Gifts from infinite-dimensional current algebras

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based on<br>[2406.02662 w/ D.M. Hofman]<br>[2310.18391 w/ J.R. Fliss]<br>| and thought: and thoughts in progress

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**Outline** 

**Motivation** 

Free fields and current algebras

Gift 1: State-operator correspondence [2406.02662 w/ D.M. Hofman]

Gift 2: Topological entanglement entropy [2310.18391 w/ J.R. Fliss]

Gift ideas (outlook)

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Gift ideas (outlook)

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IR constraints *Landau paradigm, 't Hooft anomaly matching*

Selection rules *allowed transition, decays, particles in the spectrum...*

Topological protection *topological insulators, topological superconductors, fracton phases...*

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O(700) papers<br>[IYKYK 2014–2024]

Recently vast generalisation of the notion of symmetry

One kind of generalisation: *higher-form symmetries* [Gaiotto, Kapustin, Seiberg, Willet 2014]

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*Σd−*<sup>1</sup>

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Continuous symmetry:  $\partial_{\mu}J^{\mu} = 0 \quad \Longleftrightarrow \quad d \star J_{[1]} = 0$ 

 $\implies$  codimension-one topological operator  $U(\Sigma_{d-1}) := \exp\Biggl(i\Biggr)$ Z *⋆J*[1]  $\lambda$ 



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*p*-form continuous symmetry:  $\partial_{\mu}J^{\mu\nu_1\cdots\nu_p} = 0 \iff d \star J_{[p+1]} = 0$  $\implies$  codimension-(*p* + 1) topological operator  $U(\Sigma_{d-p-1}) := \exp\left(i\frac{p^2}{2}\right)$  $\sqrt{2}$ *Σd−p−*<sup>1</sup>  $\star J_{[p+1]}$  $\lambda$ 



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Act on *p*-dimensional extended operators by linking



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Other generalisations: higher-group, non-invertible, subsystem symmetries, and more.

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Gift ideas (outlook)

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Very simple dynamics: 
$$
\partial^{\mu} J_{\mu\nu\rho\cdots} = 0
$$
 and  $\partial^{\mu} \epsilon_{\mu\nu\cdots}{}^{\alpha\beta\cdots} J_{\alpha\beta\cdots} = 0$ .

The star of the show is a free *p*-form field.

Very simple dynamics:  $d * J_{[p+1]} = 0$  and  $d J_{[p+1]} = 0$ .











generated by: 
$$
\exp\left(i \alpha \int_{M_{d-p-1}} \star J_{[p+1]}\right)
$$
  
acts on: Wilson =  $\exp\left(i \int_{\gamma_p} A_{[p]}\right)$ 

 $(d-p-2)$ -form U(1) symmetry (magnetic) generated by: exp i*α*  $\sqrt{2}$ *Mp*+<sup>1</sup> *⋆J*e[*d−p−*1]  $\lambda$ 

acts on:

't Hooft = 
$$
\exp\left(i \int_{\gamma_{d-p-2}} \widetilde{A}_{[d-p-2]}\right)
$$

There's much more to it: *infinitely many morezero-form symmetries*

*F* or any  $Λ$ <sub>[*p*]</sub>,  $Λ$ <sub>[*d−p−*2] satisfying:</sub>  $dΛ$ <sub>[*p*</sub>] + *∗*  $dΛ$ <sub>[*d*−*p*−2] = 0</sub> There's much more to it: *infinitely many morezero-form symmetries*

 $\textsf{Conserve}$  currents:  $\mathcal{J}_{\Lambda,\widetilde{\Lambda}} = \star \big( J_{[p+1]} \wedge \widetilde{\Lambda}_{[d-p-2]} + \widetilde{J}_{[p+1]} \wedge \Lambda_{[p]} \big)$ 



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To get the spectrum, turn [ , ] = 
$$
\int \cdots
$$
 into a mode algebra  $\Longrightarrow$  [ $A_n$ ,  $A_m^{\dagger}$ ] =  $E_n \delta_{nm}$ 

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 $\int_{\text{Laplacian on } \Sigma} \oint$ 

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Plus zero modes!  $\begin{cases} \mathbf{r} \in \mathbb{Z}^{b_{d-p-1}(\Sigma)} \\ \mathbf{s} \in \mathbb{Z}^{b_{p+1}(\Sigma)} \end{cases} = \begin{cases} \text{electric} \\ \text{magnetic} \end{cases}$  fluxes = higher-form charges  
Hamiltonian:  $H_{\Sigma} = k\mathbf{r}^2 + k^{-1} \mathbf{s}^2 + \sum_{n} A_n^{\dagger} A_n$ 

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#### States:

*▷ Primary states*: *|r* , *s〉*. Fixed fluxes, annihilated by all *A*<sup>n</sup> . Energy = *kr* <sup>2</sup> + *k −*1 *s* <sup>2</sup> =*· · ∆<sup>r</sup>* ,*<sup>s</sup>*

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$$
\triangleright \quad Descendants: \mathcal{A}_n^{\dagger} | \mathbf{r}, \mathbf{s} \rangle. \text{ Energy} = \Delta_{\mathbf{r}, \mathbf{s}} + E_n
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- *▷ Descendants*:  $\mathcal{A}_m^{\dagger} \mathcal{A}_n^{\dagger} | r, s$ *}*. Energy =  $\Delta_{r,s} + E_n + E_m$
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- Siegel-like Theta function  $^{\prime\prime}$ spectral eta function"  $\eta_{\,\Sigma}(q)=\prod_n \bigl(1-q^{E_n}\bigr)^{-1/2} \searrow$ very reminiscent of 2d CFTs! *⊳ Descendants:*  $\prod_n (\mathcal{A}_n^{\dagger})^{N_n} |r, s\rangle$ . Energy =  $\Delta_{r,s} + \sum_n N_n E_n$ Non-trivial check:  $\mathbf{ch}(q) = \sum_{\alpha}$ *r* ,*s*  $ch_{r,s}(q) = \sum$ *r* ,*s*  $\text{tr}\,q^{H_{\Sigma}} = \frac{\Theta_{\Sigma}(q;k)}{R_{\Sigma}(q)^2}$  $\frac{\partial \Sigma(q;k)}{\partial \Sigma(q)^2} = \mathcal{Z}\left(\mathbb{S}^1_{\beta} \times \Sigma\right)$

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Gift 2: Topological entanglement entropy [2310.18391 w/ J.R. Fliss]

#### Gift ideas (outlook)

 $\int$ <sup>*d*\**J*=0</sup> Claim:<sup>[Hofman, Iqbal 2018] A unitary CFT in  $d = 2p + 2$  with a p-form U(1) symmetry is realised by</sup> *free p-form fields.*

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very specific function [Costa, Hansen 2015]  $\mathsf{Sketch\ of\ proof}\colon \langle J(x)J(y)\rangle=\mathop{\{\!\mathrm{f}}\nolimits}_x(x-y)\implies \langle \mathrm{d} J(x)\;\mathrm{d} J(y)\rangle=0\implies \mathrm{d} J=0.$ 

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*▷ d* = 2 *standard result: free field realisation*

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- *▷ d* = 2 *standard result: free field realisation*
- *▷ d* = 4 *free photon realisation (photonisation)*

<sup>2→</sup> precise one-to-one correspondence between *states on*  $\mathcal{H}_{\Sigma}$  and *nonlocal operators* 

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precise one-to-one correspondence between *states on H<sup>Σ</sup>* and *nonlocal operators*

sets the stage for understanding higher-dimensional CFTs on compact manifolds

such a correspondence is not possible for generic CFTs [Belin, de Boer, Kruthoff 2018].





*<sup>|</sup>L〉 · ·*= Z D*A*e *<sup>−</sup>S*[*A*] *<sup>L</sup> {***0***} ×* S 1 = *L*

State-operator correspondence For illustration, take  $d = 4$  ( $p = 1$ ) and quantise on  $\Sigma = \mathbb{S}^2 \times \mathbb{S}^1$  and  $\omega$  unique flux Recover the spectrum of states by path integrals on  $\mathbb B^3 \times \mathbb S^1$  with insertions of line operators

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*A surprise*: *|*1*〉 ̸*= *|*vacuum*〉* = *|*0, 0*〉* Why? radial evolution mixes ladder operators  $\rightarrow$   $\mathcal{A}_n |\mathbb{1}\rangle \neq 0$  but  $(\#_n\mathcal{A}_n + \#_n\mathcal{A}_n^{\dagger}) |\mathbb{1}\rangle = 0$ 

State-operator correspondence unique holonomy unique flux takes care of Bogoliubov transformation  $\sim \prod_{n} \exp\left(A_n^2 + \left(A_n^{\dagger}\right)^2\right)$ For illustration, take  $d=4$  ( $p=1$ ) and quantise on  $\Sigma=\mathbb{S}^2\times\mathbb{S}^1$ Recover the spectrum of states by path integrals on  $\mathbb B^3 \times \mathbb S^1$  with insertions of line operators *<sup>|</sup>L〉 · ·*=  $\sqrt{2}$  $DAe^{-S[A]} \mathcal{L}(\{0\} \times \mathbb{S}^1) =$ *L A surprise*: *|*1*〉 ̸*= *|*vacuum*〉* = *|*0, 0*〉* Why? radial evolution mixes ladder operators  $\rightarrow$  *A*<sub>n</sub>  $|1\rangle \neq 0$  but  $(\#_{n}A_{n} + \#_{n}A_{n}^{\dagger})|1\rangle = 0 \implies$  squeezing transformation:  $|1\rangle = \mathfrak{S}|0,0\rangle$ 

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The rest is straightforward

#### The rest is straightforward

Primary states

$$
|r,s\rangle \leftrightarrow \text{Wilson-'t Hooft lines} = \mathfrak{S}^\dagger \times \exp\left(\text{if }\int_{\mathbb{S}^1} A + \text{is }\int_{\mathbb{S}^1} \widetilde{A}\right)
$$

#### The rest is straightforward

Primary states

$$
|r,s\rangle \leftrightarrow \text{Wilson--'t Hooft lines}
$$
  
 
$$
|r,s\rangle \leftrightarrow \text{dressed with squeezing operator} = \mathfrak{S}^{\dagger} \times \exp\left(i \, r \int_{\mathbb{S}^1} A + i \, s \int_{\mathbb{S}^1} \widetilde{A}\right)
$$

#### **Descendants**

Representation theory at work  $\rightsquigarrow$  sprinkle oscillators

$$
\mathcal{A}_{n}^{\dagger} |r,s\rangle \leftrightarrow \mathfrak{S}^{\dagger} \times (\#_{n}^{*} \mathcal{A}_{n}^{\dagger} + \#_{n}^{*} \mathcal{A}_{n}) \times \exp\left( i r \int_{\mathbb{S}^{1}} A + i s \int_{\mathbb{S}^{1}} \widetilde{A} \right)
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#### **Descendants**

Representation theory at work  $\rightsquigarrow$  sprinkle oscillators

$$
\prod_{n} \left(\mathcal{A}_{n}^{\dagger}\right)^{N_{n}} |r,s\rangle \leftrightarrow \mathfrak{S}^{\dagger} \times \prod_{n} \left(\#_{n}^{*}\mathcal{A}_{n}^{\dagger} + \#_{n}^{*}\mathcal{A}_{n}\right)^{N_{n}} \times \exp\left(i r \int_{\mathbb{S}^{1}} A + i s \int_{\mathbb{S}^{1}} \widetilde{A}\right)
$$

That's it. That's the entire spectrum.

**Outline** 

**Motivation** 

Gift 1: State-operator correspondence [2406.02662 w/ D.M. Hofman]

Gift 2: Topological entanglement entropy [2310.18391 w/ J.R. Fliss]

Gift ideas (outlook)

Topological order = patterns of long range entanglement

\nIn (2+1)d: 
$$
\mathcal{S}_{EE} = \frac{|\partial R|}{\varepsilon} - \gamma
$$
 [Kitaev, Preskill; Levin, Wen 2006]

\n $\text{log (total quantum dimension)}$ 



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IR effective field theories  $\rightsquigarrow$  3d TQFTs



 $\frac{1}{\sqrt{1-\frac{1}{n}}}\log\left(\frac{1}{\sqrt{1-\frac{1}{n}}}\right)$ Topological order  $=$  patterns of long range entanglement  $\ln (2+1)$ d:  $S_{EE} =$ *|R|*  $\frac{24}{\varepsilon} - \gamma$  [Kitaev, Preskill; Levin, Wen 2006]

IR effective field theories  $\rightsquigarrow$  3d TQFTs

bulk/edge correspondence

anomaly/symmetry inflow (*cf. SPTs, SymTFTs*)

entanglement spectrum = edge spectrum [Li, Haldane 2008; Chandran et. al 2011]

edge spectrum organised by Kac–Moody algebra [Elitzur, Moore, Seiberg 1989]



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Here: Consider higher-dimensional topological order and study its entanglement



$$
\text{IR TQFT:} S = \frac{k}{2\pi} \int B_{[d-p-1]} \wedge \text{d}A_{[p]}
$$

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Presence of boundary: gauge transformations that survive  $\rightsquigarrow$  global charges

Here: 
$$
Q(\alpha) = \frac{k}{2\pi} \int_{\partial \Sigma} \alpha \wedge B, \qquad \widetilde{Q}(\beta) = \frac{k}{2\pi} \int_{\partial \Sigma} \beta \wedge A
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Satisfy  $[Q(\alpha), \widetilde{Q}(\beta)] = i k \int_{\partial \Sigma} \alpha \wedge d\beta$  looks familiar?

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Edge modes is free *p*-forms + chirality condition

Edge spectrum controlled again by infinite-dimensional current algebra

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higher-d analogue of chiral boson Edge modes is free p-forms + chirality condition  $\leftarrow$ 

Edge spectrum controlled again by infinite-dimensional current algebra



**Entanglement** *ΣR*  $\partial R$ time (replica trick) Interested in  $S_{\text{EE}} = -\text{Tr}_{\mathcal{H}_R}(\rho_R \log \rho_R) \equiv \lim_{n \to 1}$ 1  $\frac{1}{1-n} \log \text{Tr}_{\mathcal{H}_R} \rho_R^n$ *R Important issues*:

 $H_{\Sigma} \neq H_R \otimes H_{\text{rest}} \rightarrow \rho_R$  needs care  $\rightarrow$  *solution*: extended Hilbert space [Buidovich, Polikarpov 2008; Donnelly 2011; Donnelly, Wall 2014]

UV issues by the introduction of  $\partial R \sim$  solution: regulate trace with quadratic Hamiltonian

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Altogether: Tr 
$$
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$$



**Entanglement** 

*R*

 $\partial R$ 

*Σ*

 $\frac{1}{1-n} \log \text{Tr}_{\mathcal{H}_R} \rho_R^n$ *R*

(replica trick)

1

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extract entanglement by modular properties of *Θ* and hard work (for *η*)

time

*Σ*

Entanglement

*R*

 $\partial R$ 

time *R*

(replica trick)

*Σ*

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extract entanglement by modular properties of *Θ* and hard work (for *η*)

$$
\mathcal{S}_{\text{EE}} = \sum_{n=1}^{\lfloor \frac{d-1}{2} \rfloor} C_{\underbrace{n}_{\text{max}}}^{(p)} \left(\frac{\ell}{\varepsilon}\right)^{d-2n} + \frac{1}{2} K^{(p)} \delta_{d,\text{even}} \log \left(\frac{\ell}{\varepsilon}\right) - \frac{1}{2} \left(b_p(\partial R) + b_{d-p-1}(\partial R)\right) \log k
$$

proportional to heat kernel coefficients

**Entanglement** 

*R*

 $\partial R$
time *R*

(replica trick)

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$$

proportional to heat

roportional to heat **subleading universal topological term**<br>kernel coefficients

*Σ*

**Entanglement** 

*R*

 $\partial R$ 

**Outline** 

**Motivation** 

Gift 1: State-operator correspondence [2406.02662 w/ D.M. Hofman]

Gift 2: Topological entanglement entropy [2310.18391 w/ J.R. Fliss]

## Gift ideas (outlook)

## Gift ideas

Non-invertible current algebras

Gauge charge conjugation *J ∼ −J* =*⇒* symmetry broken but resurrected as non-invertible [*{*others*}*, Antinucci, Galati, Rizi 2022; Aguilera-Damia,Argurio,Chaudhuri 2023]

Current algebra also:  $D(\Lambda) \otimes D(\Lambda') = e^{i k \int \Lambda \wedge d\Lambda'}\, D(\Lambda+\Lambda') \oplus e^{-i k \int \Lambda \wedge d\Lambda'}\, D(\Lambda-\Lambda')$ 

Should still fix the spectrum *TIP w/ Aguilera-Damia, Argurio, Chaudhuri*

## Current algebras in gravity

Linearised gravity enjoys "biform" symmetries with charges *R −* (traces) and *⋆R⋆ −*(traces) [[Hinterbichler et al 2022, (Hull et al 2024) $^2$ ]]

Also leads to a current algebra (depends on more parameters 2 KY tensors, 2 one-forms)

Implications: linearised spectrum, asymptotic symmetries...? *TIP w/ Mathys*

Non-linear *p*-form electrodynamics

