Gifts from infinite-dimensional current algebras

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based on [2406.02662 w/ D.M. Hofman] [2310.18391 w/ J.R. Fiiss] and thoughts in progress

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Outline

Motivation

Free fields and current algebras

Gift 1: State-operator correspondence ^[2406.02662 w/ D.M. Hofman]

Gift 2: Topological entanglement entropy ^[2310.18391 w/ J.R. Fliss]

Gift ideas (outlook)

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Gift ideas (outlook)

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fixes spectrum exactly underpins (worldsheet of) string theory further connections to quantum gravity

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Internal symmetries also enhance: Kac-Moody



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O(700) papers [IYKYK 2014–2024]

Recently vast generalisation of the notion of symmetry

One kind of generalisation: higher-form symmetries [Gaiotto, Kapustin, Seiberg, Willet 2014]

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Continuous symmetry: $\partial_{\mu}J^{\mu} = 0 \iff d \star J_{[1]} = 0$

 \implies codimension-one topological operator $U(\Sigma_{d-1}) := \exp\left(i \int_{\Sigma_{d-1}} \star J_{[1]}\right)$



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One kind of generalisation: higher-form symmetries [Gaiotto, Kapustin, Seiberg, Willet 2014]

Zero-form continuous symmetry: $\partial_{\mu} J^{\mu} = 0 \iff d \star J_{[1]} = 0$ \implies codimension-one topological operator $U(\Sigma_{d-1}) := \exp\left(i \int_{\Sigma_{d-1}} \star J_{[1]}\right)$

 $p\text{-form continuous symmetry: } \partial_{\mu} J^{\mu\nu_{1}\cdots\nu_{p}} = 0 \iff d \star J_{[p+1]} = 0$ $\implies \text{codimension-}(p+1) \text{ topological operator } U(\Sigma_{d-p-1}) := \exp\left(i \int_{\Sigma_{d-p-1}} \star J_{[p+1]}\right)$



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Other generalisations: higher-group, non-invertible, subsystem symmetries, and more.

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Very simple dynamics:
$$\partial^{\mu}J_{\mu\nu\rho\dots} = 0$$
 and $\partial^{\mu}\epsilon_{\mu\nu\dots}^{\ \ \alpha\beta\dots}J_{\alpha\beta\dots} = 0.$

The star of the show is a free p-form field.

Very simple dynamics: $d \star J_{[p+1]} = 0$ and $d J_{[p+1]} = 0$.











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Commutators: $\left[\mathcal{Q}(\Lambda), \mathcal{Q}(\Lambda')\right] = i k \int_{\Sigma} \Lambda \wedge d\Lambda'$. $k = "level" \sim$ gauge coupling
 $ln d = 2, p = 0 \rightsquigarrow free compact scalar J(z)\alpha(z) and \bar{J}(\bar{z})\bar{\alpha}(\bar{z}) conserved$
 $\left[\mathcal{Q}(\alpha), \mathcal{Q}(\alpha')\right] = i k \int \alpha d\alpha' \implies \left[Q_n, Q_m\right] = k n \delta_{n+m,m} \implies it's a Kac-Moody!$

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In higher d $[,] = \int \cdots$ is again a spectrum-generating, infinite-dimensional current algebra.

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$$[,] = \int \cdots$$
 into a mode algebra $\implies [\mathcal{A}_n, \mathcal{A}_m^{\dagger}] = E_n \delta_{nm}$

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Plus zero modes! $\begin{cases} \mathbf{r} \in \mathbb{Z}^{b_{d-p-1}(\Sigma)} \\ \mathbf{s} \in \mathbb{Z}^{b_{p+1}(\Sigma)} \end{cases} = \begin{cases} \text{electric} \\ \text{magnetic} \end{cases}$ fluxes = higher-form charges $\int \sqrt{\sum_{l=1}^{l} \sqrt{\sum_{l=1}^{l} \frac{1}{l} \sum_{l=1}^{l} \frac{1}{l} \sum_{l=$

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Hamiltonian: $H_{\Sigma} = k\mathbf{r}^{2} + k^{-1}\mathbf{s}^{2} + \sum_{n} \mathcal{A}_{n}^{\dagger}\mathcal{A}_{n}$

States:

▷ Primary states: $|r, s\rangle$. Fixed fluxes, annihilated by all A_n . Energy = $kr^2 + k^{-1}s^2 =: \Delta_{r,s}$

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$$\begin{array}{c} \triangleright \quad \text{Descendants:} \quad \prod_{n} \left(\mathcal{A}_{n}^{\dagger} \right)^{N_{n}} | \textbf{r}, \textbf{s} \rangle. \text{ Energy} = \Delta_{r,s} + \sum_{n} N_{n} E_{n} \\ \text{Non-trivial check:} \quad \textbf{ch}(q) = \sum_{r,s} \textbf{ch}_{r,s}(q) = \sum_{r,s} \operatorname{tr} q^{H_{\Sigma}} = \underbrace{\Theta_{\Sigma}(q;k)}{\eta_{\Sigma}(q)^{2}} = \mathcal{Z}\left(\mathbb{S}_{\beta}^{1} \times \Sigma \right) \\ \text{"spectral eta function"} \quad \eta_{\Sigma}(q) = \prod_{n} \left(1 - q^{E_{n}} \right)^{-1/2} \underbrace{ \begin{array}{c} \\ \end{array}} \begin{array}{c} \text{Very reminiscent of 2d CFTs!} \end{array}$$

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 \rightsquigarrow precise one-to-one correspondence between states on \mathcal{H}_{Σ} and nonlocal operators

sets the stage for understanding higher-dimensional CFTs on compact manifolds

such a correspondence is not possible for generic CFTs [Belin, de Boer, Kruthoff 2018].





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For illustration, take d = 4 (p = 1) and quantise on $\Sigma = \mathbb{S}^2 \times \mathbb{S}^1$ unique holonomy Recover the spectrum of states by path integrals on $\mathbb{B}^3 \times \mathbb{S}^1$ with insertions of line operators

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State-operator correspondence For illustration, take d = 4 (p = 1) and quantise on $\Sigma = \mathbb{S}^2 \times \mathbb{S}^1$ • unique holonomy unique flux Recover the spectrum of states by path integrals on $\mathbb{B}^3 \times \mathbb{S}^1$ with insertions of line operators $|\mathcal{L}\rangle := \int \mathrm{D}A \,\mathrm{e}^{-S[A]} \,\mathcal{L}(\{\mathbf{0}\} \times \mathbb{S}^1) = \left(\left(\mathbf{0} \right)^{1/2} \,\mathrm{d}A \,\mathrm{d}A$ takes care of **Bogoliubov** transformation $\sim \prod_{n} \exp\left(\mathcal{A}_{n}^{2} + \left(\mathcal{A}_{n}^{\dagger}\right)^{2}\right)$ A surprise: $|1\rangle \neq |vacuum\rangle = |0,0\rangle$ Why? radial evolution mixes ladder operators $\rightsquigarrow \mathcal{A}_{n} |1\rangle \neq 0$ but $(\#_{n}\mathcal{A}_{n} + \#_{n}\mathcal{A}_{n}^{\dagger}) |1\rangle = 0 \implies$ squeezing transformation: $|1\rangle = \mathfrak{S} |0,0\rangle$

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Primary states

$$|r,s\rangle \iff \frac{\text{Wilson-'t Hooft lines}}{\text{dressed with squeezing operator}} = \mathfrak{S}^{\dagger} \times \exp\left(ir \int_{\mathbb{S}^1} A + is \int_{\mathbb{S}^1} \widetilde{A}\right)$$

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Descendants

Representation theory at work ---> sprinkle oscillators

$$\mathcal{A}_{\mathsf{n}}^{\dagger} | r, s \rangle \longleftrightarrow \mathfrak{S}^{\dagger} \times \left(\#_{\mathsf{n}}^{*} \mathcal{A}_{\mathsf{n}}^{\dagger} + \#_{\mathsf{n}}^{*} \mathcal{A}_{\mathsf{n}} \right) \times \exp \left(\operatorname{i} r \int_{\mathbb{S}^{1}} A + \operatorname{i} s \int_{\mathbb{S}^{1}} \widetilde{A} \right)$$

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Descendants

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$$\prod_{n} \left(\mathcal{A}_{n}^{\dagger}\right)^{N_{n}} | r, s \rangle \nleftrightarrow \mathfrak{S}^{\dagger} \times \prod_{n} \left(\#_{n}^{*} \mathcal{A}_{n}^{\dagger} + \#_{n}^{*} \mathcal{A}_{n} \right)^{N_{n}} \times \exp\left(i r \int_{\mathbb{S}^{1}} A + i s \int_{\mathbb{S}^{1}} \widetilde{A} \right)$$

That's it. That's the entire spectrum.

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Gift ideas (outlook)





Topological order = patterns of long range entanglement
In (2+1)d:
$$S_{EE} = \frac{|\partial R|}{\varepsilon} - \gamma$$
 [Kitaev, Preskill; Levin, Wen 2006
 $\log (\text{total quantum dimension})$

IR effective field theories ~-> 3d TQFTs



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bulk/edge correspondence

anomaly/symmetry inflow (cf. SPTs, SymTFTs)

entanglement spectrum = edge spectrum [Li, Haldane 2008; Chandran et. al 2011]

edge spectrum organised by Kac-Moody algebra [Elitzur, Moore, Seiberg 1989]



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Here: Consider higher-dimensional topological order and study its entanglement



$$\mathsf{IR}\,\mathsf{TQFT}:S = \frac{k}{2\pi}\int B_{\lfloor d-p-1 \rfloor} \wedge \mathrm{d}A_{\lfloor p \rfloor}$$

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Presence of boundary: gauge transformations that survive ---> global charges

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Edge modes is free p-forms + chirality condition

Edge spectrum controlled again by infinite-dimensional current algebra

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Edge modes is free *p*-forms + chirality condition higher-d analogue of chiral boson

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Interested in $S_{EE} = -\operatorname{Tr}_{\mathcal{H}_R}(\rho_R \log \rho_R) \equiv \lim_{n \to 1} \frac{1}{1-n} \log \operatorname{Tr}_{\mathcal{H}_R} \rho_R^n$ Important issues:

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Entanglement

R

 ∂R

O_R time

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proportional to heat kernel coefficients

Entanglement

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proportional to heat kernel coefficients subleading universal topological term -

Entanglement

R

 ∂R

Outline

Motivation

Free fields and current algebras

Gift 1: State-operator correspondence [2406.02662 w/ D.M. Hofman]

Gift 2: Topological entanglement entropy ^[2310.18391 w/ J.R. Fliss]

Gift ideas (outlook)

Gift ideas

Non-invertible current algebras

Gauge charge conjugation $J \sim -J \implies$ symmetry broken but resurrected as non-invertible [{others}, Antinucci, Galati, Rizi 2022; Aguilera-Damia, Argurio, Chaudhuri 2023]

Current algebra also: $D(\Lambda) \otimes D(\Lambda') = e^{ik \int \Lambda \wedge d\Lambda'} D(\Lambda + \Lambda') \oplus e^{-ik \int \Lambda \wedge d\Lambda'} D(\Lambda - \Lambda')$

---> Should still fix the spectrum TIP w/ Aguilera-Damia, Argurio, Chaudhuri

Current algebras in gravity

Linearised gravity enjoys "biform" symmetries with charges R – (traces) and $\star R \star$ –(traces) [[Hinterbichler et al 2022, (Hull et al 2024)²]]

Also leads to a current algebra (depends on more parameters 2 KY tensors, 2 one-forms)

---> Implications: linearised spectrum, asymptotic symmetries...? TIP w/ Mathys

Non-linear p-form electrodynamics

