# On AdS Scale Separated Vacua



Irene Valenzuela

**CERN** 

IFT UAM-CSIC



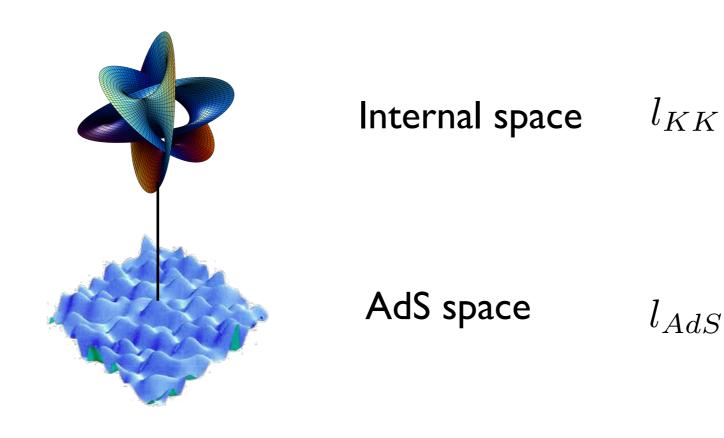
In collaboration with Miguel Montero (to appear soon)

GenHET meeting in String Theory, CERN, April 2024

# **AdS Scale Separation**

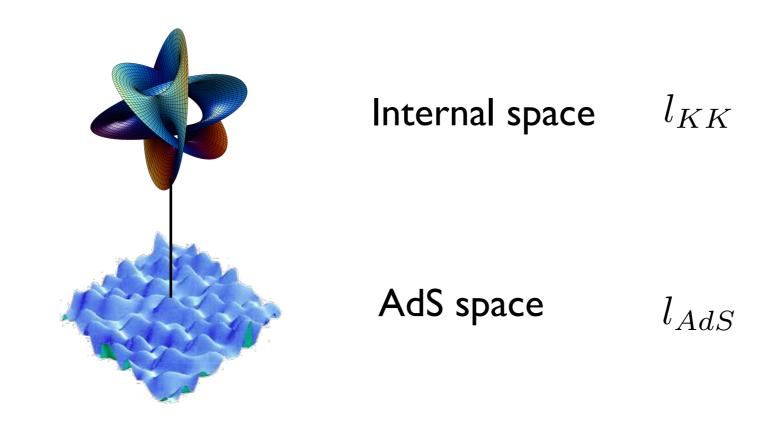
All string theory examples are of the form  $AdS_d \times X_n$ 

 $l_{AdS}$ 



# **AdS Scale Separation**

All string theory examples are of the form  $AdS_d \times X_n$ 



Can we have small extra dimensions in AdS vacua?  $l_{AdS}\gg l_{KK}$ 

Otherwise, it does not describe low dimensional physics

#### From Holography:

All known holographic AdS/CFT examples are not scale-separated Known CFTs at large N are dual to  $AdS_d \times X_n$  where the size of the internal space is of order the AdS scale

#### From Holography:

All known holographic AdS/CFT examples are not scale-separated Known CFTs at large N are dual to  $AdS_d \times X_n$  where the size of the internal space is of order the AdS scale

#### From Swampland:

The limit  $V_0 o 0$  is at infinite distance in the space of metric configurations distance  $\sim \log |V_0|$  [Luest,Palti,Vafa'19]

#### From Holography:

All known holographic AdS/CFT examples are not scale-separated Known CFTs at large N are dual to  $AdS_d \times X_n$  where the size of the internal space is of order the AdS scale

#### From Swampland:

The limit  $V_0 o 0$  is at infinite distance in the space of metric configurations distance  $\sim \log |V_0|$  [Luest,Palti,Vafa'19]

Generalization of the Distance conjecture implies the existence of a tower of states with  $m \sim \exp(-\alpha \operatorname{distance}) \sim V_0^{\alpha}$  as  $V_0 \to 0$ 

#### From Holography:

All known holographic AdS/CFT examples are not scale-separated Known CFTs at large N are dual to  $AdS_d \times X_n$  where the size of the internal space is of order the AdS scale

#### From Swampland:

The limit  $V_0 o 0$  is at infinite distance in the space of metric configurations distance  $\sim \log |V_0|$  [Luest,Palti,Vafa'19]

Generalization of the Distance conjecture implies the existence of a tower of states with  $m \sim \exp(-\alpha \operatorname{distance}) \sim V_0^{\alpha}$  as  $V_0 \to 0$ 

$$m_{\rm KK} \sim V_0^{\alpha}$$
  $\longrightarrow$   $l_{\rm KK} \sim l_{AdS}^{-2\alpha}$  If  $\alpha \geq \frac{1}{2}$  no scale separation

If we are looking for AdS scale-separated vacua, which is the most promising corner of the landscape?

If we are looking for AdS scale-separated vacua, which is the most promising corner of the landscape?

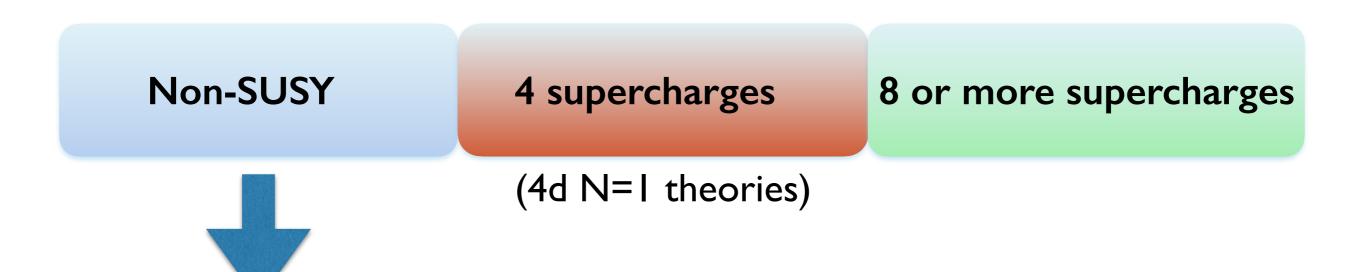
**Non-SUSY** 

4 supercharges

8 or more supercharges

(4d N=1 theories)

If we are looking for AdS scale-separated vacua, which is the most promising corner of the landscape?



Thera are examples of scale-separated vacua in string compactifications but they are unstable (they do not have a CFT dual)

If we are looking for AdS scale-separated vacua, which is the most promising corner of the landscape?

**Non-SUSY** 

4 supercharges

8 or more supercharges



(4d N=1 theories)

Thera are examples of scale-separated vacua in string compactifications but they are unstable (they do not have a CFT dual)

e.g. Scalar potential dominated by Casimir energies [De Luca et al '22]

$$V_{\mathrm{Casimir}} \sim \frac{1}{l^d} \sim m_{KK}^d$$
  $\longrightarrow l_{AdS} \sim l_{KK}^{d/2} \gg l_{KK}$  for  $d>2$ 

At the very least, they have non-perturbative bubble of nothing instabilities

If we are looking for AdS scale-separated vacua, which is the most promising corner of the landscape?



4 supercharges

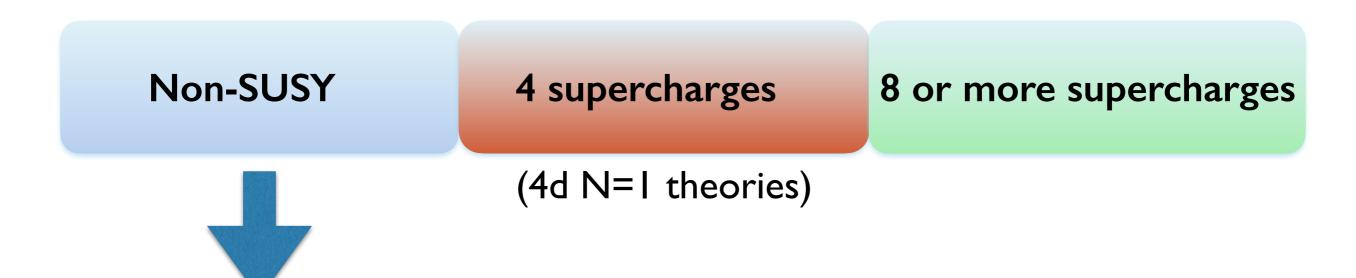
8 or more supercharges



(4d N=1 theories)

More generally, we expect from the Swampland program that all non-SUSY vacua are at best metastable

If we are looking for AdS scale-separated vacua, which is the most promising corner of the landscape?



More generally, we expect from the Swampland program that all non-SUSY vacua are at best metastable

• Bubble of nothing instabilities: always topologically allowed by [Witten'81] Cobordism conjecture (absence of global symmetries) [Garcia-Etxebarria et al'20]

If we are looking for AdS scale-separated vacua, which is the most promising corner of the landscape?



4 supercharges

8 or more supercharges



(4d N=1 theories)

More generally, we expect from the Swampland program that all non-SUSY vacua are at best metastable

- Bubble of nothing instabilities: always topologically allowed by [Witten'81] Cobordism conjecture (absence of global symmetries) [Garcia-Etxebarria et al'20]
- Brane nucleation due to the Weak Gravity Conjecture (if there are fluxes in the internal dimensions) [Ooguri-Vafa'18]

If we are looking for AdS scale-separated vacua, which is the most promising corner of the landscape?

**Non-SUSY** 

4 supercharges

8 or more supercharges

(4d N=1 theories)



There is a continuous R-symmetry

Charged BPS states have  $~m \sim q \, \ell_{AdS}^{-1}$   $~q=0,1,2,\dots$ 

If we are looking for AdS scale-separated vacua, which is the most promising corner of the landscape?

**Non-SUSY** 

4 supercharges

8 or more supercharges

(4d N=1 theories)



There is a continuous R-symmetry

Charged BPS states have 
$$~m \sim q \, \ell_{AdS}^{-1}$$
 
$$~q = 0, 1, 2, \dots$$

[Polchinski, Silverstein '09] [Alday, Perlmutter '19]

In all examples, they correspond to Kaluza-Klein modes

no scale separation

If we are looking for AdS scale-separated vacua, which is the most promising corner of the landscape?

**Non-SUSY** 

4 supercharges

8 or more supercharges

(4d N=1 theories)

Easier, but unstable

???

Difficult or impossible

If we are looking for AdS scale-separated vacua, which is the most promising corner of the landscape?

**Non-SUSY** 

4 supercharges

8 or more supercharges

(4d N=1 theories)

Easier, but unstable

???

Difficult or impossible



All proposed candidates for scale-separated (stable) vacua from the bulk perspective are here: DGKT, KKLT, ...

Not known CFT dual yet, under debate whether they are consistent top-down string constructions,...

If we are looking for AdS scale-separated vacua, which is the most promising corner of the landscape?

**Non-SUSY** 

4 supercharges

8 or more supercharges

(4d N=1 theories)

Easier, but unstable

???

Difficult or impossible



All proposed candidates for scale-separated (stable) vacua from the bulk perspective are here: DGKT, KKLT, ...

Not known CFT dual yet, under debate whether they are consistent top-down string constructions,...

Let's try to shed some light on this debate!

If we are looking for AdS scale-separated vacua, which is the most promising corner of the landscape?

**Non-SUSY** 

4 supercharges

8 or more supercharges

(4d N=1 theories)

Easier, but unstable

???

Difficult or impossible



All proposed candidates for scale-separated (stable) vacua from the bulk perspective are here: DGKT, KKLT, ...

Not known CFT dual yet, under debate whether they are consistent top-down string constructions,...

Let's try to shed some light on this debate!

[De Wolfe, Giryavets, Kachru, Taylor '05]

[Camara, Ibanez, Uranga '05]

4d N=I AdS vacuum arising from compactifying massive Type IIA on a CY3 with O6-planes and fluxes for

$$F_0, F_4, H_3$$
 AdS4xCY3

There is one unconstrained flux that does not appear on the tadpole:

$$\int_{\omega_4} F_4 = N$$

[De Wolfe, Giryavets, Kachru, Taylor '05]

[Camara, Ibanez, Uranga '05]

4d N=I AdS vacuum arising from compactifying massive Type IIA on a CY3 with O6-planes and fluxes for

$$F_0, F_4, H_3$$
 AdS4xCY3

There is one unconstrained flux that does not appear on the tadpole:

$$\int_{\omega_4} F_4 = N$$

By solving the 4d eoms, one finds a family of AdS vacua with

$$V_0 \sim N^{9/2}$$
 $m_{\rm KK}^{-2} \sim L_{\rm KK}^2 \sim N^{7/2}$ 
 $\left(\frac{\ell_{AdS}}{L_{KK}}\right)^2 \sim N$ 

So this solution is **scale-separated** in the large N limit.

The consistency of the solution is not clear because we only solved 4d equations of motion (zero mode of 10d eoms on CY3)

Lot of recent progress, everything seems fine so far, but no conclusive answer.

[Andriot, Apers, Casas, Castellano, Collins, Cribiori, Dall'Agata, De Luca, Emelin, Farakos, Graña, Herraez, Hoter, Ibañez, Junghans, Lust (x2), Marchesano, Marconnet, Montella, Morittu, Ning, Palti, Plauschinn, Prieto, Quirant, Revello, Shiu, Shukla, Tomasiello, Tonioni, Toulikas, Tringas, Tsimpis, Vafa, Van Hemelryck, Van Riet, Walcher, Wiesner, Wrasse, Xu, Yau, Zatti,...]

The consistency of the solution is not clear because we only solved 4d equations of motion (zero mode of 10d eoms on CY3)

Lot of recent progress, everything seems fine so far, but no conclusive answer.

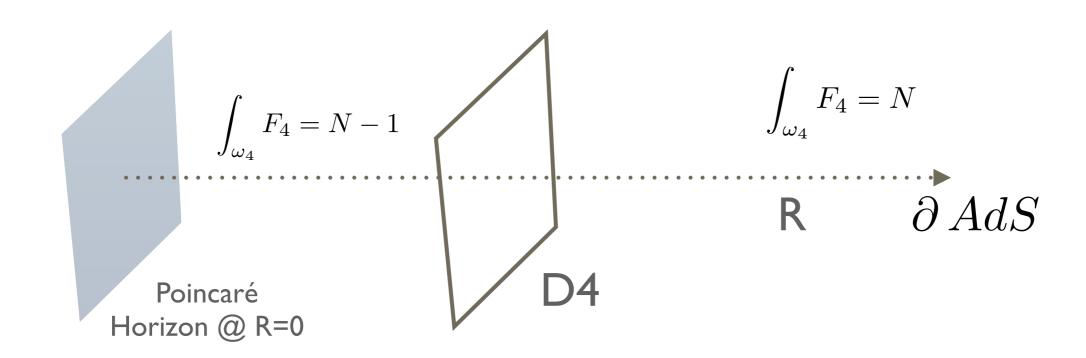
[Andriot, Apers, Casas, Castellano, Collins, Cribiori, Dall'Agata, De Luca, Emelin, Farakos, Graña, Herraez, Hoter, Ibañez, Junghans, Lust (x2), Marchesano, Marconnet, Montella, Morittu, Ning, Palti, Plauschinn, Prieto, Quirant, Revello, Shiu, Shukla, Tomasiello, Tonioni, Toulikas, Tringas, Tsimpis, Vafa, Van Hemelryck, Van Riet, Walcher, Wiesner, Wrasse, Xu, Yau, Zatti,...]

We will assume everything is OK, and study the fate of branes on DGKT vacuum, to perform a non-perturbative consistency check

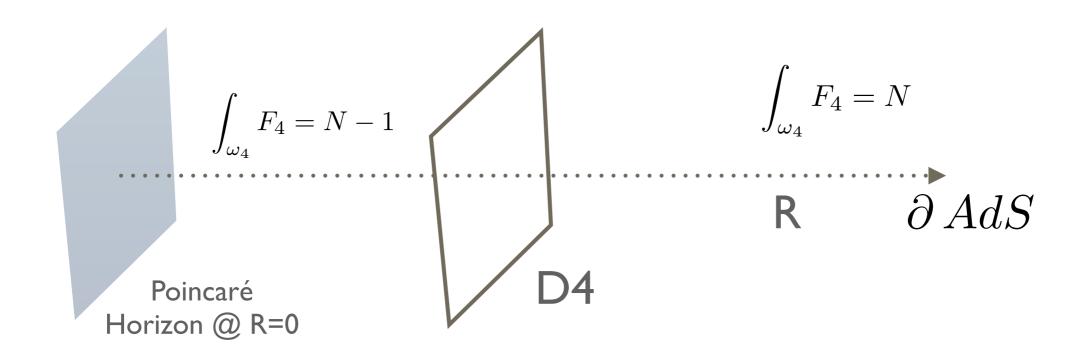
(i.e. whether it is protected against non-perturbative brane instabilities)

[Montero, Valenzuela 'ongoing]

Consider a D4-brane wrapping a holomorphic 2-cycle dual to the large N flux  $\int_{\mathcal{U}_4} F_4 = N$ 



Consider a D4-brane wrapping a holomorphic 2-cycle dual to the large N flux  $\int_{-\infty}^{\infty} F_4 = N$ 



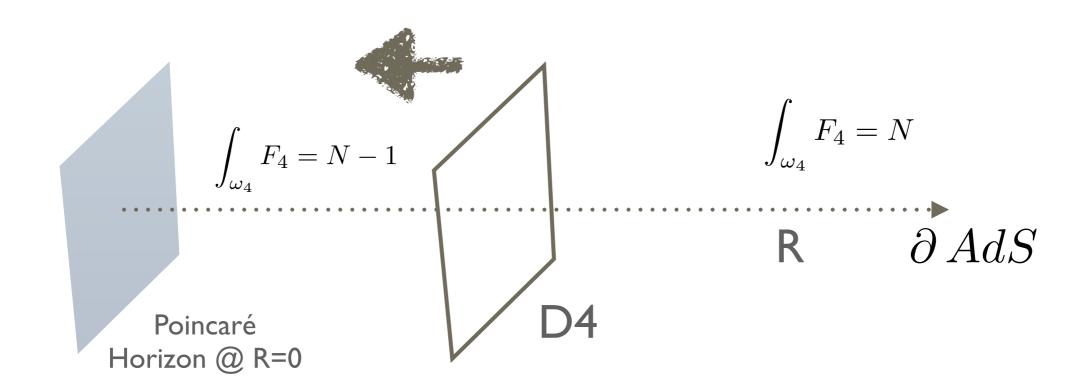
Case I) The D4-brane feels no force, so the position of the brane is a modulus and the domain wall is a static solution

--> Stable vacuum

(This is what happens e.g. in  $AdS_5 \times S^5$ )

Consider a D4-brane wrapping a holomorphic 2-cycle dual to the

large N flux  $\int_{\mathcal{U}_4} F_4 = N$ 

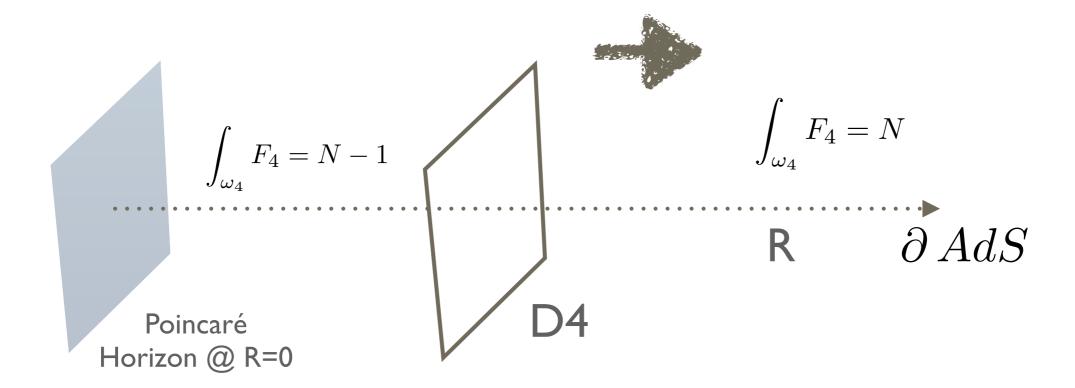


Case 2) The brane feels an attractive force towards the Poincaré Horizon

> Stable vacuum (not known example in holography yet)

Consider a D4-brane wrapping a holomorphic 2-cycle dual to the

large N flux  $\int_{\omega_4} F_4 = N$ 



Case 3) The brane feels a repulsive force from the Poincaré Horizon

It will discharge the flux —> Unstable vacuum

Consider a D4-brane wrapping a holomorphic 2-cycle dual to the large N flux  $\int_{\omega_A} F_4 = N$ 

[Arkani-Hamed et al.'06] [Ooguri, Vafa'18]

If the brane satisfies the Weak Gravity Conjecture:

$$T \leq Q$$

 $(tension) \leq (charge)$ 

Consider a D4-brane wrapping a holomorphic 2-cycle dual to the

large N flux

$$\int_{\omega_4} F_4 = N$$

[Arkani-Hamed et al.'06] [Ooguri, Vafa'18]

If the brane satisfies the Weak Gravity Conjecture:

$$T \leq Q$$

 $(tension) \leq (charge)$ 

$$rac{}{}$$
  $T=Q$ 

Consider a D4-brane wrapping a holomorphic 2-cycle dual to the large N flux  $\int_{\Gamma} F_4 = N$ 

[Arkani-Hamed et al.'06] [Ooguri, Vafa'18]

If the brane satisfies the Weak Gravity Conjecture:

$$T \leq Q$$

 $(tension) \leq (charge)$ 

static domain wall -> stable vacuum (Case I)

Consider a D4-brane wrapping a holomorphic 2-cycle dual to the large N flux  $\int_{\Gamma} F_4 = N$ 

[Arkani-Hamed et al.'06] [Ooguri, Vafa' 18]

If the brane satisfies the Weak Gravity Conjecture: T < Q

$$T \leq Q$$

 $(tension) \leq (charge)$ 

- - static domain wall -> stable vacuum (Case I)
- $\clubsuit$  T < Q the brane feels a repulsive force it describes an non-perturbative instability (Case 3)

Consider a D4-brane wrapping a holomorphic 2-cycle dual to the large N flux  $\int_{\Gamma} F_4 = N$ 

[Arkani-Hamed et al.'06] [Ooguri, Vafa'18]

If the brane satisfies the Weak Gravity Conjecture: T < Q

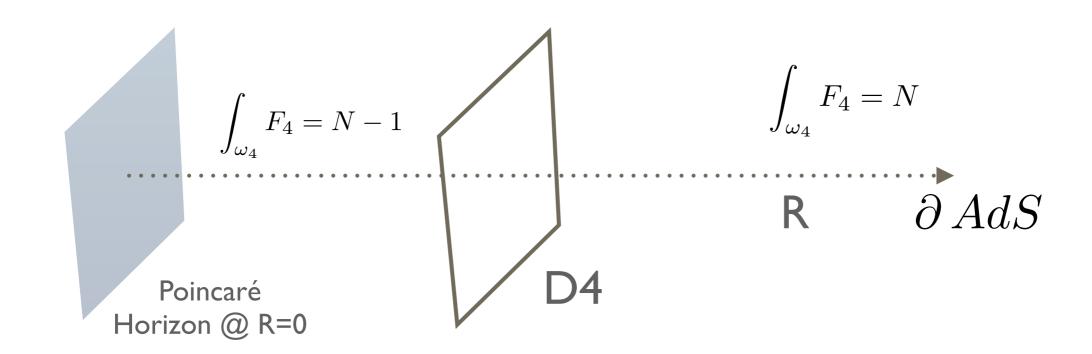
$$T' \leq Q$$
 (tension)  $\leq$  (charge)

- - static domain wall -> stable vacuum (Case I)
- $\clubsuit$  T < Q the brane feels a repulsive force it describes an non-perturbative instability (Case 3)

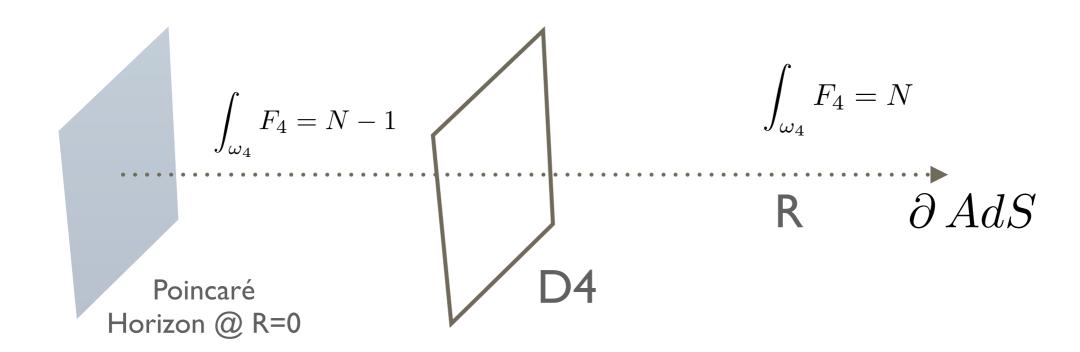
WGC suggests that non-SUSY AdS vacua with internal fluxes are mestastable due to brane nucleation that discharges the flux

(Case 2 for all branes is not allowed by WGC)

Consider a D4-brane wrapping a holomorphic 2-cycle dual to the large N flux  $\int_{\mathcal{U}_4} F_4 = N$ 



Consider a D4-brane wrapping a holomorphic 2-cycle dual to the large N flux  $\int_{\mathcal{U}} F_4 = N$ 



At the classical level, the brane is BPS, so the position of the brane is a modulus

[Aharony, Antebi, Berkooz '08]

What if we go beyond classical level and consider quantum corrections? Do they vanish by supersymmetry?

What if we go beyond classical level and consider quantum corrections? Do they vanish by supersymmetry?

Not necessarily, at low energies the worldvolume theory is  $\,3d\,\mathcal{N}=1\,$ 

What if we go beyond classical level and consider quantum corrections? Do they vanish by supersymmetry?

Not necessarily, at low energies the worldvolume theory is  $3d\mathcal{N}=1$  This is so little SUSY that there are no protected quantities

(there is no holomorphicity) [Aharony, Antebi, Berkooz '08, ... Gaiotto-Komargodski-Wu '18]

What if we go beyond classical level and consider quantum corrections?

Do they vanish by supersymmetry?

Not necessarily, at low energies the worldvolume theory is  $\,3d\,\mathcal{N}=1\,$ 

This is so little SUSY that there are no protected quantities

(there is no holomorphicity) [Aharony, Antebi, Berkooz '08, ... Gaiotto-Komargodski-Wu '18]

The superpotential can receive quantum corrections

In general: 
$$W = W_0 + \sum c_n R^n$$

$$V(R) = |W|^2$$
 the position field R will not be a modulus

What if we go beyond classical level and consider quantum corrections? Do they vanish by supersymmetry?

Not necessarily, at low energies the worldvolume theory is  $\,3d\,{\cal N}=1\,$ This is so little SUSY that there are no protected quantities

(there is no holomorphicity) [Aharony, Antebi, Berkooz '08, ... Gaiotto-Komargodski-Wu '18]

The superpotential can receive quantum corrections

In general: 
$$W = W_0 + \sum c_n R^n$$

 $V(R) = |W|^2$  the position field R will not be a modulus

the brane will feel a non-vanishing force

## 3d N=I theories

The only known way to protect the superpotential from quantum corrections is to have **Parity symmetry** 

[Gaiotto-Komargodski-Wu '18]

## 3d N=I theories

The only known way to protect the superpotential from quantum corrections is to have **Parity symmetry** 

[Gaiotto-Komargodski-Wu '18]

Action just comes from one superspace integral:  $S = \int d^2\theta \, \mathcal{W}$ 

A 3d parity symmetry  $\ \vec{x} 
ightarrow -\vec{x}$  acts as

$$\mathcal{P}(d^2\theta) = -d^2\theta$$
 so  $\mathcal{W} \to -\mathcal{W}$ 

and so a parity-even scalar R cannot generate superpotential!

## 3d N=I theories

The only known way to protect the superpotential from quantum corrections is to have **Parity symmetry** 

[Gaiotto-Komargodski-Wu '18]

Action just comes from one superspace integral:  $S = \int d^2\theta \, \mathcal{W}$ 

A 3d parity symmetry  $\ \vec{x} 
ightarrow -\vec{x} \ \ \ \ {
m acts} \ {
m as}$ 

$$\mathcal{P}(d^2\theta) = -d^2\theta$$
 so  $\mathcal{W} \to -\mathcal{W}$ 

and so a parity-even scalar R cannot generate superpotential!

e.g. 4d N=I AdS from M-theory on  $~AdS_4 imes G_2^{weak}~$  with  $~\int_{G_2^{weak}} G_7 = N$  [Forcella, Zaffaroni '09]

Preserves Pin+ symmetry of M-theory -> it has a moduli space

Could this be the case of DGKT?

Could this be the case of DGKT? No

In DGKT, the fluxes break the parity symmetries [Montero, Valenzuela 'ongoing]

Could this be the case of DGKT? No

In DGKT, the fluxes break the parity symmetries [Montero, Valenzuela 'ongoing]

By explicit computation, we show that a superpotential is generated (SUSY is broken spontaneously on the brane)



the position of D4-brane is not a modulus

Could this be the case of DGKT? No

In DGKT, the fluxes break the parity symmetries [Montero, Valenzuela 'ongoing]

By explicit computation, we show that a superpotential is generated

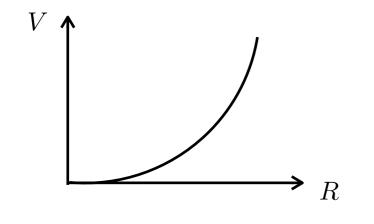
(SUSY is broken spontaneously on the brane)

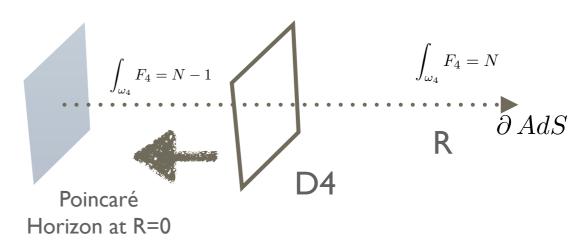


the position of D4-brane is not a modulus

Assuming the EFT of the DGKT solution is correct:  $V = |W|^2 \sim R^3$ 

We find that the brane feels an attractive force ---> stable





Could this be the case of DGKT? No

In DGKT, the fluxes break the parity symmetries [Montero, Valenzuela 'ongoing]

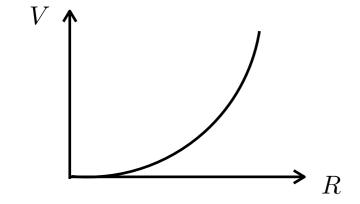
By explicit computation, we show that a superpotential is generated (SUSY is broken spontaneously on the brane)



the position of D4-brane is not a modulus

Assuming the EFT of the DGKT solution is correct:  $V=|W|^2\sim R^3$ 

We find that the brane feels an attractive force -> stable Case 2!



It violates the Weak Gravity Conjecture!!

The vacuum is in tension with the Weak Gravity Conjecture (WGC)

(there is no brane charged under F4 flux which is BPS or self-repulsive) They all have " T>Q "

The vacuum is in tension with the Weak Gravity Conjecture (WGC)

(there is no brane charged under F4 flux which is BPS or self-repulsive) They all have " T>Q "

Does this mean that the DGKT solution is inconsistent?

The vacuum is in tension with the Weak Gravity Conjecture (WGC)

(there is no brane charged under F4 flux which is BPS or self-repulsive) They all have " T>Q "

Does this mean that the DGKT solution is inconsistent?

Not clear yet, since the WGC for codimension one membranes is much less understood than for particles

The vacuum is in tension with the Weak Gravity Conjecture (WGC)

(there is no brane charged under F4 flux which is BPS or self-repulsive) They all have " T>Q "

Does this mean that the DGKT solution is inconsistent?

Not clear yet, since the WGC for codimension one membranes is much less understood than for particles

• DGKT scale-separated vacuum is inconsistent with quantum gravity

The vacuum is in tension with the Weak Gravity Conjecture (WGC)

(there is no brane charged under F4 flux which is BPS or self-repulsive) They all have " T>Q "

Does this mean that the DGKT solution is inconsistent?

Not clear yet, since the WGC for codimension one membranes is much less understood than for particles

- Either
- DGKT scale-separated vacuum is inconsistent with quantum gravity
- WGC for membranes is more subtle and needs to be understood better (no black hole argument)

The vacuum is in tension with the Weak Gravity Conjecture (WGC)

(there is no brane charged under F4 flux which is BPS or self-repulsive) They all have " T>Q "

Does this mean that the DGKT solution is inconsistent?

Not clear yet, since the WGC for codimension one membranes is much less understood than for particles

- Either
- DGKT scale-separated vacuum is inconsistent with quantum gravity
- WGC for membranes is more subtle and needs to be understood better (no black hole argument)

Attractive branes can be a hint to construct the CFT dual if it exists (first example of this type ever seen!)

### Conclusion

If we are looking for AdS scale-separated vacua, which is the most promising corner of the landscape?

**Non-SUSY** 

4 supercharges

8 or more supercharges

(4d N=1 theories)

Easier, but unstable

???

Difficult or impossible

### Conclusion

If we are looking for AdS scale-separated vacua, which is the most promising corner of the landscape?

**Non-SUSY** 

4 supercharges

8 or more supercharges

Easier, but unstable

Without parity symmetry

With parity symmetry

Difficult or impossible

One has to deal with instabilities issues

They are stable, but... scale-separated?

### Conclusion

If we are looking for AdS scale-separated vacua, which is the most promising corner of the landscape?

**Non-SUSY** 

4 supercharges

8 or more supercharges

Easier, but unstable

Without parity symmetry

With parity symmetry

Difficult or impossible

One has to deal with instabilities issues

They are stable, but... scale-separated?

Thank you!

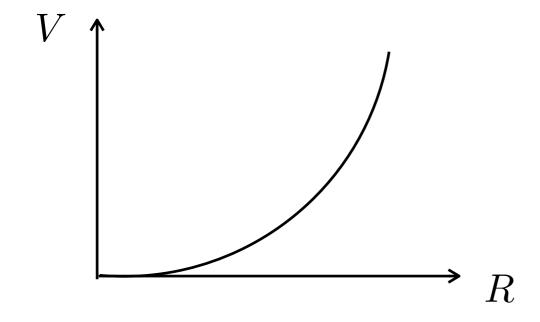
back-up slides

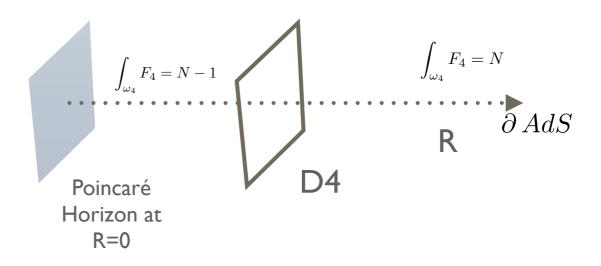
# **Worldvolume D4 theory**

At the orientifold locus, the worldvolume theory is a pure SU(2) gauge theory

$$S \propto \int d^3 \xi \left( \frac{\ell}{2R} |\nabla R|^2 + \frac{1}{2g_{YM}^2} |F|^2 \right)$$

$$V=|W|^2\sim g_{YM}^6=\left(rac{R}{\ell_{AdS}}
ight)^3$$
 in string units





It forces the position of the brane to

$$R \to 0$$

- String theory compactifications: Plethora of quantitative tests!
  - Systematic approach according to the level of supersymmetry
  - Interesting connections to mathematics

[Grimm, Palti, IV'18] [Grimm, Palti, Li'18] [Lee, Lerche, Weigand'18-19]

. . .

- String theory compactifications: Plethora of quantitative tests!
  - Systematic approach according to the level of supersymmetry
  - Interesting connections to mathematics

[Grimm, Palti, IV'18] [Grimm, Palti, Li'18] [Lee, Lerche, Weigand'18-19]

#### **♦** AdS/CFT:

[Heidenreich et al'16]

- WGC proven for AdS3 using modular invariance of the CFT [Montero et al'16]
- WGC from QI theorems and entanglement entropy [Montero'18]
- SDC formulated in terms of a CFT Distance conjecture [Perlmutter et al'20]

- String theory compactifications: Plethora of quantitative tests!
  - Systematic approach according to the level of supersymmetry
  - Interesting connections to mathematics

[Grimm, Palti, IV'18] [Grimm, Palti, Li'18] [Lee, Lerche, Weigand'18-19]

#### AdS/CFT:

[Heidenreich et al'16]

- WGC proven for AdS3 using modular invariance of the CFT [Montero et al'16]
- WGC from QI theorems and entanglement entropy [Montero'18]
- SDC formulated in terms of a CFT Distance conjecture [Perlmutter et al'20]

#### Black hole arguments:

- WGC follows from requiring black holes to decay [Arkani-Hamed et al'06]
- WGC/SDC follows from entropy bounds associated to small BHs [Hamada et al'21]
- Connection between WGC and weak cosmic censorship [Crisford et al'17]

- String theory compactifications: Plethora of quantitative tests!
  - Systematic approach according to the level of supersymmetry
  - Interesting connections to mathematics

[Grimm, Palti, IV'18] [Grimm, Palti, Li'18] [Lee, Lerche, Weigand'18-19]

#### AdS/CFT:

[Heidenreich et al'16]

- WGC proven for AdS3 using modular invariance of the CFT [Montero et al'16]
- WGC from QI theorems and entanglement entropy [Montero'18]
- SDC formulated in terms of a CFT Distance conjecture [Perlmutter et al'20]

#### Black hole arguments:

- WGC follows from requiring black holes to decay [Arkani-Hamed et al'06]
- WGC/SDC follows from entropy bounds associated to small BHs [Hamada et al'21]
- Connection between WGC and weak cosmic censorship [Crisford et al'17]
- Using positivity/unitarity bounds: lead to mild versions of the WGC

[Cheung et al'18][Hamada et al'18]...

#### Weak Gravity conjecture: [Arkani-Hamed et al'06]

Given a gauge theory coupled to gravity, there must exist an electrically charged state with:  $\mathcal{O}(1)$  factor (extremality bound of the black holes)

mass 
$$m \leq \gamma_{\mathrm{BH}} Q M_p$$

electric charge

$$Q = q \, g_{\rm YM}$$
 quantized charge  $\begin{cases} \begin{cases} \begin$ 



#### Weak Gravity conjecture: [Arkani-Hamed et al'06]

Given a gauge theory coupled to gravity, there must exist an electrically charged state with:

mass 
$$m \leq \gamma_{\rm BH}QM_p$$
  $Q = q\,g_{\rm YM}$   $Q = q\,g_{\rm YM}$  electric charge  $Q = q\,g_{\rm YM}$  quantized charge  $Q = q\,g_{\rm YM}$ 

- Independent motivation based on black hole physics
- Plethora of evidence based on string theory, AdS/CFT, scattering amplitudes,...

review: [Harlow et al'22]



#### Weak Gravity conjecture: [Arkani-Hamed et al'06]

Given a gauge theory coupled to gravity, there must exist an electrically charged state with:

mass 
$$m \leq \gamma_{\rm BH}QM_p$$
  $Q = q\,g_{\rm YM}$   $Q = q\,g_{\rm YM}$  electric charge  $Q = q\,g_{\rm YM}$  quantized charge  $Q = q\,g_{\rm YM}$ 

If applied to a p-form gauge field:

There must exist a (p-I)-brane with 
$$T \leq Q$$
 (tension)  $\leq$  (charge)

Exact equality only occurs if it is a supersymmetric BPS state

#### Weak Gravity conjecture: [Arkani-Hamed et al'06]

If applied to a p-form gauge field:

There must exist a (p-I)-brane with 
$$T \leq Q$$
 (tension)  $\leq$  (charge)

Consider an AdS vacuum supported by fluxes in the internal dimensions:

flux: 
$$f_0 \sim \int_{\Sigma_p} F_p$$
 p-form gauge field in the extra dimensions

WGC codimension-one charged brane in AdS

Exact equality only occurs if it is a supersymmetric BPS state

The cobordism group of a quantum gravity theory must be trivial:

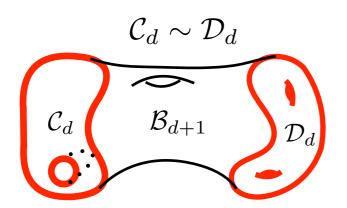
$$\Omega_k^{QG}=0$$
 [McNamara, Vafa' [9]

k: internal dimension

D: total dimension

The cobordism group of a quantum gravity theory must be trivial:

$$\Omega_k^{QG} = 0$$
 [McNamara, Vafa'19]   
  $\Sigma_k$  internal dimension  $\Sigma_k$  total dimension

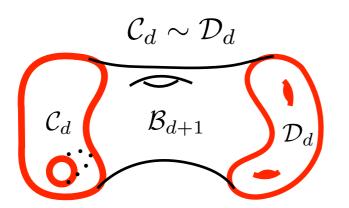


The cobordism group of a quantum gravity theory must be trivial:

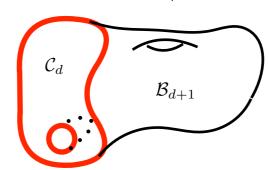
$$\Omega_k^{QG}=0$$
 [McNamara,Vafa'19]

k: internal dimension

D: total dimension



$$C_d = \partial B_{d+1}$$
 (trivial class)

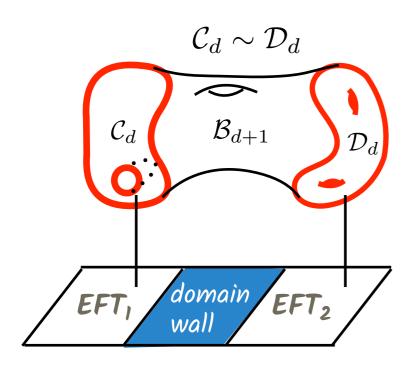


The cobordism group of a quantum gravity theory must be trivial:

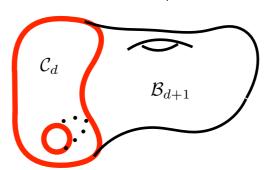
$$\Omega_k^{QG}=0$$
 [McNamara,Vafa'19]

k: internal dimension

D: total dimension



$$C_d = \partial B_{d+1}$$
 (trivial class)

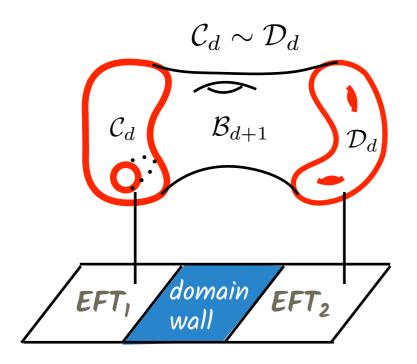


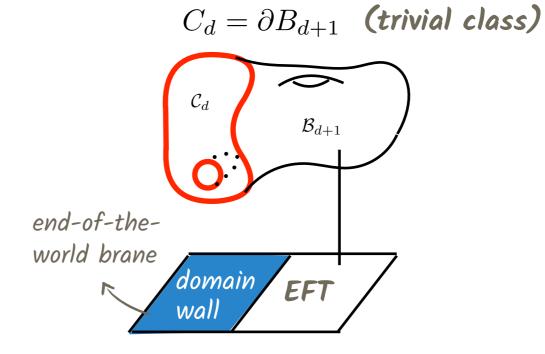
The cobordism group of a quantum gravity theory must be trivial:

$$\Omega_k^{QG}=0$$
 [McNamara,Vafa'19]

k: internal dimension

D: total dimension





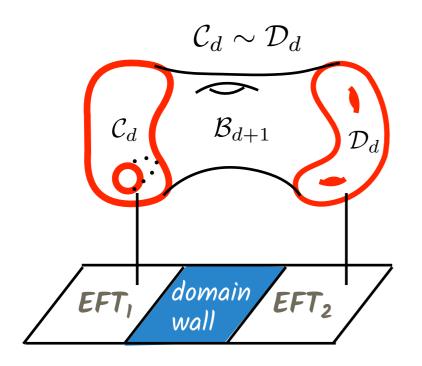
The cobordism group of a quantum gravity theory must be trivial:

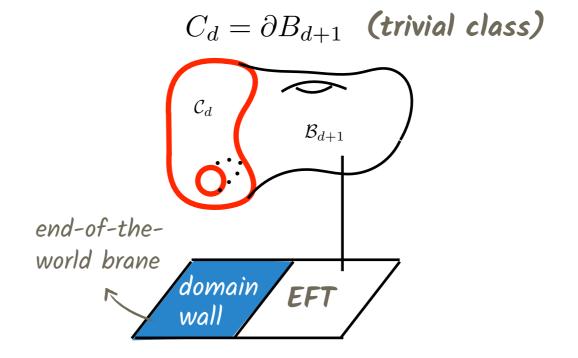
$$\Omega_k^{QG}=0$$
 [McNamara,Vafa'19]

k: internal dimension

D: total dimension

to avoid a (D-k-I)-form global symmetry with charges  $\,[M]\in\Omega_k^{QG}\,$ 





It implies all theories of same dimension are connected by finite energy domain walls, and predicts the existence of new defects in string theory!

Swampland conjecture: Any non-supersymmetric vacuum is at best metastable

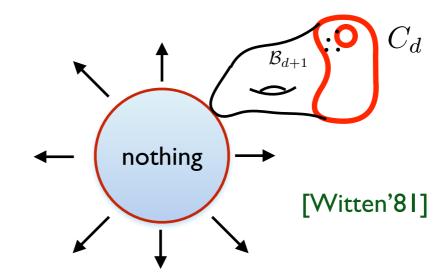
Is there any universal instability that arises when breaking supersymmetry?

Swampland conjecture: Any non-supersymmetric vacuum is at best metastable

Is there any universal instability that arises when breaking supersymmetry?

Candidate: Bubbles of nothing

Non-perturbative instability from the vacuum to nothing:

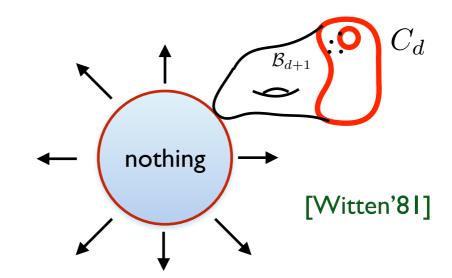


Swampland conjecture: Any non-supersymmetric vacuum is at best metastable

Is there any universal instability that arises when breaking supersymmetry?

Candidate: Bubbles of nothing

Non-perturbative instability from the vacuum to nothing:



It is topologically allowed if:

The compact space  $C_d$  shrinks to a point

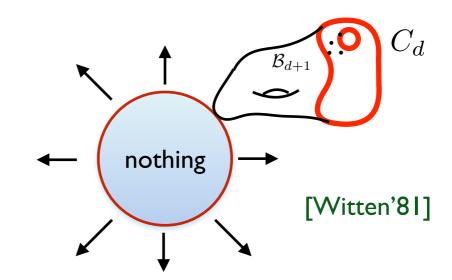
i.e. 
$$C_d = \partial B_{d+1}$$

Swampland conjecture: Any non-supersymmetric vacuum is at best metastable

Is there any universal instability that arises when breaking supersymmetry?

Candidate: Bubbles of nothing

Non-perturbative instability from the vacuum to nothing:



It is topologically allowed if:

The compact space  $C_d$  shrinks to a point

 $C_d$  belongs to trivial class of  $\Omega_d$ (in cobordism)

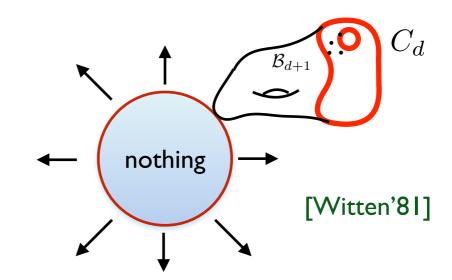
i.e. 
$$C_d = \partial B_{d+1}$$

Swampland conjecture: Any non-supersymmetric vacuum is at best metastable

Is there any universal instability that arises when breaking supersymmetry?

Candidate: Bubbles of nothing

Non-perturbative instability from the vacuum to nothing:



It is topologically allowed if:

The compact space  $C_d$  shrinks to a point

 $C_d$  belongs to trivial class of  $\Omega_d$ (in cobordism)



i.e.  $C_d = \partial B_{d+1}$ 

no topological global charges (cobordism classes)

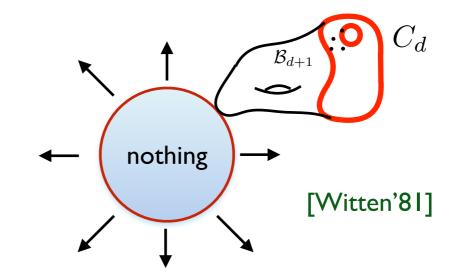
Swampland conjecture: Any non-supersymmetric vacuum is at best metastable

Is there any universal instability that arises when breaking supersymmetry?

Candidate: Bubbles of nothing

Non-perturbative instability from the vacuum to nothing:

[Garcia-Etxebarria, Montero, Sousa, IV'20]



Cobordism conjecture: (no global symmetries)

$$\Omega_k^{QG} = 0$$



No topological obstruction to have bubbles of nothing

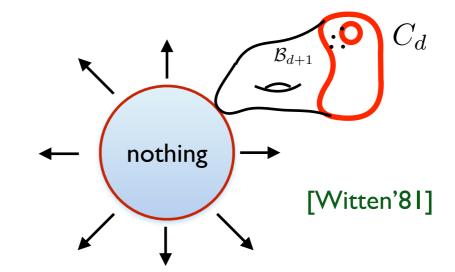
Swampland conjecture: Any non-supersymmetric vacuum is at best metastable

Is there any universal instability that arises when breaking supersymmetry?

Candidate: Bubbles of nothing

Non-perturbative instability from the vacuum to nothing:

[Garcia-Etxebarria, Montero, Sousa, IV'20]



Cobordism conjecture: (no global symmetries)

$$\Omega_k^{QG} = 0$$



No topological obstruction to have bubbles of nothing

They will expand and describe a vacuum instability if a certain energy condition (DEC) is violated semiclassically

(which can happen when supersymmetry is broken)