

QCD@LHC - FREIBURG - 8th OCTOBER 2024

MINNLO_{PS} EVENT GENERATOR: COLOUR SINGLET PLUS ONE JET EVENTS

SILVIA ZANOLI
University of Oxford

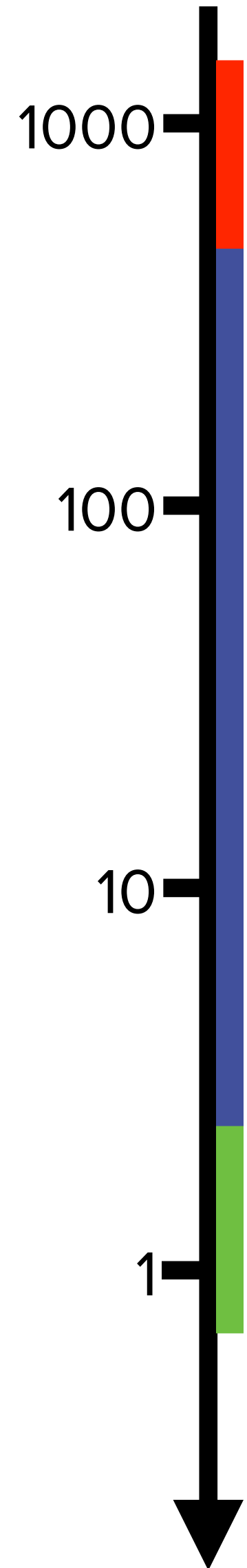
Based on **2402.00596** Ebert, Rottoli, Wiesemann, Zanderighi, **SZ**



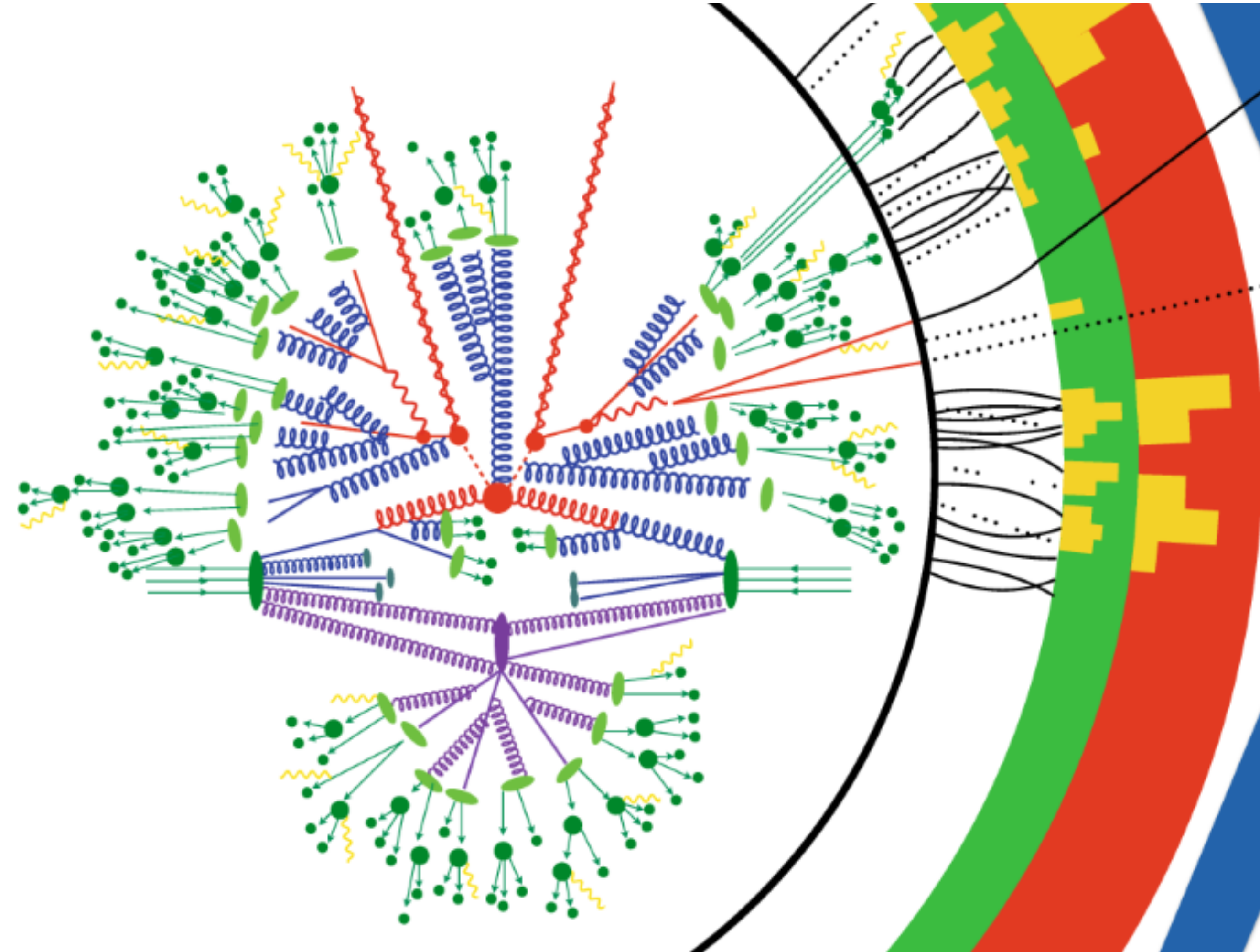
UNIVERSITY OF
OXFORD

LHC EVENT

ENERGY
SCALE [GeV]

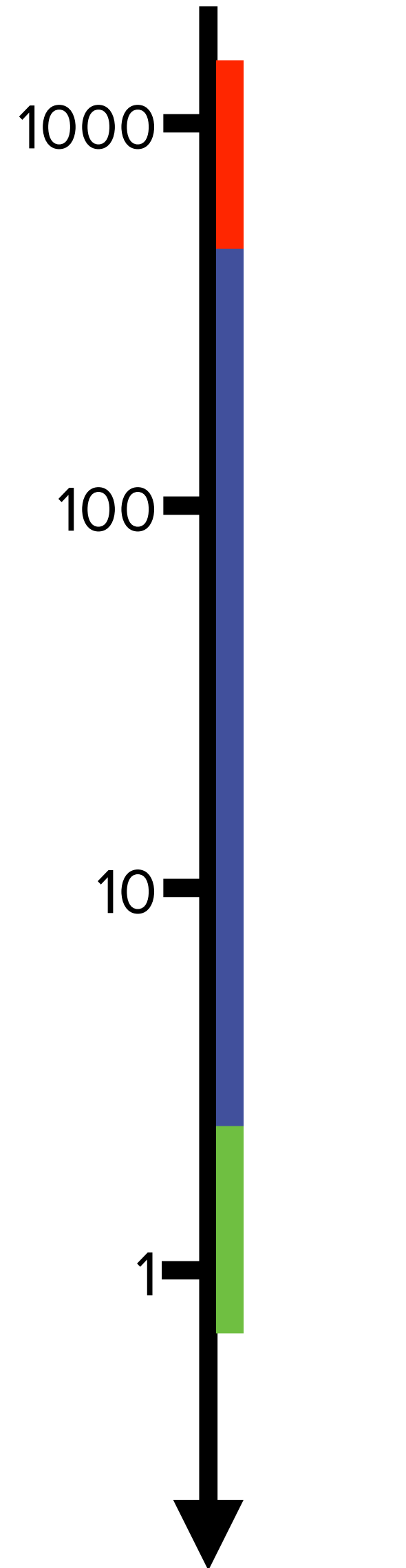


- 1. HARD SCATTERING
- 2. PARTON SHOWER
- 3. HADRONIZATION
- 4. UNDERLYING EVENT



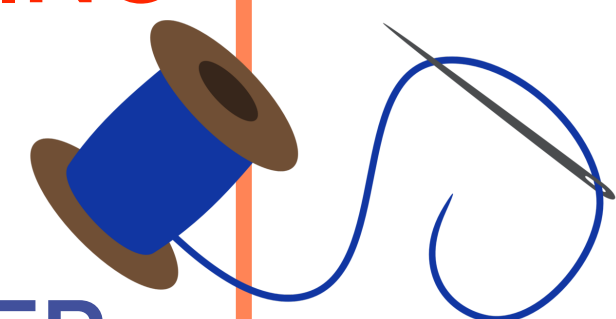
LHC EVENT

ENERGY
SCALE [GeV]



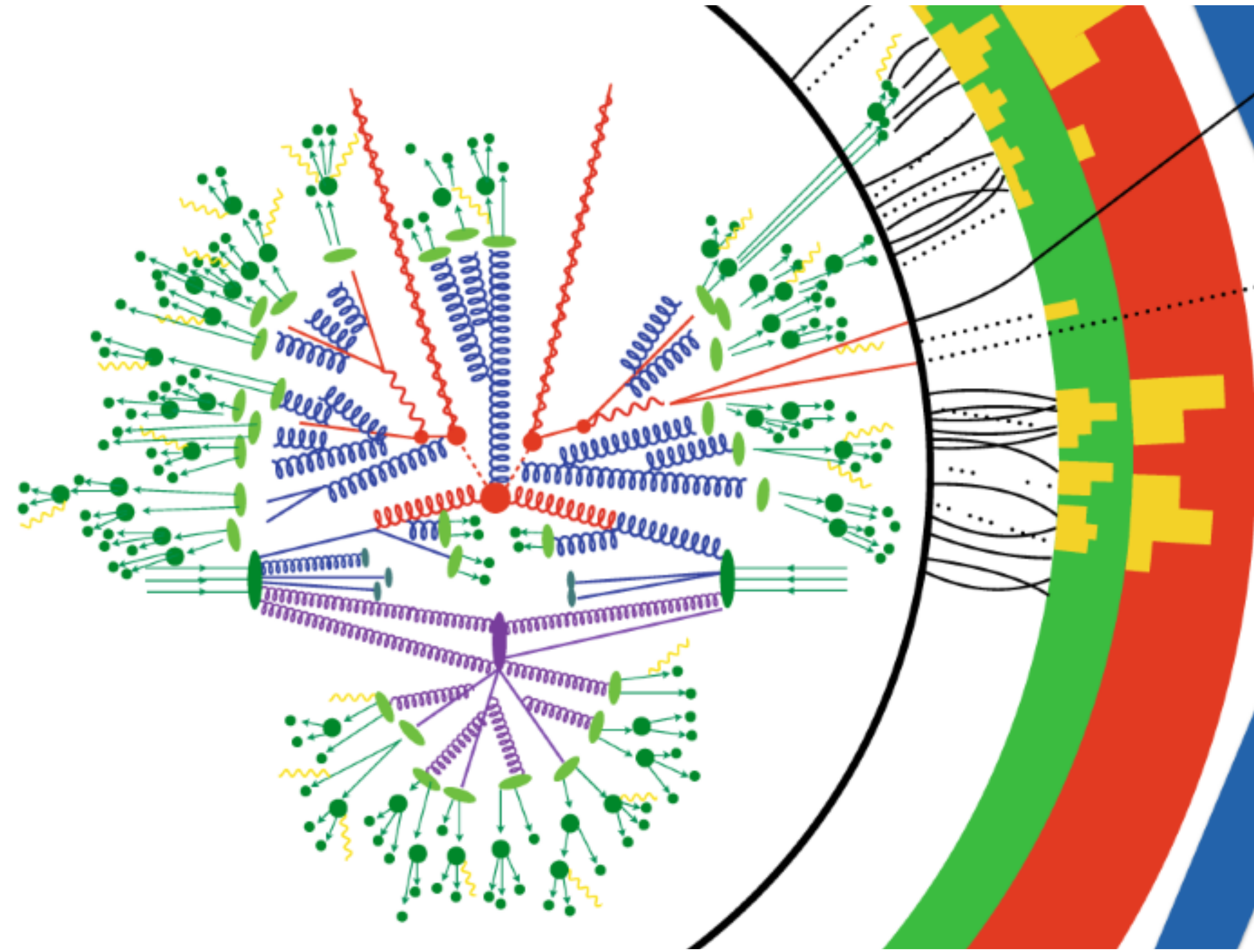
MATCHING

- 1. HARD SCATTERING
- 2. PARTON SHOWER



3. HADRONIZATION

4. UNDERLYING EVENT

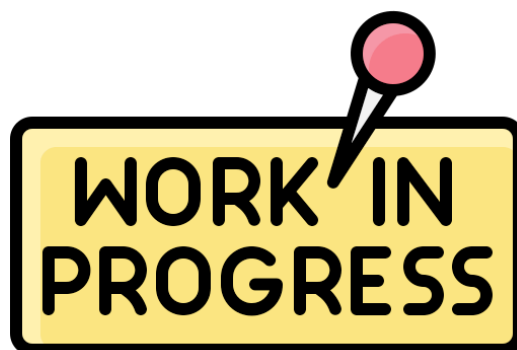


MATCHING



NLO+LL_{PS}

- A **solved problem** for long time.
- Completely understood and **fully automatized**.
- Two main approaches available: POWHEG [Nason '04; Frixione, Nason, Oleari '07; Alioli, Nason, Oleari, Re '10] and MC@NLO [Frixione, Webber '02].



NNLO+LL_{PS}

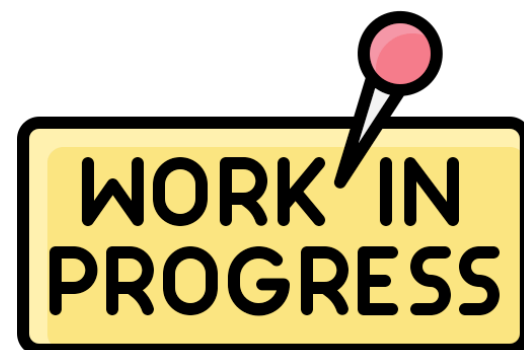
- **State-of-the-art** for precision LHC phenomenology.
- Lots of ongoing effort, **many processes already implemented**.
- Two main methods available: MiNNLO_{PS} [Monni, Nason, Re, Wiesemann, Zanderighi '19] and Geneva [Alioli, Bauer, Berggren, Tackmann, Walsh, Zuberi '13, + subsequent papers].

MATCHING



NLO+LL_{PS}

- A **solved problem** for long time.
- Completely understood and **fully automatized**.
- Two main approaches available: POWHEG [Nason '04; Frixione, Nason, Oleari '07; Alioli, Nason, Oleari, Re '10] and MC@NLO [Frixione, Webber '02].



NNLO+LL_{PS}

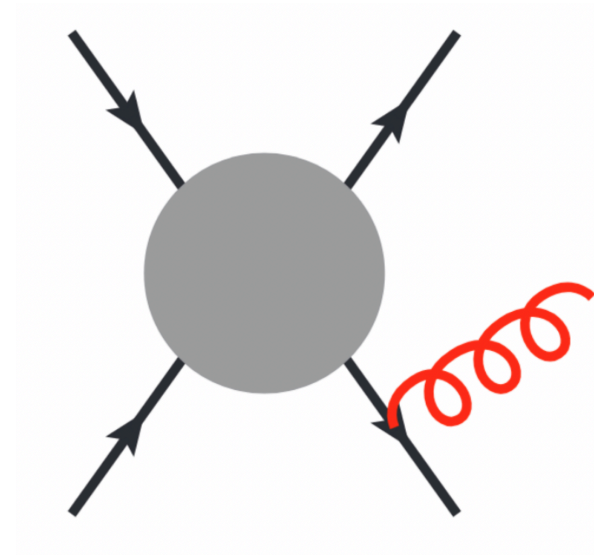
- **State-of-the-art** for precision LHC phenomenology.
- Lots of ongoing effort, **many processes already implemented**.
- Two main methods available: MiNNLO_{PS} [Monni, Nason, Re, Wiesemann, Zanderighi '19] and Geneva [Alioli, Bauer, Berggren, Tackmann, Walsh, Zuberi '13, + subsequent papers].



Logarithmically accurate showers are now available, but matching is still under investigation (at NLO). The method I am presenting today does not directly apply to matching with NLL showers!

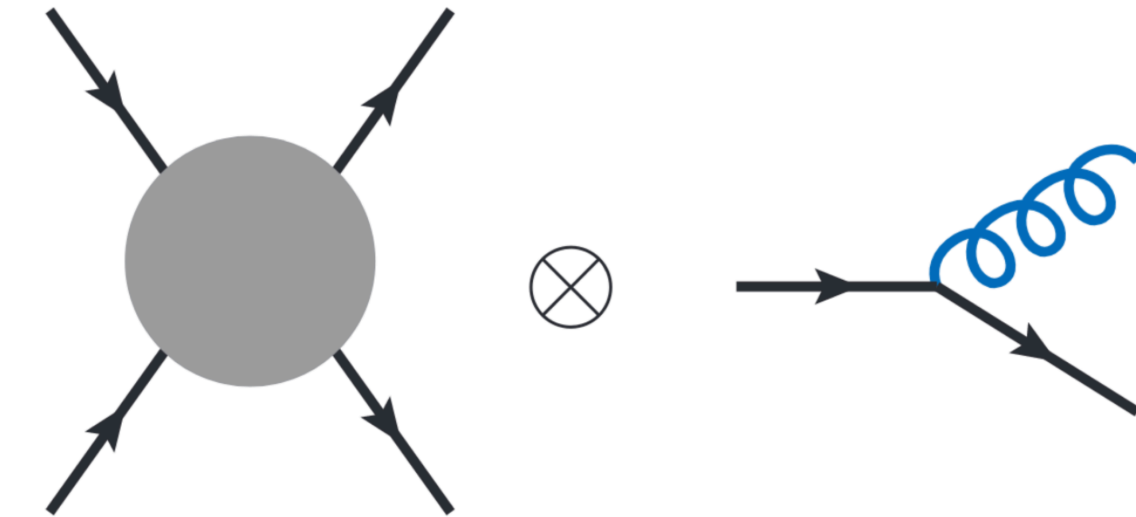
WHAT'S THE PROBLEM?

FIXED-ORDER CALCULATIONS



Correct real emission

PARTON SHOWER



Approximate real emission

!! DOUBLE COUNTING !!

THE POWHEG METHOD

[Nason '04;
Frixione, Nason, Oleari '07;
Alioli, Nason, Oleari, Re '10]

Master Formula

$$d\sigma_{\text{pwg}} = d\Phi_F \bar{B}(\Phi_F) \left\{ \Delta_{\text{pwg}}(\Lambda_{\text{pwg}}) + \int d\Phi_{\text{rad}} \Delta_{\text{pwg}}(p_{T,\text{rad}}) \frac{R(\Phi_F, \Phi_{\text{rad}})}{B(\Phi_F)} \right\}$$

NLO NORMALIZATION (= xs)

$$\bar{B}(\Phi_F) = B(\Phi_F) + V(\Phi_F) + \int d\Phi_{\text{rad}} [R(\Phi_{FJ}) - C(\Phi_{FJ})]$$

FIRST (= hardest) EMISSION
obtained with the correct matrix element R/B

$$\Delta(p_T) = \exp \left\{ - \int d\Phi'_{\text{rad}} \frac{R(\Phi_F, \Phi'_{\text{rad}})}{B(\Phi_F)} \Theta(p'_T - p_T) \right\}$$

When using a p_T -ordered shower (most common option, like PYTHIA), we apply a p_T -veto: all the emissions produced by the shower must be softer than the first emission produced by POWHEG.

THE MINNLO_{PS} METHOD

[Monni, Nason, Re, Wiesemann, Zanderighi '19]

Master Formula

$$d\sigma_F^{\text{MiNNLO}_{\text{PS}}} = d\Phi_{\text{FJ}} \bar{B}^{\text{MiNNLO}_{\text{PS}}}(\Phi_{\text{FJ}}) \left\{ \Delta_{\text{pwg}}(\Lambda_{\text{pwg}}) + \int d\Phi_{\text{rad}} \Delta_{\text{pwg}}(p_{\text{T,rad}}) \frac{R(\Phi_{\text{FJ}}, \Phi_{\text{rad}})}{B(\Phi_{\text{FJ}})} \right\}$$

NNLO NORMALIZATION (= xs)

SECOND EMISSION
obtained à la POWHEG

$$\bar{B}^{\text{MiNNLO}_{\text{PS}}}(\Phi_{\text{FJ}}) =$$

THE MINNLO_{PS} METHOD

[Monni, Nason, Re, Wiesemann, Zanderighi '19]

Master Formula

$$d\sigma_F^{\text{MiNNLO}_{\text{PS}}} = d\Phi_{\text{FJ}} \bar{B}^{\text{MiNNLO}_{\text{PS}}}(\Phi_{\text{FJ}}) \left\{ \Delta_{\text{pwg}}(\Lambda_{\text{pwg}}) + \int d\Phi_{\text{rad}} \Delta_{\text{pwg}}(p_{\text{T,rad}}) \frac{R(\Phi_{\text{FJ}}, \Phi_{\text{rad}})}{B(\Phi_{\text{FJ}})} \right\}$$

NNLO NORMALIZATION (= xs)

SECOND EMISSION
obtained à la POWHEG

$$\bar{B}^{\text{MiNNLO}_{\text{PS}}}(\Phi_{\text{FJ}}) = \left(\begin{array}{l} B(\Phi_{\text{FJ}}) \\ \downarrow \\ \text{Born FJ} \\ \text{! DIVERGENT} \end{array} + V(\Phi_{\text{FJ}}) + \int d\Phi_{\text{rad}} R(\Phi_{\text{FJJ}}) \right)$$

Virtual+Real
on FJ

	F	F+J	F+JJ	F+(>3J)
F@NNLO+PS	-	NLO	LO	PS (LL)

THE MINNLO_{PS} METHOD

[Monni, Nason, Re, Wiesemann, Zanderighi '19]

Master Formula

$$d\sigma_F^{\text{MiNNLO}_{\text{PS}}} = d\Phi_{\text{FJ}} \bar{B}^{\text{MiNNLO}_{\text{PS}}}(\Phi_{\text{FJ}}) \left\{ \Delta_{\text{pwg}}(\Lambda_{\text{pwg}}) + \int d\Phi_{\text{rad}} \Delta_{\text{pwg}}(p_{\text{T,rad}}) \frac{R(\Phi_{\text{FJ}}, \Phi_{\text{rad}})}{B(\Phi_{\text{FJ}})} \right\}$$

NNLO NORMALIZATION (= xs)

SECOND EMISSION
obtained à la POWHEG

$$\bar{B}^{\text{MiNNLO}_{\text{PS}}}(\Phi_{\text{FJ}}) = e^{-\tilde{S}(p_T)} \left(B(\Phi_{\text{FJ}}) (1 + \alpha_s \tilde{S}^{(1)}) + V(\Phi_{\text{FJ}}) + \int d\Phi_{\text{rad}} R(\Phi_{\text{FJJ}}) \right)$$

Sudakov form factor

Born FJ



~~DIVERGENT~~

Virtual+Real
on FJ

Correct NLO on FJ

	F	F+J	F+JJ	F+(>3J)
F@NNLO+PS	NLO	NLO	LO	PS (LL)

$$\tilde{S}(p_T) = \int_{p_T^2}^{Q^2} \frac{dq^2}{q^2} \left[A \log \frac{Q^2}{q^2} + B \right]$$

THE MINNLO_{PS} METHOD

[Monni, Nason, Re, Wiesemann, Zanderighi '19]

Master Formula

$$d\sigma_F^{\text{MiNNLO}_{\text{PS}}} = d\Phi_{\text{FJ}} \bar{B}^{\text{MiNNLO}_{\text{PS}}}(\Phi_{\text{FJ}}) \left\{ \Delta_{\text{pwg}}(\Lambda_{\text{pwg}}) + \int d\Phi_{\text{rad}} \Delta_{\text{pwg}}(p_{\text{T,rad}}) \frac{R(\Phi_{\text{FJ}}, \Phi_{\text{rad}})}{B(\Phi_{\text{FJ}})} \right\}$$

NNLO NORMALIZATION (= xs)

**SECOND EMISSION
obtained à la POWHEG**

$$\bar{B}^{\text{MiNNLO}_{\text{PS}}}(\Phi_{\text{FJ}}) = e^{-\tilde{S}(p_{\text{T}})} \left(B(\Phi_{\text{FJ}}) (1 + \alpha_s \tilde{S}^{(1)}) + V(\Phi_{\text{FJ}}) + \int d\Phi_{\text{rad}} R(\Phi_{\text{FJJ}}) + (D(p_{\text{T}}) - \alpha_s D^{(1)}(p_{\text{T}}) - \alpha_s^2 D^{(2)}(p_{\text{T}})) \mathcal{F} \right)$$

Sudakov form factor

$$\tilde{S}(p_{\text{T}}) = \int_{p_{\text{T}}^2}^{Q^2} \frac{dq^2}{q^2} \left[A \log \frac{Q^2}{q^2} + B \right]$$



Born FJ
~~DIVERGENT~~

Correct NLO on FJ

Virtual+Real
on FJ

α_s^3 correction needed
for NNLO normalization

spreading
 $\Phi_{\text{F}} \rightarrow \Phi_{\text{FJ}}$

	F	F+J	F+JJ	F+(>3J)
F@NNLO+PS	NNLO	NLO	LO	PS (LL)

THE MINNLO_{PS} METHOD

- Analytic all-order formula:

$$\frac{d\sigma}{d\Phi_F dp_T} = \frac{d\sigma^{\text{sing}}}{d\Phi_F dp_T} + R(p_T) = \frac{d}{dp_T} \left\{ e^{-\tilde{S}(p_T)} \mathcal{L}(p_T) \right\} + R(p_T) = e^{-\tilde{S}(p_T)} \left[D(p_T) + \frac{R(p_T)}{e^{-\tilde{S}(p_T)}} \right]$$

$D(p_T) \equiv -\frac{d\tilde{S}(p_T)}{dp_T} \mathcal{L}(p_T) + \frac{d\mathcal{L} p_T}{dp_T}$

THE MINNLO_{PS} METHOD

- Analytic all-order formula:

$$\frac{d\sigma}{d\Phi_F dp_T} = \frac{d\sigma^{\text{sing}}}{d\Phi_F dp_T} + R(p_T) = \frac{d}{dp_T} \left\{ e^{-\tilde{S}(p_T)} \mathcal{L}(p_T) \right\} + R(p_T) = e^{-\tilde{S}(p_T)} \left[D(p_T) + \frac{R(p_T)}{e^{-\tilde{S}(p_T)}} \right]$$

$$D(p_T) \equiv -\frac{d\tilde{S}(p_T)}{dp_T} \mathcal{L}(p_T) + \frac{d\mathcal{L} p_T}{dp_T}$$

- Combine with FJ fixed-order $d\sigma_{\text{FJ}}$ and expand up to α_s^3 :

$$\mu_R = \mu_F = p_T$$

$$\int \frac{dp_T}{p_T} \alpha_s^m(p_T) \ln^n \frac{p_T}{Q} e^{-\tilde{S}(p_T)} \approx \mathcal{O}\left(\alpha_s^{m-\frac{n+1}{2}}\right)$$

$$d\sigma_F = d\sigma_F^{\text{sing}} + [d\sigma_{\text{FJ}}]_{\text{f.o.}} - [d\sigma_F^{\text{sing}}]_{\text{f.o.}} = e^{-\tilde{S}(p_T)} \left\{ D + \frac{[d\sigma_{\text{FJ}}]_{\text{f.o.}}}{\underbrace{[e^{-\tilde{S}(p_T)}]_{\text{f.o.}}}_{1 - \alpha_s \tilde{S}^{(1)}}} - \frac{[d\sigma_F^{\text{sing}}]_{\text{f.o.}}}{\underbrace{[e^{-\tilde{S}(p_T)}]_{\text{f.o.}}}_{-\alpha_s D^{(1)}(p_T) - \alpha_s^2 D^{(2)}(p_T)}} \right\}$$

THE MINNLO_{PS} METHOD

- Analytic all-order formula:

$$\frac{d\sigma}{d\Phi_F dp_T} = \frac{d\sigma^{\text{sing}}}{d\Phi_F dp_T} + R(p_T) = \frac{d}{dp_T} \left\{ e^{-\tilde{S}(p_T)} \mathcal{L}(p_T) \right\} + R(p_T) = e^{-\tilde{S}(p_T)} \left[D(p_T) + \frac{R(p_T)}{e^{-\tilde{S}(p_T)}} \right]$$

$$D(p_T) \equiv -\frac{d\tilde{S}(p_T)}{dp_T} \mathcal{L}(p_T) + \frac{d\mathcal{L}(p_T)}{dp_T}$$

- Combine with FJ fixed-order $d\sigma_{\text{FJ}}$ and expand up to α_s^3 :

$$\mu_R = \mu_F = p_T$$

$$\int \frac{dp_T}{p_T} \alpha_s^m(p_T) \ln^n \frac{p_T}{Q} e^{-\tilde{S}(p_T)} \approx \mathcal{O}\left(\alpha_s^{m-\frac{n+1}{2}}\right)$$

$$d\sigma_F = d\sigma_F^{\text{sing}} + [d\sigma_{\text{FJ}}]_{\text{f.o.}} - [d\sigma_F^{\text{sing}}]_{\text{f.o.}} = e^{-\tilde{S}(p_T)} \left\{ D + \frac{[d\sigma_{\text{FJ}}]_{\text{f.o.}}}{[e^{-\tilde{S}(p_T)}]_{\text{f.o.}}} - \frac{[d\sigma_F^{\text{sing}}]_{\text{f.o.}}}{[e^{-\tilde{S}(p_T)}]_{\text{f.o.}}} \right\}$$

$$1 - \alpha_s \tilde{S}^{(1)} \quad -\alpha_s D^{(1)}(p_T) - \alpha_s^2 D^{(2)}(p_T)$$

$$\bar{B}^{\text{MINNLO}_{\text{PS}}}(\Phi_{\text{FJ}}) = e^{-\tilde{S}(p_T)} \left(B(\Phi_{\text{FJ}})(1 + \alpha_s \tilde{S}^{(1)}) + V(\Phi_{\text{FJ}}) + \int d\Phi_{\text{rad}} R(\Phi_{\text{FJJ}}) + (D(p_T) - \alpha_s D^{(1)}(p_T) - \alpha_s^2 D^{(2)}(p_T)) \mathcal{F} \right)$$

WHAT CAN WE DO WITH MiNNLO_{PS}?

2 → 1 PROCESSES

H [1908.06987, 2407.01354] ✓
Z [1908.06987] ✓
W [2006.04133] ✓
bb→H [2402.04025]

2 → 2 PROCESSES

Zγ [2010.10478] ✓
γγ [2204.12602] ✓
ZZ [2108.05337] ✓
VH (H→bb) [2112.04168]
(+SMEFT [2204.00663])
WW [2103.12077] ✓
WZ [2208.12660] ✓

QQ PRODUCTION

tt [2012.14267, 2112.12135] ✓
bb [2112.04168]

QQF PRODUCTION

bbZ [2404.08598]

INCLUSION NLO EW

WZ [2208.12660] ✓
Z ongoing

THIS TALK!

EXTENSION TO PROCESSES WITH JETS

[2402.00596]

✓ = publicly available at <https://powhegbox.mib.infn.it/>

MiNNLO_{PS} FOR PROCESSES WITH JETS

- We now aim to target NNLO+PS accuracy in processes with a jet.
- N-jettiness (τ_N) is a suitable observable to study processes with jets, and the resummation is known at a sufficient order for a NNLO+PS accurate formalism.

τ_N represents how “N-jet-like” an event looks.

$$\tau_N = \frac{2}{Q^2} \sum_k \min\{q_a \cdot p_k, q_b \cdot p_k, q_1 \cdot p_k, \dots, q_N \cdot p_k\}$$

q_a, q_b = incoming particles
 q_i = signal jets
 p_k = hadronic fs particles

$\tau_N \rightarrow 0$: exactly N jet in the event

$\tau_N \rightarrow 1$: the event has hard radiation between the N signal jets

THE PLAN

MiNNLO_{PS-p_T} : NNLO accuracy on F



MiNNLO_{PS-τ₀} : NNLO accuracy on F



MiNNLO_{PS-τ₁} : NNLO accuracy on FJ

THE PLAN

MiNNLO_{PS-p_T}: NNLO accuracy on F



MiNNLO_{PS- τ_0} : NNLO accuracy on F

- ✓ Theoretical formalism
- ✓ Implementation for H/Z production
- ✓ Phenomenological analysis



MiNNLO_{PS- τ_1} : NNLO accuracy on FJ

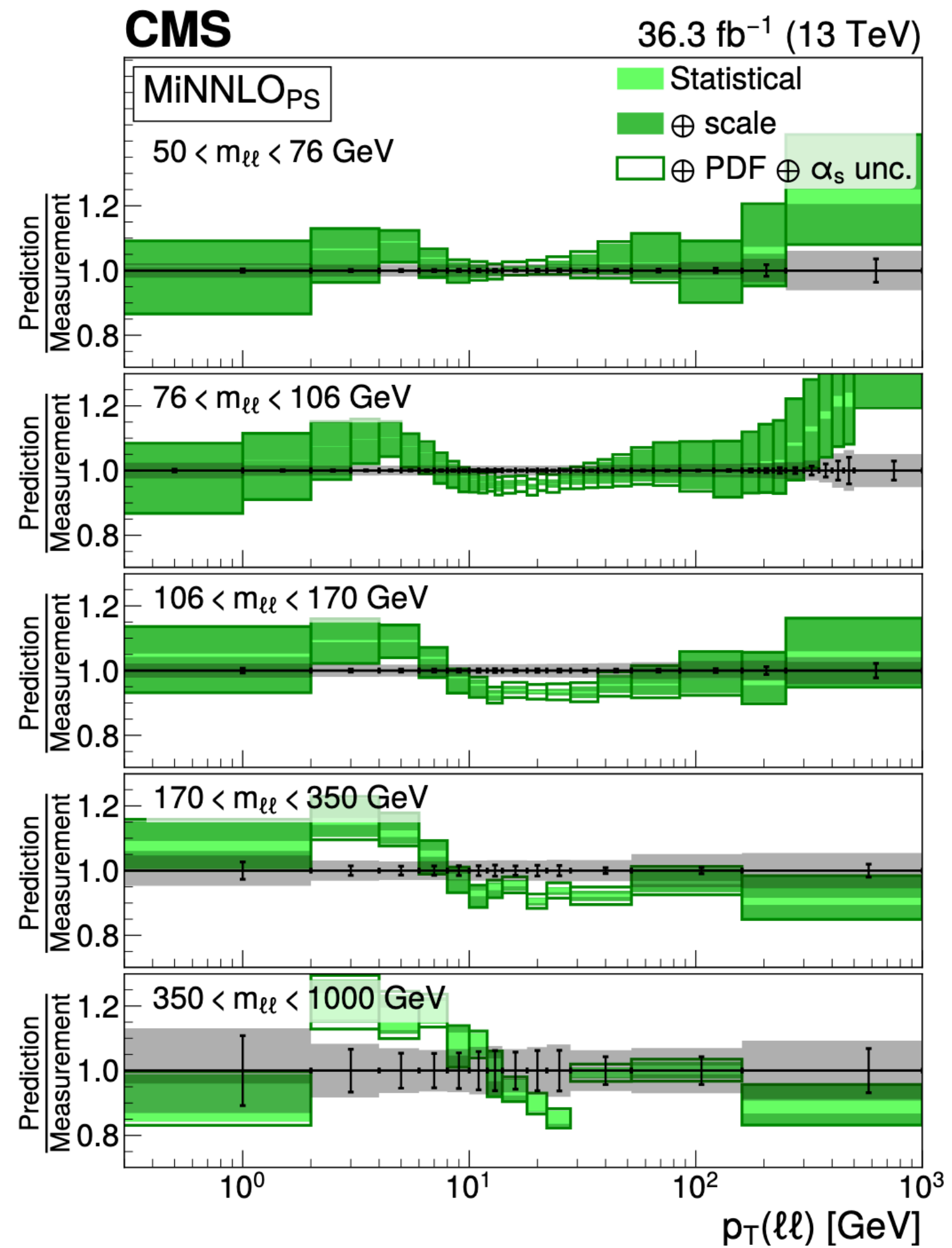
- ✓ Theoretical formalism



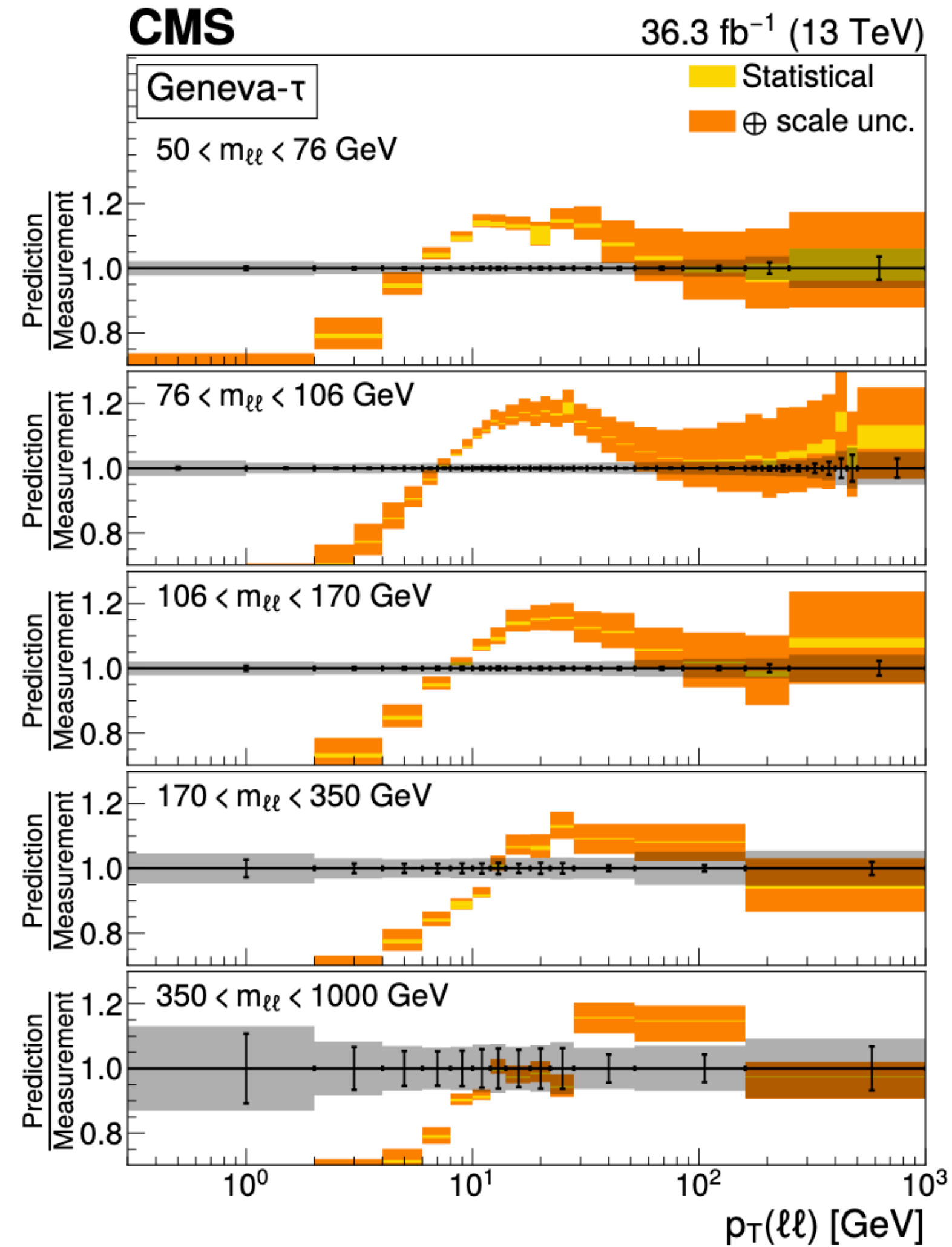
- Problems in the matching with the shower
- Large discrepancies in the 1 jet bin for H production

DY - CMS 2205.04897

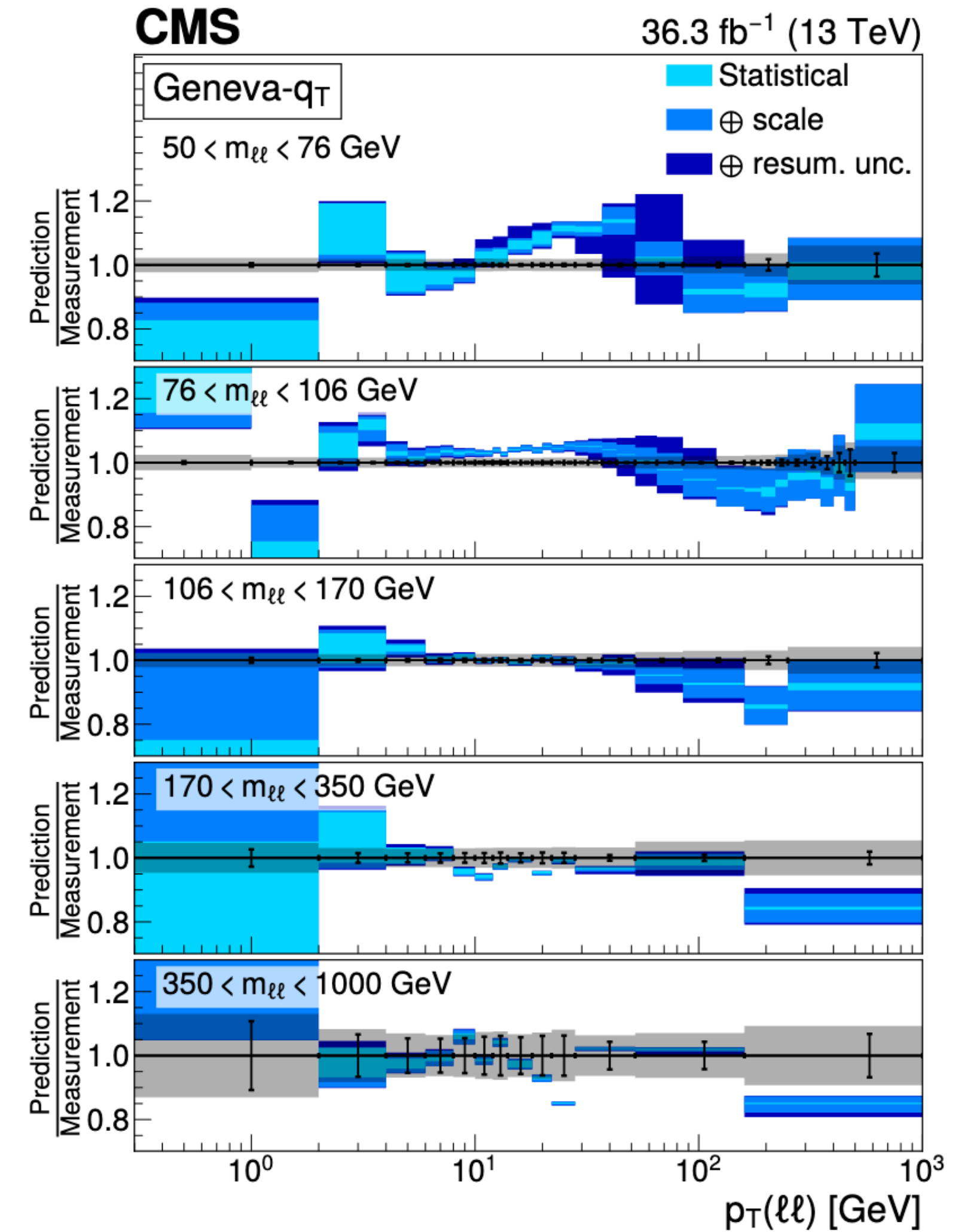
MiNNLO_{PS}-p_T



GENEVA-τ₀



GENEVA-p_T



THE MINNLO_{PS}- τ_0 METHOD

[Ebert, Rottoli, Wiesemann, Zanderighi, SZ '24]

- We start from the factorization formula of 0-jettiness at NNLL' (for 1-jettiness, include also jet function)

$$\frac{d\sigma}{d\Phi_F d\tau_0} = \frac{d\sigma^{\text{sing}}}{d\Phi_F d\tau_0} + R_f(\tau_0) \quad \frac{d\sigma^{\text{sing}}}{d\Phi_F d\tau_0} = \sum_{a,b} \frac{d|M_{a,b}|^2}{d\Phi_F} H_{a,b}(Q, \mu) \int dt_a dt_b B_a(t_a, x_a, \mu) B_b(t_b, x_b, \mu) S\left(\tau_0 - \frac{t_a}{Q} - \frac{t_b}{Q}, \mu\right)$$

- We evolve and expand all the needed ingredients in order to obtain: $\frac{d\sigma^{\text{sing}}}{d\Phi_F d\tau_0} = \frac{d}{d\tau_0} \left(e^{-\tilde{S}(\tau_0)} \mathcal{L}(\tau_0) \right)$

- We follow all the steps presented for MiNNLO_{PS}-p_T to obtain:

$$\bar{B}^{\text{MiNNLO}_{\text{PS}}}(\Phi_{\text{FJ}}) = e^{-\tilde{S}(\tau_0)} \left(B(\Phi_{\text{FJ}}) (1 + \alpha_s \tilde{S}^{(1)}) + V(\Phi_{\text{FJ}}) + \int d\Phi_{\text{rad}} R(\Phi_{\text{FJJ}}) + (D(\tau_0) - \alpha_s D^{(1)}(\tau_0) - \alpha_s^2 D^{(2)}(\tau_0)) \mathcal{F} \right)$$

THE MINNLO_{PS}- τ_0 METHOD

[Ebert, Rottoli, Wieseemann, Zanderighi, SZ '24]

- We start from the factorization formula of 0-jettiness at NNLL' (for 1-jettiness, include also jet function)

$$\frac{d\sigma}{d\Phi_F d\tau_0} = \frac{d\sigma^{\text{sing}}}{d\Phi_F d\tau_0} + R_f(\tau_0) \quad \frac{d\sigma^{\text{sing}}}{d\Phi_F d\tau_0} = \sum_{a,b} \frac{d|M_{a,b}|^2}{d\Phi_F} H_{a,b}(Q, \mu) \int dt_a dt_b B_a(t_a, x_a, \mu) B_b(t_b, x_b, \mu) S\left(\tau_0 - \frac{t_a}{Q} - \frac{t_b}{Q}, \mu\right)$$

- We evolve and expand all the needed ingredients in order to obtain: $\frac{d\sigma^{\text{sing}}}{d\Phi_F d\tau_0} = \frac{d}{d\tau_0} \left(e^{-\tilde{S}(\tau_0)} \mathcal{L}(\tau_0) \right)$

- We follow all the steps presented for MiNNLO_{PS}- p_T to obtain:

$$\bar{B}^{\text{MiNNLO}_{\text{PS}}}(\Phi_{\text{FJ}}) = e^{-\tilde{S}(\tau_0)} \left(B(\Phi_{\text{FJ}}) (1 + \alpha_s \tilde{S}^{(1)}) + V(\Phi_{\text{FJ}}) + \int d\Phi_{\text{rad}} R(\Phi_{\text{FJJ}}) + (D(\tau_0) - \alpha_s D^{(1)}(\tau_0) - \alpha_s^2 D^{(2)}(\tau_0)) \mathcal{F} \right)$$

The logarithmic structure in transverse momentum and jettiness resummation is different!
The MiNNLO_{PS} formulae in the two cases are the same in the structure, but contain different ingredients.

MATCHING WITH THE SHOWER

[Ebert, Rottoli, Wiesemann, Zanderighi, SZ '24]

The accuracy of the parton shower is not fully preserved because we rely on the POWHEG formalism.

Fix this issue requires a deep **modification of the POWHEG method**, mainly modifying the mappings and/or including truncated-vetoed showers. **We plan to implement these changes when matching with NLL accurate showers.**

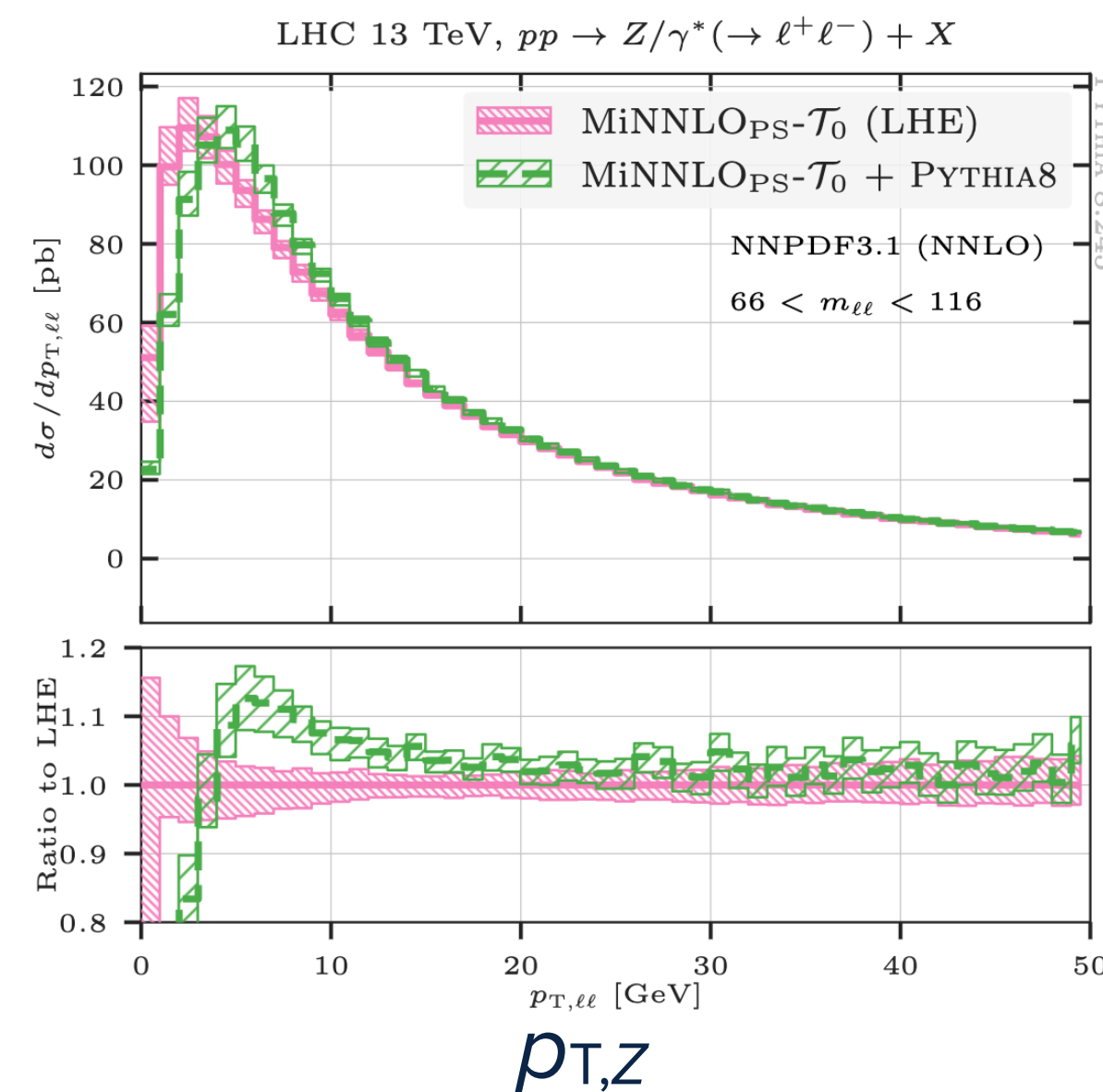
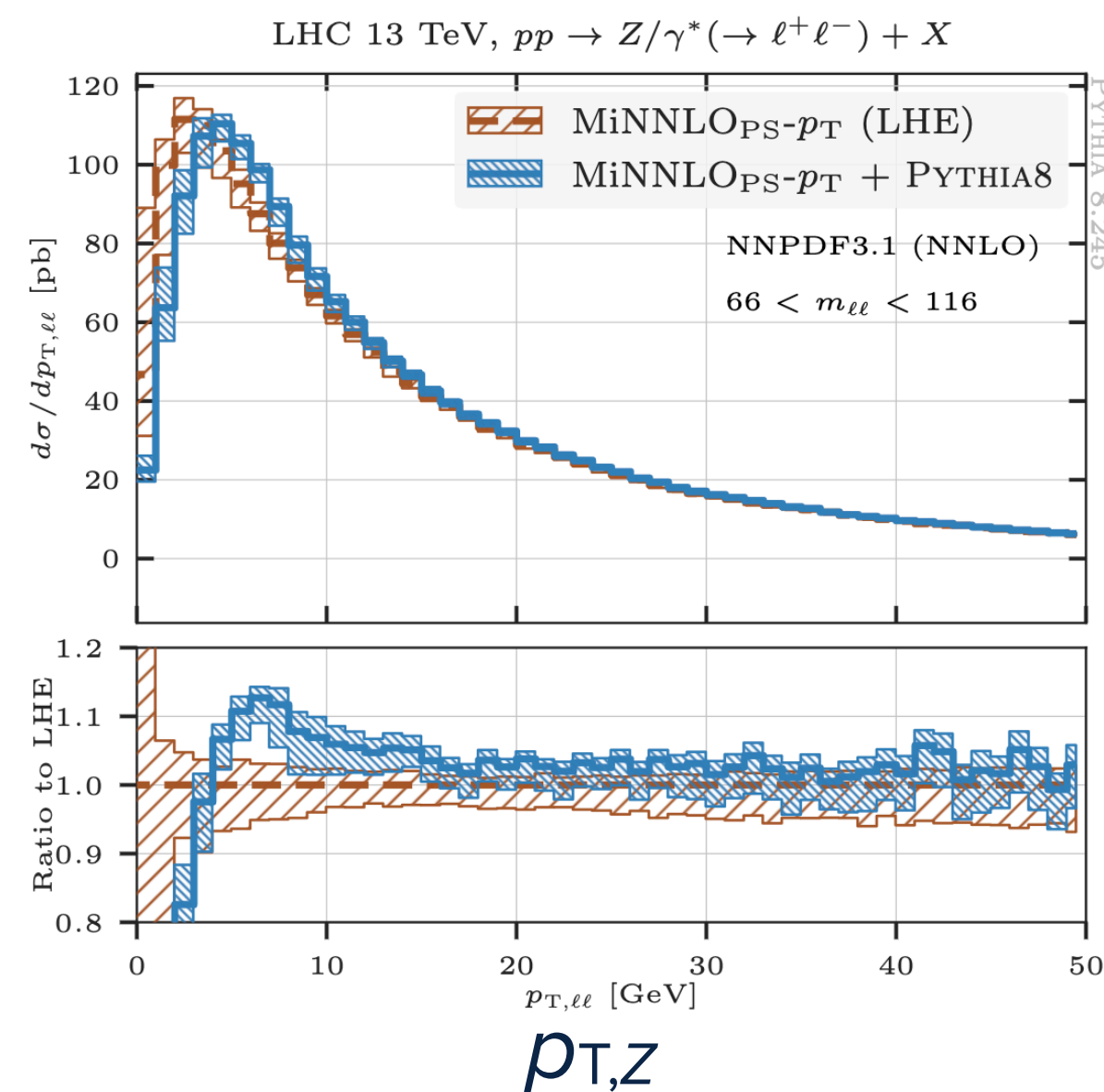
MATCHING WITH THE SHOWER

[Ebert, Rottoli, Wiesemann, Zanderighi, SZ '24]

The accuracy of the parton shower is not fully preserved because we rely on the POWHEG formalism.

Fix this issue requires a deep **modification of the POWHEG method**, mainly modifying the mappings and/or including truncated-vetoed showers. **We plan to implement these changes when matching with NLL accurate showers.**


Our results are anyway reliable because the showers effects are small and they have the same impact in $\text{MiNNLO}_{\text{PS-}p_T}$ and $\text{MiNNLO}_{\text{PS-}\tau_0}$.



- $\text{MiNNLO}_{\text{PS-}p_T}$ (LHE)
- $\text{MiNNLO}_{\text{PS-}p_T}$ (PY8)
- $\text{MiNNLO}_{\text{PS-}\tau_0}$ (LHE)
- $\text{MiNNLO}_{\text{PS-}\tau_0}$ (PY8)

CROSS SECTIONS

[Ebert, Rottoli, Wiesemann, Zanderighi, SZ '24]



	$pp \rightarrow H$ (on-shell)		$pp \rightarrow Z \rightarrow \ell^+ \ell^-$	
	σ [pb]	$\sigma/\sigma_{\text{NNLO}}$	σ [fb]	$\sigma/\sigma_{\text{NNLO}}$
NNLO	$40.32(2)^{+10.7\%}_{-10.4\%}$	1.000	$1919(1)^{+0.9\%}_{-1.1\%}$	1.000
MINNLO _{PS-p_T}	$39.33(1)^{+12.2\%}_{-11.0\%}$	0.975	$1907(2)^{+1.1\%}_{-1.2\%}$	0.994
MINNLO _{PS-\mathcal{T}_0}	$41.56(2)^{+9.4\%}_{-10.1\%}$	1.031	$1925(1)^{+1.2\%}_{-1.2\%}$	1.003

NOTE: the two MiNNLOPS descriptions differ by terms beyond accuracy, so they are expected to agree within error bands.

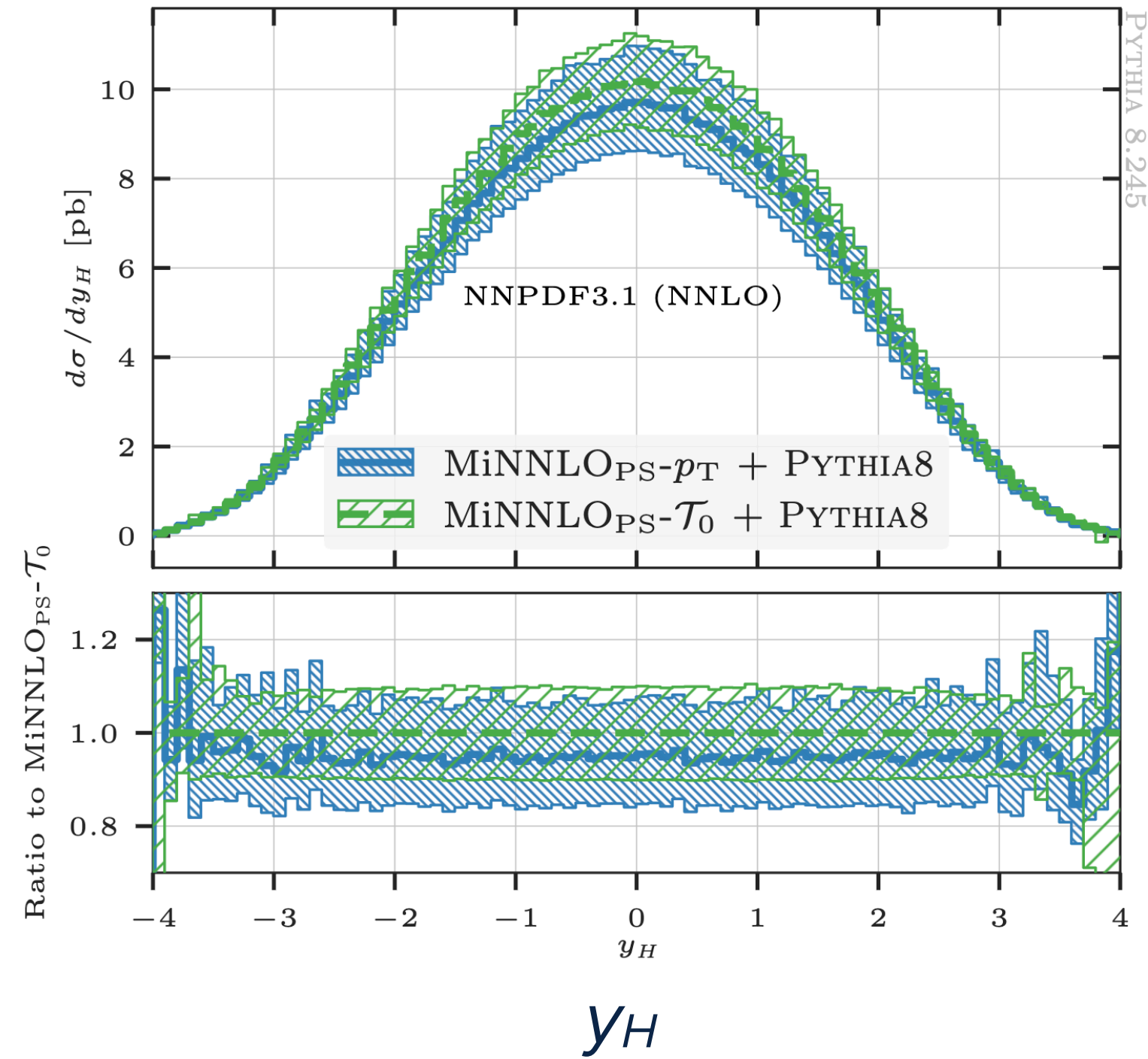
NNLO OBSERVABLES

[Ebert, Rottoli, Wiesemann, Zanderighi, SZ '24]

H PRODUCTION



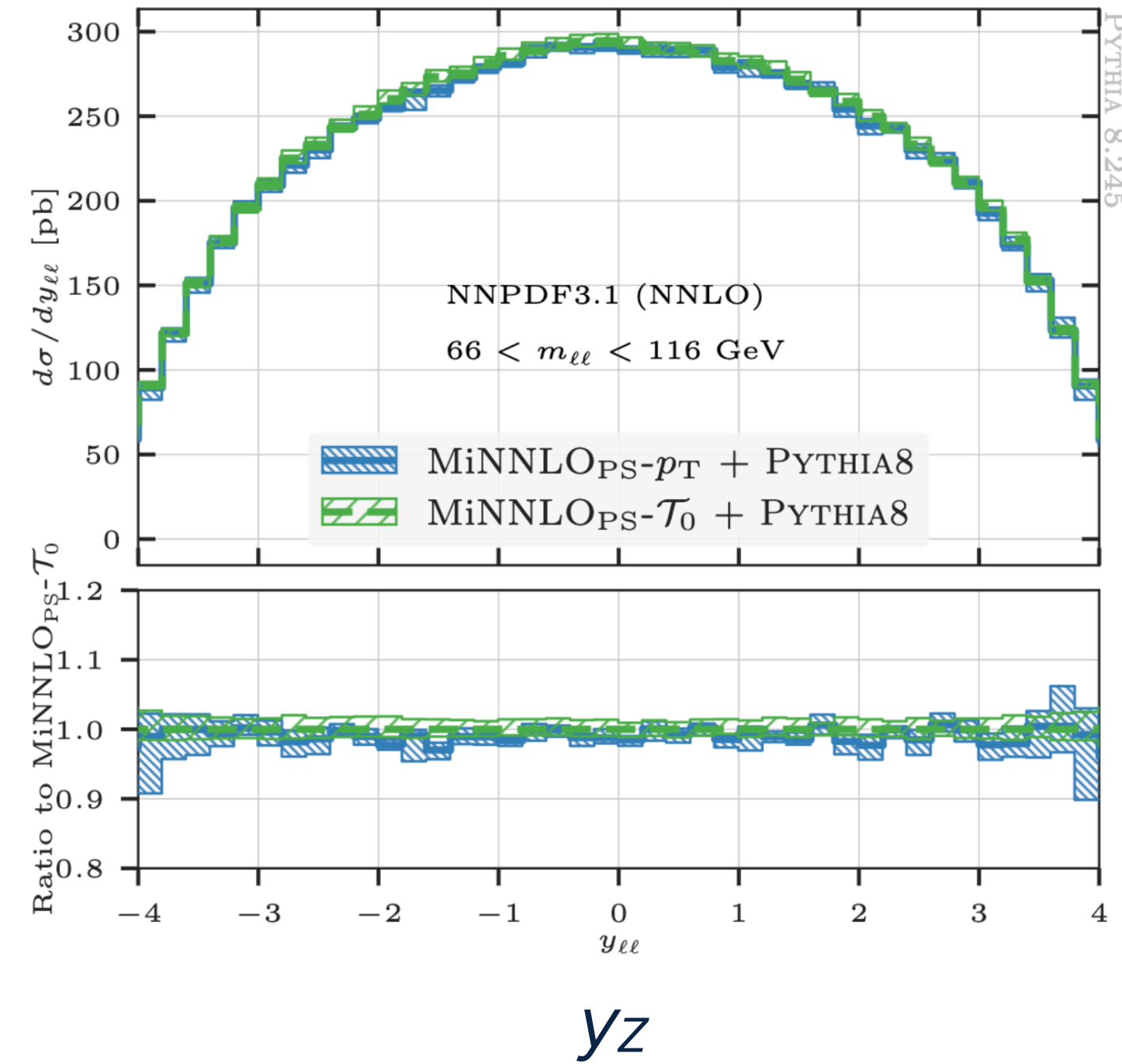
LHC 13 TeV, $pp \rightarrow H + X$



DRELL-YAN



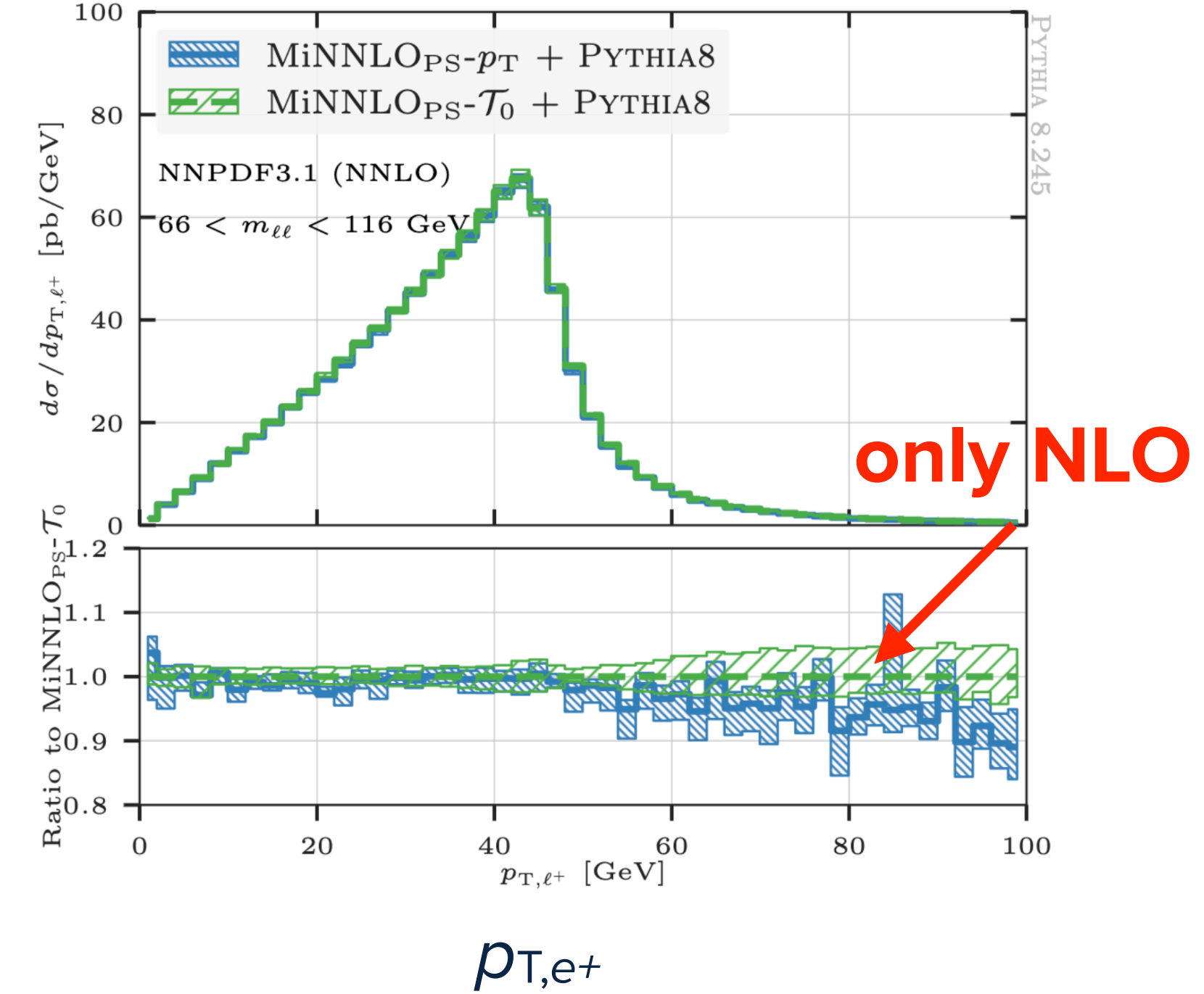
LHC 13 TeV, $pp \rightarrow Z/\gamma^*(\rightarrow \ell^+\ell^-) + X$



DRELL-YAN




LHC 13 TeV, $pp \rightarrow Z/\gamma^*(\rightarrow \ell^+\ell^-) + X$

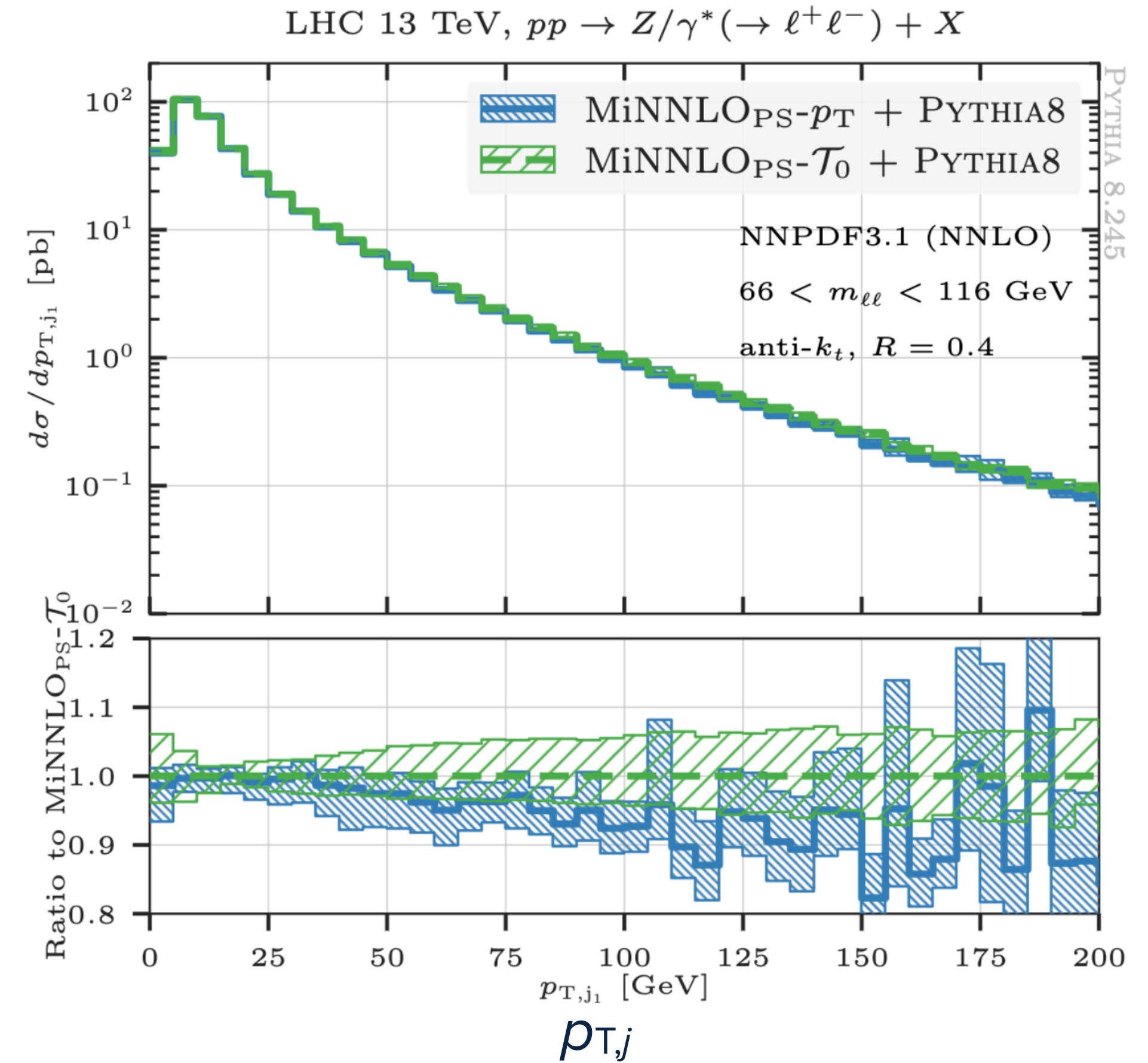
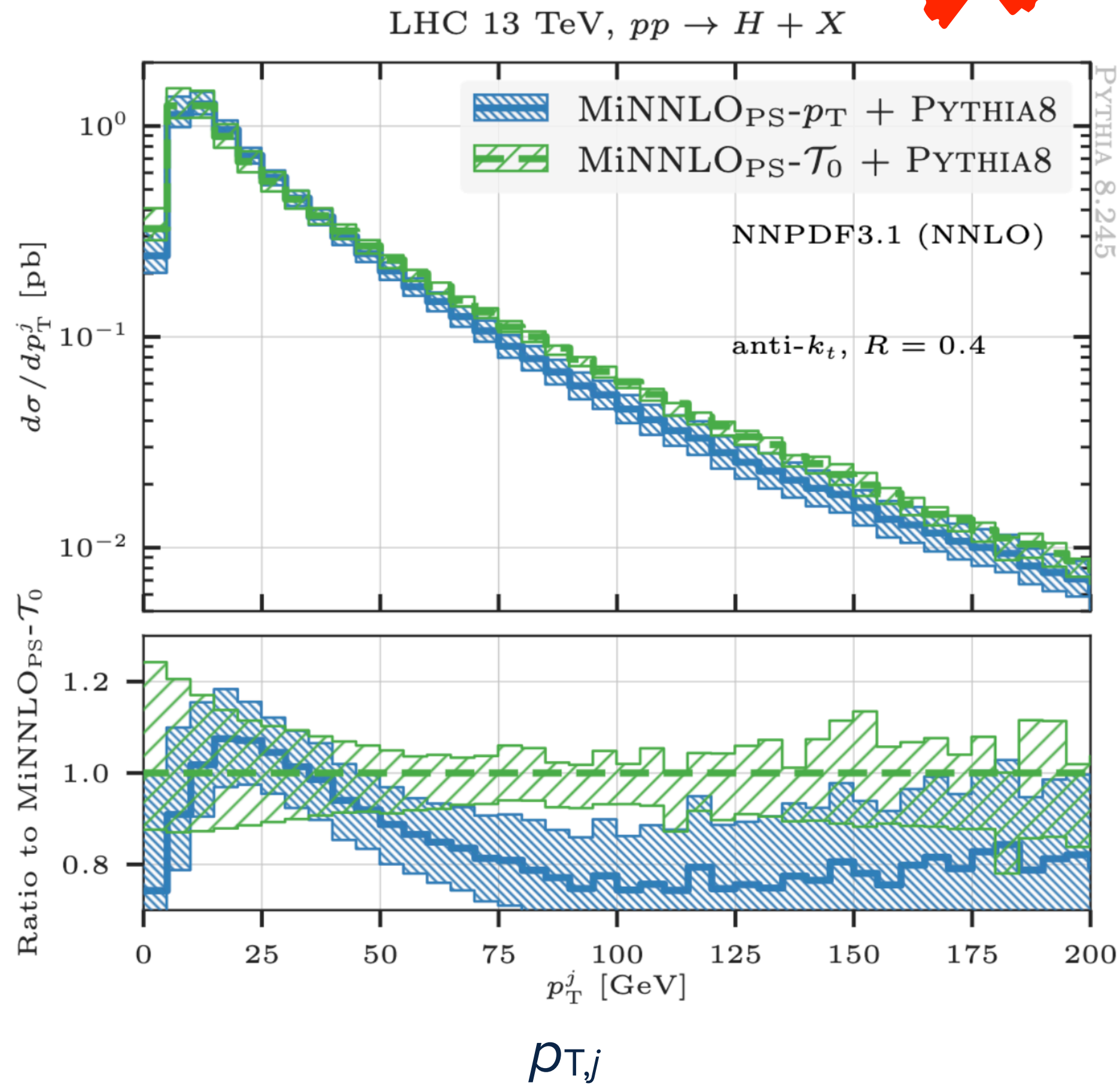


NLO OBSERVABLES

[Ebert, Rottoli, Wiesemann, Zanderighi, SZ '24]

H PRODUCTION 

DRELL-YAN 



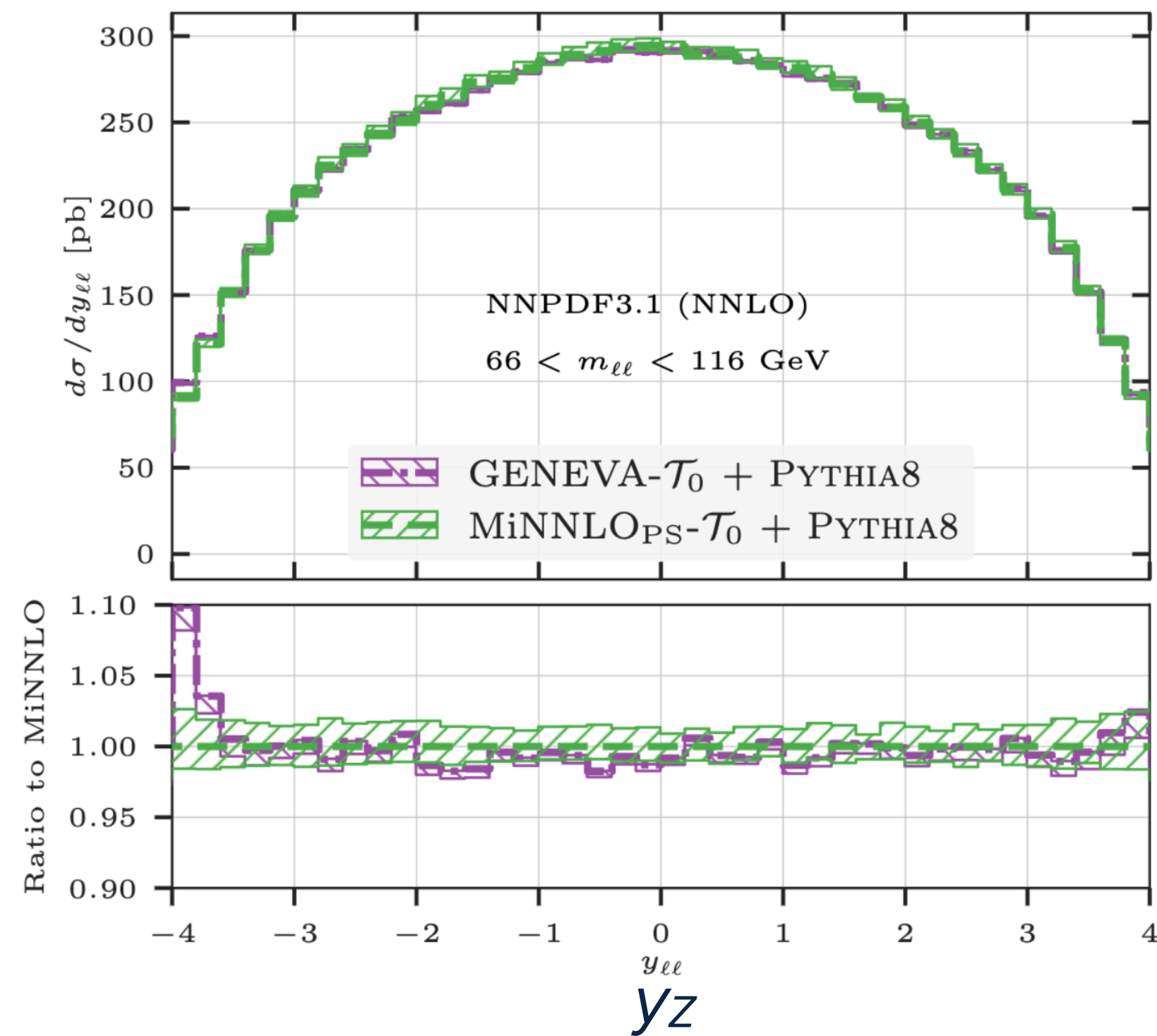
Large differences between MiNNLO- p_T and MiNNLO- τ_0 for H production **not covered by scale variation.**

COMPARISON WITH GENEVA - DY

[Ebert, Rottoli, Wiesemann, Zanderighi, SZ '24]

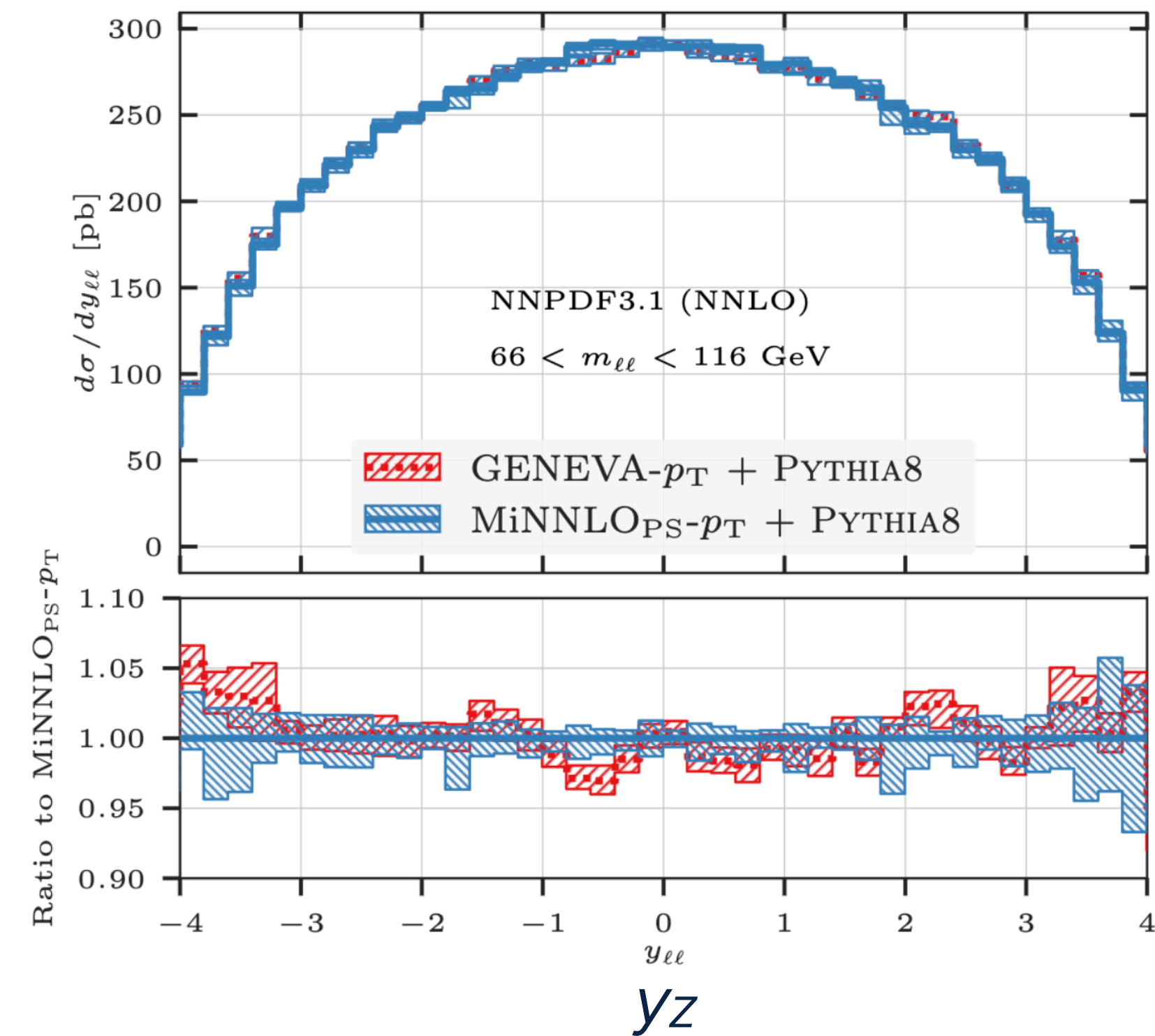
MiNNLO_{PS}- τ_0 vs GENEVA- τ_0 ✓

13 TeV, $pp \rightarrow Z/\gamma^*(\rightarrow \ell^+\ell^-) + X$



MiNNLO_{PS}- p_T vs GENEVA- p_T ✓

13 TeV, $pp \rightarrow Z/\gamma^*(\rightarrow \ell^+\ell^-) + X$

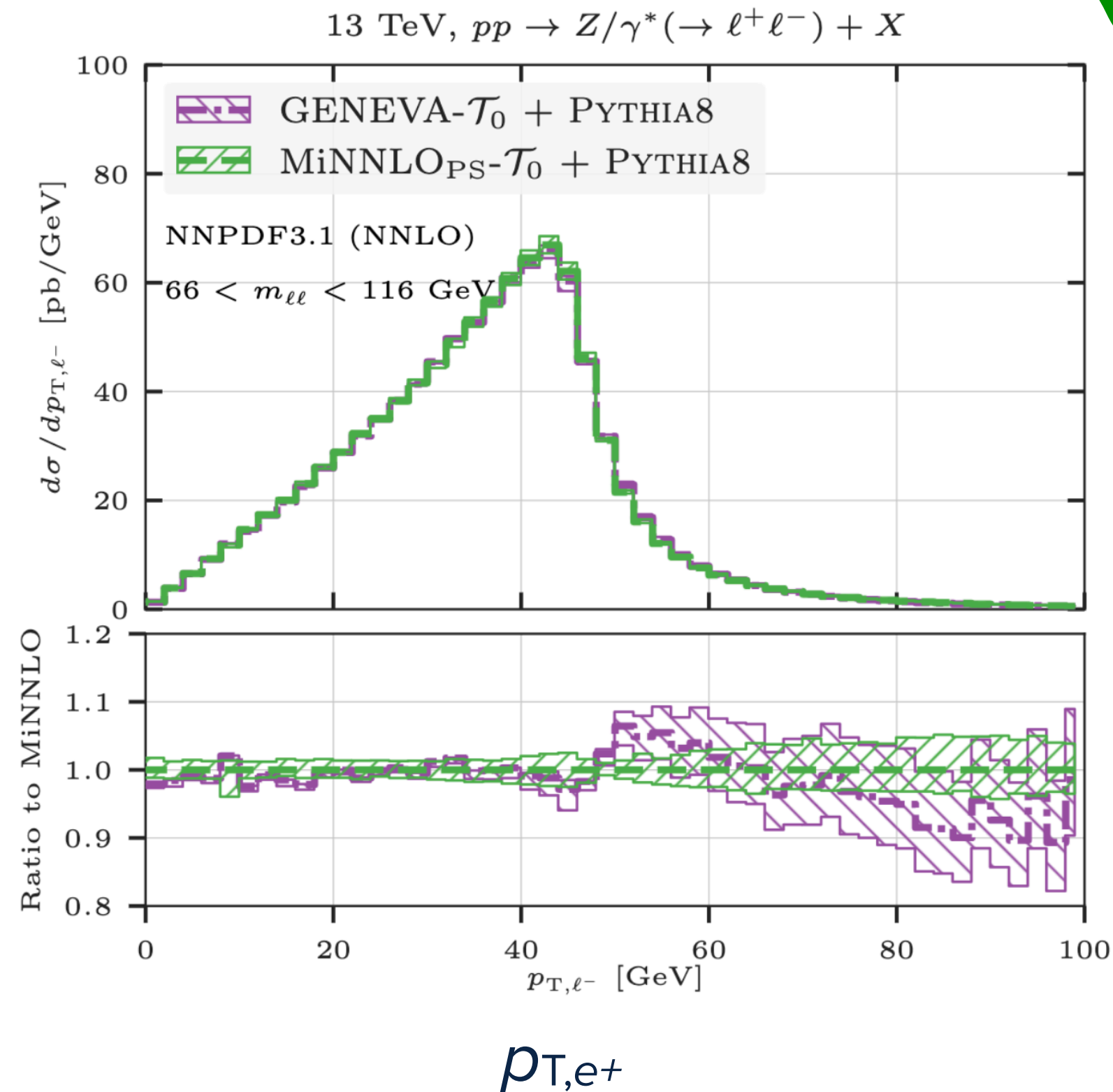


NOTE: GENEVA error bars are obtained through 3-point scale variation.

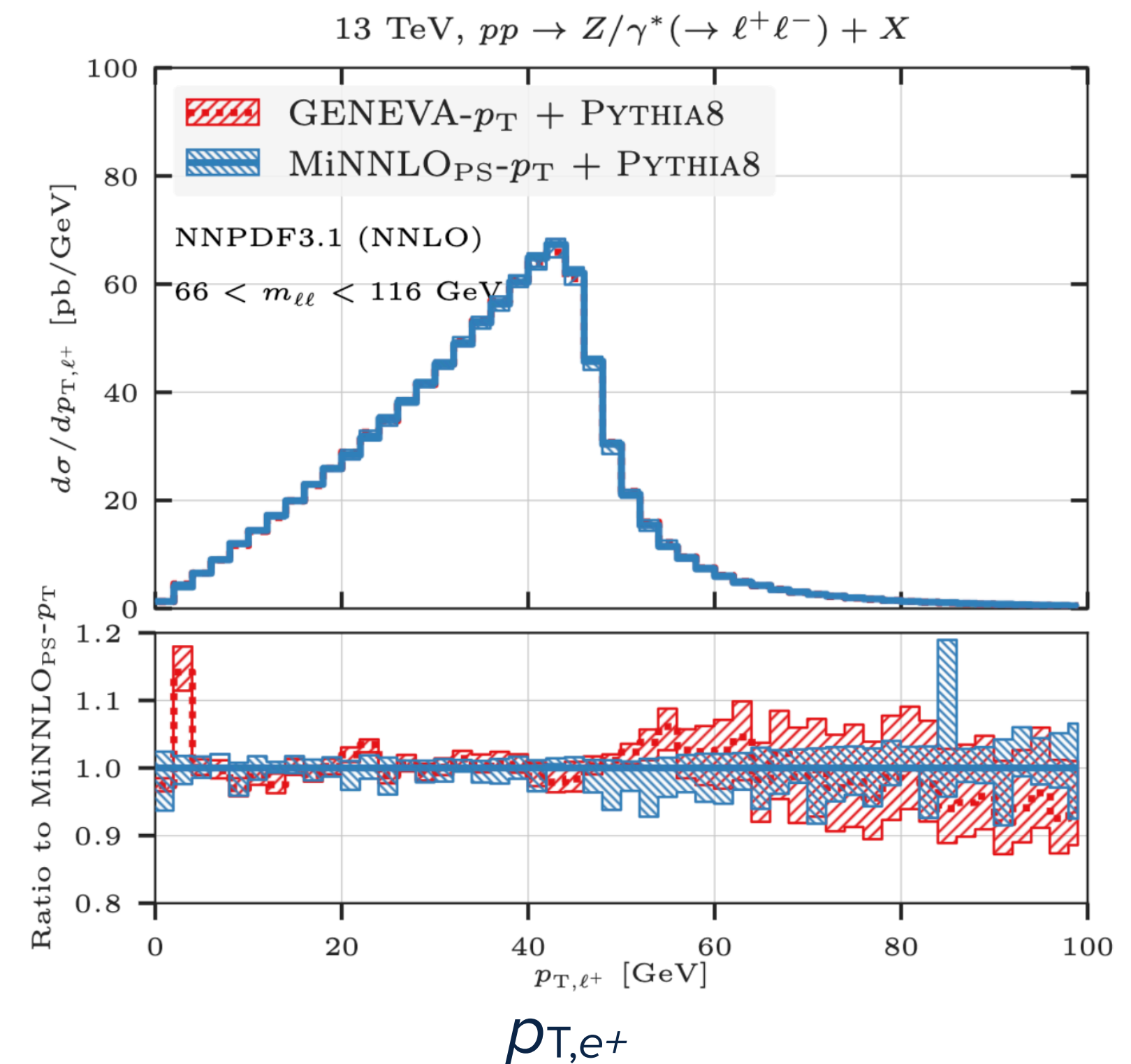
COMPARISON WITH GENEVA - DY

[Ebert, Rottoli, Wiesemann, Zanderighi, SZ '24]

MiNNLO_{PS- τ_0} vs GENEVA- τ_0



MiNNLO_{PS- p_T} vs GENEVA- p_T



NOTE: GENEVA error bars are obtained through 3-point scale variation.

COMPARISON WITH DY DATA

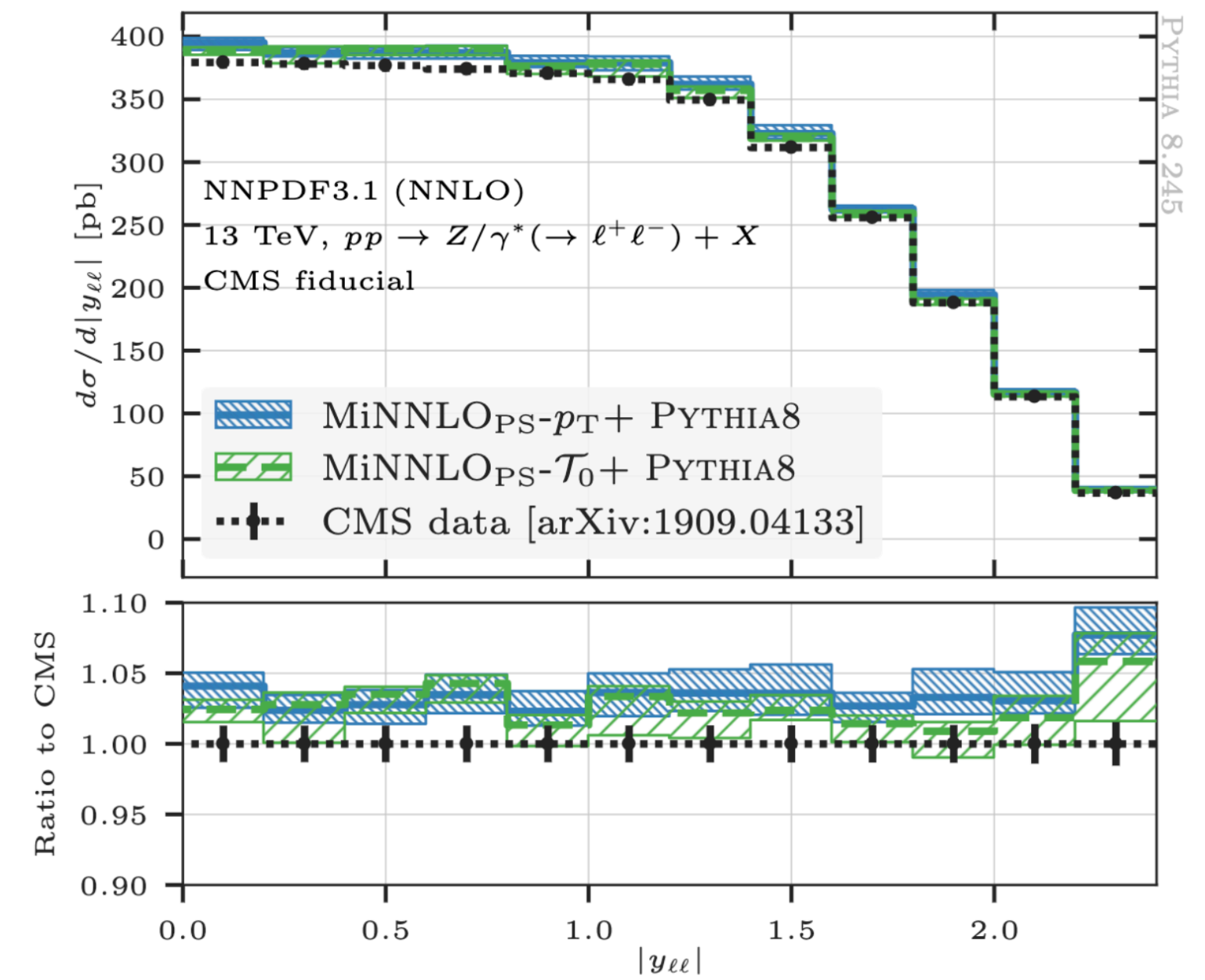
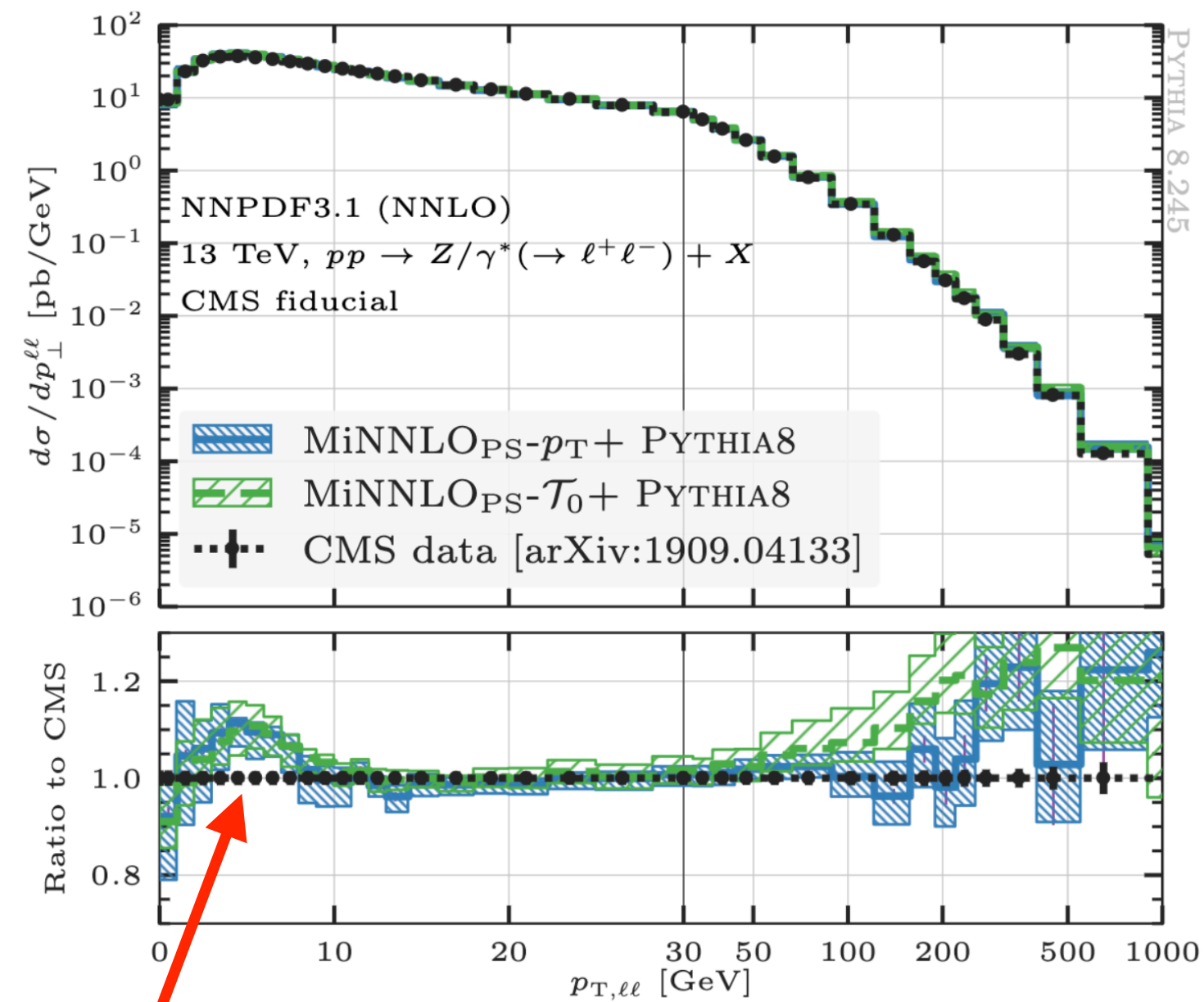
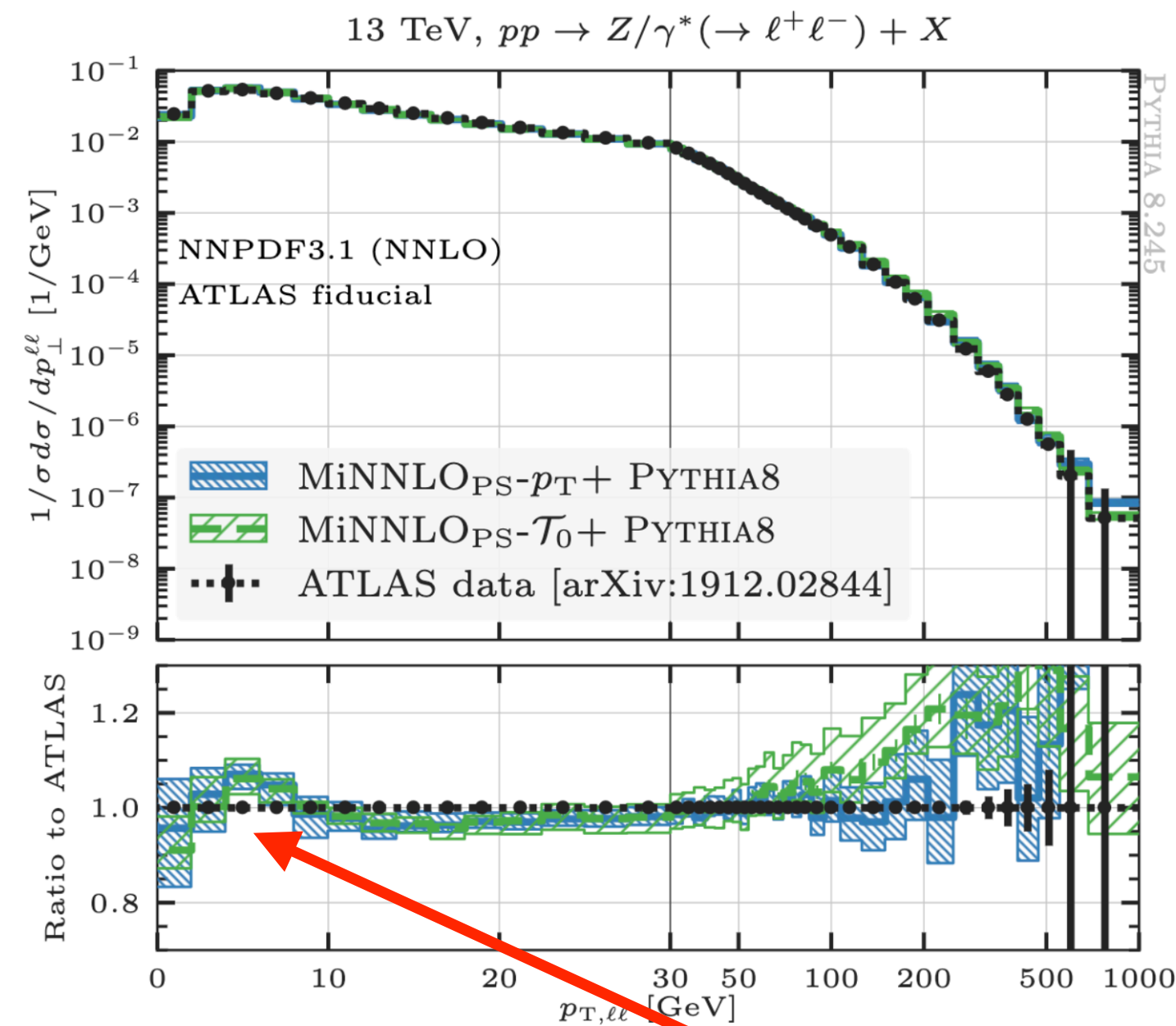


[Ebert, Rottoli, Wiesemann, Zanderighi, SZ '24]

ATLAS

CMS

CMS



$p_{T,Z}$

$p_{T,Z}$

y_Z

Accurate resummation needed

Good agreement (1-2 σ) between MiNNLO_{PS} predictions and data.

SUMMARY

- **NNLO+PS** accuracy is the **state-of-the art** for precision physics at the LHC.
- **The MiNNLO_{PS} method is a powerful framework** to reach this accuracy.
- **Extending MiNNLO_{PS} to processes with jets is not trivial.**
 - **MiNNLO_{PS} based on 0-jettiness resummation:** I presented the theoretical formalism and discussed phenomenological results for Higgs production and Drell-Yan process. We found **very good agreement for NNLO observables, while discrepancies are present for NLO ones (in the Higgs case).**
 - The **comparison with GENEVA** shows a nice agreement between the two methods, both for the p_T and the jettiness formalisms.
 - **We obtained a good agreement with ATLAS and CMS data.**
 - **MiNNLO_{PS} based on 1-jettiness resummation:** I (briefly) presented the theoretical formalism.

SUMMARY

- **NNLO+PS** accuracy is the **state-of-the art** for precision physics at the LHC.
- **The MiNNLO_{PS} method is a powerful framework** to reach this accuracy.
- **Extending MiNNLO_{PS} to processes with jets is not trivial.**
 - **MiNNLO_{PS} based on 0-jettiness resummation:** I presented the theoretical formalism and discussed phenomenological results for Higgs production and Drell-Yan process. We found **very good agreement for NNLO observables, while discrepancies are present for NLO ones (in the Higgs case).**
 - The **comparison with GENEVA** shows a nice agreement between the two methods, both for the p_T and the jettiness formalisms.
 - **We obtained a good agreement with ATLAS and CMS data.**
 - **MiNNLO_{PS} based on 1-jettiness resummation:** I (briefly) presented the theoretical formalism.

Thank you!