Theoretical predictions for  $t\bar{t}W$  and  $t\bar{t}Z$  production

Nikolaos Kidonakis

- Higher-order soft-gluon corrections
- $aN^{3}LO \ QCD + NLO \ EW$
- $t\bar{t}W$  production
- $t\bar{t}Z$  production







### Soft-gluon corrections

They are important for top-quark processes and they approximate known exact results at NLO and NNLO very well

partonic processes  $a(p_a) + b(p_b) \rightarrow t(p_t) + X$ define  $s = (p_a + p_b)^2$ ,  $t = (p_a - p_t)^2$ ,  $u = (p_b - p_t)^2$ 

For a  $2 \rightarrow n$  process with  $p_a + p_b \rightarrow p_t + p_2 + \dots + p_n$ we define the threshold variable  $s_4 = s + t + u - m_t^2 - (p_2 + \dots + p_n)^2$ Also  $s_4 = (p_2 + \dots + p_n + p_g)^2 - (p_2 + \dots + p_n)^2$  where extra gluon with  $p_g$  emitted

At partonic threshold  $p_g \to 0$  and thus  $s_4 \to 0$ 

Soft corrections  $\left[\frac{\ln^k(s_4/m_t^2)}{s_4}\right]_+$  with  $k \le 2n-1$  for the order  $\alpha_s^n$  corrections

Resum these soft corrections for the double-differential cross section

### **Soft-gluon corrections**

Note that choice of formalism and threshold variable is important two calculations may not be directly comparable even when at the same formal logarithmic accuracy if different threshold logarithms are used

Here we use single-particle-inclusive (1PI) kinematics and the threshold variable is  $s_4$ 

[another possibility of threshold variable is the invariant mass of the final state]

Finite-order expansions  $\rightarrow$  no prescription needed or used (this avoids underestimating the size of the corrections)

Approximate NNLO (aNNLO) and approximate  $N^3LO$  (aN<sup>3</sup>LO) predictions for cross sections and differential distributions

also add electroweak (EW) corrections at NLO

Thus, derive  $aN^3LO QCD + NLO EW$  predictions

### Soft-gluon Resummation

$$d\sigma_{pp \to tX} = \sum_{a,b} \int dx_a \, dx_b \, \phi_{a/p}(x_a, \mu_F) \, \phi_{b/p}(x_b, \mu_F) \, d\hat{\sigma}_{ab \to tX}(s_4, \mu_F)$$

take Laplace transforms  $d\hat{\sigma}_{ab \to tX}(N) = \int (ds_4/s) \ e^{-Ns_4/s} d\hat{\sigma}_{ab \to tX}(s_4)$ 

and  $\tilde{\phi}(N) = \int_0^1 e^{-N(1-x)} \phi(x) dx$  with transform variable N

Then

$$d\tilde{\sigma}_{ab\to tX}(N) = \tilde{\phi}_{a/a}(N_a, \mu_F) \; \tilde{\phi}_{b/b}(N_b, \mu_F) \; d\tilde{\hat{\sigma}}_{ab\to tX}(N, \mu_F)$$

Refactorization for the cross section

$$d\sigma_{ab\to tX}(N) = \tilde{\psi}_{a/a}(N_a, \mu_F) \,\tilde{\psi}_{b/b}(N_b, \mu_F) \,\tilde{J}(N, \mu_F) \,\mathrm{tr} \left\{ H_{ab\to tX}\left(\alpha_s(\mu_R)\right) \,\tilde{S}_{ab\to tX}\left(\frac{\sqrt{s}}{N\mu_F}\right) \right\}$$

 $\psi_{a/a}, \psi_{b/b} \rightarrow \text{collinear emission from incoming partons}$   $J \rightarrow \text{collinear emission from final-state gluons or massless quarks (if any)}$   $H_{ab \rightarrow tX}$  is hard function  $\rightarrow$  short distance  $S_{ab \rightarrow tX}$  is soft function  $\rightarrow$  noncollinear soft gluons

#### Thus

$$d\tilde{\hat{\sigma}}_{ab\to tX}(N) = \frac{\tilde{\psi}_{a/a}(N_a,\mu_F)\,\tilde{\psi}_{b/b}(N_b,\mu_F)}{\tilde{\phi}_{a/a}(N_a,\mu_F)\,\tilde{\phi}_{b/b}(N_b,\mu_F)}\,\tilde{J}(N,\mu_F)\,\operatorname{tr}\left\{H_{ab\to tX}\left(\alpha_s(\mu_R)\right)\,\tilde{S}_{ab\to tX}\left(\frac{\sqrt{s}}{N\mu_F}\right)\right\}$$

 $S_{ab \to tX}$  satisfies the renormalization group equation

$$\left(\mu_R \frac{\partial}{\partial \mu_R} + \beta(g_s) \frac{\partial}{\partial g_s}\right) S_{ab \to tX} = -\Gamma_{S \ ab \to tX}^{\dagger} S_{ab \to tX} - S_{ab \to tX} \Gamma_{S \ ab \to tX}$$

Soft anomalous dimension  $\Gamma_{S ab \to tX}$  controls the evolution of the soft function which gives the exponentiation of logarithms of N

Renormalization group evolution  $\rightarrow$  resummation

$$d\tilde{\sigma}_{ab \to tX}^{\text{resum}}(N) = \exp\left[\sum_{i=a,b} E_i(N_i)\right] \exp\left[\sum_{i=a,b} 2\int_{\mu_F}^{\sqrt{s}} \frac{d\mu}{\mu} \gamma_{i/i}(N_i)\right] \exp\left[E'(N)\right]$$
$$\times \operatorname{tr}\left\{H_{ab \to tX}\left(\alpha_s(\sqrt{s})\right) \bar{P} \exp\left[\int_{\sqrt{s}}^{\sqrt{s}/N} \frac{d\mu}{\mu} \Gamma_{S\ ab \to tX}^{\dagger}(\alpha_s(\mu))\right] \tilde{S}_{ab \to tX}\left(\alpha_s\left(\frac{\sqrt{s}}{N}\right)\right) P \exp\left[\int_{\sqrt{s}}^{\sqrt{s}/N} \frac{d\mu}{\mu} \Gamma_{S\ ab \to tX}(\alpha_s(\mu))\right]\right\}$$

The soft anomalous dimensions  $\Gamma_S$  and the hard and soft functions are in general matrices in the space of color exchanges in the hard scattering

Top processes studied - total and differential cross sections

```
Top pair
      aN<sup>3</sup>LO (total; top p_T, y, and double-differential; also A_{FB})
tĒ
      aN^{3}LO + NLO EW (total; top p_T, y)
t\bar{t}
t\bar{t} SMEFT aNNLO (total; top p_T)
Top-pair+X
t\bar{t}\gamma aNNLO + NLO EW (total; top p_T, y)
t\bar{t}W aN<sup>3</sup>LO + NLO EW (total; top p_T, y)
       aN^{3}LO + NLO EW (total; top p_T, y)
t\bar{t}Z
Single top
t- and s-channel aNNLO (total; top p_T) and aN<sup>3</sup>LO (total)
           aN<sup>3</sup>LO (total; p_T, y for top and W)
tW
Single-top+X
tqH aNNLO (total; top p_T, y)
tq\gamma aNNLO (total; top p_T, y)
tqZ aNNLO (total; top y)
Single-top BSM
t\gamma, tZ, tZ' aNNLO (total; top p_T, y)
tq aNNLO (total)
tH^- aNNLO (total; top p_T, y) and aN<sup>3</sup>LO (total)
```

ttW production in collaboration with Chris Foster Phys. Lett. B 854, 138708 (2024) [arXiv:2312.00861]

observation of  $t\bar{t}W$  events at 7, 8, 13 TeV collisions at the LHC measurements are significantly higher than theoretical predictions QCD corrections at NLO are large, ~ 47% at 13.6 TeV

electroweak corrections are smaller but significant

further improvement in theoretical accuracy by the inclusion of higher-order soft-gluon corrections

NLO expansions closely approximate exact NLO results for total cross sections and top-quark  $p_T$  and rapidity distributions

NNLO expansions (aNNLO) are consistent with (partial) NNLO results for total cross sections

 $aN^{3}LO QCD + NLO EW$  is state of the art

# Cross sections for $t\bar{t}W$ production



large *K*-factors

improved agreement with data at aN<sup>3</sup>LO QCD + NLO EW

# $t\bar{t}W$ cross sections

$tar{t}W$ cross sections in $pp$ collisions at the LHC					
$\sigma$ in fb	$7  \mathrm{TeV}$	8 TeV	$13  \mathrm{TeV}$	$13.6 \mathrm{TeV}$	$14  \mathrm{TeV}$
LO QCD	$128^{+39}_{-28}$	$172^{+51}_{-36}$	${}^{445}_{-84}^{+114}$	$481^{+121}_{-90}$	$506^{+126}_{-94}$
lo QCD+EW	$135^{+41}_{-29}$	$\substack{182 + 53 \\ -38}$	$467^{+119}_{-88}$	$505 + 127 \\ -94$	$531^{+132}_{-98}$
NLO QCD	$^{164}_{-17}^{+13}$	$226^{+20}_{-23}$	$646 {+83 \atop -74}$	$708^{+94}_{-82}$	$750^{egin{array}{c}+101\-88\end{array}}$
NLO QCD+EW	$175^{+12}_{-17}$	$^{239}_{-23}^{+19}$	$677^{egin{array}{c} +80 \\ -74 \end{array}}$	$741 {+90 \atop -82}$	$785^{+97}_{-88}$
aNNLO QCD	$179^{+6}_{-10}$	$^{246}_{-15}^{+9}$	$720^{+29}_{-43}$	$791 {+32 \\ -47}$	$837^{+34}_{-50}$
aNNLO QCD $+$ NLO EW	$^{190}_{-10}^{+6}$	$259^{m+9}_{m-15}$	$751^{egin{array}{c}+27\-43\end{array}}$	$824 {+29 \atop -47}$	$872^{+31}_{-50}$
aN <sup>3</sup> LO QCD	$185^{+5}_{-8}$	$253^{+7}_{-12}$	$748^{+24}_{-19}$	$822^{+26}_{-20}$	$870^{+28}_{-21}$
$aN^3LO QCD + NLO EW$	$196^{+5}_{-8}$	$266^{+7}_{-12}$	$779^{+22}_{-19}$	$855 {+23 \atop -20}$	$905^{+25}_{-21}$

At 13.6 TeV

NLO QCD corrections  $\rightarrow$  47%

aNNLO QCD corrections  $\rightarrow$  17%

 $aN^{3}LO \ QCD \ corrections \rightarrow 6\%$ 

electroweak NLO corrections  $\rightarrow$  7%

Total aN<sup>3</sup>LO QCD+NLO EW cross section is 78% bigger than LO QCD

# $t\bar{t}W^+$ and $t\bar{t}W^-$ cross sections

$tar{t}W^+$ and $tar{t}W^-$ cross sections in $pp$ collisions at the LHC						
$\sigma$ in fb	$t\bar{t}W^+$ 13 TeV	$tar{t}W^+$ 13.6 TeV	$t\bar{t}W^{-}$ 13 TeV	$tar{t}W^-$ 13.6 TeV		
LO QCD	$299^{+77}_{-57}$	$322^{+82}_{-60}$	$146 {+37 \atop -28}$	$\substack{159+40\\-30}$		
lo QCD+EW	$313 {+80 \atop -59}$	$337^{+85}_{-63}$	$^{154}_{-29}^{+39}$	$168\substack{+42\\-31}$		
NLO QCD	$431^{+54}_{-49}$	$470^{+61}_{-54}$	$215 {+29 \atop -25}$	$238 + 33 \\ -28$		
NLO QCD+EW	$450^{ig+51}_{-48}$	$490^{+58}_{-53}$	$227 {+28 \atop -25}$	$251 + 32 \\ -28$		
aNNLO QCD	$480^{+19}_{-28}$	$525^{+21}_{-31}$	$240^{ig+10}_{ig-15}$	$266^{\pm 11}_{-16}$		
aNNLO QCD $+$ NLO EW	$499^{+17}_{-28}$	$\substack{545+19\\-31}$	$252^{+10}_{-15}$	$279^{+10}_{-16}$		
aN <sup>3</sup> LO QCD	$498^{+16}_{-12}$	$545\substack{+17\\-13}$	$250 {+8 \atop -7}$	$277^{+9}_{-7}$		
$aN^3LO QCD + NLO EW$	$517 \substack{+14 \\ -12}$	$565\substack{+15\\-13}$	$262 {+8 \atop -7}$	$290 {+8 \atop -7}$		

the  $t\bar{t}W^+$  cross sections are larger than for  $t\bar{t}W^-$ 

but the corrections are slightly bigger for  $t\bar{t}W^-$ 

Comparison with 8 and 13 TeV CMS and ATLAS data NLO and even aNNLO results are not sufficient we need  $aN^3LO$  corrections to describe the data

At 8 TeV, measurements from CMS:  $382_{-102}^{+117}$  fb and from ATLAS:  $369_{-91}^{+100}$  fb Theoretical prediction is

 $aN^{3}LO \ QCD + NLO \ EW: 266^{+7}_{-12-6} \ fb$ 

At 13 TeV, CMS finds  $868 \pm 65$  fb with  $t\bar{t}W^+$   $553 \pm 42$  fb and  $t\bar{t}W^ 343 \pm 36$  fb while ATLAS finds  $880 \pm 80$  fb with  $t\bar{t}W^+$   $583 \pm 58$  fb and  $t\bar{t}W^ 296 \pm 40$  fb

Theoretical prediction is aN<sup>3</sup>LO QCD + NLO EW:  $779_{-19-13}^{+22+12}$  fb with  $t\bar{t}W^+$   $517_{-12-9}^{+14+8}$  fb and  $t\bar{t}W^ 262_{-7-4}^{+8+4}$  fb

## Top-quark $p_T$ and rapidity distributions in $t\bar{t}W$ production at 13 TeV





K-factors decrease at larger top  $p_T$ 



## Top-quark $p_T$ and rapidity distributions in $t\bar{t}W$ at 13.6 TeV







# $t\bar{t}Z$ production

#### in collaboration with Chris Foster, arXiv:2410.01214

observation of  $t\bar{t}Z$  events at 7, 8, 13 TeV collisions at the LHC

important for measuring coupling of the top quark to the Z boson

QCD corrections at NLO are large,  $\sim 32\%$  at 13.6 TeV

electroweak corrections are smaller but significant

further improvement in theoretical accuracy by the inclusion of higher-order soft-gluon corrections

NLO expansions closely approximate exact NLO results for total cross sections and top-quark  $p_T$  and rapidity distributions

 $aN^{3}LO QCD + NLO$  electroweak is state of the art

## Cross sections for $t\bar{t}Z$ production





### good agreement with data at $aN^3LO QCD+NLO EW$

# $t\bar{t}Z$ cross sections

$t \overline{t} Z$ cross sections in $pp$ collisions at the LHC					
$\sigma$ in fb	$7  \mathrm{TeV}$	$8  { m TeV}$	$13  \mathrm{TeV}$	$13.6 \mathrm{TeV}$	$14  \mathrm{TeV}$
LO QCD	$^{119}_{-30}^{+44}$	$177^{+64}_{-44}$	${}^{662}_{-152}^{+218}$	$742^{+242}_{-169}$	$797^{+258}_{-181}$
lo QCD+EW	$^{123}_{-31}^{+46}$	$183 {+67 \atop -45}$	${}^{684}_{-157}^{+226}$	$766 \substack{+250 \\ -175}$	$824 \substack{+267 \\ -187}$
NLO QCD	$^{150}_{-16}^{+9}$	${}^{225}_{-24}^{+15}$	$873^{+69}_{-93}$	$^{982}^{+79}_{-104}$	$^{1057}_{-112}^{+85}$
NLO QCD+EW	$^{155}_{-16}^{+8}$	$^{233}_{-24}^{+13}$	$889^{+61}_{-90}$	$^{999}_{-101}^{+70}$	${}^{1074}_{-107}^{+74}$
aNNLO QCD	$163^{+7}_{-10}$	${}^{245}_{-15}^{+10}$	$^{952}_{-48}^{+29}$	$^{1074}_{-54}^{+33}$	$^{1157}_{-58}^{+35}$
aNNLO QCD $+$ NLO EW	$168\substack{+6\\-10}$	$253^{+9}_{-15}$	$968 {+25 \ -46}$	$1091 {+29 \atop -51}$	$^{1174}_{-54}^{+30}$
aN <sup>3</sup> LO QCD	$168^{+5}_{-8}$	253 + 8 - 12	$^{982}^{+25}_{-28}$	$1108^{+28}_{-32}$	${}^{1194}_{-34}^{+30}$
$aN^3LO QCD + NLO EW$	$173^{+5}_{-8}$	$261^{+7}_{-12}$	$^{998}^{+21}_{-26}$	$1125 {+24 \\ -30}$	$1211 {+25 \\ -30}$

At 13.6 TeV

NLO QCD corrections  $\rightarrow$  32%

aNNLO QCD corrections  $\rightarrow$  12%

 $aN^{3}LO \ QCD \ corrections \rightarrow 5\%$ 

electroweak NLO corrections ightarrow 2%

Total aN<sup>3</sup>LO QCD+NLO EW cross section is 52% bigger than LO QCD

### Comparison with 7, 8, and 13 TeV LHC data

At 7 TeV, measurement from CMS:  $0.28^{+0.14+0.06}_{-0.11-0.03}$  pb

At 8 TeV, measurements from ATLAS:  $176^{+58}_{-52}$  fb and from CMS:  $242^{+65}_{-55}$  fb

At 13 TeV, measurements from

**CMS:**  $0.95 \pm 0.05 \pm 0.06$  **pb** 

#### and from

**ATLAS:**  $0.99 \pm 0.05 \pm 0.08$  **pb** and  $0.86 \pm 0.04 \pm 0.04$  **pb** 

aN $^3$ LO QCD + NLO EW $t\bar{t}Z$ cross section in $pp$ collisions at the LHC						
$\sigma$ in fb	$7  \mathrm{TeV}$	8 TeV	$13  \mathrm{TeV}$	13.6 TeV	$14  \mathrm{TeV}$	
MSHT20 NNLO pdf	$173^{+5}_{-8}^{+5}_{-3}^{+4}_{-3}$	$261^{+7}_{-12}^{+6}_{-4}$	$^{998}_{-26-13}^{+21+20}$	$^{1125}_{-30-14}^{+24+22}$	$^{1211}_{-30-15}^{+25+23}$	
MSHT20 aN <sup>3</sup> LO pdf	$170 {+5+4 \\ -8-4}$	$255 + 7 + 6 \\ -12 - 6$	${}^{974}_{-25-21}^{+20+19}$	$1096 {+23} {+22} \\ {-29} {-23}$	$1182 {+24} {+23} \\ {-29} {-25}$	
NNPDF4.0 aN <sup>3</sup> LO pdf	$165^{+5}_{-7}^{+2}_{-7}$	$248 {+7} {+3} \\ {-11} {-3}$	$962 {+20+6 \atop -25-6}$	$1083 {+23+7 \\ -29-7}$	$1170 {+24} {+7} \\ {-29} {-7}$	

### Top-quark $p_T$ and rapidity distributions in $t\bar{t}Z$ production at 13 TeV



K-factors decrease at larger top  $p_T$ 



# Top-quark $p_T$ and rapidity distributions in $t\bar{t}Z$ at 13.6 TeV





K-factors decrease at larger top  $p_T$  K-factors increase at larger rapidities

### Summary

- higher-order corrections for top-quark production processes
- soft-gluon resummation and aNNLO, aN<sup>3</sup>LO expansions
- $t\bar{t}W$  production
- $t\bar{t}Z$  production
- $aN^{3}LO QCD + NLO EW$  predictions
- results for total cross sections and top-quark  $p_T$  and rapidity distributions
- higher-order corrections further enhance and improve the theoretical predictions
- good agreement with LHC data