

1

# *CP* **violation in charmed baryons at LHCb**

5 March 2024

**Artur Ukleja** on behalf of the LHCb experiment AGH University of Krakow



### **Menu**



# • **The examples of the LHCb measurements**

- $\Leftrightarrow CP$  violation in  $\mathbb{E}_c^+ \to pK^-\pi^+$
- $\Diamond$  CP asymmetry difference between  $\Lambda_c^+ \to pK^-K^+$  and  $p\pi^-\pi^+$
- $\triangle$  Amplitude analysis of  $\Lambda_c^+ \to pK^-\pi^+$
- $\triangle$  Polarization of  $\Lambda_c^+ \to pK^-\pi^+$
- $\triangle$  Polarimeter vector field in  $\Lambda_c^+ \to pK^-\pi^+$

# • **Summary**

# **Why are we interested in charm physics?**

- In the Standard Model (SM), the known value of *CP* violation (CPV) is too small to explain the observed size of matter domination over antimatter in the universe
- At LHCb, we very precisely test known CPV in the SM  $\rightarrow$  finding disagreement will be indirect indication of new phenomena existence **Search for New Physics in the Flavour Sector**
- The new particles can appear in the loops



• Charm sector is very promising since the background from the SM is very small, expected CPV is only  $\leq 10^{-4} - 10^{-3}$  (much smaller than we measure in beauty meson decays) what is the scale of  $\sim$  How much different are  $\sim$ ng since the background from the SM is very sm



#### Eur. Phys. J. C80 (2020) 986  $\frac{1}{\sqrt{1-\frac{1$ EUT. PHYS. J. C80 (Z020) 986  $\sigma$





and 2012 (right) data sets after all implemented cuts: stripping, over all implemented cuts: stripping, owe cat

# The binned S<sub>CF</sub> method OCal



• In each bin a significance of a difference between  $\mathcal{Z}_c^+$  and  $\mathcal{Z}_c^-$  is calculated

$$
S_{CP}^i \equiv \frac{N_+^i - \alpha N_-^i}{\sqrt{\alpha (N_+^i + N_-^i)}} \qquad \alpha = \frac{N^+}{N^-}
$$

- To cancel global asymmetries (production  $\sigma$ , and Band proto are normalized asymmetry etc.) the Dalitz plots are normalized
- $P_{\nu}$ , atotic ticol fluotuotional • If no CPV (only statistical fluctuations) then  $S_{CP}$  is Gaussian distribution ( $\mu$ =0,  $\sigma$ =1)
- The  $\chi^2 = \Sigma S^i{}_{CP}{}^2$  test is calculated to obtain  $10^2$ p-value for the null hypothesis (no CPV) to test if  $E_c^+$  and  $E_c^-$  distributions<br>
or atoticiaally compatible and 1-10% on magnitude differences. are statistically compatible

p-value ≪ 1 in case of CPV













- Uniform and adaptive binning schemes with different bin numbers are tested
- The  $S_{CP}$  distributions agree with Gaussian
- The measured p-values are greater than 32%
- Results are consistent with no observation of *CP* asymmetry

123



- More difficult to use but can be more sensitive for small statistics
- To compare " $+$ " and " $-$ " a test statistic T is defined, which is based on the counting particles with the same sign to each event for a given number of the nearest neighbour events

$$
T = \frac{1}{n_k(n_+ + n_-)} \sum_{i=1}^{n_+ + n_-} \sum_{k=1}^{n_k} I(i, k)
$$

 $I(i,k)$  = 1 if i<sup>th</sup> event and its  $k<sup>th</sup>$  nearest neighbour have the same charge  $(+ + , -)$  $I(i,k) = 0$  if pair has opposite charge  $(+ -)$ 



- T is the mean fraction of like pairs in the pooled sample of the two datasets
- The expected distribution can be calculated using mean  $\mu$ <sub>T</sub> and variance  $\sigma$ <sub>T</sub>

$$
\mu_T = \frac{n_+(n_+ - 1) + n_-(n_- - 1)}{n(n-1)}
$$

$$
\lim_{n, n_k, D \to \infty} \sigma_T^2 = \frac{1}{n n_k} \left( \frac{n_+ n_-}{n^2} + 4 \frac{n_+^2 n_-^2}{n^4} \right)
$$

# The first searches for local  $\textbf{\textit{CPV}}$  in  $\mathbb{E}_c^+ \to pK^-\pi^+$



Eur. Phys. J. C80 (2020) 986

- Data collected in Run 1,  $\sqrt{s}$  = 7 TeV and 8 TeV, L ~ 3 fb<sup>-1</sup>
- The k-nearest neighbour method is used in regions of the Dalitz plot



- Results are consistent with *CP* symmetry but in one region of the Dalitz plot the local effect corresponds to 2.7σ
- To be continued with Run 2 data

# CP asymmetry difference between  $\Lambda_c^+ \to pK^-K^+$  and  $\,p\pi^-\pi^+$ JHEP03 (2018) 182

• <sup>Ô</sup>*<sup>s</sup>* = 7 and 8 TeV, 2011 & 2012, *<sup>L</sup>int* = 3 *fb*≠1,

• The  $\Lambda_c^+$  candidates are reconstructed from  $\Lambda_b^0 \to \Lambda_c^+ \mu^- X$  decays JHEP03 (2018) 182  $A^{\dagger}$  candidates are reconstructed from  $A^0 \rightarrow A^{\dagger} \mu^- X$  decay are reconstructed from  $\Lambda_b^0 \to \Lambda_c^+ \mu^- X$  decays

$$
A_{CP} = \frac{\Gamma(f) - \Gamma(\bar{f})}{\Gamma(f) + \Gamma(\bar{f})}, \quad A_{raw} = \frac{N(f\mu^{-}) - N(\bar{f}\mu^{+})}{N(f\mu^{-}) + N(\bar{f}\mu^{+})}
$$

$$
A_{raw}(f) = A_{CP}(f) + A_{P}^{\Lambda_{b}^{0}}(f\mu) + A_{D}^{\mu}(\mu) + A_{D}^{f}(f)
$$

• Kinematics of the  $pK^-K^+$  and  $p\pi^-\pi^+$  are and equal and data are weighted. a awet a *CP* <sup>=</sup> *<sup>A</sup>CP*(*pK* <sup>≠</sup>*<sup>K</sup>* <sup>+</sup>) <sup>≠</sup> *<sup>A</sup>wgt CP* (*pfi*≠*fi*<sup>+</sup>) ¥ *<sup>A</sup>raw* (*pK* <sup>≠</sup>*<sup>K</sup>* <sup>+</sup>) <sup>≠</sup> *<sup>A</sup>wgt*  $\begin{equation} \varDelta A^{wgt}_{CP} = A_{CP}(\rho K^-K^+) - A^{wgt}_{CP}(\rho\pi^-\pi^+) \approx A_{\textit{raw}}(\rho K^-K^+) - A_{\textit{raw}}^{wgt}(\rho\pi^-\pi^+) \end{equation}$ • Kinematics of the  $pK^-K^+$  and  $p\pi^-\pi^+$  are and equal and data and  $\Delta A_{CP}^{wgt} = A_{CP}(pK^-K^+) - A_{CP}^{wgt}(p\pi^-\pi^+) \approx A_{raw}(pK^-K^+) - A_{CP}^{wgt}$ 



**LHCh** 



# JHEP03 (2018) 182 JHEP03 (2018) 182

$$
A_{raw}(pK^-K^+) = (3.72 \pm 0.78)\%
$$
  

$$
A_{raw}^{wgt}(p\pi^-\pi^+) = (3.42 \pm 0.47)\%
$$

$$
\Delta A_{\rm CP}^{\rm wgt}=(0.30\pm 0.91\pm 0.61)\%
$$

- Central value is measured to be **be consistent with zero. in** *»*<sup>+</sup> consistent with zero **CONSISTENT VALUE ZETO** 
	- **•** The first measurement of the CP violation parameter in  $\varLambda_c^+$  decays







- $\Lambda_c^+$  from semileptonic beauty hadron decays  $\Lambda_b^0 \to \Lambda_c^+ l^- \bar{\nu}$
- $\Lambda_c^+ \to pK^-\pi^+$  has a complex resonant structure with multiple overlapping states in the  $K^-\pi^+$ ,  $pK^-$  and  $p\pi^+$ systems
- A full amplitude analysis determines the composition of the decay amplitude
- The knowledge of the resonant structure is useful in searches for *CP* symmetry violation, which is still unobserved in baryon decays





- LHCb  $2016: -1.7$  fb<sup>-1</sup>
- The signal region chosen within 15 MeV of the known  $\Lambda_c^+$  mass, containing about 95% of the signal candidates
- $~1$  ~400 000 candidates
- The fraction of background in the signal region is 1.69%





The Dalitz plot displays a rich structure with resonant contributions from all possible pairs of final state particles:  $A \to pK^-$ ,  $K^* \to K^-\pi^+$ ,  $\Delta^{++} \to p\pi^+$ 



## are needed for the simultaneous measurement of  $12$  resonances

 $\mathbf{F}^{\text{max}}_{\text{max}}$ 





Figure 3: Distributions for selected candidates together with amplitude fit projections in the *lab* resonances have not been previously been reported projections. polarization system in Eq. (20), *|*1*/*2*, m*⇤<sup>+</sup> *<sup>c</sup>* i, to the decay plane system in Eq. (21), *l* cyloff (*A*(2000)), writere clear A significant contribution from a resonant in *m(pK)*~2GeV region ( $\Lambda$ (2000)), where clear method for matching proton spin states among di↵erent decay chains of Ref. [38]. The resonances have not been previously been reported



The  $\Lambda_c^+\to pK^-\pi^+$  amplitude model can be employed to measure the polarization of the  $\Lambda_c^+$  baryon



A large averaged polarization is measured 65% with absolute uncertainty of order 1% with dominating transverse component 60% and a smaller longitudinal component -25%

#### Including polarization component in the laboratory Figure 4: Distributions for selected candidates together with amplitude fit projections in the *B*˜ including polarization component in the laboratory

#### Measuring polarimeter vector field in  $\varLambda_c^+\to pK^-\pi^+$  decays gives the relation between the rotation of the transition amplitude (spin-1*/*2 representation **of the SU(2) group) and the supplication of the rotation of the three-dimensions** (spin-1  $\alpha$  representation of the three-dimensional vector  $\pi$  representation of the three-dimensional vector  $\pi$  representation of the



The averaged squared matrix element:

LHCb-PAPER-2022-044, JHEP 2023(2023) 228

$$
|\mathcal{M}(\phi,\theta,\chi,\kappa)|^2 = I_0(\kappa) \left(1 + \sum_{i,j} P_i R_{ij}(\phi,\theta,\chi)\alpha_j(\kappa)\right)
$$



:hree-dimensional rotation matrix implementing the Euler transformation to a physical vector that describe the orientation of the decay products in space  $\overline{\mathcal{L}}$ a physical vector that describe the orientation of the decay products in space  $R_{ij}(\phi,\theta,\chi)$  – three-dimensional rotation matrix implementing the Euler transformation to

#### The left years that is the momenta of the momenta of the momenta of the production production production produc plane contains the momenta of the momenta of the decay reaction. The central blue plane is defined by the decay reaction. The central blue plane is defined by the decay reaction. The central blue plane is defined by the ce  $\alpha_j(\kappa)$  polarimeter vector

where the components of the *averaged aligned polarimeter vector* ~

Z

 $\frac{1}{2}$ 



LHCb-PAPER-2022-044, JHEP 2023(2023) 228

 $\vec{\alpha}$  – polarimeter vector with respect to the decay plane it gives a model-agnostic representation for polarization dependence of the decay rate



 $R_{\mathsf y}$  is alignment ratio: who is z axis (the sum of the pion and kaon momenta defines the positive z direction) Puzzles of Heavy Baryons April 21st 2023 26 / 30  $\pm$  2023 26  $\pm$  30  $\pm$  3

### LHCb-PAPER-2022-044, JHEP 2023(2023) 228

- The length of the polarimeter vector (shown by the colour) changes from point to point
- It is greater than 0.5 for most of the kinematic domain
- The structures are driven by resonances in different subsystems and their interferences
- The  $\bar{\alpha}$  field gives information needed to determine the  $\Lambda_c^+$ polarization in future analyses





- Charm physics measurements are providing a wealth of interesting results over the last few years
- So far, *CP* violation in the charm sector is confirmed in mesons only
- In all other charm decays, results are consistent with *CP* symmetry
	- statistical uncertainties dominate
	- increasing data statistics will allow to test the SM in more details
- We are running Run 3
- The goal is to reach  $\sim$  23 fb<sup>-1</sup> (Run 3) and  $\sim$  50 fb<sup>-1</sup> (Run 4)  $(Run 1+2: -9 fb^{-1})$





## Back up





- 1. In the mixing (only neutral particles)  $P^0 \rightarrow$  anti- $P^0 \neq$  anti- $P^0 \rightarrow P^0$
- 2. In the amplitudes of direct decays (neutral and charge particles)  $P^{\pm} \rightarrow f \neq$  anti- $P^{\pm} \rightarrow$  anti-f
- 3. In the interference between direct decays and decays via mixing (only neutral particles)







*P0* = *K0, B0, B0 s, D0 , D0 s P*<sup>±</sup> = *K*<sup>±</sup>*, B*<sup>±</sup>*, B*<sup>±</sup> *s, D*<sup>±</sup> , *D*<sup>±</sup> *s,* <sup>L</sup><sup>±</sup> *b,* L<sup>±</sup> *c,* X<sup>±</sup> *<sup>c</sup>* … Federico's & Jolanta's talks This talk



## The single-arm forward spectrometer (a new concept for HEP experiments)

10 < <sup>q</sup> < 300 mrad (2<h<5) (*b*¯*b*) = 75*.*<sup>3</sup> *<sup>±</sup>* <sup>5</sup>*.*<sup>4</sup> *<sup>±</sup>* <sup>13</sup>*.*<sup>0</sup> *µb* Phys.Lett.B694 (2010) 209-216

 $\sigma(c\bar{c}) = 1419 \pm 12 \pm 116 \mu b \sim 20 \times \sigma(b\bar{b})$ 

Nucl.Phys.B871 (2013) 1

Int. J. Mod. Phys. A 30 (2015) 1530022





Singly Cabibbo−suppressed decays (SCS):

- the only place for *CP* violation in the Standard Model
- both: tree and penguin diagrams



To observe *CP* violation, at least two amplitudes must interfere with different weak phases AND DIFFERENT STRONG PHASES

 $\lambda = 0.22$ 



LHCb 2016-2018 dataset at 13TeV and 5.4 fb<sup>-1</sup> (Prompt production) Two-dimensional unbinned extended maximum likelihood fits are performed to the mass and  $log \chi^2$  distributions



The results confirm the charmed-hadron lifetime hierarchy, improve the precision of the previous  $\Omega_c$  lifetime by a factor of two