie et al., JHEP 11 (2021) 024]

- Non-perturbative scale evolution of pseudo distributions:
 - lattice scale much different than scale for light-cone PDFs
 - addresses the subtle z² behavior of matrix elements





[H. Dutrieux et al. (HadStruc), JHEP 04 (2024) 061]

Evolution of vector operator much larger than anticipated

Gluon PDFs for the proton









Elimination of Mixing with Quark Singlet PDFs





Elimination of Mixing with Quark Singlet PDFs



Gluon Helicity PDF

- Neural network analysis of lattice calculation disfavors negative gluon polarizability



[T. Khan et al., PRD 108, 074502]



Gluon Helicity PDF

Without Lattice

- Neural network analysis of lattice calculation disfavors negative gluon polarizability



[T. Khan et al., PRD 108, 074502]

0.5 $\mathcal{M}(
u,z_3^2)$ 0.0 2.55.0 7.5 2.57.5 0.0 5.0 0.0 ν ν $\Delta g > 0$ $\mu^2 = 10 \text{ GeV}^2$ $\Delta g < 0$ 0.2 0.2 $\pm |g|$ $x \Delta g$ 0.0 0.0 -0.2-0.20.2 0.6 0.8 0.2 0.4 0.4 0.6 0.8 \boldsymbol{x} \boldsymbol{x}

[J. Karpie et al., PRD 109 (2024) 3, 036031]

Including Lattice

LQCD: Hint for a nonzero gluon spin (proton) JAM analysis: No positivity constraint $(\Delta g > |g|$ for some regions of x)



Gluon Helicity PDF

Without Lattice

- Neural network analysis of lattice calculation disfavors negative gluon polarizability



[T. Khan et al., PRD 108, 074502]

0.5 $\mathcal{M}(
u,z_3^2)$ 0.0 2.55.0 7.5 2.55.0 7.5 0.0 0.0 ν ν $\Delta g > 0$ $\mu^2 = 10 \text{ GeV}^2$ $\Delta g < 0$ 0.2 0.2 $\pm |g|$ $x \Delta g$ 0.0 0.0 -0.20.2 0.2 0.4 0.6 0.8 0.6 0.80.4 \boldsymbol{x}

[J. Karpie et al., PRD 109 (2024) 3, 036031]

Including Lattice

LQCD: Hint for a nonzero gluon spin (proton) JAM analysis: No positivity constraint $(\Delta g > |g|$ for some regions of x)



New Developments

★ Twist-3 PDFs

★ GPDs



Twist-classification of PDFs, GPDs, TMDs $f_i = f_i^{(0)} + \frac{f_i^{(1)}}{Q} + \frac{f_i^{(2)}}{Q^2} \cdots$

★ Twist: The order in Q^{-1} entering factorization



(Selected) Twist-3 $(f_i^{(1)})$

() Nucleon	γ^j	$\gamma^j \gamma^5$	σ^{jk}
U	G_1, G_2 G_3, G_4		
L		$\widetilde{G}_1, \widetilde{G}_2$ $\widetilde{G}_3, \widetilde{G}_4$	
т			$H_2'(x,\xi,t)$ $E_2'(x,\xi,t)$



Twist-classification of PDFs, GPDs, TMDs $f_i = f_i^{(0)} + \frac{f_i^{(1)}}{O} + \frac{f_i^{(2)}}{O^2} \cdots$

 \star Twist: The order in Q^{-1} entering factorization



- **Twist-2**: probabilistic densities a wealth of information exists (mostly on PDFs)
- \star Twist-3: poorly known, but very important and have physical interpretation: - as sizable as twist-2
 - contain information about quark-gluon correlations inside hadrons
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The extraction of twist-3 is very challenges both experimentally and theoretically

★ Correlation functions in coordinate space

$$F^{[\Gamma]}(x,\Delta;P^3) = \frac{1}{2} \int \frac{dz^3}{2\pi} e^{ik \cdot z} \langle p_f, \lambda' | \bar{\psi}(-\frac{z}{2}) \Gamma \mathcal{W}(-\frac{z}{2},\frac{z}{2}) \psi(\frac{z}{2}) | p_i, \lambda \rangle \Big|_{z^0 = 0, \vec{z}_\perp = \vec{0}_\perp}$$

★ Parametrization of coordinate-space correlation functions

[D. Kiptily and M. Polyakov, Eur. Phys. J. C37 (2004) [F. Aslan et a., Phys. Rev. D 98, 014038 (2018)]

$$\begin{split} F^{[\gamma^{\mu}\gamma_{5}]}(x,\Delta;P^{3}) &= \frac{1}{2P^{3}} \bar{u}(p_{f},\lambda') \bigg[P^{\mu} \frac{\gamma^{3}\gamma_{5}}{P^{0}} F_{\widetilde{H}}(x,\xi,t;P^{3}) + P^{\mu} \frac{\Delta^{3}\gamma_{5}}{2mP^{0}} F_{\widetilde{E}}(x,\xi,t;P^{3}) \\ &+ \Delta^{\mu}_{\perp} \frac{\gamma_{5}}{2m} F_{\widetilde{E}+\widetilde{G}_{1}}(x,\xi,t;P^{3}) + \gamma^{\mu}_{\perp}\gamma_{5} F_{\widetilde{H}+\widetilde{G}_{2}}(x,\xi,t;P^{3}) \\ &+ \Delta^{\mu}_{\perp} \frac{\gamma^{3}\gamma_{5}}{P^{3}} F_{\widetilde{G}_{3}}(x,\xi,t;P^{3}) + i\varepsilon^{\mu\nu}_{\perp}\Delta_{\nu} \frac{\gamma^{3}}{P^{3}} F_{\widetilde{G}_{4}}(x,\xi,t;P^{3}) \bigg] u(p_{i},\lambda) \end{split}$$



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Pion mass:	260 MeV	
Lattice spacing:	0.093 fm	
Volume:	32 ³ x 64	

Twist-3 counterpart as sizable as twist-2

Burkhardt-Cottingham sum rule important check

 $\int_{-1}^{1} dx g_1(x) - \int_{-1}^{1} dx g_T(x) = 0$

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WW approximation

[S. Bhattacharya et al., PRD 102 (2020) 11 (Editors Selection)]

WW approximation:

twist-3 $g_T(x)$ determined by the twist-2 $g_1(x)$

 $g_T^{WW}(x) = \int_{-\infty}^{1} \frac{dy}{y} g_1(y)$

- $g_T(x)$ agrees with $g_T^{WW}(x)$ for x < 0.5(violations up to 30-40% possible)
- Violations of 15-40% expected from experimental data

 [A. Accardi et al., JHEP 11 (2009) 093]





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Twist-3 h_L(x) PDF

[S. Bhattacharya et al., PRD 104 (2021) 11, 114510]



- h_L^u dominant tension between h_L & h_L^{WW}
- h_L^d <0 and decays faster then h_L^u

Proton GPDs

 Tomographic imaging of proton has central role in the science
 program of EIC
 GPDs, FFs, GFFs, TMDs, ...
 [R. Abdul Khalek et al.,
 EIC Yellow Report 2021, arXiv:2103.05419]



★ GPDs are not well-constrained experimentally

$$\langle N(P')|O_V^{\mu}(x)|N(P)\rangle = \overline{U}(P') \left\{ \gamma^{\mu}H(x,\xi,t) + \frac{\imath\sigma^{\mu\nu}\Delta_{\nu}}{2m_N}E(x,\xi,t) \right\} U(P) + \mathrm{ht}$$

$$\langle N(P')|O_A^{\mu}(x)|N(P)\rangle = \overline{U}(P') \left\{ \gamma^{\mu}\gamma_5\widetilde{H}(x,\xi,t) + \frac{\gamma_5\Delta^{\mu}}{2m_N}\widetilde{E}(x,\xi,t) \right\} U(P) + \mathrm{ht}$$

★ Can be accessed also at the twist-3 level

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Light-cone GPDs





★ Direct access to \widetilde{E} -GPD not possible for zero skewness $P^{\mu} \frac{\Delta^{3} \gamma_{5}}{2mP^{0}} F_{\widetilde{E}}(x,\xi,t;P^{3})$

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Synergy/Complementarity of lattice and phenomenology





Incorporating lattice PDFs in global analyses

Synergy between lattice and phenomenology

 Lattice and experimental data sets data within the same global analysis (JAM framework)
 [J. Bringewatt et al., PRD 103 (2021) 016003, arXiv:2010.00548]





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- Significant impact for helicity PDF



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★ Other efforts within NNPDF framework

[K. Cichy et al., JHEP 10 (2019) 137, arXiv:1907.06037] [L. Del Debbio et al., JHEP 02 (2021) 138, 2010.03996] ★ Interest in applying similar approach to quantities that are more challenging to extract experimentally (GPDs, twist-3 distributions, ...)







★ Lattice data may be incorporated in global analysis of experimental data and may influence parametrization of t and ξ dependence



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- 1. Theoretical studies of high-momentum transfer processes using perturbative QCD methods and study of GPDs properties
- 2. Lattice QCD calculations of GPDs and related structures
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Other GPD global analysis efforts:

- Gepard [https://gepard.phy.hr/]
- PARTONS [https://partons.cea.fr]
- EXCLAIM [https://exclaimcollab.github.io/web.github.io/#/]



Synergies: constraints & predictive power of lattice QCD



M. Constantinou, QCD@LHC 2024

Concluding remarks



Concluding Remarks

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- ★ New Developments in several promising directions:



- **Extensive programs in Gluon PDFs**
- ★ Synergy with phenomenology has the potential to enhance the impact of lattice QCD data and complement data sets



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